



This is a repository copy of *Diffusion of dextran within poly(methacrylic acid) hydrogels*.

White Rose Research Online URL for this paper:  
<http://eprints.whiterose.ac.uk/79454/>

Version: Accepted Version

---

**Article:**

AL-Baradi, A.M., Mears, M., Jones, R.A.L. et al. (1 more author) (2012) Diffusion of dextran within poly(methacrylic acid) hydrogels. *Journal of Polymer Science Part B: Polymer Physics*, 50 (18). 1286 - 1292. ISSN 0887-6266

<https://doi.org/10.1002/polb.23120>

---

**Reuse**

Unless indicated otherwise, fulltext items are protected by copyright with all rights reserved. The copyright exception in section 29 of the Copyright, Designs and Patents Act 1988 allows the making of a single copy solely for the purpose of non-commercial research or private study within the limits of fair dealing. The publisher or other rights-holder may allow further reproduction and re-use of this version - refer to the White Rose Research Online record for this item. Where records identify the publisher as the copyright holder, users can verify any specific terms of use on the publisher's website.

**Takedown**

If you consider content in White Rose Research Online to be in breach of UK law, please notify us by emailing [eprints@whiterose.ac.uk](mailto:eprints@whiterose.ac.uk) including the URL of the record and the reason for the withdrawal request.



[eprints@whiterose.ac.uk](mailto:eprints@whiterose.ac.uk)  
<https://eprints.whiterose.ac.uk/>



## The diffusion of dextran within poly(methacrylic acid) hydrogels

Journal:	<i>Journal of Polymer Science Part B: Polymer Physics</i>
Manuscript ID:	12-02-0085.R1
Wiley - Manuscript type:	Full Papers
Date Submitted by the Author:	n/a
Complete List of Authors:	Al-Baradi, Ateyyah; The University of Sheffield, Department of Physics and Astronomy Mears, Matthew; The University Of Sheffield, Physics And Astronomy Geoghegan, Mark; University of Sheffield, Dept. of Physics and Astronomy; Jones, Richard; University of Sheffield, Department of Physics and Astronomy
Keywords:	gels, diffusion, diffusion control, fluorescence, networks, swelling, viscosity

SCHOLARONE™  
Manuscripts

Review

## The diffusion of dextran within poly(methacrylic acid) hydrogels

Ateyyah M. AL-Baradi, Matthew Mears, Richard A. L. Jones, Mark Geoghegan

Department of Physics and Astronomy, The University of Sheffield, Hounsfield Road, Sheffield, S3 7RH, UK

Correspondence to: Matthew Mears (m.mears@sheffield.ac.uk)

### ABSTRACT

We describe a fluorescence correlation spectroscopy investigation into the diffusion of fluorescein-tagged dextran (FDEX) in a poly(methacrylic acid) (PMAA) hydrogel. The temperature dependence of FDEX diffusion is shown to follow Zimm behaviour in pure water, and the decrease in the diffusion coefficient when in the PMAA hydrogel has been modelled. The addition of acid and alkali (HCl and NaOH respectively) not only control the swelling and collapse of the hydrogel but also reveal a strong pH dependence of the dextran diffusion coefficient, which shows a (non-monotonic) increase with pH. The addition of NaCl and CaCl<sub>2</sub> salts similarly showed evidence of network swelling, most notably at low salt concentration, but also that the diffusion coefficient within the gel at these low concentrations is larger than in the equivalent solution without the hydrogel, indicating that the combination of hydrogel and salt works to increase the diffusion coefficient above that in pure water.

**KEYWORDS** Diffusion, hydrogel, FCS

### INTRODUCTION

The theory of polymer diffusion through a molten polymer network has already been well established by Doi and Edwards,<sup>1</sup> showing that the diffusion coefficient depends upon the distance between crosslinks known as mesh size. However a combination of this theory with the responsive behaviour of polymer hydrogel systems is of considerable practical importance because controlling the diffusion of polymers in a hydrogel network has applications in drug delivery systems<sup>2,3</sup> as well as in responsive adhesives,<sup>4</sup> biosensors,<sup>5</sup> and contact lens development.<sup>6</sup> The underlying phenomena of gelation and swelling have been extensively studied,<sup>7,8</sup> with more recent developments utilising fluorescence<sup>9,10</sup> and light-scattering techniques<sup>11</sup> to investigate the swelling properties of hydrogels.

1  
2  
3 Diffusion through heterogeneous media is important in biological systems, ranging from  
4  
5 examining the absorption of doxylamine in mice<sup>12</sup> to the controlled delivery of drugs from  
6  
7 biocompatible responsive hydrogels such as hydroxypropyl methylcellulose,<sup>13</sup> or more  
8  
9 complicated systems involving hydrogels impregnated with biodegradable microcarriers.<sup>14</sup>  
10  
11 Theoretical work has also been published that provides some understanding of such network  
12  
13 systems and highlights the parameters that can be modified to provide a hydrogel system that  
14  
15 is specifically designed for a certain task.<sup>15,16</sup>  
16  
17

18  
19 One of the drawbacks to studying the diffusion of molecules within such systems is that they  
20  
21 either require large and expensive experimental setups (such as pulsed-field-gradient NMR<sup>17,18</sup>)  
22  
23 or rely on measurements of molecules as they diffuse across a gel boundary rather than within  
24  
25 the gel itself; details of such techniques can be found in the review by Pasch.<sup>19</sup>  
26  
27  
28  
29  
30

## 31 32 **EXPERIMENTAL**

### 33 34 **Gel Synthesis**

35  
36 A solution of 16.7% (by weight) methacrylic acid (MAA), 0.04% methylene-bis-acrylamide (MBA),  
37  
38 0.02% 2,2'-azobis (2-methylpropionamidine) dihydrochloride (AMPA), and the remainder HPLC  
39  
40 grade water was prepared in a sealed container and flooded with nitrogen for 30 min. The  
41  
42 solution was then placed in a vacuum oven at 65°C for 8h to ensure complete polymerisation of  
43  
44 the MAA.  
45  
46

47  
48 The resulting gel was then immersed in an excess of distilled water, such that the swelling of  
49  
50 the gel was not limited by the environmental volume nor quantity of solvent available. When  
51  
52 swelling had reached equilibrium (a constant mass of hydrated gel measured over 3 periods  
53  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 separated by at least 4h), pieces of the gel were subsequently removed for analysis in the  
4  
5 confocal microscope.  
6  
7

8 Dextran is a branched, neutral glucan widely used due to its biocompatibility and ease of  
9  
10 conjugation with labels and functional groups.<sup>20-22</sup> In order to introduce fluorescein-tagged  
11  
12 dextran molecules (FDEX) ( $M_w = 20$  kDa) into the gel, a nanomolar concentration in either pure  
13  
14 water, HCl acid, NaOH base, or salt solution (depending upon the experiment) was produced  
15  
16 and a small volume of gel was allowed to rest in each solution for sufficient time that the probe  
17  
18 penetrated the network and any new structural equilibrium had been reached. In order to  
19  
20 exclude the possibility of free-dye diffusion competing with that of the dextran during the  
21  
22 experiments, FDEX solutions were dialysed and measurements performed with this solution  
23  
24 were compared to those using undialysed solutions. No change in the diffusion coefficients was  
25  
26 found between dialysed and undialysed solutions.  
27  
28  
29  
30  
31  
32  
33  
34  
35

### 36 **Fluorescence Correlation Spectroscopy**

37  
38 FCS measurements were made using a ConfoCor2 FCS Module fitted to an LSM510 inverted  
39  
40 confocal microscope (Zeiss) and temperature was controlled using a Linkam heating stage  
41  
42 (Linkam Scientific Instruments Ltd, Surrey, UK) with TMS94 heat controller and LNP-1 nitrogen  
43  
44 flow control.  
45  
46  
47

