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Rheology and tribology of dextran/ polyethylene oxide-based water-in-water emulsions

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2 Highlights

- 3 ● Dextran-in-polyethylene oxide (D/P) water-in-water emulsion was studied
- 4 ● D/P emulsions exhibit composition-dependent viscosity behaviour
- 5 ● Higher P increases viscosity, more D content intensifies shear-thinning behaviour
- 6 ● Emulsions exhibit speed-independent mixed lubrication ($\mu < 0.01$) before EHL
- 7 ● Droplet deformation-induced shear thinning may influence lubrication
- 8 performance

Abstract

This study investigated the microstructural, rheological and tribological properties of model W/W emulsions composed of dextran (D) and poly(ethylene oxide) (P) at a fundamental level. Rheological analysis revealed that increasing the P concentration, [P], resulted in increased viscosity (η), whilst increasing the D concentration, [D], intensified shear-thinning behaviour, likely due to changes in the size of D-based droplets. Confocal laser scanning microscopy (CLSM) demonstrated a significant increase in the average droplet size with higher [D] or [P]. A striking tribological result was that the W/W emulsions demonstrated an unusual speed-independent regime, with a coefficient of friction (μ) < 0.01 over a considerable range of sliding contact speed (~ 10 to 100 mm s^{-1} , of physiological relevance) before the onset of the elastohydrodynamic lubrication (EHL) regime. This was not observed for solutions of the individual polymers on their own. Such composition-dependent behaviour may be due to W/W emulsion droplets entering the tribological gap, flattening and reducing the η of the entrained lubricants, thus delaying the formation of a fluid film. Overall, this detailed study shows how fabrication of W/W emulsions via phase-separating polymers can

offer unique lubrication characteristics that could provide advantageous aqueous lubricants for biomedical applications.

9 **1.Introduction**

10 Water-in-water (W/W) emulsions are intriguing colloidal dispersions composed of two
11 immiscible aqueous phases, in which one phase — rich in polymer A — forms droplets that
12 are dispersed within a continuous aqueous phase rich in polymer B [1, 2]. Unlike oil–water
13 interfaces, the interfacial tension of water–water systems is extremely low (~ 0.01 mN/m) [3].
14 Without stabilizers, the droplet phase can persist for a considerable period (up to a few hours)
15 due to the slow diffusion and sedimentation/creaming of the droplets — especially when the
16 continuous phase is viscous [1, 4]. In these systems, the clean droplets with their ultralow
17 interfacial tension and the high mobility of this interface are easily deformed and often form
18 string-like structures under certain flow conditions [5, 6]. Additionally, these emulsions exhibit
19 volume fraction-dependent shear-thinning behaviour even when both constituent polymers
20 are Newtonian fluids [7]. The viscosity (η) of a lubricant is a critical parameter in tribology
21 and most studies have focused on Newtonian fluids, which makes investigating the
22 tribological properties of these non-Newtonian biphasic systems particularly compelling [8].
23 We are aware that investigations into the soft tribological properties of W/W emulsions are
24 scarce, with only three published studies [9-11], but the mechanism of their tribological
25 effects has not been assessed in any great detail.

26 Aqueous mixtures of dextran (D) and polyethylene oxide (P) is one of the most well-known
27 polymer pairs exhibiting segregative phase separation [12, 13]. There is no issue of
28 electrostatic interactions between D and P, and their phase behaviour has been extensively
29 characterized in the literature, making them ideal model systems for studying W/W emulsion
30 lubrication performance [14]. Aqueous mixtures of dextran (D) and polyethylene oxide (P) is

31 one of the most well-known polymer pairs exhibiting segregative phase separation [12, 13].
32 There is no issue of electrostatic interactions between D and P, and their phase behaviour
33 has been extensively characterized in the literature, making them ideal model systems for
34 studying W/W emulsion lubrication performance [14]. Furthermore, W/W emulsions
35 composed of D and P are fully biocompatible dispersions, and their applications extend
36 beyond theoretical studies. For example, they have been used to encapsulate adherent cells
37 in Dulbecco's Modified Eagle Medium to generate tissue spheroids [15]. Unilamellar
38 liposomes-stabilized D/ polyethylene glycol (PEG) emulsions have been applied as
39 microreactors for the synthesis of CaCO_3 materials in biomimetic mineralization processes
40 [16]. More recently, Wang et al. [17] reported the lubrication performance of collagen
41 nanofibril-stabilized PEG/D emulsions, these emulsions shows promising potential as bio-
42 lubricants for the treatment of osteoarthritis.

43 Particles are effective stabilizers for W/W emulsions by adsorbing to the interface [18, 19].
44 Although the interfacial tension in W/W systems is extremely low, the accumulation of
45 particles at the interface can still lead to a reduction in free energy. However, it is important
46 to note that this energy reduction is much smaller compared to O-W interfaces. To explain
47 this phenomenon, Firoozmand et al. [20] proposed that the primary driving force for particle
48 accumulation at the W/W interface is depletion-induced attraction, whereby particles are
49 expelled from the polymer-rich phases and migrate toward the more solvent-rich interface.
50 In addition to colloidal particles, triblock copolymers—comprising segments with distinct
51 affinities for each of the two aqueous phases—can also assemble at the W/W interface to
52 prevent droplet coalescence. Nicolai and Murray [18] suggested that such block copolymers

53 form polymeric micelles that adsorb at the interface and stabilize the emulsion similarly to
54 colloidal particles. However, the presence of stabilizers of the emulsions can make the
55 rheology and tribology much more complicated, because the deformation of droplets
56 covered with a stabilizing layer is likely to change [21]. In this study, we use a simple system
57 without stabilizers to isolate the intrinsic effects of droplet deformation on the overall
58 rheological and tribological performance of the W/W emulsions.

59 For W/W emulsions, the viscosity ratio (λ) of the two phases, defined as the η of
60 continuous phase divided by that of dispersed phase, is a critical parameter determining the
61 extent of deformation of droplets [7]. It has been reported that the formation of string phases
62 in shear flow occurs more easily when $\lambda < 1$ [7]. Therefore, we selected a formulation in
63 which an aqueous solution of a high molecular weight P, with high η , forms the continuous
64 phase, whilst solutions of a low molecular weight D (low η) forms the droplet phase. By
65 limiting our scope to this single dextran-in-polyethylene oxide (D/P) emulsion system, we
66 aim to understand if droplet deformation and string-phase formation also occurs in this
67 system within the tribological contact zone and how this and the low viscosity droplet phase
68 within the more viscous continuous phase influences the overall viscosity and lubrication
69 performance.

70 **2. Materials and methods**

71 **2.1 Materials**

72 Dextran (D) and poly(ethylene oxide) (P) with average molecular weights of
73 35,000~45,000 and 200,000 Da, respectively, were purchased from Sigma-Aldrich, Dorset,

74 UK. Fluorescein isothiocyanate-dextran (FITC-D) with average molecular weight of 40,000
75 Da was also purchased from Sigma-Aldrich, Dorset, UK. MilliQ water purified by a Milli-Q
76 apparatus (Millipore, Bedford, UK), with an electrical resistivity not less than 18.2 MΩ.cm
77 was used as the solvent throughout the experiments. Sodium azide was used as a
78 bactericide to prevent microbial growth.

