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Role of microstructure and surface defects on the dissolution kinetics of CeO_2 , a UO_2 fuel analogue.

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Abstract

The release of radionuclides from spent fuel in a geological disposal facility is controlled by the surface mediated dissolution of UO₂ in groundwater. In this study we investigate the influence of reactive surface sites on the dissolution of a synthesised CeO₂ analogue for UO₂ fuel. Dissolution was performed on: CeO₂ annealed at high temperature, which eliminated intrinsic surface defects (point defects and dislocations); CeO_{2-x} annealed in inert and reducing atmospheres to induce oxygen vacancy defects; and on crushed CeO_2 particles of different size fractions. BET surface area measurements were used as an indicator of reactive surface site concentration. Cerium stoichiometry, determined using X-ray Photoelectron Spectroscopy (XPS) and supported by X-ray Diffraction (XRD) analysis, was used to determine oxygen vacancy concentration. Upon dissolution in nitric acid medium at 90°C, a quantifiable relationship was established between the concentration of high energy surface sites and CeO_2 dissolution rate; the greater the proportion of intrinsic defects and oxygen vacancies, the higher the dissolution rate. Dissolution of oxygen vacancy-containing CeO_{2-x} gave rise to rates that were an order of magnitude greater than for CeO₂ with fewer oxygen vacancies. While enhanced solubility of Ce³⁺ influenced the dissolution, it was shown that replacement of vacancy sites by oxygen significantly affected the dissolution mechanism due to changes in the lattice volume and strain upon dissolution and concurrent grain boundary decohesion. These results highlight the significant influence of defect sites and grain boundaries on the dissolution kinetics of UO₂ fuel analogues and reduce uncertainty in the long-term performance of spent fuel in geological disposal.

Key words: Nuclear fuel; dissolution; defects; oxygen vacancies; grain boundaries

1. INTRODUCTION

Spent nuclear fuel is a heterogeneous ceramic material composed primarily of UO₂, with a minor component of actinides and fission products. The internationally supported strategy for the safe disposal of this nuclear waste material is within a geological disposal facility¹, where the release of radionuclides to the environment will be dominated by the interaction of the UO₂ matrix with groundwater. To assess the long-term performance of the geological disposal site towards the containment of radionuclides, a safety case is being prepared, which assesses the mechanism and kinetics of fuel dissolution; accordingly, a large number of laboratory experiments have been focussed on determining the dissolution kinetics of UO₂ under a wide variety of temperature and redox conditions^{2,3}. However, dissolution rates determined from natural uraninite (UO₂) ore weathering, indicate that rates measured in the laboratory are orders of magnitude greater than in nature^{4,5}. This phenomenon has been widely considered, and is attributed to a number of factors including artefacts of specimen preparation, solution saturation state and changes in surface area with time⁶. Clearly, it is important to the development of a robust post-closure safety case for spent fuel disposal, that such uncertainty is reduced.

In the current study, we consider the role of "energetically reactive sites", which contain fewer, and weaker, chemical bonds than those on defect-free surfaces or in the bulk, in the determination of laboratory dissolution rates of spent UO₂ fuel. These sites, which may include natural crystal defects (point defects and dislocations), grain boundaries or sample preparation-induced artefacts (e.g. polishing scratches or sharp edges on crushed particles), comprise atoms that have greater reactivity than those in the defect-free surfaces or in the bulk, thus may dissolve more rapidly^{7,8}. For example, the dissolution of UO₂ and a CeO₂ analogue for UO₂ have been found to be influenced by sample crushing^{9,10} and polishing¹¹; dissolution rates were overestimated by 3 orders of magnitude due to the reactive surface sites induced through these sample preparation techniques¹¹. Natural defects, especially randomly distributed oxygen vacancy defects, dominate the lattice structure of spent UO₂ fuel, due to the incorporation of trivalent rare earth fission products in the UO₂ lattice¹². It has been hypothesised that the presence of such energetically reactive defect sites may influence the dissolution rate¹³⁻¹⁸. However, the extent of such an effect on the overall

dissolution rate is difficult to observe when other factors, such as the enhanced solubility of U(VI) over U(IV), strongly influence the dissolution behaviour. Recent work by the current authors has shown that these seemingly minor features may significantly influence the *initial* dissolution rate of UO₂ fuel analogues^{11,19}, therefore should be taken into consideration alongside other important factors (notably oxidation, solubility and radiolysis) when understanding dissolution mechanisms of spent fuel.

In an attempt to preclude the dominant effects of U(IV) oxidation, the high solubility of U(VI) and radiolysis on the dissolution kinetics, we here report the role of energetically reactive features on the dissolution of a UO₂ analogue, CeO₂, that was developed to closely resemble the microstructure of UO₂ fuel²⁰. Both UO₂ and CeO₂ crystallise in the same fluorite-type structure (Fm-3m) and can be engineered to have similar grain sizes and crystallographically randomly orientated grains. We describe an evaluation of how energetically reactive sites, including intrinsic surface defects and oxygen vacancies, influence dissolution kinetics. Through quantification of these effects, we attempt to improve our understanding of the uncertainties associated with dissolution rates of spent fuel in geological disposal conditions.