48 Autocorrelation curves were fitted using the model described by Widengren and colleagues<sup>23</sup>  
49  
50 using pro Fit v6.1.8 (Quantum Soft, Switzerland). A Monte Carlo algorithm was initially used to  
51  
52 determine the starting parameters, and a Levenberg-Marquardt routine was used to find the  
53  
54 final best fit parameters. The autocorrelation function is given by  
55  
56  
57  
58  
59  
60

$$G(\tau) = 1 + \frac{1}{N_V} \left[ \frac{1}{1 + \frac{\tau}{\tau_D}} \frac{1}{\sqrt{1 + \frac{\tau}{\Gamma^2 \tau_D}}} \left[ 1 + \left( \frac{P_t}{1 - P_t} \right) \exp\left(-\frac{\tau}{\tau_t}\right) \right] \right] \quad (1)$$

and depends upon the number of fluorescently labelled molecules  $N_V$  within the detection volume, the diffusion time  $\tau_D$ , the confocal volume structure parameter  $\Gamma$ , and the fraction of excitations to the triplet state  $P_t$  that has a decay time  $\tau_t$ . An example of fitted autocorrelation functions for typical data obtained in this work are shown in figure 1. The Widengren equation determines the diffusion time of a polymer species (that is, the average time taken to cross the detection volume), and therefore it is necessary to convert this into a diffusion coefficient  $D$  for use with the theory described in the following sections. For this purpose, the diffusion coefficient can be obtained from the diffusion time using the relationship defined by Varma and colleagues<sup>24</sup>

$$D = \frac{a^2}{4\tau_D} \quad (2)$$

where  $a$  is the average confocal width (300 nm for the experimental setup used in this work).

## THEORY

### Zimm Model

The Stokes-Einstein relationship is the standard model for characterising the diffusion of a particle through a solvent. The diffusing species is treated as a sphere of radius  $R$  moving through a solvent of viscosity  $\eta_S$  at absolute temperature  $T$ , with a diffusion coefficient given by

$$D_{SE} = \frac{k_B T}{6\pi R \eta_S} \quad (3)$$

where  $k_B$  is the Boltzmann constant.

Whilst equation 3 has been shown to hold true for the diffusion of spherical or nearly-spherical

particles, a polymer chain requires modifications to the theory. By approximating the polymer chain as a series of balls connected by springs, and by considering the drag effect of each monomer ball on the surrounding solvent, the Zimm model<sup>25</sup> gives

$$D_{\text{Zimm}} = \frac{k_{\text{B}}T}{\zeta_{\text{Zimm}}} = \frac{0.196k_{\text{B}}T}{\eta_{\text{s}}bN^{\nu}} \quad (4)$$

where  $\zeta_{\text{Zimm}}$  is the polymer friction coefficient,  $b$  is the Kuhn monomer length and  $N$  is the number of monomers. The numerical prefactor is found for an ideal chain, and the polymer size is given by

$$R = bN^{\nu}$$

in which the exponent  $\nu$  depends upon the solvent quality.

### Swelling of Ionic Gels

The swelling of polymer networks comprising of charged constituents has been extensively studied. Flory<sup>26</sup> showed that, for dilute polymers, the swelling of an ionic gel can be expressed as

$$Q_{\text{m ion}}^{5/3} = Q_{\text{m non}}^{5/3} + \frac{V_0}{4\nu_e\nu_u^2} \frac{i^2}{S^*} \quad (5)$$

where  $Q_{\text{m ion}}$  and  $Q_{\text{m non}}$  are the equilibrium swelling ratios of the ionic and non-ionic gels respectively,  $V_0$  is unswollen gel volume,  $\nu_e$  the effective number of chains in the network,  $\nu_u$  is the molar volume of a structural unit,  $i$  is the degree of ionisation per unit of polymer, and  $S^*$  is the ionic strength of the surrounding medium. Assuming that no structural changes occur within the gel upon swelling, it follows that an ionic gel will swell more than a non-ionic gel, and subsequently collapse back towards the non-ionic system with increasing ionic concentration.

### RESULTS

### Temperature dependence of dextran diffusion

For FDEX in pure water, the diffusion coefficient increases with temperature as expected from equation 4, after allowing for changes in the viscosity of water based upon a theory first proposed<sup>27</sup> in 1971 and subsequently expanded to account for the atypical increase in fluidity at higher temperatures due to a disruption in the hydrogen bond network with increased thermal energy.<sup>28, 29</sup>

When the FDEX solution is within a hydrogel the overall viscosity changes. Figure 2 shows that the network introduces an additional effect that overwhelms the normal thermal motion of diffusion (equation 4) and causes a decrease in the diffusion coefficient at higher temperatures. By consideration of the Zimm equation it follows that the viscosity of the system must contribute more than the effect of any change in temperature.

The change in gel viscosity can be characterised by considering the correlation length of the network, known also as the average mesh size  $\xi$ . De Gennes and colleagues proposed<sup>30</sup> that the diffusion of a molecule of diameter  $d$  through a network with this average mesh size follows an exponential relationship

$$D(T) = D_0 \exp \left[ -\beta \left( \frac{d}{\xi(T)} \right)^\delta \right] \quad (6)$$

where  $D_0$  is the diffusion coefficient at a reference point,  $\delta = 2.5$  for cross-linked networks and  $\beta$  is a constant of order unity.

The mesh size is a length scale of the network that describes the average length between crosslink sites, given by

$$\xi = Q^{1/3} (C_n n)^{1/2} l \quad (7)$$

1  
2  
3 where  $Q$  is the equilibrium swelling ratio (determined experimentally through swelling  
4 measurements),  $C_n$  is a characteristic ratio of the polymer (14.6 for PMAA,<sup>31</sup>  $n$  is the number of  
5 crosslinks per chain, and  $l$  is the C-C bond length (0.154 nm). This expression for the mesh size  
6 can be substituted into equation 6 and the model applied to the data is presented in figure 2.  
7  
8  
9  
10  
11  
12  
13  
14  
15

16 Whilst the model described above provides a good description of the data in figure 2, it is  
17 instructive to consider alternative explanations of our results. If we consider the viscosity of the  
18 system to be described by changes in the concentration of PMAA when the hydrogel undergoes  
19 swelling, we also can fit the data above and is presented here. Assuming that no polymer is lost  
20 from the hydrogel, any decrease of solvent mass (per unit volume) increases the polymer  
21 concentration,  $c$ . Substitution of the Huggins equation<sup>25</sup>  
22  
23  
24  
25  
26  
27  
28  
29  
30

$$\frac{\eta - \eta_s}{\eta_s c} = [\eta] + k_H [\eta]^2 c + K \quad (8)$$

31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44

45 into equation 4 produces a relationship that describes the data in figure 2. The change in  
46 polymer mass concentration can be inferred from the mass loss of the swollen network as given  
47 in figure 3.  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

Our use of the Huggins equation, which requires the radius of gyration of the diffusing molecule to be comparable to the mesh size as shown in figure 4 indicates that changes in hydrogel structure can be characterised by a change in the net viscosity. It is worth noting that increasing temperature also affects diffusion coefficient due to a change in the excluded volume of the dextran molecule within the gel. Separation of the chain size variation from the structural changes of the gel requires a thermal characterisation of FDEX molecules, which is not covered

1  
2  
3 in this work but has been extensively covered for other molecules such as polystyrene.<sup>30</sup> The  
4  
5  
6 change in polymer radius can be inferred from the measurements of dextran diffusion in pure  
7  
8  
9 water (figure 2).

10  
11  
12  
13 It is found that this method also produces a good fit to the data in figure 2 although  
14  
15  
16 macroscopic swelling and how individual strands respond to changes in solvent concentration  
17  
18  
19 as shown in figure 3 must be considered. These are not affine networks described in this work;  
20  
21  
22 locally within the gel a great deal of molecular reorganisation can take place that is not  
23  
24  
25 mimicked on the macroscopic scale. This is a key consideration when attempting to produce  
26  
27  
28 hydrogels that control the release of absorbed molecules.