79 **2.2 Sample preparation**

80 D or P were dissolved in MilliQ water and stirred overnight to obtain stock solutions. The
81 P solution contained a small amount of insoluble particles that were removed by
82 centrifugation at 5,000 g for 30 min. The phase diagram was determined by mixing P and D
83 solutions at different compositions and checking for phase separation. Water-in-water (W/W)
84 emulsion samples used for investigation were prepared by mixing required amount of the
85 stock solutions using a vortex mixer for 1 min. **Table 1** shows the initial compositions of the
86 pure polymer solutions and selected mixed D/P systems that phase separated into W/W
87 emulsions as well as the volume fraction of D phase, along with the measured volume
88 fraction of the dextran-rich phase. The phase volume fractions were determined after one
89 week to ensure complete macroscopic phase separation. Unlike previous literature focusing
90 purely on rheological [22] and stability assessments [12], and selecting samples along the
91 same tie-line, W/W emulsion samples in this study were chosen within a narrow high shear
92 viscosity window for tribological assessment.

93 **Table 1.** Nomenclature and composition of the samples.

Sample	Nomenclature	D (wt%)	P (wt%)	Volume fraction of D
Dextran (D)	4D	4	0	-
	6D	6	0	-
	10D	10	0	-
Polyethylene oxide (P)	6P	0	6	-
	8P	0	8	-
	10P	0	10	-
Water-in-water (W/W) emulsions of dextran in poly(ethylene) oxide (D/P)	1D-8P	1	8	1.8
	2D-8P	2	8	6.0
	3D-8P	3	8	12.6
	4D-4P	4	4	12.6
	4D-6P	4	6	15.2
	4D-8P	4	8	16.7
	6D-8P	6	8	23.8

94 **2.3 Confocal laser scanning microscopy (CLSM)**

95 A Zeiss LSM880 inverted confocal microscope (Carl Zeiss Microscopy GmbH, Germany)

96 was used to visualize the emulsion droplets and images were analysed using ZEN Black

97 software. A small amount (0.01 wt%) of FITC-D was added to the D stock solutions. FITC

98 was excited at wavelength (λ) of 488 nm and the fluorescence was collected between $\lambda \approx$

99 480 to 800 nm. Dark areas were therefore assumed to be unlabelled P-rich regions. The

100 freshly prepared emulsions or pure solutions were pipetted onto a concave microscopic slide

101 and covered with a glass coverslip before imaging. ImageJ software (version 1.54g, National

102 Institute of Health, Bethesda, USA) was used to determine the diameter of the emulsion

103 droplets, and the mean droplet size was calculated using at least 100 droplets in multiple

104 images. The volume-weighted mean diameter ($D_{4,3}$) was calculated according to the

105 following equation:

106
$$D_{4,3} = \frac{\sum n_i d_i^4}{\sum n_i d_i^3}$$

107 where n_i represents the number of droplets with diameter d_i .

108 **2.4 Apparent viscosity**

109 The viscosity of samples was measured using an MCR-302 modular compact rheometer
110 (Anton Paar, Austria). Cone and plate geometry (CP50-1) was used: cone diameter = 50
111 mm and cone angle = 1°. The shear rates were varied from 0.1 to 1000 s⁻¹. All experiments
112 were carried out within 2 h of W/W emulsion formation, by which time no visible phase
113 separation had occurred. For each measurement, 2 mL of the sample was pipetted onto the
114 plate and a temperature-controlled cover was used to prevent evaporation and maintain the
115 temperature at 37 ± 0.1 °C. Samples were left on the plate for 1 min to achieve thermal
116 equilibrium before rheological measurements commenced.

117 **2.5 Tribology**

118 Tribological data were obtained using an MTM2 Mini Traction Machine (PCS instruments,
119 UK) equipped with a tribo-pair consisting of smooth polydimethylsiloxane (PDMS) ball
120 (diameter = 19 mm, Young modulus ~2.6 MPa) and disc (diameter = 46 mm). All
121 experiments were carried out 5 minutes after W/W emulsion formation, during which time
122 no visible phase separation occurred. A normal load (W) of 2 N, equivalent to a maximum
123 Hertzian contact pressure of ~200 kPa [23] was used. The temperature was controlled at 37
124 ± 1 °C to mimic physiological conditions. The coefficient of friction (μ) was recorded for all

125 samples as a function of the entrainment speed (U) between 1 to 1000 mm s⁻¹ with $U = (u_D$
126 $+ u_S)/2$, where u_D and u_S are the disc and ball speeds, respectively. The sliding rolling ratio
127 (SRR), defined as the ratio of the *absolute* value of sliding speed $|u_D - u_S|$ to U , was kept
128 constant at 0.5.

129 **2.6 Statistical analysis**

130 All results presented were performed in triplicate ($n = 3 \times 3$), with samples prepared on
131 different days and mean values plotted with the standard deviation as the error bar, unless
132 otherwise stated. Statistical analyses were carried out using one-way ANOVA and multiple
133 comparison test via SPSS software and differences between samples were deemed
134 significantly different at $p < 0.05$ via Tukey's test.

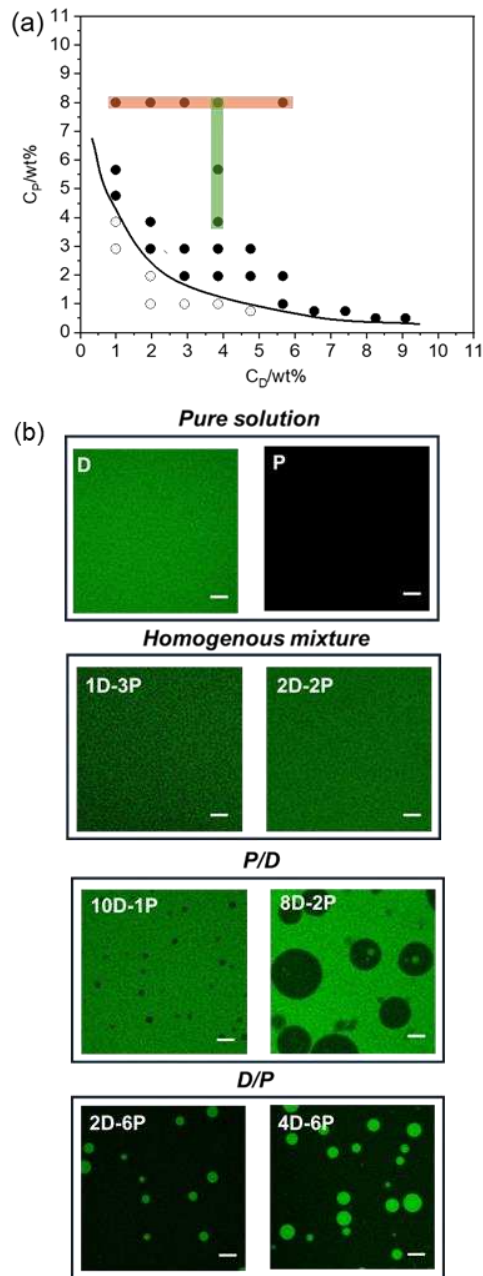
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136 3. Results and discussion

