

Tailoring ultrathin magnetic multilayers at terraced topologically insulating interfaces for perpendicularly magnetized domains

Benjamin A. Brereton^{1,2,3}, Soumyarup Hait¹, Ahmet Yagmur¹, Christy J. Kinane², Francesco Maccherozzi³, Michele Conroy⁴, Satoshi Sasaki¹, Thomas A. Moore¹, Sarnjeet S. Dhesi³, Sean Langridge², and Christopher H. Marrows^{1,*}

¹*School of Physics and Astronomy, University of Leeds, Leeds LS2 9JT, United Kingdom*

²*ISIS Neutron and Muon Source, STFC Rutherford Appleton Laboratory, Harwell Science and Innovation Campus, Chilton, Oxfordshire OX11 0QX, United Kingdom*

³*Diamond Light Source, Harwell Science and Innovation Campus, Chilton, Oxfordshire OX11 0DE, United Kingdom*

⁴*Department of Materials, Imperial College London, Exhibition Road, London SW7 2AZ, United Kingdom*



(Received 10 February 2026; accepted 20 May 2026; published 24 June 2026)

Topological insulators and skyrmion-hosting, chiral magnetic multilayers are two well-explored areas of modern condensed matter physics, each offering unique advantages for spintronics applications. In this paper, we demonstrate the optimization process for the growth of a $\text{Bi}_2\text{Se}_3/\text{buffer}/[\text{Pt}/\text{CoB}/\text{Ru}]_{\times N}$ heterostructure that combines these two material classes: the Bi_2Se_3 epilayer was grown by molecular beam epitaxy before transfer under ultrahigh vacuum to a separate growth chamber where the polycrystalline metallic multilayer was sputter deposited. The structure of the samples was characterized by cuffed x-ray and polarized neutron reflectometry measurements and scanning transmission electron microscopy. Polarized neutron models and standard magnetometry show that a buffer layer exceeding a critical thickness is required to obtain the desired uniform, perpendicular magnetic anisotropy in every magnetic layer in the multilayer. Samples with both Ta and Mo buffers were used requiring thicknesses of 1.5 and 0.9 nm, respectively. In minimizing the Bi_2Se_3 terracing, buffered samples yield well-defined, out-of-plane, magnetic domains suitable for spin-orbit torque-induced manipulation as determined by x-ray photoemission electron microscopy.

DOI: [10.1103/hp66-tldh](https://doi.org/10.1103/hp66-tldh)

I. INTRODUCTION

Topological insulators (TIs) have long been of interest in next-generation spintronic devices because of their highly efficient charge-to-spin conversion, analogous to the spin Hall angles of heavy metals: $\theta_{\text{SH}} = J_s/J_c$. In such systems θ_{SH} can be much greater than unity [1,2], arising from spin-momentum-locked surface states. Such devices utilize the large spin-orbit torques (SOTs) generated by this effect, and often are based on switching thin ferromagnetic films with perpendicular magnetic anisotropy (PMA) for low-current magnetic random access memory applications [3–6]. Magnetic multilayers (MMLs) comprising ultrathin layers of magnetic and heavy metals are known to host multitudes of exotic, chiral spin textures in systems of broken inversion symmetry due to a nonvanishing interfacial Dzyaloshinskii-Moriya interaction (DMI): most notably Néel-type magnetic skyrmions [7–12]. Skyrmions are topologically nontrivial with a resultant quantized Berry phase for transport electrons:

their topology grants them relative stability against perturbations and admits SOT-driven motion and manipulation from spin currents such as those generated by the spin-Hall effect (SHE) in heavy metals [9,13–17]. These properties have not gone overlooked for potential applications in energy efficient data storage and computation applications, most recently in reservoir computing [18–21]. TI/chiral magnetic multilayer heterostructures present an opportunity to utilize the strengths of both materials to further reduce necessary current densities and enhance energy efficiency. Often chiral multilayers rely on the interfacial hybridization between Co and Pt. This yields both strong PMA and DMI effects. The total effective anisotropy, K_{eff} , can be expressed as a function of ferromagnet (FM) thickness, t_{FM} , as

$$K_{\text{eff}} = \frac{K_i}{t_{\text{FM}}} + \frac{1}{2}\mu_0 M_s H_D, \quad (1)$$

where K_i is the interfacial PMA term, M_s the saturation magnetization, and H_D the demagnetization field. In ultrathin films the demagnetization field arising from shape anisotropy, $H_D = -M_s$. The quality of the Pt/Co interface plays a key role in determining the magnitudes of the DMI vector and K_i [22–24], which increasingly dominates over the other volumetric contributions at low t_{FM} . The DMI contribution competes against a standard Heisenberg-like exchange interaction producing noncollinear and incommensurate Néel-type spin canting over a characteristic length scale, Δ . Previous PMA TI/FM

*Contact author: c.h.marrows@leeds.ac.uk

Published by the American Physical Society under the terms of the Creative Commons Attribution 4.0 International license. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI.

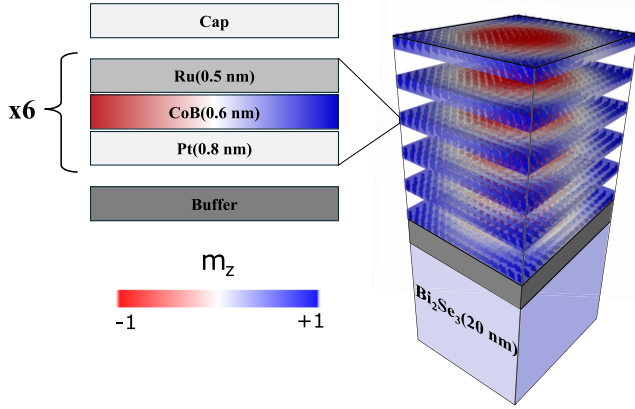


FIG. 1. Schematic representing standard sample structure grown on Bi_2Se_3 with illustrated skyrmion spin texture in the CoB layers where m_z lies along the out-of-plane axis.

switching studies [25,26] note their use of buffer layers deposited on the TI surface before the MML is grown for out-of-plane (OOP) switching geometries, and, when buffer layers are removed, FM easy axes tend to lie in-plane (IP). Inserting a buffer layer over a TI will introduce spin-flip scattering causing the spin-signal to be lost over a material dependent spin-diffusion length l_{SD} according to $J_s(t) = J_{s,0} e^{-t/l_{\text{SD}}}$; this is in addition to the additional scattering introduced as a result of introducing a second interface [27]. This problem is particularly acute for heavy, refractory metals such as Ta, which possess large spin-orbit coupling (SOC) and hence spin diffusion lengths are short, $l_{\text{SD}}^{\text{Ta}} \sim 2$ nm [28], yet are the standard materials for buffers for this type of MML stack. Typical TIs belonging to the commonly used Bi_2X_3 class of materials are known to form in quintuple layers weakly bound by van der Waals forces [29–31]. This layering forms distinct triangular terracing at the surface of an TI surface, which raises concerns for magnetic pinning sites in TI/MML heterostructures. As such this study aims to tailor a buffer to ensure a high-quality multilayer structure and optimize for magnetic domains decoupled from the TI surface topography while minimizing the loss of spin current between the TI and MML.

II. HETEROSTRUCTURE GROWTH

The heterostructures that we studied combine 20 nm epitaxial layers of Bi_2Se_3 , a well-known TI [31–35], and a $[\text{Pt}(0.8 \text{ nm})/\text{CoB}(0.6 \text{ nm})/\text{Ru}(0.5 \text{ nm})] \times 6$ MML on a Ta or Mo buffer layer as shown in Fig. 1. The TI was grown by molecular beam epitaxy (MBE), and the multilayer was grown by dc magnetron sputtering at room temperature. The Bi_2Se_3 (20 nm) epilayers were prepared in a solid-source molecular beam epitaxy (MBE) system that operated at a base pressure of 7×10^{-10} mbar. The films were grown by codepositing evaporated bismuth and selenium, both supplied from standard dual-filament Knudsen cells, on c-plane (0001) sapphire substrates. To effectively suppress the formation of selenium vacancies, the growth was carried out under Se-rich conditions, with the selenium flux at least 20 times higher than the bismuth flux. Two different samples were prepared.

The first sample, characterized by narrow surface terraces ~ 50 nm wide, was grown at a temperature of 235°C . The second, wide-terrace sample was deposited using a two-step method, with a nucleation layer of ~ 2 nm grown at 130°C and the remainder of the layer grown at 300°C increasing average terrace width to ~ 100 nm. Connected MBE and sputtering chambers allowed the entire heterostructure to be deposited without breaking ultra-high vacuum, thereby minimizing degradation of the TI surface and keeping the interface between the materials as clean as possible. After transfer to the sputtering chamber, any buffer layer, followed by the MML with a 4.0 nm Pt capping layer, was deposited at room temperature. Here we have studied the buffer layer materials Ta and Mo, a $5d$ and $4d$ refractory metal, respectively. On reference MML samples grown on thermally oxidized $\text{Si}(100)/\text{SiO}_2(100 \text{ nm})$ substrates and 5.0 nm Ta buffers, PMA could be reliably obtained down to a minimum nominal Pt thickness of 0.8 nm; however, the PMA strength is noticeably reduced when reducing Pt thickness beneath 1 nm [27,36–38]. To maximize the effects of interfacial PMA against shape anisotropy, the FM CoB layer was thinned to the minimum thickness to be reliably continuous, which was found here to be 0.6 nm. This optimization yields a multilayer structure with strong PMA through minimized use of heavy metals between the FM and TI layers. Samples which were not required for imaging were grown with an additional 4 nm Pt cap to prevent oxidation.