2. EXPERIMENTAL METHODS

2.1. CeO₂ preparation. Cerium dioxide monoliths were prepared according to Stennett et al.²⁰, at a sintering temperature of 1700 °C. Monoliths were crushed using a percussion mortar and ball mill, sieved to four different size fractions; $25 - 50 \mu m$, $75 - 150 \mu m$, $300 - 600 \mu m$ and 2 - 4 mm, and washed, all according to ASTM C1285 (the Product Consistency Test)²¹. Prior to use, the particles were inspected by Scanning Electron Microscopy (SEM) to ensure no fine particles <1 μm in size remained adhered to particle surfaces. Particles in the size fraction 25 – 50 μm were subjected to further heat treatment; annealing in air was performed in a standard muffle furnace heated to 300°C, 400°C, 600°C, 800°C, 1000°C and 1250°C at a ramp rate of 5°C min⁻¹ and held for 4 hours. Annealing in inert (Ar) and reducing (5% H₂ / 95% N₂) atmospheres was performed in a standard tube furnace, sealed and purged with gas for 2 hours prior to

heating to 1000°C at a ramp rate of 5°C min⁻¹ and held for 4 hours. Following dissolution, particles were dried by patting with tissue and placed in an inert-atmosphere desiccator prior to analysis.

2.2. Particle Characterisation. The surface area of each crushed size fraction and the 25 – 50 μ m heat-treated particles was determined using the Brunauer, Emmett and Teller (BET) method, using a nitrogen adsorbate. Particles were degassed for 3 hours at 200°C and analysed using a Beckman Coulter SA3100 instrument. The surface oxidation state of annealed particles was determined using a Kratos Axis Ultra X-ray Photoelectron Spectrometer (XPS) with a monochromatic Al K α source. The spectral intensities recorded were converted to surface oxidation state concentrations by first subtracting a Shirley-type background²² and then obtaining accurate peak positions by fitting peaks using a mixed Gaussian / Lorenzian line shape. All photoelectron binding energies were referenced to C 1s adventitious contamination peaks at a binding energy of 285 eV.

Powder X-ray Diffraction (XRD) patterns were acquired using a Bruker D2 Phaser diffractometer operating in transmission mode at 40 kV and 30 mA with Cu K α_1 radiation. Data were collected between 20 < 2 θ < 60° at 2° min⁻¹ and a step size of 0.02°. Rietveld analysis of phases identified in the diffraction patterns was performed using the refinement software, Topas (Bruker), and strain was calculated using the Double-Voigt approach²³. Secondary electron images of particles were acquired using a Hitachi TM3030 SEM, operating at an accelerating voltage of 15 kV and a working distance of 18 mm.

2.3. CeO₂ Dissolution. Samples of 0.1 g of each size fraction of CeO₂ particles were placed in 50 mL PTFE vessels, which had been cleaned according to ASTM C1285²¹. Each vessel was filled with 40 mL of 0.01 M HNO₃ (pH 2.0), sealed and placed within a Carbolite oven. For 25 – 50 μ m CeO₂ particles with no annealing treatment, experiments were performed at 40 °C, 70 °C and 90°C (± 0.5 °C). All other experiments were performed at 90 (± 0.5) °C. Each experiment was performed in triplicate with duplicate blanks, and sampling was conducted at regular intervals from 0 to 35 days. During each sampling, an aliquot (1.2 mL) of the supernatant liquid in each vessel was removed and filtered (0.22 µm) prior to analysis by Inductively

Coupled Plasma – Mass Spectroscopy (ICP-MS) (Agilent 4500 Spectrometer). Solution data expressed as the normalised weight loss (N_L , expressed in g m⁻²) were calculated from the Ce concentrations after normalisation by the surface area S (m^2) of the sample in contact with the solution (assuming that all the surface is "reactive") and by the mass loading of the element considered in the solid (i.e. f_i is expressed as the ratio between the mass of the considered element and the overall mass of the leached sample), according to Eqn. 1 ^{24,25}:

$$N_L(i) = \frac{m_i}{f_i \times S} \tag{1}$$

where m_i (g) corresponds to the total amount of the element *i* measured in the solution.

Derivation of Eqn. 1 as a function of dissolution time gives the normalised dissolution rate of Ce ($R_L(Ce)$), expressed in g m⁻² d⁻¹) ²⁶:

$$R_L(Ce) = \frac{dN_L(Ce)}{dt} = \frac{1}{f_{Ce} \times S} \times \frac{dm_{Ce}}{dt} = \frac{1}{f_{Ce} \times S} \times \frac{d}{dt} (C_{Ce} \times V) \approx \frac{V}{f_{Ce} \times S} \times \frac{dC_{Ce}}{dt}$$
(2)

where C_{ce} (g L⁻¹) is the concentration of Ce in solution and V (L) is volume of leachate. In this expression, the f_{Ce} and S values are usually assumed to remain almost constant during the leaching experiments, which can be considered as a valid assumption at the beginning of the dissolution tests for powdered samples.