137 3.1 Phase diagrams

138 It is generally known that polymers P and D become incompatible at higher concentrations
139 mainly caused by steric interactions [13, 24]. Figure 1a shows the phase diagram of the
140 polymers used in this study, *i.e.*, D (35,000~45,000 Da) and P (200,000 Da). The curved line
141 represents the estimated binodal curve that separates the one-phase region from the
142 biphasic region. A homogeneous single-phase solution at the macroscopic scale is achieved
143 when the concentrations of both components lie below this curve, as demonstrated by the
144 CLSM images of 1D-3P and 2D-2P (Fig 1b). For compositions above this curve, the system
145 undergoes phase separation, resulting in a polymer A-enriched phase that forms droplets
146 dispersed in a continuous phase enriched in polymer B, as visualized in the CLSM images
147 of P/D or D/P in Fig. 1b.

148 In this study, we selected phase-separated D/P with various ratios, dividing them into two
149 groups: one group with a fixed concentration of D [D] of 4.0 wt% (4D-4P, 4D-6P, 4D-8P, as
150 highlighted by the green line in the Fig. 1a) and the other with a fixed concentration of P [P]
151 of 8.0 wt% (1D-8P, 2D-8P, 3D-8P, 4D-8P, 6D-8P, as highlighted by the orange line in the Fig.
152 1a) resulting in different overall η and volume fractions of dispersed phase, to investigate
153 their microstructure, rheological and tribological properties.



154

155 **Figure 1.** (a) Phase diagram of poly(ethylene oxide) (P) and dextran (D). Open circles
 156 represent homogeneous samples, whilst closed circles indicate phase-separated samples.
 157 The solid line represents the binodal curve. (b) Confocal laser scanning micrographs (CLSM)
 158 of the pure D and pure P solutions, homogenous mixtures, P/D and D/P emulsions. The
 159 scale bar = 20 μm .

3.2 Microstructure of W/W emulsions

The microstructure of these emulsions is in principle time dependent. All the images were captured 5 min after the preparation of the emulsions. Based on the CLSM images (Figs. 2 and 3), all the emulsions investigated are of the D/P type. Both the concentrations of D and P influence the size of droplets. One anomaly that can see is that whilst the D phase volume increases when the ratio D/P is increased (in Figure 3), as expected, in Figure 2 when more P is added keeping the D concentration constant (decreasing the D/P ratio), the volume fraction of D appears to increase. We cannot account for this other than of course sedimentation/creaming and also possibly coalescence of the droplet phase may result in different apparent volume fractions even though we tried to keep the observation times and handling of the samples the same in all instances.

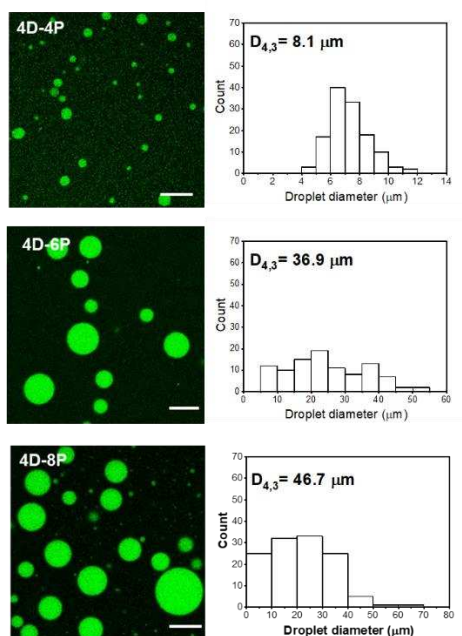
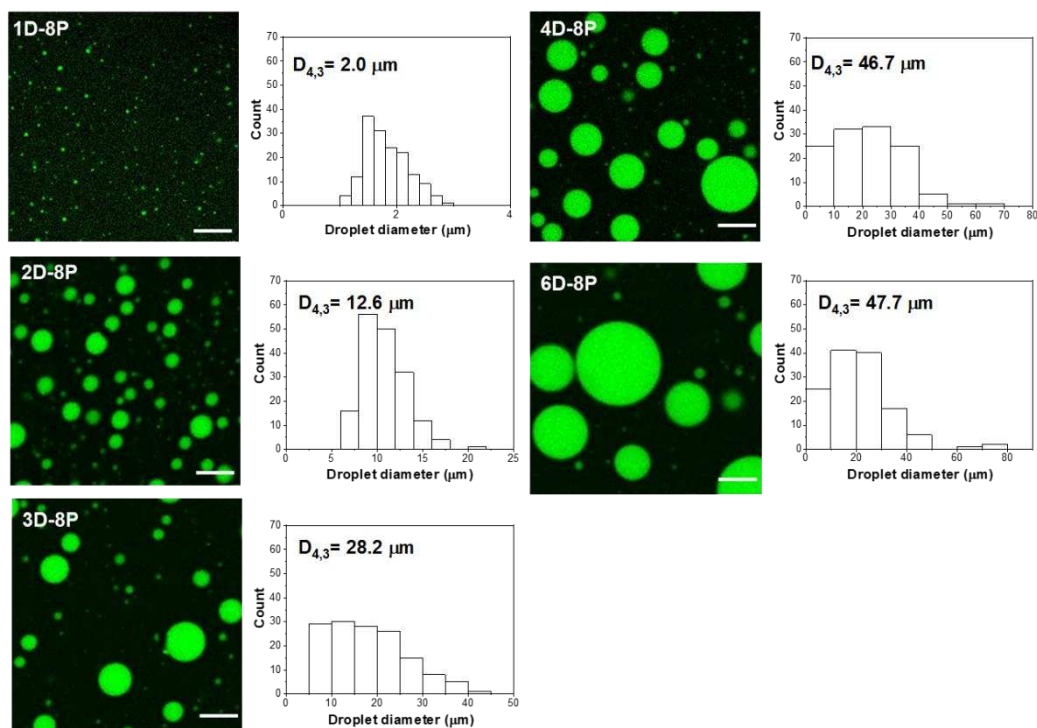


Figure 2. CLSM micrographs of W/W emulsions with a fixed concentration of D (4 wt.%) plus histograms of the droplet size distributions and corresponding $D_{4,3}$ values. The scale bar = 20 μm .



175

176 **Figure 3.** CLSM micrographs of W/W emulsions with fixed concentration of P (8 wt.%) plus
 177 histograms of droplet size distributions and corresponding D_{43} values. The scale bar is 20
 178 μm .