III. MULTILAYER STRUCTURE ON TERRACED Bi_2Se_3

Epilayers of Bi_2Se_3 grow with a terraced surface as a result of the lateral expansion and step flow growth of quintuple layers from defects during MBE growth. The terracing represents a unique challenge for Bi_2Se_3 -based devices featuring such thin layers since a single quintuple layer, determined to be 0.96 ± 0.03 nm from x-ray diffraction (XRD) measurements, is equivalent to half a single trilayer thickness of the MML (1.9 nm). As such the average terrace width is an important parameter to optimize for in film growth: wider terraces will mean fewer potentially disruptive terrace edges. From the atomic force microscopy (AFM) of complete $\text{Bi}_2\text{Se}_3/\text{Ta}/\text{MML}$ heterostructures shown in Figs. 2(a) and 2(b), the Bi_2Se_3 surface topography can be seen to impart through the overlying multilayer. The result of different Bi_2Se_3 growth recipes on average island size and terrace width is evident. The terrace widths were detected using watershed grain detection, using the open-source software GWYDDION [39]. The distribution in equivalent disk radii, r_T where $\pi r_T^2 = A_T$ for terrace area A_T can be shown to be well approximated by Gaussian distributions with mean radii of $\langle r_T \rangle = 51 \pm 1$ and 120 ± 4 nm for narrow and wide terraced Bi_2Se_3 , respectively, as shown in Fig. 2(c). The fitted standard deviation of distributions is higher for the wide terrace recipe at 100 nm versus 20 nm for narrower terraces indicating a greater variety of island sizes and terrace widths. The terracing is evident in the sample height distributions in Fig. 2(d) with most of both sample surfaces being found within ± 1 terrace steps of the dominant terrace height. Through slab-model fitting of x-ray reflectometry (XRR) measurements, our MMLs exhibit monolayer-level precision with interfacial roughnesses

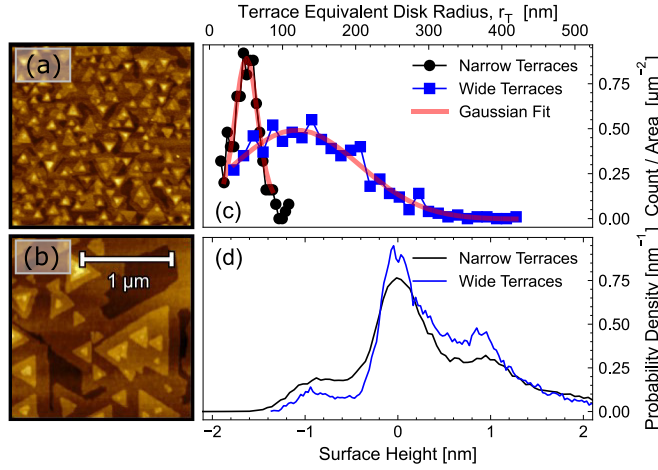


FIG. 2. $2\ \mu\text{m} \times 2\ \mu\text{m}$ AFM images of $\text{Bi}_2\text{Se}_3/\text{Ta}/\text{MML}$ samples grown with recipes leading to different terrace widths: narrow (a) and wide (b). The terrace widths are detected using watershed grain detection, the areas were converted into equivalent disk radii r_T and fitted with Gaussian distributions shown in (c). Resultant distributions in sample height are shown in (d).

on the order of $2\text{--}3\ \text{\AA}$ in corresponding fitted slab models shown in Figs. 3(a)–3(f). Measurements were taken using a Cu $K\alpha$ source of $\lambda = 1.541\text{\AA}$ with a coherence length of $L_c \sim \lambda^2/\Delta\lambda \approx 300\ \text{nm}$. This characteristic length scale is sufficiently similar to the Bi_2Se_3 terracing for the effect of widening the average terrace width to be apparent on the reflectivity. The extracted multilayer periodicity of $1.9\ \text{nm}$ corresponds to a multilayer Bragg peak at $2\theta \sim 4.8^\circ$. For the narrow terraces where $\langle r_T \rangle \ll L_c$, the effective multilayer scattering length density (SLD) is averaged over multiple terraces, and so multilayer features are suppressed. When the terraces become wider such that $\langle r_T \rangle \gtrsim L_c$, the effective medium approximation (EMA) begins to break down and coherent reflection from wider terraces allows for a broad characteristic Bragg peak to emerge. This is reproducible on both Mo and Ta buffer layers shown in Figs. 3(a) and 3(b) and confirms that the trilayer structure grown on Bi_2Se_3 is resolvable. Difficulties fitting high-angle data for samples grown on both terrace widths is attributed to the effect of terracing on the multilayer Bragg peak in addition to the Laue oscillations of the lowest order Bi_2Se_3 XRD peak, which appears at $2\theta = 9.3^\circ$. Hence these XRR curves are fitted only

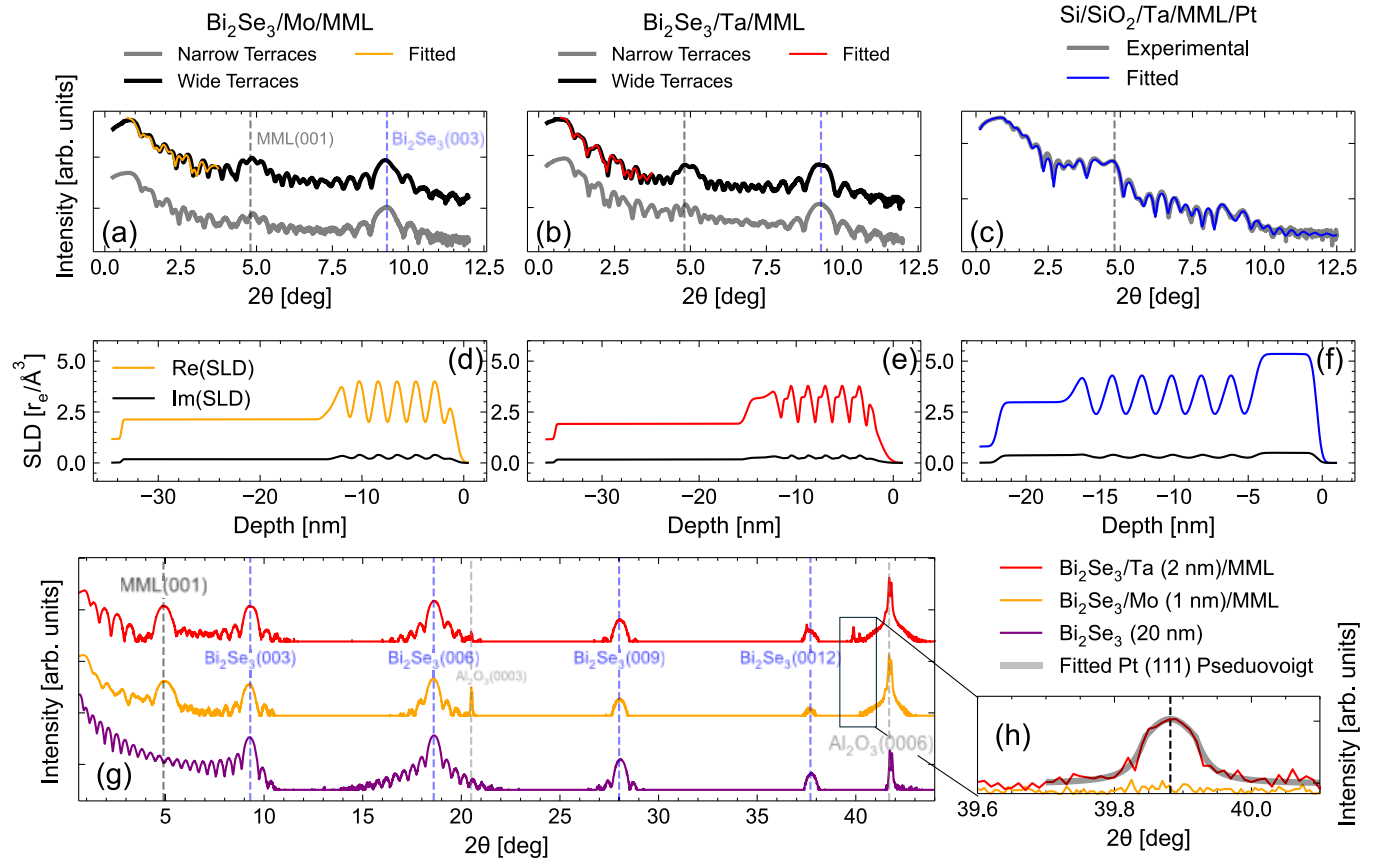


FIG. 3. X-ray characterization of sample structure. (a)–(c) XRR curves of the main $[\text{Pt}(0.8\ \text{nm})/\text{CoB}(0.6\ \text{nm})/\text{Ru}(0.5\ \text{nm})]_6$ MML stack as grown on Bi_2Se_3 with (a) Mo and (b) Ta buffer layers as well as (c) on a thermally oxidized silicon substrate with $5.0\ \text{nm}$ Ta buffer and $4.0\ \text{nm}$ Pt cap. Multilayer and Bi_2Se_3 peaks are marked with gray and blue dashed lines, respectively. Panels (a)–(b) show curves using narrow and wide terraced Bi_2Se_3 ; however, only the wide terraced samples are fitted. The range over which the slab model is fitted is shown in a solid line. Axes beneath (d)–(f) show the fitted SLD distribution corresponding to these fits. The underlying slab models from (d) and (e) are shared with those shown in Figs. 6 and 7. (g) XRD patterns of wide-terrace Bi_2Se_3 grown both with and without overlying MMLs using both buffers. Relevant peaks are labeled. (h) Zoomed axes of (g) showing the Pt (111) peak. All intensity scales in this figure are logarithmic.