3. RESULTS

3.1. Characterisation of defect structures in CeO₂. Particles from the 25 – 50 μ m size fraction were annealed in air at temperatures ranging from 300 – 1250°C. One sample was analysed without annealing for comparison. High temperature annealing in air lead to a reduction in surface area, as shown in Figure 1. For example, annealing at 600°C gave a surface area of 0.20 ± 0.01 m² g⁻¹, while at 1000°C the surface area was 0.11 ± 0.01 m² g⁻¹. We hypothesise that the surface area decrease reflects the removal of 'sharp edges' and other intrinsic surface defects in the CeO₂ particles, during heating. When the annealing temperature was increased to the range 700°C to 1250°C, a sharp drop in surface area was observed. This may be due to

agglomeration of particles at these higher temperatures. Dissolution rate investigation was performed on samples annealed at 600°C and 1000°C; both were representative of the removal of 'sharp edges' and intrinsic surface defects, but samples annealed at 1000°C additionally represented particles with a significantly lower surface area.



Figure 1. Surface area analysis of CeO₂ 25 – 50 μ m particles as a function of annealing temperature, as determined by the BET method. Errors represent 1 σ from triplicate samples.

An attempt to generate oxygen vacancies in particles of CeO_2 in the 25 – 50 µm size fraction was made by annealing at 1000°C in the absence of oxygen (inert or reducing environments). The creation of an oxygen vacancy in CeO_2 is accompanied by the release of two electrons from lattice O^{2-} , resulting in the reduction of Ce^{4+} ions to Ce^{3+} ²⁷. Reduction from CeO_2 to Ce_2O_3 occurs through the formation of CeO_{2-x} intermediate phases²⁸; to confirm the presence of oxygen vacancies, the oxidation state of Ce was determined using XPS, shown in Figure 2. Trivalent Ce was identified in the particle surfaces (top ~10 nm), confirming that Ce^{4+} reduction to Ce^{3+} had occurred. The Ce 3d XPS data were resolved using ten components²⁹; Table 1 describes each of the components found in the current study according to the accepted Burroughs nomenclature³⁰. These data are in excellent agreement with previously published XPS data for partially reduced CeO_2 ^{31,32} and show that samples annealed in an inert or reducing environment were not fully reduced to Ce_2O_3 , but contained a combination of identifiable contributions from both Ce^{3+} and Ce^{4+} oxidation states at the sample surface. This indicates the formation of partially reduced CeO_2 with a stoichiometry of CeO_{2-x} , and the generation of oxygen vacancies. These defects have been previously characterised as Frenkel-type defects³³.



Figure 2. X-ray photoelectron oxidation state analysis of CeO_2 annealed under reducing conditions, showing Ce 3d spectra for CeO_{2-x} annealed in a reducing (H₂ / N₂) atmosphere. Standard spectra for Ce^{3+} and Ce^{4+} -containing materials are shown for comparison. Peak deconvolution was performed and labelled according to the accepted Burroughs nomenclature³⁰ for Ce 3d as described in Table 1.

Table 1. Experimental binding energies and FWHM for Ce 3d peaks derived from a CePO₄ standard, and CeO₂ and CeO_{2-x}

Ce 3d peak	Oxidation state	Binding energy	FWHM	
assignment	assignment	(eV) (± 0.1)	(eV) (± 0.05)	
V ⁰	3+	881.90	3.0	
V	3+	882.40	1.8	
V	4+	885.30	4.3	
V''	4+	888.40	4.2	
V'''	4+	898.50	2.3	
u ⁰	3+	899.95	3.0	
U	3+	901.10	1.8	
u'	4+	903.00	4.3	
u''	4+	907.60	4.2	
u'''	4+	916.80	2.3	

particles investigated in the current study.

The fraction of Ce^{3+} (denoted α) was determined from Ce 3d peak area analysis, according to Equation 3. On the basis that half an oxygen vacancy is formed per Ce^{3+} created, the deviation, *x*, from the ideal oxygen stoichiometry was determined for each sample, as shown in Table 2.

$$\alpha = \frac{Ce^{3+}}{\Sigma Ce} \tag{3}$$

Annealing in a reducing (H₂ / N₂) atmosphere resulted in CeO_{2-x} with the greatest Ce³⁺ fraction ($\alpha = 0.370 \pm 0.019$) and a stoichiometry of CeO_{1.82}, while annealing in an inert (Ar) atmosphere gave a higher Ce³⁺ fraction than annealing in air under the same conditions ($\alpha = 0.221 \pm 0.011$ and $\alpha = 0.080 \pm 0.004$, respectively). As the Ce³⁺ fraction is directly related to the concentration of oxygen vacancies generated, it can be assumed that the oxygen vacancy concentration is proportional to the Ce³⁺ fraction of each sample.

Table 2. Calculated surface stoichiometry and Ce³⁺ fraction (α) of samples and standards investigated in this study, as determined from Ce 3d XPS data. Errors quoted for α are standard deviations based upon duplicate XPS scans on the same sample.