### 29 **Diffusion of dextran in salt solutions**

30  
31 The effect of added salt on PMAA swelling is similar to the effect of pH in that  $Q$  can be made to  
32  
33  
34 change by over two orders of magnitude, although in this instance the most pronounced effect  
35  
36  
37 is at very dilute salt concentrations (figure 5).

38  
39 As expected from inverse relation between swelling ratio and ion concentration given in  
40  
41  
42 equation 5, the diffusion coefficient drops rapidly over the dilute salt concentration range (up  
43  
44  
45 to  $0.5 \text{ M L}^{-1}$ ) as the swelling of the hydrogel decreases towards the non-ionic system for  
46  
47  
48 increasing salt concentrations; what may be less obvious is that the gel-free system also follows  
49  
50  
51 this trend of a sharp decrease in  $D$  at low salt concentrations, although the overall decrease is  
52  
53  
54 less in magnitude when compared to the gel system. Reconsideration of equation 4 suggests  
55  
56  
57 that the dextran diffusion in the absence of gel may be affected by added salt due to an  
58  
59  
60 increase in viscosity or radius of hydration caused by a change in the solvent quality exponent

v.

The ionic concentration dependence on viscosity has already been well investigated; Falkenhagen<sup>33</sup> derived an expression for symmetric electrolytes which was subsequently extended by Onsager and Fuoss to include asymmetric electrolytes.<sup>34</sup> Both derivations lead to the ionic concentration dependent viscosity relationship:

$$\eta(c) = \eta_0 + \frac{\kappa_D \zeta_0}{480\pi} \quad (9)$$

where  $\eta_0$  is viscosity of pure solvent,  $\zeta_0$  is the friction coefficient of an ion in a solution of infinitesimal concentration, and  $\kappa_D$  is the inverse Debye screening length which scales as the square root of the ion concentration and linearly with ion valency. It should be noted however that the Debye screening length provides a poor model of the charge behaviour around a polymer chain especially in low salt concentrations: it has been demonstrated that the distribution of charges around a polymer chain is very different to that around free ions for which the Debye length model was originally derived.<sup>35</sup>

Another point of interest is that at the lowest salt concentrations, the diffusion coefficient of the dextran in the gel,  $D_{\text{gel}}$ , is greater than that in the absence of a gel,  $D_{\text{water}}$  (figure 7). This means that whilst there is extensive swelling of the gel (as predicted in equation 5 for low concentration), there is an additional effect causing the dextran molecules to diffuse faster than in pure water; it is expected that, for maximum swelling, the system will tend towards (but not exceed) pure water conditions – exceeding this would imply there is an additional component of the salt-gel-FDEX system that is causing the dextran to diffuse faster through the gel than in the equivalent situation without a gel. A possible explanation for this increase in dextran diffusion lies with the interaction between any additional charge and the dextran, either from

1  
2  
3 the salt solution itself or from any free charge released by the network. Indeed, PMAA has been  
4  
5 found to ionise in the presence of salt,<sup>36-38</sup> which may provide some justification for the non-  
6  
7 monotonic swelling behaviour of the gel, but does not suitably explain why the diffusion in the  
8  
9 gel exceeds that in salt solution, nor the non-monotonic behaviour of dextran in the gel-free  
10  
11 system.  
12  
13

14  
15 The effect of electrolyte concentration on the radius of dextran molecules is not yet  
16  
17 understood; it is not unreasonable to assume that the conformation of the dextran molecule  
18  
19 will also affect the viscosity term described above, and therefore a theoretical model to fit the  
20  
21 data presented in figures 6 and 7 would require a characterisation of the effect of ionic  
22  
23 concentration on the labelled dextran molecule.  
24  
25

26  
27 Finally, the presence of salt may affect or dissociate the fluorescein probe from the dextran  
28  
29 molecule, but this is not considered to be a contributing factor to the observation that  $D_{\text{gel}} >$   
30  
31  $D_{\text{water}}$  as the diffusion through the gel is normalised to the corresponding diffusion in salt  
32  
33 solution and consequently any effect on the dye due to the salt will be present in both the gel  
34  
35 and gel-free diffusion measurements.  
36  
37  
38  
39

#### 40 41 **Variation of pH**

42  
43 Any change in hydrogel structure can again be determined by measuring  $Q$  for all pH values,  
44  
45 which is shown in figure 8. Unlike the effect of temperature upon  $Q$  (figure 3) that shows very  
46  
47 little swelling ( $Q \approx 0$ ), the swelling ratio in alkaline pH can be several hundreds indicating a  
48  
49 significant degree of swelling.  
50  
51

52  
53 The introduction of acid or alkali has an effect not only upon the structure of the charged  
54  
55 hydrogel but also on the dextran molecule itself. Figure 9 shows that the diffusion of dextran in  
56  
57  
58  
59  
60

1  
2  
3 a gel-free environment is dependent upon the charge concentration. This result may be  
4  
5 somewhat surprising considering that previous work<sup>39</sup> shows that the radius of hydration  $R_H$  for  
6  
7 dextran in water and 0.5 mM NaOH solution follows the same relationship for different  
8  
9 molecular weights of polymer. These authors also found that the intrinsic viscosity of dextran is  
10  
11 comparable when in water and NaOH, which is in agreement with figure 9 in the basic regime.  
12  
13 The apparent anomaly between pH 6 and 7 can easily be explained when the method used to  
14  
15 prepare the pH solutions is examined. HPLC-grade water of pH 7 was initially used, into which  
16  
17 different masses of HCl or NaOH were dissolved into the water to reach the required pH. Whilst  
18  
19 the concentration of hydrogen ions was measured, the additional atoms from the original  
20  
21 solutes form a solution of additional ions that will influence the entire system in a similar  
22  
23 manner to that expressed in the previous section.  
24  
25

26  
27 Even with this consideration in mind, it is now possible to examine the overall effect of charge  
28  
29 upon the structure of the gel itself. As figure 9 shows the diffusion relationship between 'free'  
30  
31 dextran and that confined in PMAA gel follow significantly different trends. More acidic  
32  
33 environments result in a decrease in the diffusion coefficient, and given that the diffusion of  
34  
35 'free' dextran is maximised at these low pH values we can ascertain that the hydrogel itself has  
36  
37 undergone a structural change, collapsing and thus impeding the motion of dextran, in  
38  
39 agreement with the results shown in figure 8. As with the temperature effect, this pH-induced  
40  
41 decrease in the diffusion coefficient can be interpreted as an increase in the hydrogel  
42  
43 concentration as described by equation 8. Such a collapse has already been documented,<sup>40</sup> and  
44  
45 indeed put into practical use.<sup>2,4,41</sup> Dextran diffusion in magnetite impregnated gels (ferrogels)  
46  
47 has been shown to be controlled by the mesh size of the hydrogel,<sup>42</sup> but in those experiments  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

1  
2  
3 the swelling ratio changed by a factor of two compared to approximately 350 in these  
4  
5  
6 experiments.