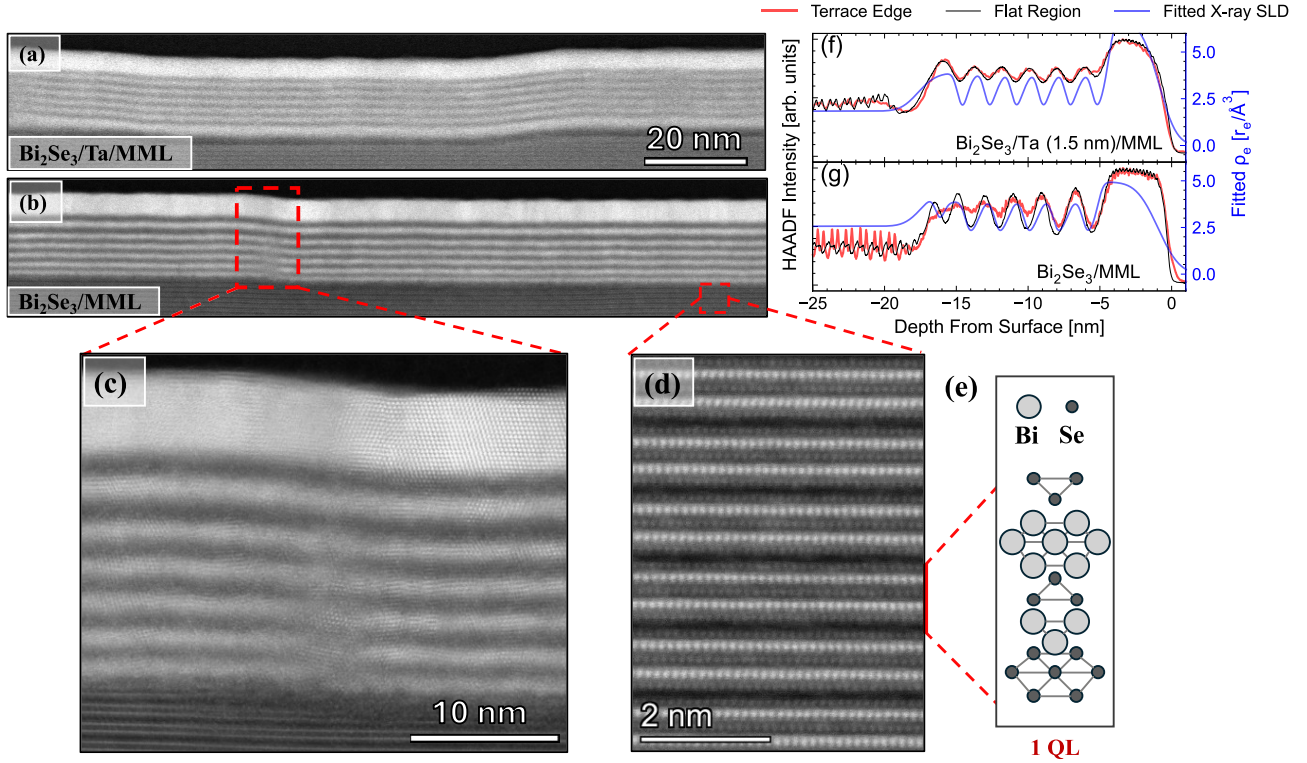


FIG. 4. Cross sections prepared by FIB of $\text{Bi}_2\text{Se}_3/\text{buffer}/\text{MML}/\text{Pt}(4.0 \text{ nm})$ heterostructures as seen under HAADF-STEM imaging. (a), (b) Cross sections showing MML overlaying multiple terrace edges (a) with and (b) without a Ta (1.5 nm) buffer on narrow terrace Bi_2Se_3 . (c) Zoomed-in section of (b) showing terrace edge with resultant layer intermixing; (d) observed quintuple layering of Bi_2Se_3 with corresponding schematic of a single quintuple layer Bi_2Se_3 shown in (e); (f), (g) integrated line profiles of HAADF gray count at and far away from the terrace edge with (f) and without (g) buffer. Overlaid in (f)–(g) are the respective, cofitted x-ray SLDs corresponding to models used in Fig. 7.

in the low angle regime. The full XRD patterns shown in Fig. 3(g) show distinct diffraction peaks matching the Bi_2Se_3 crystal structure. The large number of Laue oscillations are seen indicating good sample quality. Notably, these are seen before and after MML deposition, indicating little to no damage to the TI throughout the sputtering process. The choice between Mo and Ta buffers creates a noticeable difference in the XRD spectra. When grown without the 4 nm Pt capping layer, multilayers grown upon a Ta (2.0 nm) buffer display a diffraction peak seen at 39.88° , characteristic of Pt (111) ordering within the multilayer structure as shown in Fig. 3(h). Pt grown on Mo (1.0 nm) does not show evidence of such Pt(111) ordering in XRD above the noise floor. This is potentially supported by subtle differences in x-ray SLD distributions where the fitted $\text{Bi}_2\text{Se}_3/\text{Ta}/\text{MML}$ profile is the only sample to show three distinct metallic layers in the multilayer repeat unit, indicating sharper interfaces. We can directly observe the structural effect of TI surface terracing on the MMLs using high angle annular dark field scanning transmission electron microscopy (HAADF-STEM), shown in Fig. 4. The multilayer structure over the terrace edge is observed both [Fig. 4(a)] with and [Fig. 4(b)] without a Ta (1.5 nm) buffer, and in both cases can be seen to closely follow the profile of the underlying Bi_2Se_3 surface. The quintuple layering of Bi_2Se_3 which leads to terracing is evident under closer imaging as shown in Fig. 4(d). Notably no Se vacancies are found on the atomic scale. Taking an averaged vertical line profile,

the layers in both samples can be seen to be well defined when away from terrace edges shown in Figs. 4(f) and 4(g). The features in the HAADF count broadly match the fitted x-ray SLD distributions used throughout. Some slight modulation in stack thickness is observed in the nonbuffered sample in Fig. 4(f) with layers being thicker towards the top of the stack; however, this is attributed to the sputter targets heating up during growth leading to growth rate drift. At the terrace edge, the multilayer without any buffer is seen to break down due to increased layer-to-layer intermixing which worsens closer to the Bi_2Se_3 interface as shown in Figs. 4(c) and 4(g). When the Ta buffer is inserted, however, this effect is not seen and the line profile is unchanged at the terrace edge with all layers remaining resolvable as shown in Figs. 4(a) and 4(f). This continuous multilayer structure can be expected to reduce the effect of pinning of domain walls and topological spin textures at terrace edges, essential for any future device applications.

IV. ACHIEVING FULL PERPENDICULAR MAGNETIC ANISOTROPY VIA BUFFER INSERTION

A. Magnetometry

To determine the magnetic properties, the samples were measured by superconducting quantum interference device vibrating sample magnetometry (SQUID-VSM) from which key magnetic parameters such as the stack-averaged, effective anisotropy field H_K , the here-named “PMA ratio,” and