Sample treatment	Surface Stoichiometry	α	
Ce(III)PO ₄ standard		0.952 ± 0.051	
H_2/N_2 atmosphere 1000°C	CeO _{1.82}	0.370 ± 0.019	
Ar atmosphere 1000°C	CeO _{1.89}	0.221 ± 0.011	
Air 1000°C	CeO _{1.96}	0.080 ± 0.004	
Air 600°C	CeO _{1.93}	0.143 ± 0.007	
Air - no annealing	CeO _{1.99}	0.021 ± 0.008	

The crystal structure of non-stoichiometric $CeO_{2\cdot x}$ was determined using X-ray diffraction. Because the samples were found to be chemically unstable upon exposure to air, they were loaded into the diffractometer within minutes of being removed from the annealing furnace. By comparison to the known correlation of the fluorite lattice parameter with the oxygen stoichiometry in the $CeO_{2\cdot x}$ system³⁴, the three main diffraction peaks identified in Figure 3 were assigned to $CeO_{1.75}$, $CeO_{1.81}$ and CeO_2 . The second phase is in reasonable agreement with the XPS analysis of this sample, which gave a stoichiometry of $CeO_{1.82}$ (Table 2). The $CeO_{1.75}$ phase was not identified by XPS; this may be due to the rapid oxidation, as described below. The spatial distribution of the phases within the samples was not investigated, however, since Ce^{3+} species have previously been found to accumulate at grain boundaries within partially reduced CeO_2^{-35} , grain boundaries may host $CeO_{1.75}$ and $CeO_{1.81}$. The identification of these phases following annealing in a reducing atmosphere is consistent with the formation of oxygen vacancies³⁴.



Figure 3. X-ray diffraction data for CeO_{2-x} formed by annealing CeO_2 at 1000°C under a reducing atmosphere (H₂ / N₂). XRD patterns were recorded, in air, at regular periods from 0.76 h to 200 h after annealing. Correlation of the fluorite lattice parameter with the known oxygen stoichiometry in the CeO_{2-x} system³¹ allowed diffraction peaks to be identified as $CeO_{1.75}$, $CeO_{1.81}$ and CeO_2 . Inset shows full diffraction patterns recorded at 1h and 200h, indexing peaks for CeO_2 .

With prolonged exposure to ambient conditions (air, for up to 200 hours), the relative intensity of the peaks assigned to $CeO_{1.75}$, $CeO_{1.81}$ and CeO_2 changed (Fig. 3). Quantitative Rietveld analysis of the distribution of phases as a function of time are shown in Figure 4. Initially, the sample was composed of ~15% $CeO_{1.75}$, ~68% $CeO_{1.81}$ and ~16% CeO_2 (± 1.0 %). After 56 hours, $CeO_{1.75}$ was no longer observed and the distribution of phases remained at 4.5 ± 0.9 % $CeO_{1.81}$ and 95.2 ± 1.3 % CeO_2 (Fig. 4) until 200 hours. The increase in the content of oxygen, corresponding to a decrease in oxygen vacancies, was found to occur at a rate of (1.40 ± 0.14) × 10⁻³ mol h⁻¹ during the first 24 hours of exposure to ambient conditions. This process was associated with a change in the volume of the fluorite unit cell, as calculated from Rietveld-refined lattice parameters. The fluorite unit cell was found to contract, from 167.62 ± 3.4 Å³ for $CeO_{1.75}$ to 164.97 ± 3.30 Å³ for $CeO_{1.81}$ and to 157.58 ± 3.15 Å³ for CeO_2 . This volume reduction is consistent with the oxidation of Ce^{3+} to

 Ce^{4+} ; Ce^{3+} has a larger ionic radius than Ce^{4+} (1.03 Å and 0.97 Å, respectively)³⁶. Overall, the sample annealed in a H₂/N₂ atmosphere underwent a 6% volume contraction within several hours of exposure to oxygen under ambient atmospheric conditions. The lattice strain, as calculated using Rietveld analysis of the CeO_2 phase, was found to significantly decrease during oxidation, from a value of $\mathcal{E} = 0.092 \pm 0.01$ at 2 hours to $\mathcal{E} =$ 0.034 ± 0.01 after 200 hours, consistent with the observed change in lattice volume. Due to the rapid nature of these chemical and structural changes, all dissolution experiments were performed within 30 minutes of sample annealing to minimise the effects of air oxidation.



Figure 4. Weight fraction of CeO_2 and CeO_{2-x} phases present in a sample of CeO_2 annealed at 1000°C in a H₂ / N₂ atmosphere as a function of time, as determined by quantitative Rietveld analysis of powder XRD patterns. After 56 h, there was no further change in the distribution of phases up to 200 h of observation. Errors are ± 3% based on Rietveld analysis methods.

3.2. CeO₂ dissolution. Cerium dioxide dissolution experiments were performed under aggressive conditions (0.01 M HNO₃, 90°C) to attain quantitative dissolution kinetic data within several weeks. It has been suggested that reduction of Ce⁴⁺ to Ce³⁺ may be induced by nitrous acid (HNO₂) formed through nitric acid instability at high acidity (> 0.5 M)¹³, however, this is not expected to influence the dissolution reaction in the current experiments, which were at an acidity much less than 0.5 M.