179 3.3 Rheological characteristics

180 Figures 4a and 4b show viscosity (η) versus shear rate ($\dot{\gamma}$) for pure D and P solutions at
 181 different concentrations. The pure D solutions exhibited almost Newtonian η that increased
 182 with [D], a behavior observed elsewhere [25]. The pure P solution exhibited shear-thinning
 183 behavior at high $\dot{\gamma}$ ($> 100 \text{ s}^{-1}$), and its η was at least one order of magnitude higher than that
 184 of the D solutions of the same concentration. In Figures 5a and 5b, η versus $\dot{\gamma}$ is shown for
 185 the D/P emulsions with fixed [D] and [P], respectively. D comprises the droplet phase and P
 186 comprises the continuous phase for all these emulsions.

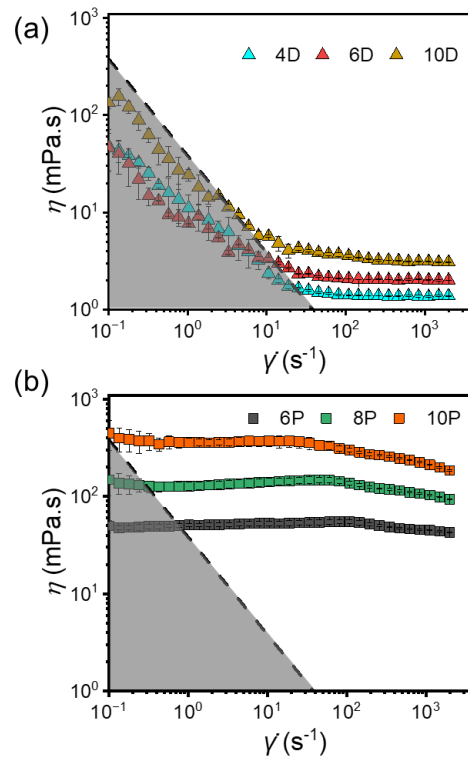
187 Of course we do not know exactly the volume fractions of the dextran droplets in the P
 188 phase, nor the exact concentrations of D or P in each phase, which will vary depending on

189 the exact location of the mixed composition in the phase diagram, according to the tie-lines
190 (though it is probably safe to assume that most of the D resides in the droplet phase and
191 most of the P in the continuous phase). Accordingly, the η of the biphasic systems will
192 depend on the exact concentration of P in the continuous phase and the volume fraction of
193 droplets of D solution. If the discontinuous phase consisted of monodisperse hard spheres
194 it might be possible to predict, on theoretical grounds, the additional contribution that these
195 would make to the overall η of the continuous phase, but of course they are neither, rather
196 highly deformable fluid objects of varying sizes. At the same time, the interfacial tension will
197 also vary depending on the composition, which will determine the degree of deformation of
198 the droplets in a given shear field. (The tension will be lower the closer the composition is to
199 the binodal). However, the interfacial tension for such systems is difficult to measure and
200 indeed the exact shear rate conditions as a function of speed are also not known and our
201 objective in presenting the actual, measured η was to see how far this could account for the
202 observed tribological changes.

203 Increasing [D] and [P] increased emulsion shear-thinning behaviour, but to a varying
204 extent depending on the ratio between D and P. The shear-thinning effect is attributed to the
205 stretching out and aligning of droplets in the direction of the flow [22]. 1D-8P showed almost
206 Newtonian behaviour for $\dot{\gamma} > 60 \text{ s}^{-1}$, indicating that the continuous phase dominated the
207 overall bulk η . The mean droplet radius of 1D-8P was low, with a $D_{4,3}$ of $2.05 \text{ }\mu\text{m}$, which is
208 too small to influence the structure of the emulsion under flow. Wolf and Frith (2003) reported
209 that no string phase formed in those emulsions with only 10% droplet phase volume, and
210 the flow behaviour of this polymer blend is dominated by the continuous phase [7]. In our

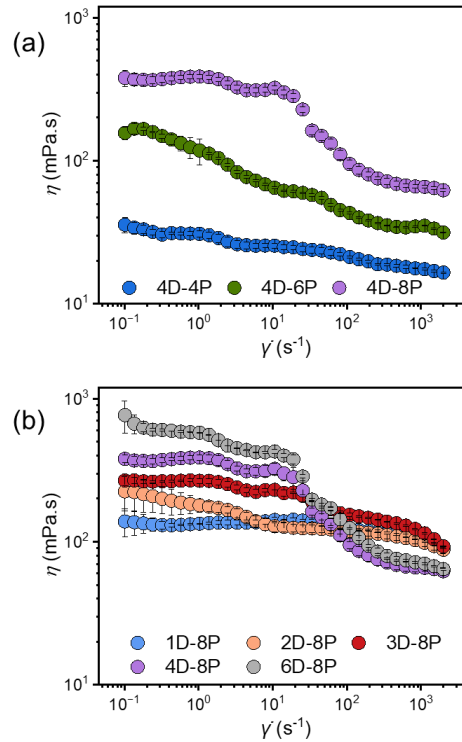
211 system, the droplet phase (D) is much less viscous than the continuous phase (P), meaning
212 that at a fixed $\dot{\gamma}$, the high η of the continuous phase imposes substantial viscous stresses
213 on the droplets, promoting deformation, elongation or break up. Furthermore, given the ultra-
214 low interfacial tensions typical of W/W emulsions ($\sim 10^{-5}$ N/m) [26], even low $\dot{\gamma}$ generate high
215 capillary numbers, leading to strong deformation [27].

216 For emulsions with constant [D] = 4 wt% and increasing [P], this resulted in intensified
217 shear-thinning and increased η . The more pronounced shear-thinning can be attributed to
218 the increased [P], shifting the system further from the critical point, leading to stronger phase
219 separation and larger droplets. Larger droplets start to deform at lower $\dot{\gamma}$ keeping the
220 interfacial tension constant. Emulsions containing a fixed [P] of 8 wt% exhibited a low shear
221 rate plateau in their flow curves, consistent with the shear behaviour observed in the pure
222 8P solution (Fig. 4b), indicating that the η of the D/P emulsions was dominated by the
223 continuous phase under at low $\dot{\gamma}$. The low shear rate plateau in η of these emulsions is
224 higher than that of the pure 8P solution and increases significantly in emulsions with
225 increasing [D], even though the η of pure D is very low compared to that of P. It was expected
226 that increasing [D] results in a higher effective [P] within the continuous phase, thereby
227 increasing the η , although the exact concentration cannot be measured in this dynamic
228 system. Emulsions with higher [D] exhibit more pronounced shear-thinning behaviour over
229 $\dot{\gamma} = 10$ to 100 s^{-1} which is probably related to the droplet deformation and alignment. At high
230 $\dot{\gamma}$ ($> 100 \text{ s}^{-1}$), increasing [D] led to lower η , indicating that the D droplets predominantly
231 govern the shear-thinning response at these higher $\dot{\gamma}$.



232

233 **Figure 4.** Apparent viscosity (η) versus shear rate (γ) for (a) pure D solutions and (b) pure
 234 P solutions. Data represents means and standard deviations of triplicate samples ($n = 3 \times$
 235 3). The dashed line in (a) is the lowest measurable shear viscosity corresponding to $80\times$
 236 the minimum measurable torque (10^{-9} N m) of our shear rheometer and the shaded region
 237 denotes the invalid data range due to the geometrical limitations of the rheometer.