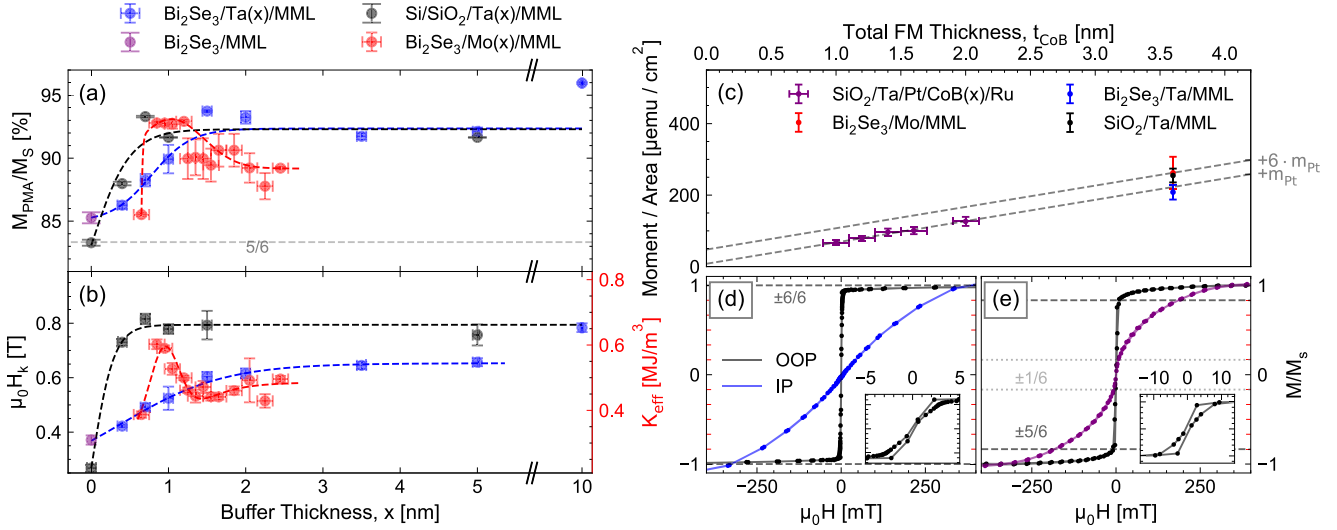


FIG. 5. Buffer thickness dependence of MML magnetic properties, extracted from SQUID-VSM measurements for Ta and Mo buffer thickness series on narrow-terraced Bi_2Se_3 : (a) PMA fraction taken as the fraction of M_s ; (b) the effective anisotropy field, $\mu_0 H_K$, or corresponding K_{eff} assuming constant M_s . Both (a) and (b) include fitted sigmoid + Gaussian functions included to highlight the trend shape as well as a reference series on $\text{Si}/\text{SiO}_2/\text{Ta}(x)$. (c) Moment per unit area versus total t_{CoB} in the stack. A linear relation is fitted to simple $\text{Pt}/\text{CoB}(x)/\text{Ru}$ trilayers to calculate $M_s = 600 \pm 20$ emu/cc; a second is included to account for six proximity magnetized Pt interfaces in the MML stacks. Included are the averaged M_s values for Mo and Ta buffer thickness series shown in (a)–(b); (d), (e) SQUID-VSM loops MMLs on Bi_2Se_3 on a Ta(10 nm) buffer (d) or no buffer (e) taken in both IP and OOP geometries. Inset axes show zoomed OOP data at low field. Easily magnetizable fractions in both geometries are highlighted corresponding to a signal equivalent to an integer number of layers magnetized ($n/6$).

saturation magnetization M_s could be extracted for each sample. The anisotropy field is defined as $H_K = 2K_{\text{eff}}/\mu_0 M_s$ and where $K_{\text{eff}} > 0$ corresponds to an out-of-plane easy axis. The stack-averaged effective anisotropy K_{eff} was calculated by integrating the difference between in-plane and out-of-plane $M(H)$ curves. We define the PMA ratio as M_{PMA}/M_s , the magnetization fraction of the sample magnetized under a small 100 Oe bias field out-of-plane as a fraction of the total saturation magnetization M_s . This small field is needed to saturate out any domains that exist at zero field and is used in place of the more usual squareness ratio which would be near-zero for samples with wasp-waisted hysteresis loops, despite possessing strong PMA. As shown in Figs. 5(a) and 5(e), samples grown without a buffer between the TI and MML show easily magnetizable fractions of 5/6 and 1/6 for the out-of-plane and in-plane directions, respectively, as well as a reduced H_K . This suggests that five of the six CoB layers in the MML possess PMA, but one does not and is instead in-plane magnetized at low fields. Samples were then grown with a range of Mo and Ta buffer layers thicknesses x to find the minimum thickness needed to recover full PMA in every layer. K_{eff} , the PMA fraction, and M_s were then plotted as a function of buffer material and thickness. These values are shown in Fig. 5(a), 5(b), and 5(c), respectively. Inserting a Ta buffer between the TI and the multilayer recovers the PMA fraction to $>90\%$ once a thickness of $x \sim 1.5$ nm is exceeded, as shown in Fig. 5(a), with H_K continuing to rise up to a thickness of 2 nm where H_K reaches a plateau at $\mu_0 H_K = 0.65 \pm 0.02$ T. This is a similar process to that of Ta on natively oxidized Si, where the PMA fraction increases from $\sim 5/6$ without any buffer, plateauing once a critical thickness is surpassed. Differences to note are that on Si, the PMA ratio and H_K recover

faster, plateauing at the same point at $x \sim 0.8$ nm and to a higher H_K value of $\mu_0 H_K = 0.80 \pm 0.01$ T or the equivalent of $K_{\text{eff}} = 0.66 \pm 0.01$ MJ/m³. To obtain a H_K value comparable to multilayers grown on natively oxidized Si on Bi_2Se_3 , the buffer must be increased further, seen in Fig. 5(b), to agree with Si at $x = 10$ nm. At this large thickness of Ta we can expect that the multilayer would be effectively shielded from any spin signal from the Bi_2Se_3 leaving it redundant. Repeating the same experiment with Mo instead of Ta, a different behavior was observed, with both the H_K and the PMA ratio peaking at a low thickness of 0.94 ± 0.06 nm to an increased $\mu_0 H_K = 0.72 \pm 0.05$ T. Increasing the buffer thickness further leads to a sharp decrease in H_K and small decrease in the PMA ratio. This shows there is sufficient Pt (111)/Co hybridization for PMA despite the Pt (111) peak not being visible in XRD measurements in Fig. 3(d). We can hypothesize that there is increased intermixing between Pt and CoB when grown on Mo versus Ta leading to insufficient long range ordering for the Pt (111) peak to be visible. The drop observed in H_K could be explained as this effect worsening with increased Mo thickness. We can therefore conclude that depositing at least 1.5 nm of Ta or ~ 0.9 nm of Mo on the Bi_2Se_3 epilayer surface as a buffer layer recovers the PMA in the layer that is in-plane magnetized with no buffer. The choice of buffer comes with additional subtle changes in the MML's magnetic characteristics. Multilayers grown on both Ta and Mo buffers show M_s values consistent with a trilayer $\text{Pt}(1.5 \text{ nm})/\text{CoB}(x)/\text{Ru}$ thickness series as grown on Ta (5 nm) on natively oxidized Si, as shown in Fig. 5(c), where M_s for $\text{Co}_{68}\text{B}_{32}$ is 600 ± 20 emu/cc. This indicates no substantial magnetic dead layers forming as a result of the Bi_2Se_3 or of the reduced Pt thickness. Proximity-induced magnetism

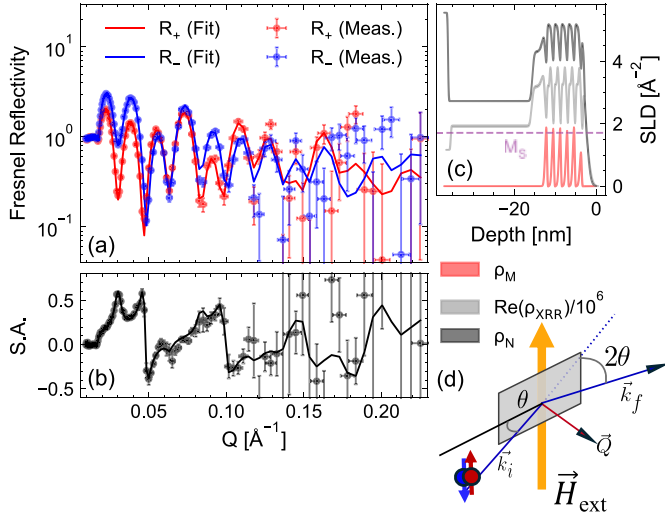


FIG. 6. Example polarized neutron reflectometry data for the $\text{Bi}_2\text{Se}_3/\text{Ta}(2.0\text{ nm})/\text{MML}$ sample with wide terraces: (a) shows Fresnel reflectivity curves for both neutron polarizations: the resultant SA is plotted in (b). Simulated results from a fitted slab model are shown in (a) and (b): the corresponding model's 1σ confidence intervals for magnetic and nonmagnetic neutron scattering distributions are shown in (c) as well as the co-fitted x-ray SLD from 3(e). The equivalent ρ_M from the calculated CoB M_s value taken from Fig. 5(c) is shown as reference; (d) shows the sample geometry with respect to the neutron beam and applied field.

(PIM) in Pt is indicated by a nonzero intercept of $(m/A)_{\text{Pt}} = 8 \pm 3 \mu\text{emu}/\text{cm}^2$ per Pt interface.