The dissolution of CeO₂ can be described by two distinct dissolution regimes, as shown in Figure 5; the first, rapid regime is far from solution saturation, and thus represents the kinetically-controlled dissolution rate, R_k . The second regime is representative of near-solution saturation, controlled by the solubility of CeO₂ and concentration of Ce(OH)₂^{2*} in solution. This is less rapid due to thermodynamic effects that occur close to equilibrium, and is denoted R_t . To understand the role of defects in the dissolution of CeO₂, it was necessary to focus only on the kinetic regime of dissolution, therefore we herein report R_k only. Generally, the change between the kinetic and the thermodynamically-controlled dissolution regimes occurred at ~7 days, unless stated. Figure 5 shows the normalised mass loss of Ce (N_L , g m⁻²) from non-annealed 25 – 50 µm particles of CeO₂ as a function of time at 40°C, 70°C and 90°C. The normalised dissolution rate of Ce was $R_k = (1.90 \pm 0.74) \times 10^{-4}$ g m⁻² d⁻¹ at 90°C. The dependence of dissolution on temperature gave an activation energy of 30.4 ± 4 kJ mol⁻¹, assuming an Arrhenius-type relationship. This is consistent with previous work on CeO₂ and other fluorite-type dioxides such as ThO₂ and PuO₂ (E_a = 20 - 37 kJ mol⁻¹)^{13,25,37} and is indicative of a surface-controlled dissolution mechanism.



Figure 5. Concentration of Ce, normalised to surface area (g m⁻²) for non-annealed CeO₂ particles $25 - 50 \mu m$ in size, reacted in 0.01M HNO₃ at 40°C, 70°C and 90°C, as a function of time. Graph depicts the two main dissolution rates observed, kinetically controlled dissolution (R_k) and thermodynamic effect-controlled dissolution (R_t). Errors represent 1σ from triplicate samples.

3.3. Effect of defect annealing / surface area on CeO₂ dissolution kinetics. The dissolution rates of $25 - 50 \mu$ m particles of CeO₂ annealed in air, and non-annealed CeO₂ of four different particle sizes, are shown in Table 3 and Figure 6.

The surface area normalised dissolution rates were observed to increase with increasing surface area. This behaviour is unexpected, since normalisation to the surface area should result in similar dissolution rates. Samples of CeO₂ that were crushed and sieved to different size fractions clearly showed this correlation, with dissolution rates increasing from $R_k = (4.17 \pm 1.13) \times 10^{-5} \text{ g} \text{ m}^{-2} \text{ d}^{-1}$ for 2-4 mm particles to $R_k = (1.38 \pm 0.28) \times 10^{-4} \text{ g} \text{ m}^{-2} \text{ d}^{-1}$ for particles in the 75 – 150 µm size fraction. The exposure of grain boundaries during particle crushing may be responsible for enhancing the reactive surface area, resulting in the observed trend, as follows: each particle is composed of a number of grains, with a grain size of 5 – 30 µm. As the particle size decreases, from 2 – 4mm to 75 – 150 µm, the ratio of particle size to the number of grains also decreases. Because the grain size remains constant, the shrinking size of the particle exposes more grain boundaries to the dissolution medium. Grain boundaries are known to preferentially dissolve due to the presence of a greater number of energetically reactive surface sites (e.g. defects) compared to the surfaces¹¹. Once the smallest particle size is reached, each particle comprises only a few grains, with a high number of grain boundaries exposed to solution, which leads to the most rapid dissolution rates (Fig. 6).



Figure 6. Normalised dissolution rates of Ce, R_k (g m⁻² d⁻¹) as a function of specific surface area and particle size.

Similarly, thermally annealed particles of CeO₂ with the same particle size ($25 - 50 \mu$ m) exhibited a factor of seven decrease in their dissolution rate compared with non-annealed particles ($R_k = (2.98 \pm 0.50) \times 10^{-5} \text{ g m}^{-2} \text{ d}^{-1}$ and (2.67 ± 0.30) $\times 10^{-5} \text{ g m}^{-2} \text{ d}^{-1}$, for 600°C and 1000°C respectively), as shown in Figure 7. Because the reduction in surface area in these samples was associated with the high temperature annealing of intrinsic defects and sharp edges generated through sample crushing, it can be concluded that these energetically reactive features influence the dissolution kinetics. This is in agreement with the work of Claparede *et al.*¹³ who found that increasing the sintering temperature of Th_{1-x}Ce_xO₂ powders led to a reduction in dissolution rate, attributed to the elimination of crystal defects and amorphous domains. The current work extends this hypothesis to samples prepared by crushing, generating surface defects in the form of "sharp edges", as proposed by Corkhill et al.¹¹



Figure 7. Summary of the normalised dissolution rates of Ce, R_k (g m⁻² d⁻¹) as a function of particle size and annealing in air, Ar and H_2/N_2 atmospheres.

Sample treatment	Surface area (m ² g ⁻¹)	α	R_k (g m ⁻² d ⁻¹)		
25 – 50 μm samples					
H_2/N_2 atmosphere 1000°C*	0.53 ± 0.01	0.370 ± 0.019	-		
H_2/N_2 atmosphere 1000°C**	2.49 ± 0.02	0.090 ± 0.003	(4.06 ± 0.75) x 10 ⁻⁴		
Ar atmosphere 1000°C	0.20 ± 0.01	0.221 ± 0.011	(4.62 ± 0.90) x 10 ⁻⁵		
Air 1000°C	0.11 ± 0.01	0.080 ± 0.004	(2.67 ± 0.30) x 10 ⁻⁵		
Air 600°C	0.20 ± 0.01	0.143 ± 0.007	(2.98 ± 0.50) x 10 ⁻⁵		
No annealing	0.30 ± 0.01	0.021 ± 0.008	(1.90 ± 0.74) x 10 ⁻⁴		
Other particle sizes, no annealing					
75 – 150 μm	0.07 ± 0.01	0.022 ± 0.007	(1.38 ± 0.28) x 10 ⁻⁴		
300 – 600 μm	0.02 ± 0.01	0.021 ± 0.008	(1.02 ± 0.26) x 10 ⁻⁵		
2 – 4 mm	0.01 ± 0.01	0.021 ± 0.004	(4.17 ± 1.13) x 10 ⁻⁵		