238

239 **Figure 5.** Apparent viscosity (η) versus shear rate (γ) for W/W emulsions with (a) fixed
 240 concentration of D and (b) fixed concentration of P. Data represents means and standard
 241 deviations of triplicate samples ($n = 3 \times 3$).

242 3.4 Tribological characteristics

243 3.4.1 W/W emulsions with fixed [D]

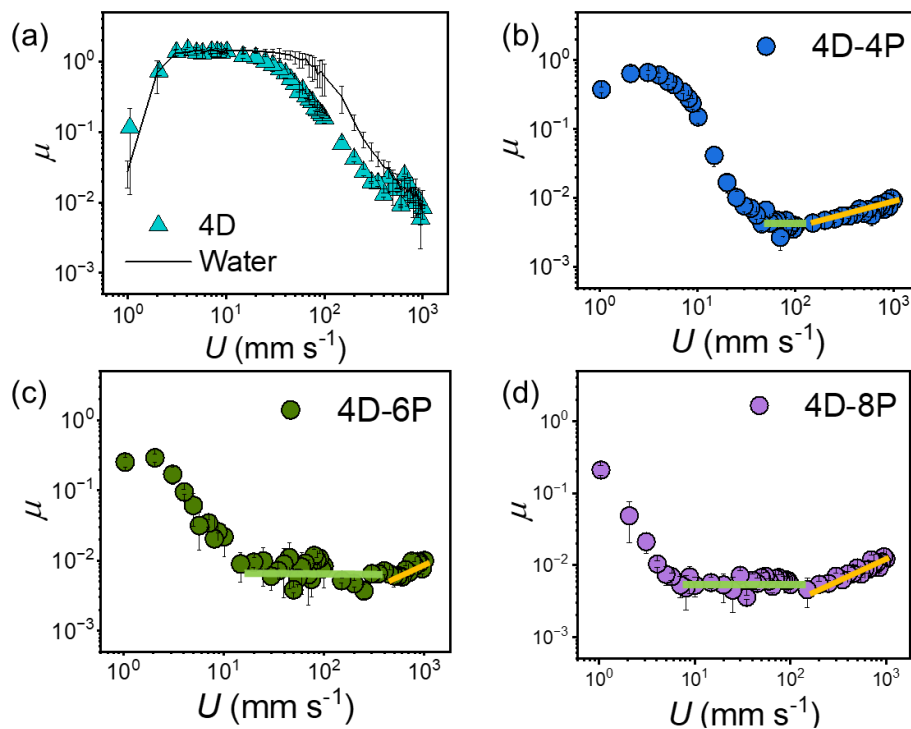
244 Figure 6 shows the friction results for pure D, water and D/P emulsions with fixed [D]. Fig.
 245 6a shows that pure 4D exhibited nearly identical coefficient of friction (μ) vs entrainment
 246 speed (U) as water alone for the entire speed regime studied. This indicates that D
 247 molecules rarely adsorb onto PDMS surface and do not contribute to further friction lowering
 248 compared to water. Although μ decreases rapidly from approximately 30 mm/s to about 10^{-2}
 249 at 1000 mm/s, indicating a transition into a mixed lubrication regime, a further transition to
 250 an elasto-hydrodynamic lubrication (EHL) was never achieved. This behavior is attributed

251 to the low η of pure 4D. For 4D-4P, 4D-6P and 4D-8P emulsions, the transition to a mixed
252 lubrication regime was observed from gradually lower speeds (within 1–10 mm s⁻¹) with
253 increasing [P], and 4D-8P did not exhibit a distinct boundary regime. This can be attributed
254 to the bulk viscosity; lubricants with higher η reduce friction at lower speeds because viscous
255 fluids are more resistant to being squeezed out of the contact zone than less viscous fluids
256 [8].

257 In sharp contrast, the emulsion samples displayed an unusually speed-independent
258 mixed lubrication regime (highlighted by the green line in Fig. 6) with constant $\mu < 0.01$ over
259 a substantial range of sliding speeds (~ 10 to 100 mm s⁻¹, which is physiologically relevant)
260 before transitioning to the EHL regime. The entrainment rate of the lubricants depends on
261 both U and η of fluids at the contact inlet. As U increases and the contacting gap widens due
262 to fluid pressurization, droplets in the W/W emulsions will more easily enter the gap, where
263 they probably flatten and reduce the effective η of the entrained lubricant, a phenomenon
264 described as "inlet shear-thinning." [28]. To support this, the theoretical film thickness in the
265 full-fluid film hydrodynamic regime was calculated, showing values in the range of ca. 1 to 5
266 μm for 4D-4P and 4D-6P, and ca. 3 to 10.5 μm for 4D-8P (Fig. S1). The droplet diameters
267 observed in the emulsions (Fig. 2), ranging from 8.1 to 46.7 μm , are generally larger than
268 the calculated lubricating film thickness. However, we propose that these droplets could still
269 be entrained into the contact and contribute to lubrication, as they are highly deformable and
270 can be stretched into string-like structures. This can further contribute to shear-thinning of
271 the lubricant.

272 Additionally, the γ' under the tribological contacts were estimated at different viscosities in

273 the calculation of the theoretical μ values. Detailed calculation procedure and results are
 274 presented in Fig. S2 and relevant discussion therein. Briefly, these calculated γ' values (ca.
 275 10^4 to 10^5 s $^{-1}$, shown in Fig. S2c and Table S1) are orders of magnitude higher than the
 276 maximum γ' measured using the rheometer (ca. 10^3 s $^{-1}$, Fig. 5), primarily due to the much
 277 thinner film thickness under the tribological contacts compared to the 1.03 mm gap applied
 278 in rheological measurement. Thus, shear-thinning of the lubricants tends to delay the
 279 formation of a full fluid film. As [P] was increased, the speed range of this regime extended,
 280 which might be attributed to the more intense shear-thinning of the systems. This special
 281 mixed regime has been observed elsewhere for shear-thinning xanthan solutions in PDMS-
 282 PDMS ball-on-three-plates contact [29]. However, the authors focused on the EHL of the
 283 fluids and did not discuss this regime.



284
 285 **Figure 6.** Coefficient of friction (μ) versus entrainment speed (U) for W/W emulsions with
 286 fixed concentration of D in PDMS-PDMS contact for 3 different mixtures: (b) 4D-4P, (c) 4D-

287 6P, (d) 4D-8P. The results for 4D alone and water are plotted in (a). The green line and
288 yellow line in (b-d) show the linear fits to the mixed and hydrodynamic regimes, respectively,
289 for the 4D-4P, 4D-6P and 4D-8P system. Data represents means and standard deviations
290 of triplicate samples ($n = 3 \times 3$).