B. Polarized neutron reflectometry

While the conventional magnetometry measurements show that full PMA can be recovered in these ultrathin films on Bi_2Se_3 , they still highlight that the magnetic anisotropy can vary greatly throughout the stack. This could be problematic when spin torques are predominantly exerted on the bottom-most FM layer where we cannot easily directly image. A solution to this is to verify the multilayer stack is magnetically homogenous such that we can be confident magnetic textures seen through surface-sensitive measurements correspond to the full texture throughout the stack through magnetostatic coupling. Polarized neutron reflectivity (PNR) was used to obtain a depth resolved magnetization profile of multilayers grown on the various buffers. In this way we can identify the layer that lacks PMA when no buffer is present as well as any large changes in anisotropy. Spin-polarized neutron reflectivities R_+ and R_- were taken as a function of wave vector transfer Q through time-of-flight detection on the POLREF instrument at ISIS, and the spin asymmetry (SA), defined as $\text{SA} = (R_+ - R_-)/(R_+ + R_-)$, between them was calculated. An example PNR measurement is shown in Fig. 6, including the measurement geometry. SAs were measured at applied in-plane fields from full or near saturation ($\mu_0 H = 700\text{ mT}$) down to the minimum applicable field to prevent beam depolarization ($\mu_0 H = 10\text{ mT}$) referred to as near-remnance. Since PNR in the geometry we used [see Fig. 6(d)] is sensitive to the in-plane component of magnetization, our experiment

was designed to measure that component as it increases while the in-plane applied field forces the magnetization into the magnetically hard film plane. When the PNR data were fitted (Figs. 6(a) and 6(b)) using the REFL1D code [40], the resulting magnetic SLD distribution appears as a spike for each of the six magnetic CoB layers [Fig. 6(c)] with an amplitude proportional to that in-plane component. The PNR confirms the SQUID-VSM result that both Ta (2 nm) and Mo (1 nm) buffers yield multilayers with zero in-plane remanance, demonstrated by a vanishing spin-asymmetry under a 10 mT field, which shows that there is no in-plane component of magnetization present. The lack of any notable change in shape of the spin asymmetry with field, shown in Figs. 7(e) and 7(f), suggests magnetic profiles changing in magnitude and not substantially in shape. We can conclude that magnetic anisotropy is uniform throughout the multilayer stack, which possesses full PMA. Measured neutron reflectivity curves were fitted to nuclear and magnetic SLD profiles (ρ_N and ρ_M) using slab model fitting. Due to the low thickness of the individual layers, neutron data were co-fitted with the x-ray data and models shown in Figs. 3(a), 3(b), 3(d), and 3(e) to extend our range in Q . Our magnetic slab models assume a phenomenological, Langevin-type $M(H)$ dependence and fit each CoB with a individual effective anisotropy, $K_{\text{eff},n}$, where the magnetic SLD in the n th CoB layer follows

$$\rho_{M,n}(H) = \mu_0 \alpha M_s \mathcal{L}(a_n \mu_0 H), \quad (2)$$

where $\mathcal{L}(a_n H)$ is the Langevin function and $\alpha = 2.32 \text{ \AA}^{-2} \text{ T}^{-1}$ is a conversion factor [41]. This allows for field-dependent simulations based on a single scaling parameter a_n per layer which can be related to $K_{\text{eff},n}$ and has units of T^{-1} . The final fits are shown in Figs. 7(a) and 7(b) and were obtained using the open source library REFL1D [40]. Integrating $\frac{1}{\mu_0 \alpha} \int \rho_M dz$ yields the moment per unit area which can be shown to agree with the measured SQUID-VSM data shown in Figs. 7(c) and 7(d). In fitted ρ_M distributions at saturation where $H_K < H_{\text{ext}}$, all but the topmost layers can be shown to agree with the M_s values of CoB as shown in Fig. 6(c). Discrepancies in the topmost layer can be discounted as being due to the TI and are more likely the result of slight oxidation due to the thin cap required for subsequent synchrotron imaging (see next section) or the result of removing the overlying Pt layer. Extending over all fields, the model achieves goodness of fit of $\chi^2 = 2.84$ and 2.87 for Ta (2.0 nm) and Mo (1.0 nm) buffers, respectively. A more detailed description of the fitting process can be found in the Supplemental Material [42]. Integrated ρ_M distributions give an estimation for the dipole moment per unit area, which is shown to closely agree with values extracted from SQUID-VSM measurements in Figs. 7(c) and 7(d). The extracted, layer-resolved anisotropy fields $H_{K,n}$ shown in Fig. 7(h) are broadly distributed around the stack averaged values determined by SQUID-VSM in the $\mu_0 H_{K,n} \sim 0.5\text{--}1.0\text{ T}$ range. Averaging over these layer-resolved anisotropies for the Ta (2.0 nm) buffer yields $\frac{1}{6} \sum_{n=1}^6 \mu_0 H_{K,n} = 0.63 \pm 0.03\text{ T}$, agreeing with the corresponding value from SQUID-VSM. Likewise on the Mo (1.0 nm) buffer, this method slightly overestimates H_K compared to the SQUID-VSM-derived value with $\frac{1}{6} \sum_{n=1}^6 \mu_0 H_{K,n} = 0.84 \pm 0.05\text{ T}$. However, due to the

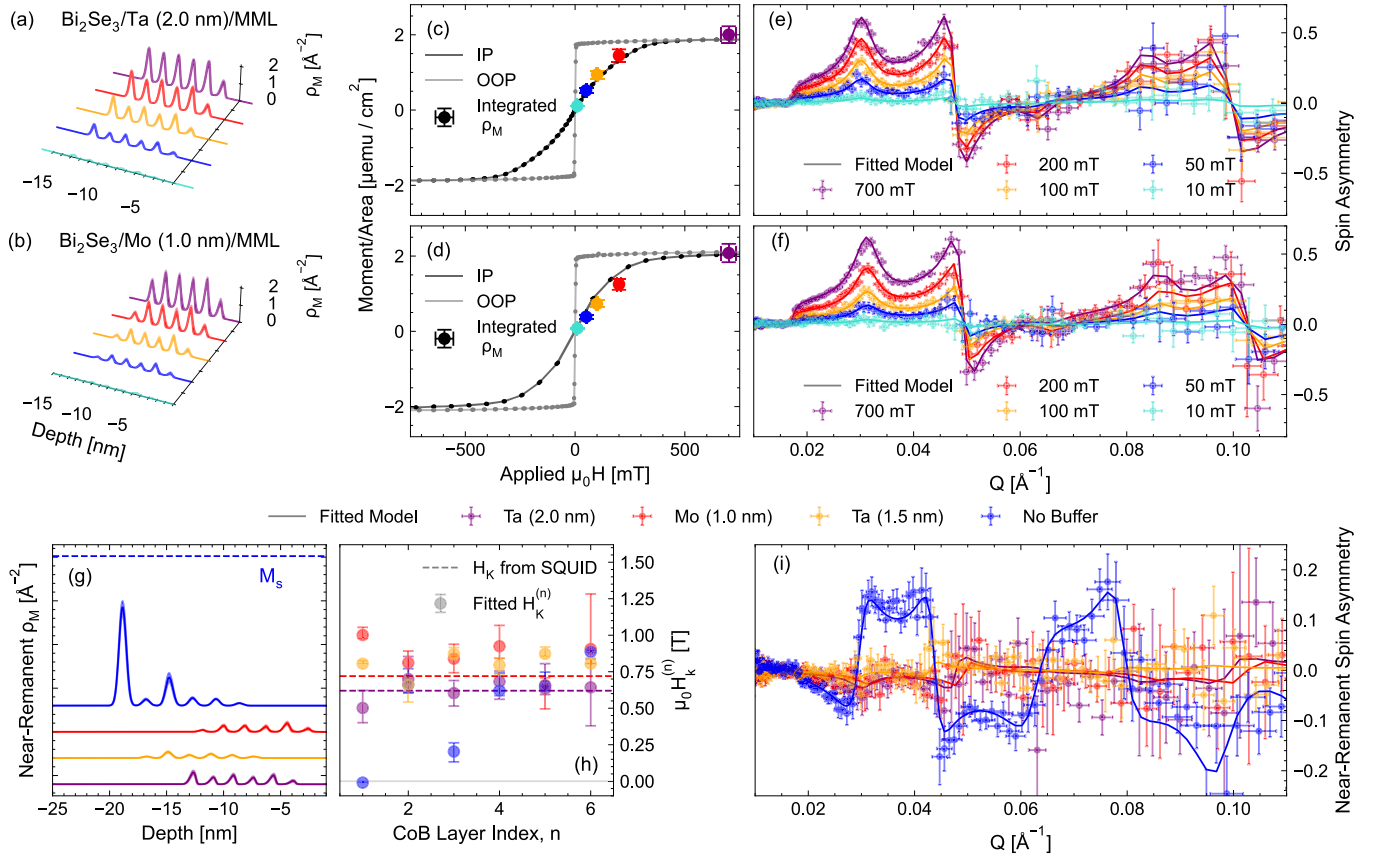


FIG. 7. PNR as a function of external applied field. (a), (b) Fitted magnetic scattering length density profiles, $\rho_M(z)$ for $[\text{Pt}/\text{CoB}/\text{Ru}]_{\times 6}$ multilayers grown on Ta (2 nm) (a) and Mo (1 nm) (b) buffer layers on wide-terrace Bi_2Se_3 . (c), (d) Corresponding magnetometry data for MMLs on Ta (2 nm) and Mo (1 nm) buffers, respectively. Out-of-plane and in-plane SQUID-VSM measurements are shown over the integrated ρ_M profiles from (a) and (b). (e), (f) Fitted spin-asymmetry profiles for Ta and Mo buffers, respectively, with simulated data corresponding to the $\rho_M(z)$ profiles shown in (a) and (b). (g) Fitted SLD distributions for various buffers taken at the minimum 10 mT alignment field. Note that layer depths differ due to differences in the capping layer used. Opaque and translucent profiles show the best obtained fit and the 1σ confidence range, respectively. The equivalent saturation magnetization M_s relative to the nonbuffered sample is shown for reference. (h) Individual, layer-resolved $H_{K,n}$ values obtained from the model fittings for various buffer types. Stack-averaged H_K values extracted from SQUID measurements in Fig. 1(b) are shown for reference. (i) Fitted low-field (10 mT) spin asymmetries corresponding to the slab models in (g).