## 8 CONCLUSIONS

10 It has been shown that FCS is a viable technique for measuring the diffusion of tagged  
11  
12 molecules within a hydrogel. The expected Zimm diffusion of FDEX was observed in pure water  
13  
14 over a range of temperatures, and that with suitable modification using the Huggins equation it  
15  
16 is possible to model the change in viscosity of a hydrogel and related this to the change in  
17  
18 diffusion of the probe molecule.  
19

21 Furthermore we have shown that PMAA undergoes structural changes with the addition of  
22  
23 charge into the surrounding solvent. Interestingly, the addition of salt into the hydrogel will  
24  
25 induce swelling of the network as predicted by Flory,<sup>26</sup> but rather than bringing the in-gel  
26  
27 diffusion closer to that in pure water we have shown that the diffusion exceeds that in pure  
28  
29 water. Whether this is due to changes in the network itself or a consequence of salt-probe  
30  
31 interactions is yet to be determined.  
32  
33  
34  
35  
36

## 38 ACKNOWLEDGEMENTS

39 We thank the EPSRC and the Cultural Bureau of the Embassy of the Kingdom of Saudi Arabia for  
40  
41 financial support. Helpful discussions from Professor Anthony Ryan, Professor John Torkelson,  
42  
43 and Dr Drew Tarmey are acknowledged.  
44  
45  
46

## 48 REFERENCES AND NOTES

- 49 1. Doi, M; Edwards, S. F. In *The Theory Of Polymer Dynamics*; Oxford University Press,  
50 **1992**.
- 51 2. Kozlovskaya, V.; Kharlampieva, E.; Mansfield, M. L.; Sukhishvili, S. A. *Chem. Mater.* **2006**,  
52 *18*, 328–336.
- 53 3. Katz, J. S.; Burdick, J. A. *Wiley Interdisp. Rev. Nanomed. Nanobiotechnol.* **2009**, *1*, 128–  
54 139.
- 55 4. La Spina, R.; Tomlinson, M. R.; Ruiz Perez, L.; Chiche, A.; Langridge, S; Geoghegan, M.  
56  
57  
58  
59  
60

- 1  
2  
3  
4  
5  
6  
7  
8  
9  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60
- Angew. Chem. Int. Ed.*, **2007**, *46*, 6460–6463.
5. Endo, T.; Ikeda, R.; Yanagida, Y.; Hatsuzawa, T. *Anal. Chim. Acta*, **2008**, *611*, 205–211.
6. May, C.; Nazar, L.; Jones, L.; Simpson, T. *Contact Lens and Anterior Eye*. **2002**, *25*, 147–156.
7. Colby, R. H.; Rubinstein, M.; Gillmor, J. R.; Mourey, T. H. *Macromolecules*, **1992**, *25*, 7180–7187.
8. Rubinstein, M.; Colby, R. H. *Macromolecules*, **1994**, *27*, 3184–3190.
9. Gianelli, M.; Beines, P. W.; Roskamp, R. F.; Koynov, K.; Fytas, G.; Knoll, W. *J. Phys. Chem. C*, **2007**, *111*, 13205–13211.
10. Jia, P.; Yang, Q.; Gong, Y.; Zhao, J. *J. Chem. Phys.*, **2012**, *136*, 084904–084910
11. Gianelli, M.; Roskamp, R. F.; Jonas, U.; Loppinet, B.; Fytas, G.; Knoll, W. *Soft Matter*, **2008**, *4*, 1443–1447.
12. Venter, J. P.; Muller, D. J.; du Plessis, J.; Goosen, C. *Eur. J. Pharm. Sci.* **2001**, *13*, 169–177.
13. Siepman, J.; Kranz, H.; Bodmeier, R.; Peppas, N. A. *Pharm. Res.* **1999**, *16*, 1748–1756.
14. Zalfen, A. M.; Nizet, D.; Jerome, C.; Jerome, R.; Frankenne, F.; Foidart, J. M.; Maquet, V.; Lecomte, F.; Hubert, P.; Evrard, B. *Acta Biomaterialia*. **2008**, *4*, 1788–1796.
15. Bernik, D. L.; Zubiri, D.; Monge, M. E.; Negri, R. M. *Colloids and Surfaces A: Physicochem. Eng. Aspects*, **2006**, *273*, 165–173.
16. Barbeiro, S.; Ferreira, J. A. *Comput. Meth. Appl. Mech. Eng.* **2009**, *198*, 2078–2086.
17. Gao, P.; Fagerness, P. E. *Pharm. Res.* **1995**, *12*, 955–964.
18. Gao, P.; Meury, R. H. *J. Pharm. Sci.* **1996**, *85*, 725–731.
19. Pasch, H. *Adv. Polym. Sci.* **2000**, *150*, 1–66.
20. Ansorge, W.; Pepperkok, R. *J. Biochem. Biophys. Methods* **1988**, *16*, 283.
21. Graziadei, L.; Burfeind, P.; Bar-Sagi, D. *Anal. Biochem.* **1991**, *194*, 198.
22. Olson, D. J.; Christian, J. L.; Moon, R. T. *Science*, **1991**, *252*, 1173.
23. Widengren, J.; Mets, Ü.; Rigler, R. *J. Phys. Chem.* **1995**, *99*, 13368–13379.
24. Varma, B. K.; Fujita, Y.; Takahashi, M.; Nose, T. *J. Polym. Sci. B; Polym. Phys.* **1984**, *22*, 1781–1798.
25. Rubinstein, M.; Colby, R. H. In *Polymer Physics*; Oxford University Press, **2006**.
26. Flory, P. J. In *Principles of Polymer Chemistry* 7<sup>th</sup> edition; Cornell University Press, **1971**.
27. Zwolinski, B. J.; Eicher, L. D. *J. Phys. Chem.* **1971**, *75*, 2016–2024.
28. Stanley, H. E.; Teixeira, J. *J. Chem. Phys.* **1980**, *73*, 3404–3422.
29. Errington, J. R.; Debenedetti, P. G. *Nature*. **2001**, *409*, 318–320.
30. Langevin, D.; Rondelez, F. *Polymer*. **1978**, *19*, 875–882.
31. Bell, C. L.; Peppas, N. A. *J. Control. Release*. **1996**, *39*, 201–207.
32. Berry, G. C. *J. Chem. Phys.* **1966**, *44*, 4550–4564.
33. Falkenhagen, H. *Phys. Z.* **1931**, *32*, 745–764.
34. Onsager, L.; Fuoss, R. M. *J. Phys. Chem.* **1932**, *36*, 2689–2778.
35. Valteau, J. P. *Chem. Phys.* **1989**, *129*, 163–174.
36. Oth, A.; Doty, P. *J. Phys. Chem.* **1952**, *56*, 43–50.
37. Koňák, Č.; Bansil, R. *Polymer*, **1989**, *30*, 677–680.
38. Pohlmeier, A.; Haber-Pohlmeier, S. *J. Colloid Interface Sci.* **2004**, *273*, 369–380.
39. Ioan, C. E.; Aberle, T.; Burchard, W. *Macromolecules*. **2000**, *33*, 5730–5739.
40. Tanaka, T.; Fillmore, D.; Sun, S.-T.; Nishio, I.; Swislow, G.; Shah, A. *Phys. Rev. Lett.* **1980**,

45, 1636–1639.

41. Sukhishvili, S.; Granick, S. *Macromolecules*. **2002**, *35*, 301.

42. Al-Baradi, A. M.; Mykhaylyk, O. O.; Blythe, H. J.; Geoghegan, M. *J. Chem. Phys.* **2011**, *134*, 094901.

FIGURE 1 Autocorrelation functions for one dataset of the diffusion of FDEX in pure water at different temperatures. The solid lines represent the fit of equation 1 to the data.

FIGURE 2 Temperature effect on diffusion of free dextran (open circles) and dextran in PMAA gel (crosses). The dotted line presents a fit using the Zimm model (equation 4) with temperature-dependent viscosity,<sup>19-21</sup> whereas the dashed line also utilises a Huggins relationship for viscosity (equation 6) and the solid line presents a fit of the mesh-size relation given by equation 7. The same data are shown in the lower figure but expanded for the gel data and two fitted models.

FIGURE 3 Swelling ratio (given as a percentage change in mass) and mesh size of swollen PMAA at different temperatures. Note that at the highest temperature the mass decreases only by 1% compared to the initial value at  $T = 283$  K.  $Q$  values were determined experimentally whereas mesh size was determined from FCS measurements and modelled via equation 8. The data have been fitted to

empirical functions  $Q = A \left( \frac{1}{1 + \exp\left(\frac{T - T_0}{K}\right)} \right)$  and  $\xi = B(T + T_1)^\sigma + C$  to provide a guide to the eye,

where  $A$ ,  $T_0$ ,  $K$ ,  $B$ ,  $T_1$ ,  $\sigma$ , and  $C$  are empirical constants.

FIGURE 4 Comparison of the mesh size (determined using equation 8) and the radius of gyration (determined using equation 3 from the diffusion of dextran through pure water) plotted as a function of temperature. Note that these length scales are comparable over all temperature ranges supporting the use of the Huggins theory of viscosity (equation 8) in modelling diffusion through PMAA hydrogels.