Table 3. Calculated dissolution rates (R_k) for CeO₂ and CeO_{2-x} reacted in 0.01M HNO₃ at 90°C between 0 and 35 days of

dissolution, with the	corresponding measu	ured surface area and	l oxygen vacancy content (Ce	^{2*} fraction, α)
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* Initial surface area and Ce³⁺ fraction, measured within minutes of removal of furnace

**Surface area and Ce^{3+} fraction, measured 1 day after dissolution.

3.4. Effect of oxygen vacancies on dissolution kinetics. For CeO₂ samples annealed in different atmospheres at 1000°C, the samples with the greatest oxygen vacancy content were found to dissolve at the most rapid rates. For example, the dissolution rate of CeO₂ annealed in Ar was slightly enhanced compared with CeO₂ annealed in air at the same temperature, by a factor of two (Table 3). The dissolution of CeO_{1.82} (annealed in H₂ / N₂) showed a significantly enhanced dissolution rate (Table 3, Fig. 7). Samples of this material were analysed by SEM at several time points during dissolution; in contrast to particles of stoichiometric CeO₂, which did not change their shape or morphology upon 1 or 35 days of dissolution, (Figure 8a – c), examination of the CeO_{1.82} particles revealed a dramatic transformation. Prior to dissolution, these particles were 25 – 50 µm in size (Fig. 6d), but within 1 day they were observed to break apart into smaller fragments ~5 µm in size, with a morphology suggestive of individual grains (Fig. 8e). After 35 days of dissolution, some larger particles remained, but most of the sample had decomposed into single grain-sized fragments (Fig. 8f). Post-dissolution BET analysis of these samples after 1 day of dissolution gave a surface area of 2.49 ± 0.01 m² g⁻¹, an increase by a factor of 5 compared with the initial surface area (Table 3).

Normalisation of the dissolution rate data to this surface area value gave a rate of $R_k = (4.06 \pm 0.75) \times 10^{-4}$ g m⁻² d⁻¹, 15 times greater than for air or Ar annealed samples of the same particle size (Fig. 7). Due to the extremely rapid dissolution of this material, it was necessary to calculate the R_k between 0 and 1 day of dissolution, after which time the normalised release of Ce was constant. When analysed by XPS after 1 day of dissolution, these samples gave a surface stoichiometry of CeO_{1.96}, which indicates that a rapid oxidation of CeO_{2-x} occurred, concurrent with a significant reduction in the oxygen vacancy content.



25 µm

25 μm

25 µm

Figure 8. SEM images of $25 - 50 \mu$ m particles of non-annealed CeO₂ (a) prior to dissolution; and after dissolution in 0.01 M HNO₃ at 90°C for (b) 1 day and (c) 35 days; and H₂ / N₂ annealed CeO_{2-x} (d) prior to dissolution; and after dissolution in 0.01 M HNO₃ at 90°C for (e) 1 hr and (f) 35 days.

4. DISCUSSION

4.1. Microstructure and surface defect – dissolution rate relationship in CeO₂ and CeO_{2-x}. A first-

order expression that describes the concentration of Ce (C) as a function of time (t) has been considered:

$$C = A \left(1 - e^{-bt} \right) \tag{4}$$

where A is the solubility limit of CeO_2 and b is the rate constant. Expanding Eqn. 4 to account for the factors that are expected to contribute to the dissolution kinetics (surface area, oxygen vacancy content and the different relative solubilities of Ce^{3+} and Ce^{4+}) a first order dissolution rate expression can be derived for the CeO₂ samples investigated in this study:

$$C = \frac{s}{s_0} [\alpha A_3 (1 - e^{-b_3 t}) + (1 - \alpha) A_4 (1 - e^{-b_4 t})]$$
(5)

where S is the measured surface area (after dissolution in the case of CeO_{2-x}), S₀ is a reference surface area of 1 m² g⁻¹, α is the fraction of Ce³⁺ (representing the oxygen vacancy content) at the particle surface as determined by XPS, A₃ and A₄ correspond to the solubility limits of Ce³⁺ and Ce⁴⁺ respectively, and b₃ and b₄ are the rate constants for Ce³⁺ and Ce⁴⁺.