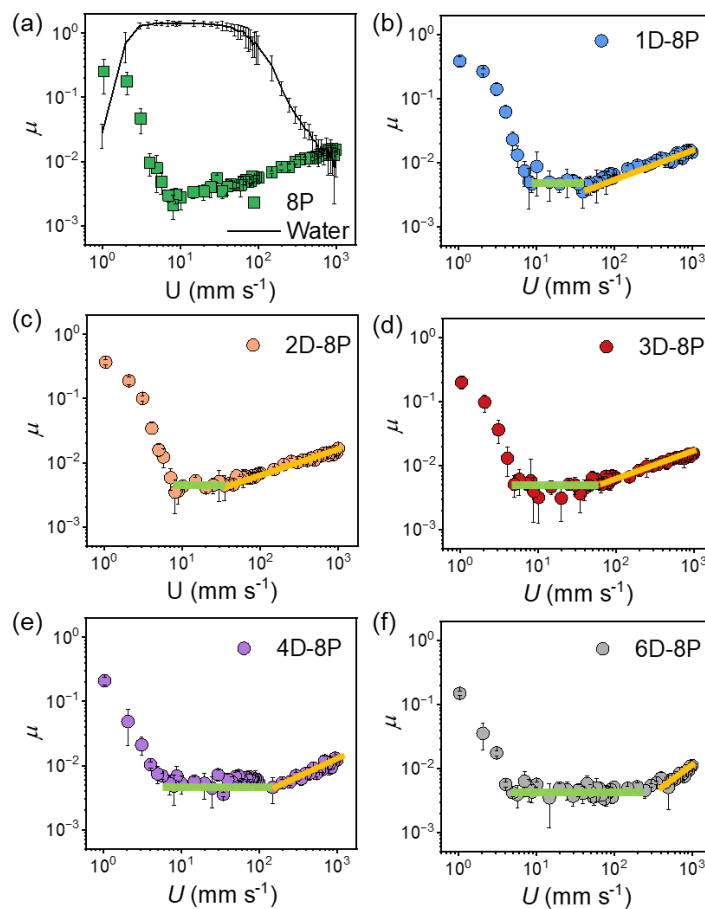
291 3.4.2 W/W emulsions with fixed [P]

292 Figure 7 presents the friction data for all D/P emulsions with fixed [P], along with pure 8P
293 and water as a control, plotted as μ versus U . Fig. 7a shows that pure 8P exhibits a classic
294 Stribeck curve, with μ decreasing as speed increases to a minimum value, then rising again
295 at higher speeds. The different tribological behaviour of 8P compared to 4D (Fig. 6a) is
296 attributed to the much higher viscosity of 8P (Fig. 4). All XD-8P emulsions ($X = 1, 2, 3, 4,$
297 and 6) exhibited a downward trend in μ — indicative of the mixed regime — from the lowest
298 speeds, and no distinct boundary regime was observed. Moreover, as [D] increased, the μ
299 values in the mixed regime gradually decreased (Fig. 8a). Stribeck curves were also
300 constructed by multiplying the value of U by the η value in the first plateau region (specifically
301 at $\dot{\gamma} = 3.3 \text{ s}^{-1}$) to ‘eliminate’ the influence of η (Fig. 8b). All samples collapsed onto a Master
302 curve within the mixed regime, indicating that this regime is predominantly governed by the
303 bulk phase η [30].

304 Similar to the 4D-XP emulsions (Fig. 5), a speed-independent mixed lubrication regime
305 (highlighted by the green line) was also observed for XD-8P emulsions. However, with
306 increasing [D] this special regime, characterized by invariant μ values, starts at lower U and
307 extends over a broader range, causing the transition to the EHL regime to occur at higher

308 *U*. It has been reported that, under shear, droplets deform into string-like structures, as
309 evidenced by images captured during shear experiments [22, 31, 32]. We propose that
310 deformation of droplets to string phases also occurred in the D/P emulsion systems under
311 the tribological shear. According to our film thickness calculation in Fig. S1, the calculated
312 film thickness (ca. 2 to 5 μm for 6D-8P and 2~3 to 16.5 μm for the other D/P emulsions,
313 Table S1) was 2-3 times higher than those of 4D-XP emulsions. Therefore, the likelihood of
314 the droplets being entrained to comprise the lubricating films, although in a deformed string
315 shape, and dominate the lubrication is even higher. Tromp and De Hoog reported that string
316 formation in systems with higher dispersed phase volumes occurred at lower shear rates
317 [33]. In emulsions with higher [D], the droplets are more prone to deformation and entry into
318 the contact zone, which would cause the 'flat' mixed lubrication regime to be initiated at
319 lower speeds and persist over a more extended range. In contrast, the 1D-8P sample
320 displayed a μ value comparable to that of pure 8P over the entire speed range, which is
321 consistent with the flow curve results, indicating that the small fraction of droplets had a
322 limited effect. In the EHL regime (indicated by the orange line), the 4D-8P and 6D-8P
323 samples exhibited significantly lower μ values than the emulsions with lower [D] *ie.*, 1D-8P,
324 2D-8P and 3D-8P, which could be attributed to their lower η at high $\dot{\gamma}$. However, within the
325 range of shear rates we tested (up to 2000 s^{-1}), the second plateau was not reached—the
326 samples continued to shear thin at higher $\dot{\gamma}$ (corresponding to shear conditions in the EHL
327 regime within the tribological contact). Therefore, scaling the Stribeck curve using the high
328 shear η measured at 2000 s^{-1} is not appropriate. Alternatively, this phenomenon could be
329 interpreted by the calculated μ in Fig. S2b, the η best fit the tribological data of 6D-8P and

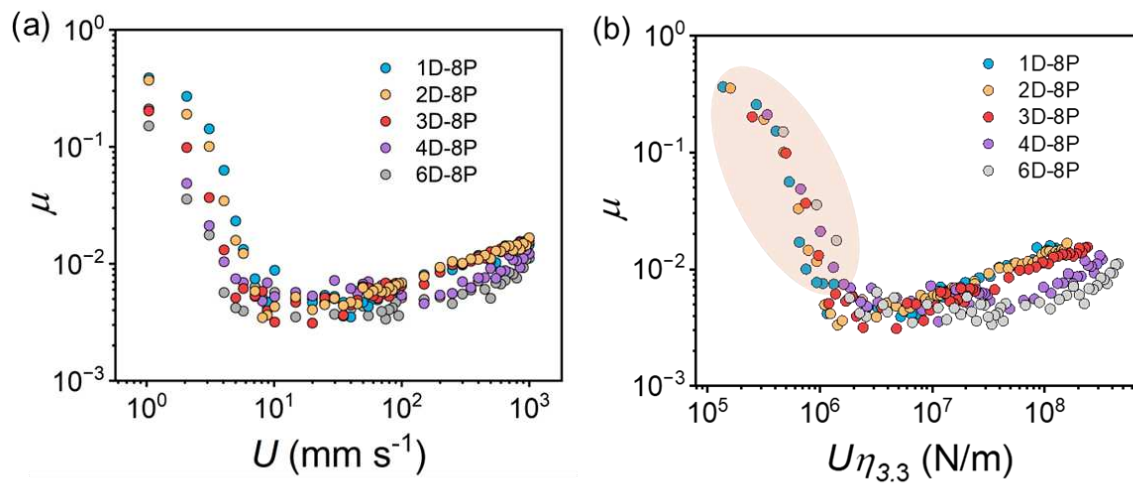
330 4D-8P (10 mPa.s) are lower than those of 1D-8P, 2D-8P and 3D-8P (60 mPa.s), suggesting
 331 that droplet deformation and entrainment contribute additionally to shear thinning behaviour
 332 of the lubricants. To further clarify the tribological effect of droplets, a pure P solution (6P)
 333 with a η at high-shear rates comparable to that of 4D-8P was selected for tribological
 334 comparison (Fig. S3). 6P and 4D-8P show comparable μ values in the EHL regime due to
 335 their similar high shear rate η , whereas 4D-8P exhibits a markedly lower μ in the mixed
 336 regime. No 'flat' mixed-lubrication regime was observed for 6P. This indicates that the unique
 337 mixed-lubrication regime in D/P emulsions can be attributed to droplet deformation and
 338 entrainment within the contact.