FIGURE 5 Swelling ratio of PMAA for concentrations of NaCl and CaCl<sub>2</sub> over the concentration range used in the present study. Whilst the maximum swelling ratio is not as high as that for which pH was varied (figure 8) one can still see that very low concentrations of salt induce a dramatic structural change in the gel resulting in up to a factor of 60 increase in mass. The inset figure provides the same data on a log-linear plot but does not include the value for zero salt concentration.

FIGURE 6 Diffusion of FDEX in NaCl and CaCl<sub>2</sub> salt solutions, both with and without PMAA gel. The sharp increase in diffusion at low salt concentrations is evident in all cases, with the system without gel reaching a plateau at higher concentrations. The lower and unequal diffusion coefficient within the gels arises from different degrees of collapse due to the different valencies between the salts.

FIGURE 7 Normalised diffusion coefficients in NaCl and CaCl<sub>2</sub>.  $D_N$  is the normalised diffusion coefficient, and is defined as  $D_N = D_{\text{gel}} / D_{\text{water}}$ . The dashed line at  $D_N = 1$  represents a diffusion coefficient equal to that in water at the same salt concentration. Values of  $D_N > 1$  indicate that the hydrogel is causing the diffusion of dextran to be greater than when it is not present.

FIGURE 8 Swelling ratio of swollen PMAA for the pH range. The two plots show the time given for the

1  
2  
3 hydrogel to test for equilibrium conditions. The effect of pH is the most pronounced, when compared to  
4 the other systems studied within this work, with a factor of 350 change in mass.  
5

6  
7 FIGURE 9 Diffusion of FDEX in different pH environments. It is worth noting that both the acid and alkali  
8 regions were produced by the addition of HCl and NaOH respectively and thus also include contributions  
9 from the additional ions ( $\text{Cl}^-$  and  $\text{Na}^+$ ).  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60

For Peer Review

**GRAPHICAL ABSTRACT**

## AUTHOR NAMES

Ateyyah M. AL-Baradi, Matthew Mears, Richard A. L. Jones, and Mark Geoghegan

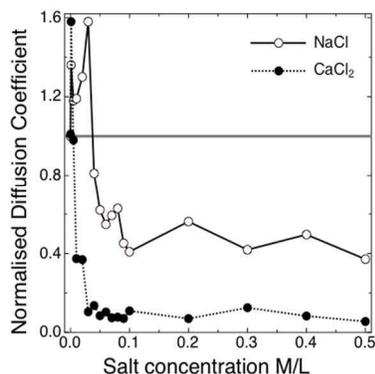
## TITLE

The diffusion of dextran within poly(methacrylic acid) hydrogels

TEXT ((up to 75 words, present tense, written for a non-specialist, see recent issue for examples))

Responsive hydrogels allow the diffusive properties of molecules within their network to be controlled. Here we show that poly(methacrylic acid) hydrogels respond to temperature, salt, and pH to alter the diffusion coefficient of fluorescently labelled dextran. The presence of ions in the surrounding environment cause significant swelling of the gel, with low ion concentrations providing the largest effect on both the mesh size of the gel and the diffusion of the dextran molecules.

GRAPHICAL ABSTRACT FIGURE ((50 mm wide by 50 mm high))



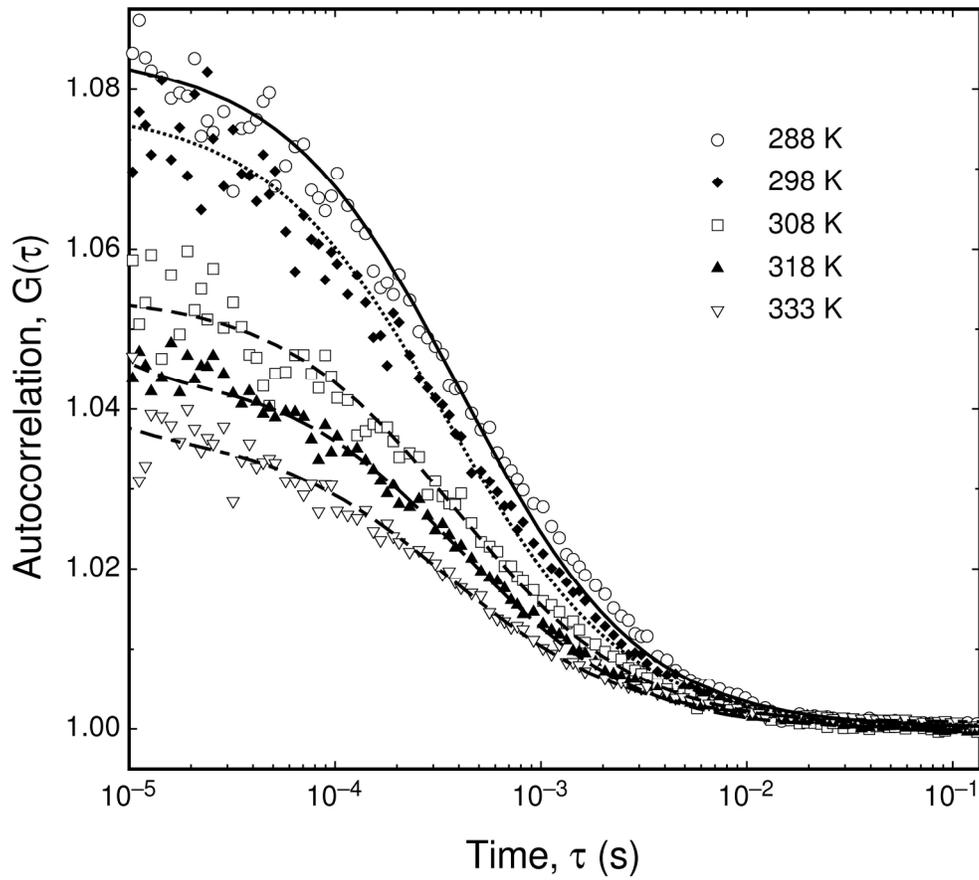


FIGURE 1 Autocorrelation functions for one dataset of the diffusion of FDEX in pure water at different temperatures. The solid lines represent the fit of equation 1 to the data.

200x181mm (300 x 300 DPI)