A genetic algorithm, which can be applied to solve constrained optimisation problems³⁸, was used to fit the individual parameters in Eqn. 5, providing optimum solubility and rate constant values for Ce³⁺ and Ce⁴⁺ based on the dissolution data obtained in this study; the data and model fits are shown in Figure 9. The modelled Ce³⁺ and Ce⁴⁺ solubilities, A_3 and A_4 , gave best-fit values of $(5.10 \pm 0.54) \times 10^{-4}$ mol dm⁻³ and $(1.59 \pm 0.08) \times 10^{-6}$ mol dm⁻³, respectively (errors represent the 1 σ confidence interval). These are in reasonable agreement with geochemically modelled solubility values, for example, Ce⁴⁺ solubility was calculated using PHREEQC (LLNL database) to be 4.01×10^{-6} mol dm⁻³. When considering only CeO₂ annealed at various temperatures in air and CeO₂ particles of different sizes (shown in Figs. 9a and b, respectively), the rate constant for Ce³⁺ was found to be $b_3 = (0.245 \pm 0.06) d^{-3}$. For Ce⁴⁺ the scatter in the experimental data, combined with the low absolute solubility limit, resulted in poorly-defined estimates for b_4 ; a value of the order $1.00 \times 10^{-7} d^{-1}$ is tentatively suggested. Because the model to data fit was in agreement for CeO₂ (Fig. 9a and b), and the solubilities of the different Ce oxidation states were similar to geochemically modelled solubility values, it can be surmised that the parameters in Eqn. 5 adequately describe the dissolution rate of stoichiometric CeO₂.

Figure 9c shows the dissolution data and model fit for CeO_{2-x} samples annealed in Ar and H_2 / N_2 atmospheres. The model fit to the parameters in Eqn. 5 agrees somewhat with samples annealed in Ar,

however, the model fit to the parameters is not in agreement with data from samples annealed in H_2 / N_2 , despite normalisation to the post-dissolution surface area and consideration for the elevated content of the more highly soluble Ce³⁺ species in this sample (Eqn. 5). Surface area data in Table 3 show that during dissolution, the surface area of CeO_{2-x} annealed in a H_2 / N_2 atmosphere increased by a factor of 5 after 1 day of dissolution. The first order dissolution model described above predicts that a surface area increase of a factor of 50 would be required to create a good fit to the parameters in Eqn. 5, assuming all other values are correct. Specific surface area, as measured by BET, is routinely accepted to be the same as reactive surface area. However, recent studies have discussed that this may not necessarily be the case due to different activation energies of different high energy surface sites, e.g., grain boundaries and oxygen vacancies³⁹. It is possible that the BET surface area measurements taken here, 1 day after dissolution, underestimated the reactive surface area of these samples during the initial stages of dissolution due to the presence of a high proportion of exposed grain boundaries, which are known to be significantly more reactive than the bulk.



Figure 9. Ce concentrations (mol dm⁻³) derived from dissolution at 90°C in 0.01M HNO₃ (points) compared with modelled fit to Eqn. 5 (lines) for: (a) $25 - 50 \mu m$ CeO₂ subject to different annealing temperatures in air; (b) CeO₂ crushed to different size fractions, with no subsequent annealing; and (c) $25 - 50 \mu m$ CeO₂ subject to annealing in different atmospheres at 1000°C. The parameters in Eqn. 5 do not fit the data for H₂ / N₂ annealed samples shown in (c), suggesting another factor may contribute to the dissolution rate.

4.2. Physical effects of oxidation on the kinetics of dissolution. As described above, not only the enhanced chemical solubility of Ce^{3+} influenced the dissolution kinetics of CeO_{2-x} , but the physical effects leading to a change in the sample surface area also played a significant role. The mechanism through which this is postulated to occur for CeO_{2-x} investigated in the current study is described below.

Upon contact with air or an oxic solution, the Ce^{3+} ions in CeO_{2-x} , present as a result of charge compensation for oxygen vacancies, were rapidly oxidised to Ce^{4+} . Due to the difference in ionic radius between Ce^{3+} and Ce⁴⁺ (with Ce⁴⁺ the smaller cation), this oxidation process was associated with a contraction of the fluorite unit cell volume, on the order of $\Delta V = -3.25 \pm 0.06 \text{ Å}^3$ according to Rietveld analysis of XRD data (a 6.0 ± 0.1 % volume decrease relative to the starting volume). Ceramic materials that have undergone volume changerelated phase transition exhibit significant internal lattice strain; according to Rietveld analysis, the lattice strain was initially high in CeO_{2-x} , reducing by a factor of 2.6 during oxidation to CeO_2 . Additionally, the presence of oxygen vacancies results in enhanced oxygen mobility, especially at grain boundaries where it is expected that a higher proportion of Ce³⁺ exists³⁵; this facilitates preferential dissolution of grain boundaries¹¹. In this study, during the dissolution of CeO_{1.82} to CeO_{1.96}, the combined influence of lattice strain and enhanced oxygen mobility, created by the removal of oxygen vacancies, resulted in the disintegration of particles, preferentially along grain boundaries, producing individual grain-sized fragments after 1 day of dissolution (Fig. 8e - f). Importantly, this resulted in an increase in reactive surface area, which exposed fresh surfaces for dissolution and enhanced the overall dissolution rate. After 1 day, the dissolution rate was observed to be more or less constant, likely because there is reduced oxygen atom mobility within the nearly stoichiometric CeO_{1.96}. Hence, the combined roles of oxygen vacancies and grain boundaries should be considered significant in defining the dissolution kinetics of non-stoichiometric CeO₂ and other fluorite-type materials.