339

340 **Figure 7.** Coefficient of friction (μ) versus entrainment speed (U) for W/W emulsions with

341 fixed concentration of P (8 wt.%) in PDMS-PDMS contact: (b) 1D-8P, (c) 2D-8P, (d) 3D-8P,
 342 (e) 4D-8P, (f) 6D-8P. The results for 8P alone and water are plotted in (a) as control. The
 343 green line and yellow line in (b-f) show the linear fits to the mixed and hydrodynamic regimes
 344 for the 1D-8P, 2D-8P, 3D-8P, 4D-8P and 6D-8P system, respectively. Data represents means
 345 and standard deviations of triplicate samples ($n = 3 \times 3$).



346
 347 **Fig. 8** Stribeck curves for D/P emulsions with [P] = 8 wt.% and variable [D] (1, 2, 3, 4 and 6
 348 wt%) (a) as a function of entrainment speed (U) and (b) U multiplied by the viscosity
 349 measured at $\dot{\gamma} = 3.3 \text{ s}^{-1}$ ($U\eta_{3.3}$).

350

351 4. Conclusions

352 This study demonstrates that the microstructural, rheological, and tribological properties
 353 of W/W emulsions composed of D and P are highly sensitive to their composition. Increasing
 354 [P] primarily enhances the overall bulk viscosity, while higher [D] intensifies shear-thinning
 355 behavior by promoting larger and more deformable droplets. It is supposed that these
 356 droplets, under shear, can transform into string-like structures that enter the tribological

357 contact zone, thereby reducing the effective bulk viscosity and delaying the formation of a
358 full fluid film. Most notably, the emulsions exhibited an unusually speed-independent mixed
359 lubrication regime, with friction coefficients below 0.01 over a wide range of entrainment
360 speed, a behavior distinctly different from that of the individual polymer solutions. The lower
361 friction observed in emulsions with higher [D] in the EHL regime further underscores the
362 critical role of droplet deformation in governing lubrication performance of these W/W
363 systems. Future studies should also focus on various volume fractions of the droplets along
364 the tie-line to see how rheology affects tribology in these biphasic systems. Overall, these
365 findings highlight the potential of tailoring W/W emulsions via phase-separating polymers to
366 achieve unique lubrication characteristics, offering promising avenues for the development
367 of advanced, oil-free lubricants for biomedical applications.

368

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379 Methodology, Project administration; Writing- Reviewing & Editing, Supervision.

380

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Rheology and tribology of dextran/ polyethylene oxide-based water-in-water emulsions

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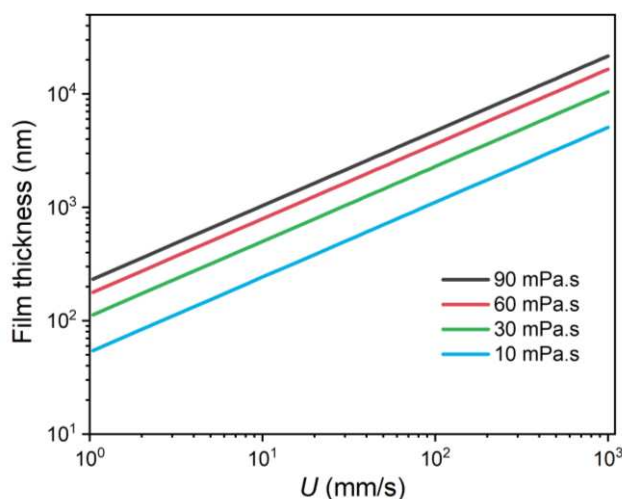
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484 Numerical model to predict the lubricating film thickness

485 Based on the materials characteristics of the tribopairs and lubricants, we have firstly
 486 identified that the EHL in this study belongs to isoviscous-elastic, namely 'soft-EHL' regime
 487 [1,2]. Therefore, the lubricating film thickness could be estimated according to Equation 1.
 488 [1,2].

$$489 \quad h_c = 3.28(U^{0.66}\eta^{0.66}E'^{-0.45}R^{0.76}W^{-0.21}) \quad a)$$

490 where η is the viscosity of the lubricant, U is the entrainment speed, E' is the reduced elastic
 491 modulus, R is the reduced radius in the entrainment direction and W is applied normal load.
 492 While the E' , R , and W are measured parameters, the viscosities, 10, 30, 60, 90 mPa.s, are
 493 estimated parameters from fitting the theoretical friction coefficient, μ , to the experimentally
 494 obtained μ vs U plots (Figure 6 & 7). Here, theoretical μ values represent the friction within
 495 the fluid in full-fluid film hydrodynamic regime (elaborated below).



496
 497 **Figure S1.** Calculated film thickness for PDMS-PDMS contact by applying equation 1 and
 498 effective viscosities of 10, 30, 60, 90 mPa.s.

499 500 Calculation of theoretical friction coefficient (μ) in full-fluid film hydrodynamic regime 501 for PDMS- PDMS contact.

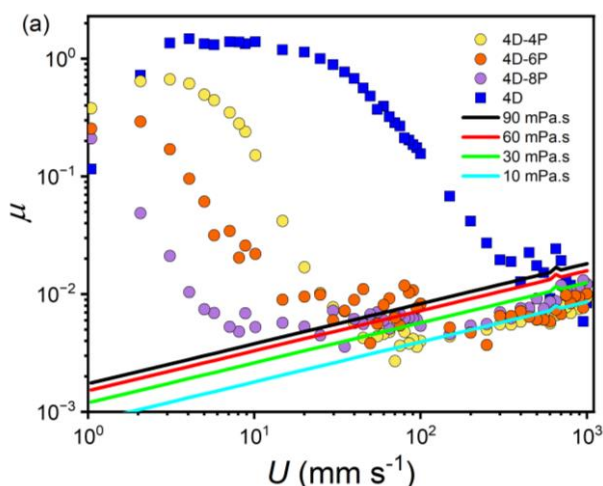
502 The μ in the hydrodynamic regime could be estimated according to the Equation 2 (where
 503 A is the contact area calculated with Hertzian contact equation) [2].

$$504 \quad \mu = \frac{U^{0.34}\eta^{0.34}E'^{0.45}A}{3.28W^{0.79}R^{0.76}} \quad b)$$