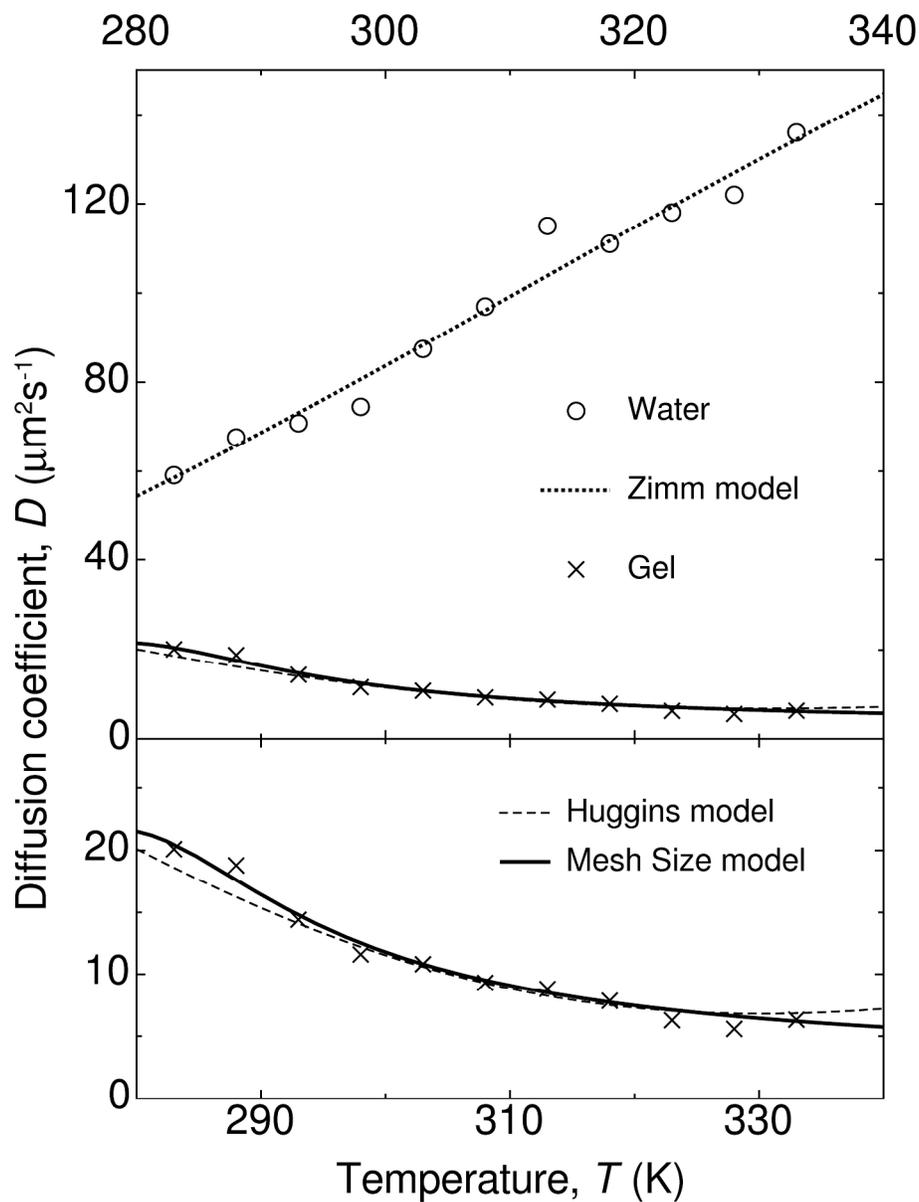


FIGURE 2 Temperature effect on diffusion of free dextran (open circles) and dextran in PMAA gel (crosses). The dotted line presents a fit using the Zimm model (equation 4) with temperature-dependent viscosity,<sup>19-21</sup> whereas the dashed line also utilises a Huggins relationship for viscosity (equation 6) and the solid line presents a fit of the mesh-size relation given by equation 7. The same data are shown in the lower figure but expanded for the gel data and two fitted models.

242x319mm (300 x 300 DPI)

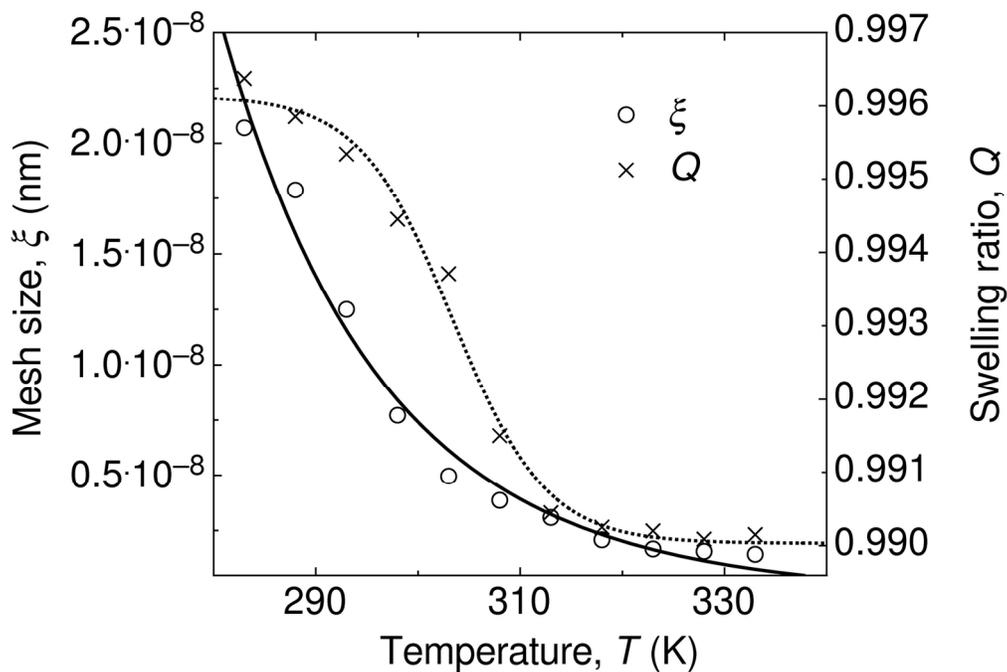


FIGURE 3 Swelling ratio (given as a percentage change in mass) and mesh size of swollen PMAA at different temperatures. Note that at the highest temperature the mass decreases only by 1% compared to the initial value at  $T = 283$  K.  $Q$  values were determined experimentally whereas mesh size was determined from FCS measurements and modelled via equation 8. The data have been fitted to empirical functions and to provide a guide to the eye, where  $A$ ,  $T_0$ ,  $K$ ,  $B$ ,  $T_1$ ,  $\sigma$ , and  $C$  are empirical constants.

149x100mm (300 x 300 DPI)

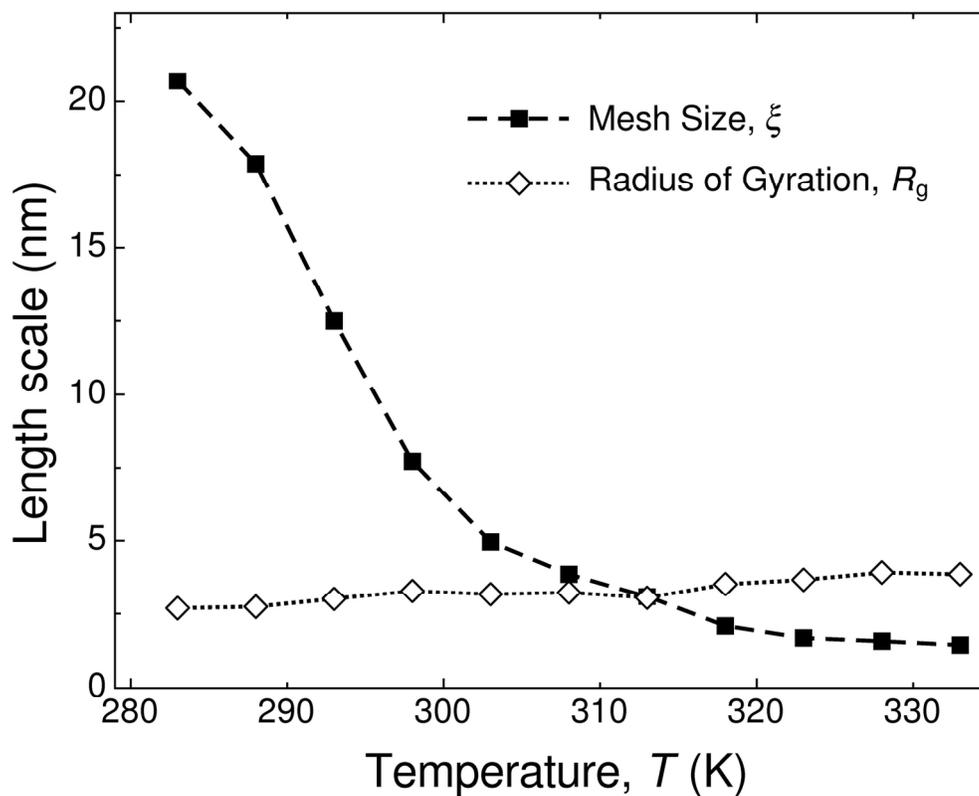


FIGURE 4 Comparison of the mesh size (determined using equation 8) and the radius of gyration (determined using equation 3 from the diffusion of dextran through pure water) plotted as a function of temperature. Note that these length scales are comparable over all temperature ranges supporting the use of the Huggins theory of viscosity (equation 8) in modelling diffusion through PMAA hydrogels.