The dual role of oxygen vacancies and grain boundaries in dissolution has been observed in other fluoritetype materials, albeit those comprisising a solid solution of different chemical species. For example, Horlait et al.¹⁷ reported the 'crumbling' of $Ce_{1-x}Ln_xO_{2-x/2}$ materials along grain boundaries during dissolution. This was associated with an increase in reactive surface area by up to a factor of 5. It was postulated that the incorporation of Ln³⁺ played a significant role in the dissolution mechanism through changes to the lattice volume^{17,40}. Similarly, Finkeldei et al.⁴¹ found that defect fluorite and pyrochlore structures in the ZrO₂-Nd₂O₃ group preferentially dissolved along grain boundaries and at triple junctions, leading to disintegration. Substitution of elements into a fluorite lattice results in a volume change due to the difference in ionic radius of the substituting element³⁶; the current study highlights the equally important role of oxygen vacancies on lattice volume change when only a single chemical species is present (CeO₂), confirming the importance of this defect feature in dissolution rate kinetics.

The association of defect structures with grain boundaries and their influence on the dissolution of CeO₂ and ThO₂ analogues for spent fuel was recently hypothesised by Corkhill et al.¹¹, who showed that the dissolution of grain boundaries was influenced by the crystallographic orientation of the grains and the corresponding mean misorientation angle of the grain boundaries. Boundaries with higher mean misorientation angles were found to dissolve at faster rates than those with low mean misorientation angles; it was hypothesised that higher misorientation angle grain boundaries contained a greater proportion of defects, with a high reactive surface area, which led to the enhanced dissolution rates. This study and others^{17,18}, have highlighted the important role of grain boundaries, and particularly their association with defect structures (e.g. oxygen vacancies) on dissolution processes in spent fuel analogues. This underlines the requirement for new methodologies capable of linking grain boundary structure with dissolution kinetics, to fully understand their role in the dissolution of spent fuel materials⁴².

4.3. Implications for the dissolution of spent fuels. It was recently demonstrated for a suite of fluorite type materials, including CeO₂, CaF₂ and UO₂, that the surface stability is correlated with the proportion of high energy surfaces, even when different chemical processes of dissolution occur, for example for CeO₂ and CaF₂⁸. ThO₂ has also been shown to confirm this hypothesis^{11,19}. Together with the resulted presented herein, these findings highlight the potential for energetically reactive surface sites to play a role in UO₂ dissolution.

Spent UO_2 fuel forms a complicated defect structure during fission within a reactor and it is not yet clear the extent to which different types of reactive surface sites may influence the dissolution kinetics of spent fuel, however studies on SIMFUEL have shown that doping with trivalent rare elements (to simulate fission products) resulted in lattice contraction¹² and a loss of cubic symmetry. On the basis of the results presented in the current study, we can hypothesise that such a lattice modification may lead to grain boundary decohesion. Evidence supporting this hypothesis includes a study of UO₂ fuel stored in aerated, moist environments, which was found to degrade along its grain boundaries and become heavily fractured⁴³. The effects of lattice volume change on the integrity of grain boundaries in these materials, and the extent to which it may affect the dissolution kinetics, requires further exploration. Recent work has shown that Mixed Oxide Fuels (MOx), containing Pu^{3+} , with a PuO_{2-x} stoichiometry of $PuO_{1.61}^{44}$, may also be prone to accommodating oxygen vacancies⁴⁵, although the distribution of such defects between grains and the grain boundaries is currently not known. In light of the current limited knowledge of the atomic-scale structure of spent fuels, further work is required on inactive analogues, UO₂ and spent fuel materials to ascertain the effects of reactive surface sites on dissolution kinetics, as distinct from oxidative dissolution or radiolysis effects, to reduce the uncertainty associated with the prediction of spent fuel dissolution within a geological disposal facility.

5. CONCLUSIONS

We present an investigation of the role of energetically reactive surface sites on the initial dissolution rate of CeO_2 to ascertain whether these features may result in over-estimation of dissolution rates in the laboratory. Reactive surface sites and 'sharp edges' were simulated by crushing samples to different size fractions and through high temperature annealing. Oxygen vacancy defects in CeO_{2-x} were generated by high temperature annealing in oxygen-free conditions. Surface area was used as a measurement of 'sharp edges' and intrinsic defects, and the total Ce^{3+} fraction measured by XPS was used as a measurement of the oxygen vacancy concentration. The relative dissolution rates of these materials showed that there was a quantifiable relationship between energetically reactive surface sites and dissolution rate, which was dependent upon reactive surface site concentration and the solubilities of Ce^{3*} and Ce^{4*} ; the greater the concentration of reactive surface sites, the higher the dissolution rate. The elimination of oxygen vacancies in CeO_{2*x} was found to significantly influence the dissolution rate; this was associated with changes in the lattice volume and lattice strain, manifested as grain boundary decohesion and increased surface area. Thus, it is clear that reactive surface sites, including artefacts of sample preparation, intrinsic material defects and grain boundaries, have the potential to significantly enhance the dissolution of spent oxide fuel. We show that such effects can be quantified, thus it is possible to reduce the uncertainty associated with laboratory measurements of UO_2 dissolution, albeit in a simplified analogue system. It will be essential to develop further careful studies designed to interrogate the effect of reactive surface features on the durability and dissolution of materials such as UO_2 and MOx, as distinct from solubility, surface area and radiation effects.

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