maximum applied field $0.7 \text{ T} < \mu_0 H_K$ for the Mo (1.0 nm) buffer, unlike for the Ta (2.0 nm), this may limit the accuracy of the final fitted values. Upon removal of the buffer layer, the near-remanent spin asymmetry greatly increases as seen in Fig. 7(i). Furthermore as the shape of the SA profile changes, we can infer the magnetic profile changed in shape compared to saturation and therefore is rendered inhomogeneous throughout the stack. Through slab fitting shown in Fig. 7(g) this can be shown to be localized to an in-plane magnetization equivalent to greater than half of the saturation magnetization in the bottom of the multilayer stack. Here the interfacial layer's modeled $K_{\text{eff},1}$ drops to zero, consistent with this single CoB layer losing PMA and consistent with the missing 1/6 from the PMA fraction seen in Fig. 5(a). This is consistent with previous literature where single ferromagnetic films grown directly onto TIs did not possess PMA [25,26]. We see here where a buffer is not used, this first interfacial layer acts as sufficient buffer material such that the following five layers form above with PMA.

V. MAGNETIC DOMAIN DEPENDENCE ON BUFFER SELECTION

The realization of surface state-mediated spin torques on complex magnetic textures and domains will require tuning not only of the anisotropy but also of the nucleated domain structure, raising again the question of the dependence underlying topography. The effect of the average terrace width on magnetic hysteresis in polar magneto-optical Kerr effect (MOKE) measurements can be seen in Fig. 8(a). Decreasing terrace width on a Mo (1.0 nm) buffer results in a loss of the linear region near zero field indicative of domain pinning increasingly disrupting the formation of demagnetization-induced labyrinthine domains, causing the loop to become squarer when the TI surface terraces are narrow. This can be similarly seen on a Ta (1.5 nm) buffer where the OOP hysteresis becomes almost fully square when grown on narrow Bi_2Se_3 terraces. The effect on the domain structure can be seen directly using photoemission electron microscopy (PEEM) with x-ray magnetic circular dichroism

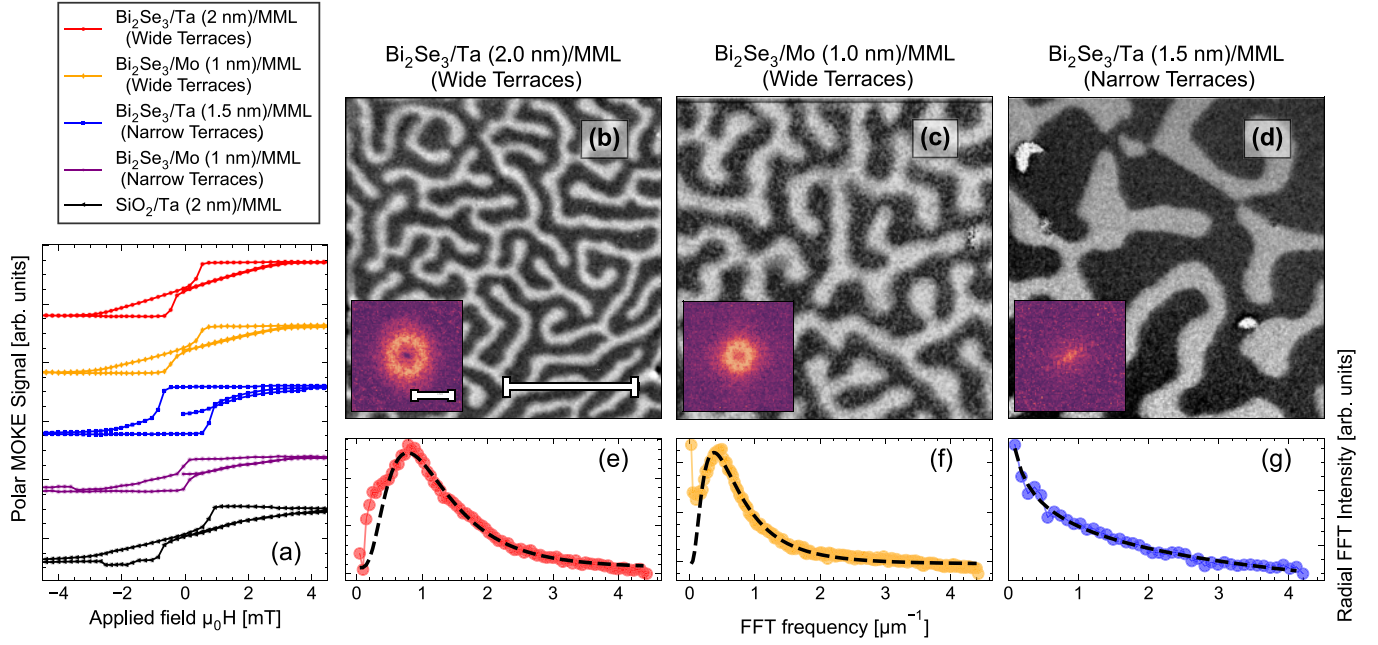


FIG. 8. Domain dependence of MMLs grown on Bi₂Se₃. (a) Polar MOKE hysteresis of MMLs with virgin curves as grown on different buffer layers, TI surface terrace width, and compared against identical MMLs grown on thermally oxidized Si/Ta(5.0 nm). (b)–(d) XMCD-PEEM difference images of magnetic domains in samples with (b) Ta (2.0 nm), (c) Mo (1.0 nm), and (d) Ta (1.5 nm) buffers with either wide (b), (c) or narrow (d) terraces. Scale bars represent 2 μm . (e)–(g) Azimuthally integrated, radial FFT intensities with fitted log-normal distributions shown in black for the same three images. Corresponding FFTs are shown as insets. The included scale bar represents 2 μm^{-1} .

(XMCD) magnetic contrast, carried out at the I06 beamline at Diamond Light Source. The XMCD-PEEM images, taken with the photon energy tuned to the Co L_3 edge, are shown in Figs. 8(b)–8(d). These are difference images, taken by subtracting and normalizing images taken with left- and right-handed circularly polarized x-ray photons. Photoemissive techniques are very surface-sensitive, and so the signal is dominated by the top magnetic layer. As the PNR measurements show that the layer-to-layer anisotropy is rather uniform, we can assume domains do not possess substantial depth dependence and couple through standard magnetostatic effects so as to appear in every layer in the stack. This is supported by basic micromagnetic simulations using a pristine, defectless model with values extracted from magnetometry [42] performed using the MUMAX3 code library [43]. Our XMCD-PEEM results show that, in samples without proper optimization of the Bi₂Se₃ surface morphology that hence possess terraces that are narrow, magnetic domains lack a clear periodicity when demagnetized, as shown in Figs. 8(d) and 8(g), and appear highly pinned to the Bi₂Se₃ surface in Fig. 8(d). On samples with wide TI surface terraces, MMLs on both Mo and Ta buffers demonstrate labyrinthine domain patterns, as shown in Figs. 8(b) and 8(c), with well-defined periodicities seen in fast Fourier transformed (FFT) images [see insets of Figs. 8(e) and 8(f)] and appear to be fully isotropic in-plane. Azimuthal averages of these FFTs were fitted with log-normal profiles, yielding peaks corresponding to domain widths in real space of $\Delta^{\text{Ta}} = 600 \pm 200$ nm and $\Delta^{\text{Mo}} = 1.1 \pm 0.6$ μm for Ta (2.0 nm) and Mo (1.0 nm) buffers, respectively. Identical multilayers grown on Si wafers with 100 nm thermal oxide using Ta (5.0 nm) buffers show an average domain width of $\Delta^{\text{SiO}_2/\text{Ta}} = 400 \pm 100$ nm extracted from

MOKE microscopy [42]. These images show that growing multilayers on Bi₂Se₃, regardless of buffer, leads to changes in the domain structure due to the QL surface terraces; however, the narrower distribution in Δ seen in the Ta (2.0 nm) buffer compared to Mo (1.0 nm) seems indicative of a reduced effect of pinning from the underlying Bi₂Se₃. As the periodicity is characteristic of competition between exchange, the DMI, the anisotropy, and magnetostatic effects the widened and less labyrinthine domain structure on the Mo (1.0 nm) buffer may be attributed to its increased effective anisotropy and apparent differences in Pt ordering. Due to the constraining nature of the molybdenum buffer thickness K_{eff} dependence, there is little room available for optimization by altering Mo buffer thickness. Despite this, as domain structure is greatly improved on the Mo on wide terraces compared to Ta on narrow terraces, we conclude that the Bi₂Se₃ optimization is the greater factor in domain optimization than the buffer choice.