152x125mm (300 x 300 DPI)

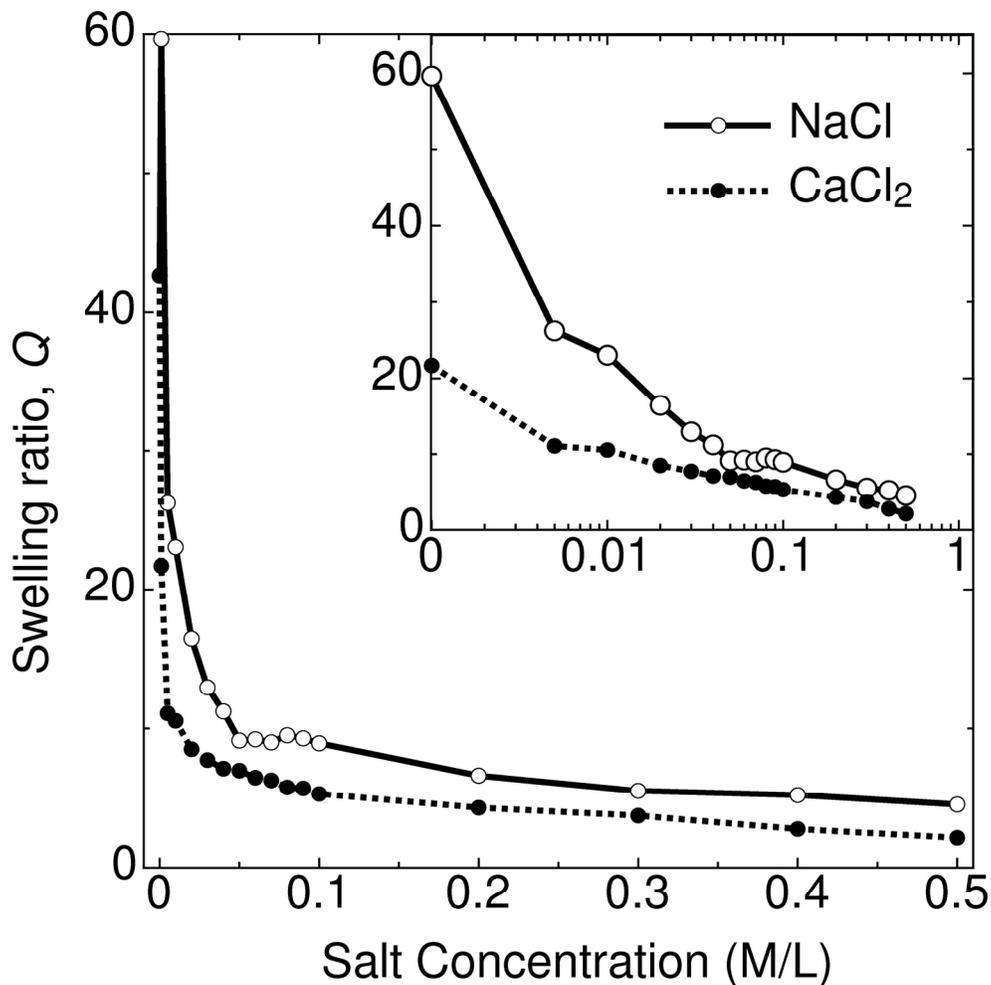


FIGURE 5 Swelling ratio of PMAA for concentrations of NaCl and CaCl<sub>2</sub> over the concentration range used in the present study. Whilst the maximum swelling ratio is not as high as that for which pH was varied (figure 8) one can still see that very low concentrations of salt induce a dramatic structural change in the gel resulting in up to a factor of 60 increase in mass. The inset figure provides the same data on a log-linear plot but does not include the value for zero salt concentration.

177x177mm (300 x 300 DPI)

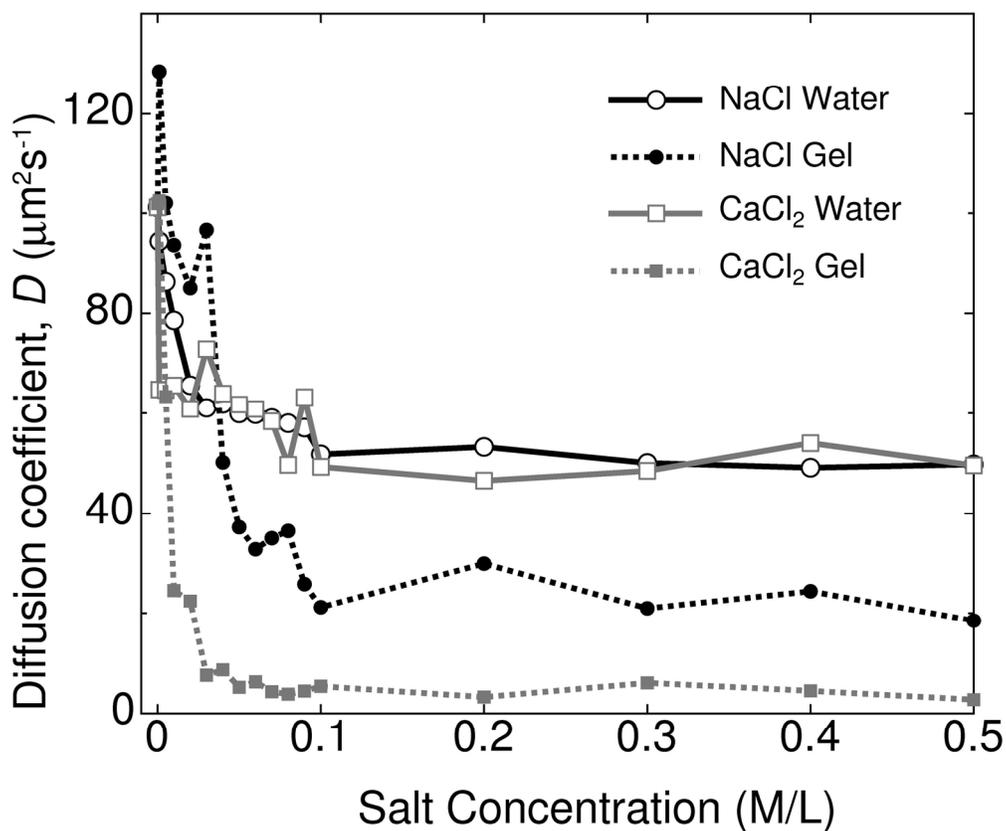


FIGURE 6 Diffusion of FDEX in NaCl and CaCl<sub>2</sub> salt solutions, both with and without PMAA gel. The sharp increase in diffusion at low salt concentrations is evident in all cases, with the system without gel reaching a plateau at higher concentrations. The lower and unequal diffusion coefficient within the gels arises from different degrees of collapse due to the different valencies between the salts.

163x140mm (300 x 300 DPI)