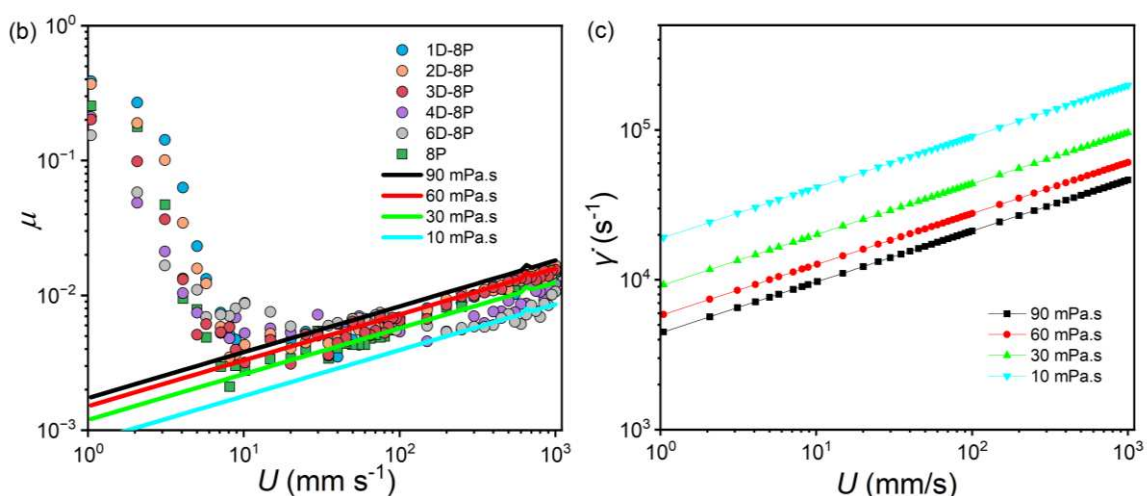
505 Thus, μ should be linearly proportional to U in log-log scale. Figures S2(a) and S2(b) present

506 an excellent agreement between theoretical μ values and experimental values. In turn, this
 507 fitting led us to estimate effective viscosities (90, 60, 30, and 10 mPa·s) for different solution
 508 samples, including pure D, pure P solutions, and D/P emulsions with constant [D] and [P].
 509 Lastly, based on the estimated effective viscosities, we could estimate the shear rate, $\dot{\gamma}$, in
 510 the tribological contact (Figure S2c), by applying the relation: $\dot{\gamma}$ = entrainment speed (U)/film
 511 thickness (h_c).

512



513



514 **Figure S2.** Stribeck curves for (a) D/P emulsions with constant [D] and pure D, and (b) D/P
 515 emulsions with constant [P] and pure P, overlaid with theoretically calculated friction
 516 coefficients (μ) based on hydrodynamic lubrication at viscosities of 10, 30, 60, and 90 mPa·s.
 517 (c) Theoretically calculated shear rates ($\dot{\gamma}$) within the PDMS–PDMS contact as a function of
 518 entrainment speed (U). The shaded region denotes the invalid data range before the full-
 519 fluid film hydrodynamic regime.

520 It should be noted that the film thickness calculated according to the Equation 1 is valid *in*
 521 *practice* only in the speed range where EHL regime is activated. This can be readily judged
 522 from the transition of μ to a linear increase in μ vs U plots in log-log scale (Figure 6, 7 or

Figure S1, S2). At lower speeds, the two surfaces are in direct contacts that the calculated film thickness is meaningless. The threshold speed that this transition occurs is different from sample to sample. Thus, the range of speed for full-fluid films for each sample and respective estimated film thickness at the lowest and highest speeds, along with the estimated viscosities, are presented in the Table S1.

528

Table S1. The list of fitted viscosities, η , the minimum, U_{mini} , and maximum, U_{max} , speeds for the full-fluid hydrodynamic regimes, and the calculated film thickness at the respective speeds for all the emulsion samples.

Samples	Fitted η (mPa.s)	U_{mini} of full- fluid hydrodynamic regime (mm/s)	Cal. film thickness at U_{mini} (μm)	U_{max} of full- fluid hydrodynamic regime (mm/s)	Cal. film thickness at U_{max} (μm)
4D	/	/	/	/	/
4D-4P	10	100	1.11	1000	5.06
4D-6P	10	100	1.11	1000	5.06
4D-8P	30	150	2.99	1000	10.45
1D-8P	60	40	1.97	1000	16.50
2D-8P	60	45	2.14	1000	16.50
3D-8P	60	80	3.12	1000	16.50
6D-8P	10	250	2.03	1000	5.06
8P	60	8	0.69	1000	16.50

532

It should be also noted that the calculated film thickness could be smaller than *real* film thickness due to surface roughness of the tribopairs. However, the surface roughness (R_q) of PDMS disc is known to be less than 2 nm and that of ball is ca. 120 nm [3], whereas the film thickness is typically a few μm as shown in Table S1. Thus, the influence of surface roughness is not likely substantial.

Overall, the lubricating film thickness (Table S1) is generally smaller than the diameters of the droplets in the emulsion (8.1 to 46.7 μm , Figure 2). However, as the droplets are highly deformable, it is still possible that they are entrained into the contact zone in a deformed, string-like shape to comprise the lubricating film, especially with increasing speed.

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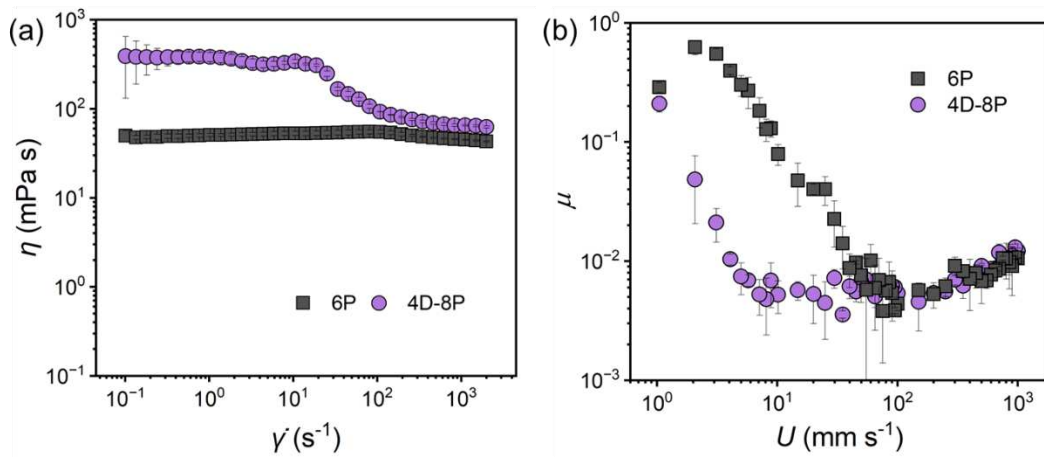


Figure S3. Apparent viscosity (η) (a) as a function of shear rate ($\dot{\gamma}$) and coefficient of friction (μ) (b) as a function of entrainment speed (U) for pure polymer (6P) and W/W emulsion (4D-8P) with comparable viscosity.

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