VI. CONCLUSIONS

The insertion of a thin refractory metal buffer between a Bi₂Se₃ epilayer and a subsequently deposited MML has been shown to recover full stack-averaged PMA for multilayers grown in the ultrathin limit for suitably chosen buffer thicknesses. We achieve a similar degree of sample quality of magnetic films as those grown on standard Si/SiO₂ substrates using Ta and Mo buffer layers, highlighting their different thickness dependences, and indicating with PNR that these multilayers are magnetically uniform. This PMA is lost for the first magnetic layer if this buffer is omitted, and HAADF-STEM imaging showing increased intermixing

at terrace sites leads to a local breakdown in multilayer structure rendering the system unsuitable for any potential SOT-driven spin-texture motion. Inserting a buffer layer is shown to ensure the multilayer remains continuous and well defined as it overlays the Bi_2Se_3 epilayer surface topography. Through optimization of Bi_2Se_3 growth to control the surface morphology of the TI epilayer, we can achieve a labyrinthine domain ground state at zero field—promising for hosting skyrmion textures. While the main variable to control appears to be the average Bi_2Se_3 surface terrace width, the choice of buffer is additionally shown to result in changes to the domain structure. We have shown that most well-defined domain ordering is obtained when the MML is grown on Ta (2.0 nm) as opposed to the lighter, thinner Mo (1.0 nm) buffer resulting in some trade-off between terrace-induced pinning and transparency to the spincurrent generated in the TI. Future transport measurements will provide a quantitative characterisation of the exerted spin torques on the MML through the buffer layers. We conclude that the SOT-driven spin texture manipulation, now well established in heavy metal/multilayer films, may well be achievable using the Bi_2X_3 class of topological insulators as the SOT source through correct buffering. To this end both tantalum and molybdenum show promise for

use in such devices so long as their respective strengths are considered.

ACKNOWLEDGMENTS

S.H., A.Y., M.C., S.S., T.A.M., and C.H.M. acknowledge support from the EPSRC Program Grants “CAMIE” (EP/X027074/1) and “NAME” (EP/V001914/1). The samples were grown in the Royce Deposition system at the University of Leeds, which is supported by the Henry Royce Institute, United Kingdom, through Grants No. EP/P022464/1 and No. EP/R00661X/1. Experiments at the ISIS Neutron and Muon Source were supported by beamtime allocations RB2410236 and RB2510555 from the Science and Technology Facilities Council. XMCD-PEEM imaging was undertaken at Diamond Light Source under sessions MM37770-3 and MM38770-1.

DATA AVAILABILITY

The data associated with this paper are openly available at the Research Data Leeds repository [48]. Raw data associated with the ISIS experiments are available at [49,50].

-
- [1] A. R. Mellnik, J. S. Lee, A. Richardella, J. L. Grab, P. J. Mintun, M. H. Fischer, A. Vaezi, A. Manchon, E.-A. Kim, N. Samarth, and D. C. Ralph, Spin-transfer torque generated by a topological insulator, *Nature (London)* **511**, 449 (2014).
 - [2] Y. Fan, X. Kou, P. Upadhyaya, Q. Shao, L. Pan, M. Lang, X. Che, J. Tang, M. Montazeri, K. Murata, L.-T. Chang, M. Akyol, G. Yu, T. Nie, K. L. Wong, J. Liu, Y. Wang, Y. Tserkovnyak, and K. L. Wang, Electric-field control of spin-orbit torque in a magnetically doped topological insulator, *Nat. Nanotechnol.* **11**, 352 (2016).
 - [3] X. Zhang, A. Chen, Y. Zhang, Z. Zeng, Y. Guo, D. Zheng, B. Cui, C. Wu, W. Song, S. Yang, Z. Luo, J. Li, G. Gubbiotti, X. Han, J. Zhao, P. Yan, X. Kou, X. Zhang, and H. Wu, Ultralow electrical current driven field-free spin-orbit torque switching of magnetic tunnel junctions by topological insulators, *Adv. Electron. Mater.* **11**, 2500022 (2025).
 - [4] Y. Fan, P. Upadhyaya, X. Kou, M. Lang, S. Takei, Z. Wang, J. Tang, L. He, L. T. Chang, M. Montazeri, G. Yu, W. Jiang, T. Nie, R. N. Schwartz, Y. Tserkovnyak, and K. L. Wang, Magnetization switching through giant spin-orbit torque in a magnetically doped topological insulator heterostructure, *Nat. Mater.* **13**, 699 (2014).
 - [5] Y. Wang, D. Zhu, Y. Wu, Y. Yang, J. Yu, R. Ramaswamy, R. Mishra, S. Shi, M. Elyasi, K. L. Teo, Y. Wu, and H. Yang, Room temperature magnetization switching in topological insulator-ferromagnet heterostructures by spin-orbit torques, *Nat. Commun.* **8**, 1364 (2017).
 - [6] J. Han, A. Richardella, S. A. Siddiqui, J. Finley, N. Samarth, and L. Liu, Room-temperature spin-orbit torque switching induced by a topological insulator, *Phys. Rev. Lett.* **119**, 077702 (2017).
 - [7] S. Heinze, K. V. Bergmann, M. Menzel, J. Brede, A. Kubetzka, R. Wiesendanger, G. Bihlmayer, and S. Blügel, Spontaneous atomic-scale magnetic skyrmion lattice in two dimensions, *Nat. Phys.* **7**, 713 (2011).
 - [8] C. Moreau-Luchaire, C. Moutafis, N. Reyren, J. Sampaio, C. A. Vaz, N. V. Horne, K. Bouzehouane, K. Garcia, C. Deranlot, P. Warnicke, P. Wohlhüter, J. M. George, M. Weigand, J. Raabe, V. Cros, and A. Fert, Additive interfacial chiral interaction in multilayers for stabilization of small individual skyrmions at room temperature, *Nat. Nanotechnol.* **11**, 444 (2016).
 - [9] S. Woo, K. Litzius, B. Krüger, M. Y. Im, L. Caretta, K. Richter, M. Mann, A. Krone, R. M. Reeve, M. Weigand, P. Agrawal, I. Lemesch, M. A. Mawass, P. Fischer, M. Kläui, and G. S. Beach, Observation of room-temperature magnetic skyrmions and their current-driven dynamics in ultrathin metallic ferromagnets, *Nat. Mater.* **15**, 501 (2016).
 - [10] G. Chen, A. Mascaraque, A. T. N'Diaye, and A. K. Schmid, Room temperature skyrmion ground state stabilized through interlayer exchange coupling, *Appl. Phys. Lett.* **106**, 242404 (2015).
 - [11] A. Soumyanarayanan, M. Raju, A. L. G. Oyarce, A. K. C. Tan, M.-Y. Im, A. Petrović, P. Ho, K. H. Khoo, M. Tran, C. K. Gan, F. Ernult, and C. Panagopoulos, Tunable room-temperature magnetic skyrmions in Ir/Fe/Co/Pt multilayers, *Nat. Mater.* **16**, 898 (2017).
 - [12] W. Jiang, P. Upadhyaya, W. Zhang, G. Yu, M. B. Jungfleisch, F. Y. Fradin, J. E. Pearson, Y. Tserkovnyak, K. L. Wang, O. Heinonen, S. G. E. te Velthuis, and A. Hoffmann, Blowing magnetic skyrmion bubbles, *Science* **349**, 283 (2015).
 - [13] K. Hamamoto, M. Ezawa, and N. Nagaosa, Quantized topological Hall effect in skyrmion crystal, *Phys. Rev. B* **92**, 115417 (2015).