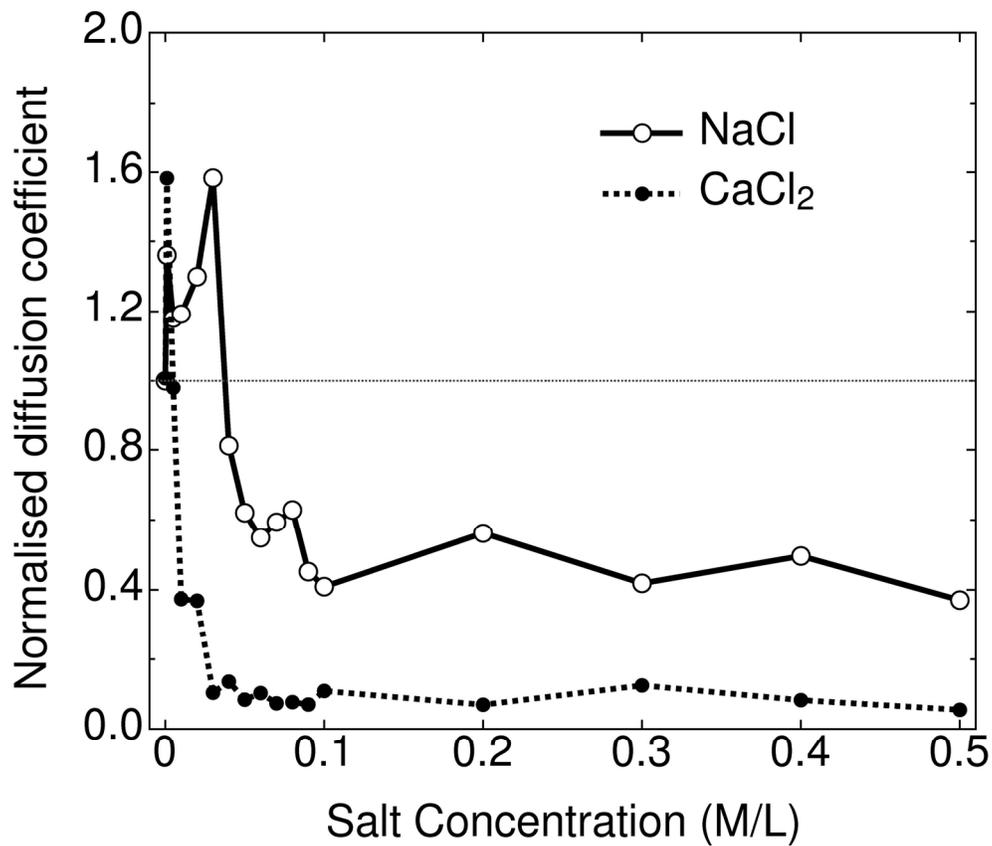


FIGURE 7 Normalised diffusion coefficients in NaCl and CaCl<sub>2</sub>. DN is the normalised diffusion coefficient, and is defined as  $DN = D_{gel} / D_{water}$ . The dashed line at  $DN = 1$  represents a diffusion coefficient equal to that in water at the same salt concentration. Values of  $DN > 1$  indicate that the hydrogel is causing the diffusion of dextran to be greater than when it is not present.

163x139mm (300 x 300 DPI)

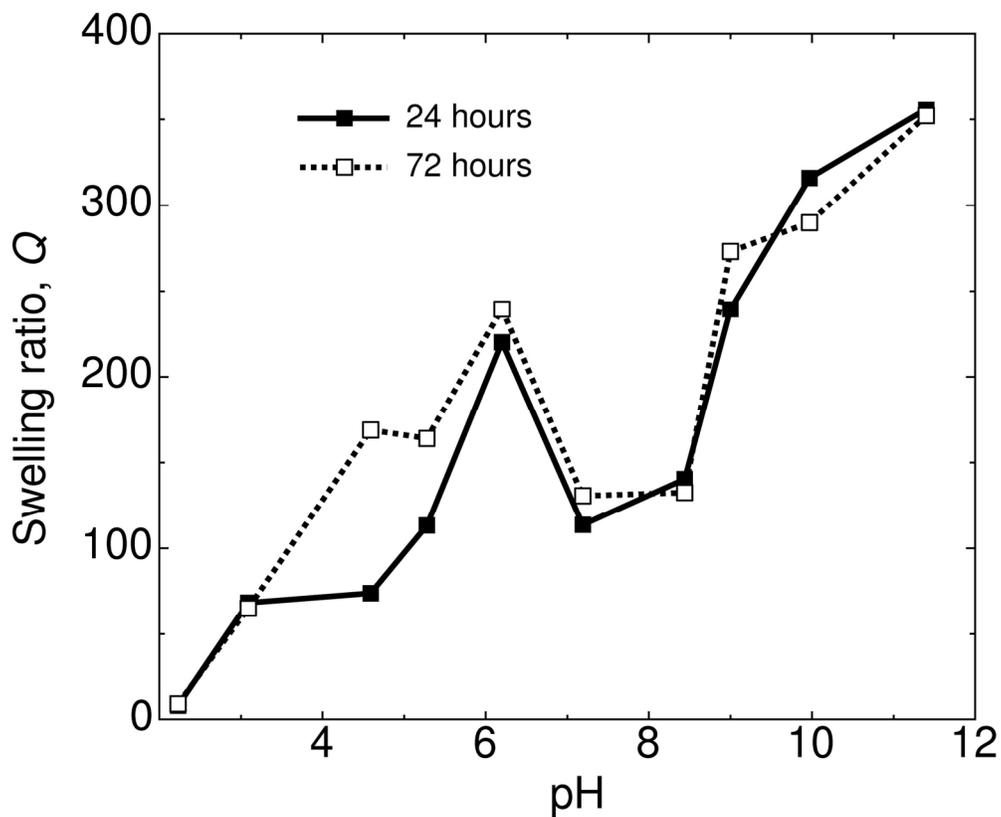


FIGURE 8 Swelling ratio of swollen PMAA for the pH range. The two plots show the time given for the hydrogel to test for equilibrium conditions. The effect of pH is the most pronounced, when compared to the other systems studied within this work, with a factor of 350 change in mass.  
160x131mm (300 x 300 DPI)

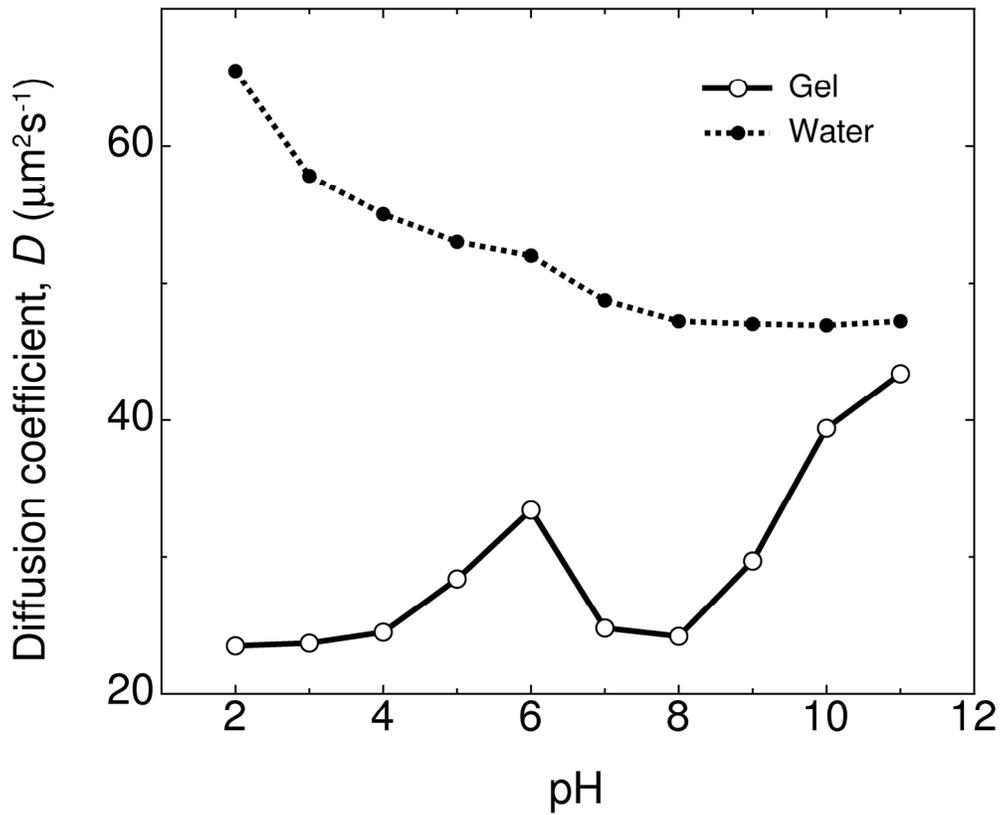


FIGURE 9 Diffusion of FDEX in different pH environments. It is worth noting that both the acid and alkali regions were produced by the addition of HCl and NaOH respectively and thus also include contributions from the additional ions ( $\text{Cl}^-$  and  $\text{Na}^+$ ).  
164x138mm (300 x 300 DPI)