- [14] N. Romming, C. Hanneken, M. Menzel, J. E. Bickel, B. Wolter, K. von Bergmann, A. Kubetzka, and R. Wiesendanger, Writing and deleting single magnetic skyrmions, *Science* **341**, 636 (2013).
- [15] W. Jiang, X. Zhang, G. Yu, W. Zhang, X. Wang, M. Benjamin Jungfleisch, J. Pearson, X. Cheng, O. Heinonen, K. L. Wang, Y. Zhou, A. Hoffmann, and S. te Velthuis, Direct observation of the skyrmion Hall effect, *Nat. Phys.* **13**, 162 (2017).
- [16] K. Zeissler, S. Finizio, C. Barton, A. J. Huxtable, J. Massey, J. Raabe, A. V. Sadovnikov, S. A. Nikitov, R. Brearton, T. Hesjedal, G. van der Laan, M. C. Rosamond, E. H. Linfield, G. Burnell, and C. H. Marrows, Diameter-independent skyrmion Hall angle observed in chiral magnetic multilayers, *Nat. Commun.* **11**, 428 (2020).
- [17] F. Büttner, I. Lemesch, M. Schneider, B. Pfau, C. M. Günther, P. Hensing, J. Geilhufe, L. Caretta, D. Engel, B. Krüger, J. Viehhaus, S. Eisebitt, and G. S. D. Beach, Field-free deterministic ultrafast creation of magnetic skyrmions by spin-orbit torques, *Nat. Nanotechnol.* **12**, 1040 (2017).
- [18] A. Fert, V. Cros, and J. Sampaio, Skyrmions on the track, *Nat. Nanotechnol.* **8**, 152 (2013).
- [19] A. Fert, N. Reyren, and V. Cros, Magnetic skyrmions: Advances in physics and potential applications, *Nat. Rev. Mater.* **2**, 17031 (2017).
- [20] D. Pinna, G. Bourianoff, and K. Everschor-Sitte, Reservoir computing with random skyrmion textures, *Phys. Rev. Appl.* **14**, 054020 (2020).
- [21] T. Yokouchi, S. Sugimoto, B. Rana, S. Seki, N. Ogawa, Y. Shiomi, S. Kasai, and Y. Otani, Pattern recognition with neuromorphic computing using magnetic field-induced dynamics of skyrmions, *Sci. Adv.* **8**, eabq5652 (2022).
- [22] M. Bersweiler, K. Dumesnil, D. Lacour, and M. Hehn, Impact of buffer layer and Pt thickness on the interface structure and magnetic properties in (Co/Pt) multilayers, *J. Phys.: Condens. Matter* **28**, 336005 (2016).
- [23] A. W. J. Wells, P. M. Shepley, C. H. Marrows, and T. A. Moore, Effect of interfacial intermixing on the Dzyaloshinskii-Moriya interaction in Pt/Co/Pt, *Phys. Rev. B* **95**, 054428 (2017).
- [24] A. Hrabec, N. A. Porter, A. Wells, M. J. Benitez, G. Burnell, S. McVitie, D. McGrouther, T. A. Moore, and C. H. Marrows, Measuring and tailoring the Dzyaloshinskii-Moriya interaction in perpendicularly magnetized thin films, *Phys. Rev. B* **90**, 020402 (2014).
- [25] Q. Pan, Y. Liu, H. Wu, P. Zhang, H. Huang, C. Eckberg, X. Che, Y. Wu, B. Dai, Q. Shao, and K. L. Wang, Efficient spin-orbit torque switching of perpendicular magnetization using topological insulators with high thermal tolerance, *Adv. Electron. Mater.* **8**, 2200003 (2022).
- [26] M. DC, R. Grassi, J. Y. Chen, M. Jamali, D. R. Hickey, D. Zhang, Z. Zhao, H. Li, P. Quarterman, Y. Lv, M. Li, A. Manchon, K. A. Mkhoyan, T. Low, and J. P. Wang, Room-temperature high spin-orbit torque due to quantum confinement in sputtered $\text{Bi}_x\text{Se}_{(1-x)}$ films, *Nat. Mater.* **17**, 800 (2018).
- [27] J.-C. Rojas-Sánchez, N. Reyren, P. Laczkowski, W. Savero, J.-P. Attané, C. Deranlot, M. Jamet, J.-M. George, L. Vila, and H. Jaffrès, Spin pumping and inverse spin Hall effect in platinum: The essential role of spin-memory loss at metallic interfaces, *Phys. Rev. Lett.* **112**, 106602 (2014).
- [28] G. Zahnd, L. Vila, V. T. Pham, M. Cosset-Cheneau, W. Lim, A. Brenac, P. Laczkowski, A. Marty, and J. P. Attané, Spin diffusion length and polarization of ferromagnetic metals measured by the spin-absorption technique in lateral spin valves, *Phys. Rev. B* **98**, 174414 (2018).
- [29] L. He, F. Xiu, Y. Wang, A. V. Fedorov, G. Huang, X. Kou, M. Lang, W. P. Beyermann, J. Zou, and K. L. Wang, Epitaxial growth of Bi_2Se_3 topological insulator thin films on Si (111), *J. Appl. Phys.* **109**, 103702 (2011).
- [30] N. Bansal, Y. S. Kim, E. Edrey, M. Brahlek, Y. Horibe, K. Iida, M. Tanimura, G.-H. Li, T. Feng, H.-D. Lee, T. Gustafsson, E. Andrei, and S. Oh, Epitaxial growth of topological insulator Bi_2Se_3 film on Si(111) with atomically sharp interface, *Thin Solid Films* **520**, 224 (2011).
- [31] G. Zhang, H. Qin, J. Teng, J. Guo, Q. Guo, X. Dai, Z. Fang, and K. Wu, Quintuple-layer epitaxy of thin films of topological insulator Bi_2Se_3 , *Appl. Phys. Lett.* **95**, 053114 (2009).
- [32] H. Zhang, C.-X. Liu, X.-L. Qi, X. Dai, Z. Fang, and S.-C. Zhang, Topological insulators in Bi_2Se_3 , Bi_2Te_3 and Sb_2Te_3 with a single Dirac cone on the surface, *Nat. Phys.* **5**, 438 (2009).
- [33] Y. Xia, D. Qian, D. Hsieh, L. Wray, A. Pal, H. Lin, A. Bansil, D. Grauer, Y. S. Hor, R. J. Cava, and M. Z. Hasan, Observation of a large-gap topological-insulator class with a single Dirac cone on the surface, *Nat. Phys.* **5**, 398 (2009).
- [34] C. H. Li, O. M. J. van 't Erve, J. T. Robinson, Y. Liu, L. Li, and B. T. Jonker, Electrical detection of charge-current-induced spin polarization due to spin-momentum locking in Bi_2Se_3 , *Nat. Nanotechnol.* **9**, 218 (2014).
- [35] Y. Wang, P. Deorani, K. Banerjee, N. Koirala, M. Brahlek, S. Oh, and H. Yang, Topological surface states originated spin-orbit torques in Bi_2Se_3 , *Phys. Rev. Lett.* **114**, 257202 (2015).
- [36] M.-H. Nguyen, D. C. Ralph, and R. A. Buhrman, Spin torque study of the spin Hall conductivity and spin diffusion length in platinum thin films with varying resistivity, *Phys. Rev. Lett.* **116**, 126601 (2016).
- [37] H. Nakayama, K. Ando, K. Harii, T. Yoshino, R. Takahashi, Y. Kajiwara, K. Uchida, Y. Fujikawa, and E. Saitoh, Geometry dependence on inverse spin Hall effect induced by spin pumping in $\text{Ni}_{81}\text{Fe}_{19}/\text{Pt}$ films, *Phys. Rev. B* **85**, 144408 (2012).
- [38] W. Zhang, V. Vlaminc, J. E. Pearson, R. Divan, S. D. Bader, and A. Hoffmann, Determination of the Pt spin diffusion length by spin-pumping and spin Hall effect, *Appl. Phys. Lett.* **103**, 242414 (2013).
- [39] D. Nečas and P. Klapetek, Gwyddion: An open-source software for SPM data analysis, *Cent. Eur. J. Phys.* **10**, 181 (2012).
- [40] B. B. Maranville, P. Kienzle, R. Sheridan, A. J. Caruana, J. Borreguero, A. Nelson, D. P. Hoogerheide, A. Book, M. Blackman, P. P. Balakrishnan, and R. Ghose, Reflectometry/refl1d: v1.0.2a0 (v1.0.2a0), Zenodo, 2026, <https://doi.org/10.5281/zenodo.19851918>.
- [41] B. B. Maranville, B. J. Kirby, A. J. Grutter, P. A. Kienzle, C. F. Majkrzak, Y. Liu, and C. L. Dennis, Measurement and modeling of polarized specular neutron reflectivity in large magnetic fields, *J. Appl. Crystallogr.* **49**, 1121 (2016).
- [42] See Supplemental Material at <http://link.aps.org/supplemental/10.1103/hp66-tldh> for details, which includes Refs. [43–47].
- [43] A. Vansteenkiste, J. Leliaert, M. Dvornik, M. Helsen, F. Garcia-Sanchez, and B. Van Waeyenberge, The design and verification of Mumax3, *AIP Adv.* **4**, 107133 (2014).

- [44] A. Glavic and M. Björck, Genx 3: The latest generation of an established tool, *J. Appl. Cryst.* **55**, 1063 (2022).
- [45] D. Foreman-Mackey, corner.py: Scatterplot matrices in Python, *J. Open Source Softw.* **1**, 24 (2016).
- [46] L. Huang, J. Barker, L. Kailas, S. Hait, S. D. Connell, G. Burnell, and C. H. Marrows, Temperature gradient driven motion of magnetic domains in a magnetic metal multilayer by entropic forces, *Phys. Rev. B* **113**, 014428 (2026).
- [47] J. Mulkers, B. Van Waeyenberge, and M. V. Milošević, Effects of spatially-engineered Dzyaloshinskii-Moriya interaction in ferromagnetic films, *Phys. Rev. B* **95**, 144401 (2017).
- [48] B. A. Brereton, S. Hait, A. Yagmur, C. J. Kinane, F. Maccherozzi, M. Conroy, S. Sasaki, T. A. Moore, S. S. Dhesi, S. Langridge and C. H. Marrows (2025), Research Data Leeds, <https://doi.org/10.5518/1847>.
- [49] <https://doi.org/10.5286/ISIS.E.RB2410236>.
- [50] <https://doi.org/10.5286/ISIS.E.RB2510555>.