High-resolution radio observations of massive protostars in the SARAO MeerKAT Galactic Plane Survey

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Accepted 2025 April 22. Received 2025 April 19; in original form 2025 February 27

ABSTRACT

We present high-resolution observations made with the Australia Telescope Compact Array (ATCA) at 5.5 and 9.0 GHz of a sample of 28 massive protostars. These protostars were initially detected at radio wavelengths in the MeerKAT Galactic Plane Survey at 1.3 GHz, where they were unresolved with an angular resolution of approximately 8 arcsec. The resolution of the ATCA observations at 5.5 and 9.0 GHz is significantly higher at 2.0 and 1.2 arcsec, respectively. The highest angular resolution corresponds to a linear resolution of 1920 au for our nearest object. This improvement in resolution enabled us to resolve the components of extended emission and to more accurately determine the nature of the objects. Approximately 80 per cent of the detections were classified as jets, with 68 per cent of the jets found to be associated with non-thermal emission.

Key words: radiation mechanisms: non-thermal-stars: jets-stars: massive-stars: protostars.

1 INTRODUCTION

Massive protostars form through the accretion of materials that are often accompanied by the ejection of a fraction of the accreted material, closely linking the jet-driving phase to star formation processes (Carrasco-González et al. 2021). This ejection has been proposed to be either magnetically (Shu et al. 1994b; Pudritz et al. 2007; Carrasco-González et al. 2010) or radiatively (Carrasco-González et al. 2015; Kuiper, Yorke & Turner 2015) launched, accelerated, and collimated. The ejections play a crucial role as they drive mass away from the accretion discs, thereby conserving angular momentum and facilitating further accretion on to the protostar (Lovelace, Romanova & Bisnovatyi-Kogan 1995). Despite their astrophysical significance, the fundamental processes underlying the ejection of mass to form these jets remain unresolved. This is mainly due to the limitations in current observations, which lack the angular resolution required to resolve the inner 100 au of the jets, where launching and collimation occur (Carrasco-González et al. 2021). Indeed, the most current observational representation of a massive protostellar jet is that of the Cep A HW2 radio jet (Carrasco-González et al. 2021), where the outflowing material is composed of two components: a wide-angle spherical wind launched from the protostellar disc, and a highly collimated jet emanating from approximately 20-30 au from the protostar.

The protostars form in dusty environments where they are typically deeply embedded in their natal clouds. Their jets serve as the earliest indicators of massive star formation at radio wavelengths. These radio jets are now known to have components that may exhibit positive or negative spectral indices, depending on the position of the radio source relative to the central driving source (Garay et al. 2003; Guzmán et al. 2012; Moscadelli et al. 2016; Purser et al. 2016; Anglada, Rodríguez & Carrasco-González 2018; Obonyo et al. 2019). The spectral index, α , is defined by the relation $S_{\nu} \propto \nu^{\alpha}$, where S_{ν} is the flux density at frequency ν . Generally, the region associated with a luminous infrared (IR) source exhibit a positive index, while the more separated lobes may display either positive or negative spectral indices (Purser et al. 2016; Obonyo et al. 2019, 2021; Sanna et al. 2019) at centimetre wavelengths. The thermal component of the jet that is associated with an IR source is known as a core, while the separated parts are called lobes (Purser et al. 2016; Obonyo et al. 2019; Sanna et al. 2019).

Moreover, it is increasingly becoming evident that about 40– 50 per cent massive protostellar jets are associated with non-thermal emission, with the rest emitting free–free emission. This split could indicate an evolutionary sequence, consistent with the current theory where the envelope of an evolving massive protostar undergoes a convective phase before becoming fully radiative (Hosokawa & Omukai 2009). While theories of massive star formation typically suggest the formation of jets before H II regions (Tanaka, Tan & Zhang 2016), the observation by Guzmán et al. (2016) manifested a contemporaneous jet and H II region in the same object.

Advances in telescope sensitivity and resolution have significantly improved the detection of massive protostars, including the phases exhibiting weak radio emissions, such as the jet-driving phase (Purser et al. 2016; Obonyo et al. 2019; Sanna et al. 2019). None the less, previous Galactic Plane surveys e.g. Wood & Churchwell (1989), Urquhart et al. (2009), and Irabor et al. (2023) with typical sensitivities $\simeq 0.32$, 0.2, and 0.11 mJy beam⁻¹, respectively, were not sensitive enough to detect large numbers of the massive young stellar object (MYSO) jets. Notably, the SARAO MeerKAT Galactic Plane

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Table 1. Table showing object names, associated IRAS objects, positions (Goedhart et al. 2024) where observation pointing centres were set to, distance from the Sun, and Bolometric luminosities that were taken from Lumsden et al. (2013) catalogue. Local rms noise levels of the maps and the integrated fluxes of the sources in the MeerKAT's 1.3 GHz observation are also shown.

RMS name	Associated IRAS Object	RA (J2000)	Dec. (J2000)	D (kpc)	L_{Bol} (L_{\odot})	$rms~(5.5~GHz) \\ (\mu Jy~beam^{-1})$	$rms (9.0 GHz) \\ (\mu Jy beam^{-1})$	Int. flux (mJy) (1.3 GHz)
G270.8247-01.1112	IRAS 09088-4929	09 ^h 10'30.89"	-49°41′29.8″	6.3	7100	15	21	$0.09 {\pm} 0.03$
G271.2225-01.7712	IRAS 09075-5013	$09^{h}09'11.16''$	$-50^{\circ}25'55.3''$	8.6	4800	25	24	$0.80 {\pm} 0.08$
G281.2206-01.2556	IRAS 09598-5632	10 ^h 01'33.78"	$-56^{\circ}47'18.2''$	7.0	3900	25	24	$0.08 {\pm} 0.02$
G282.8969-01.2727	-	10 ^h 11'31.60"	$-57^{\circ}47'03.7''$	7.0	17 000	38	27	$0.75 {\pm} 0.05$
G287.3716+00.6444	IRAS 10460-5811	$10^{h}48'04.55''$	$-58^{\circ}27'01.5''$	4.5	18 000	27	37	4.22 ± 0.38
G287.8768-01.3618	IRAS 10423-6011	$10^{h}44'17.93''$	$-60^{\circ}27'46.0''$	6.1	19 000	28	26	17.4 ± 0.3
G289.0543+00.0223	-	$10^{h}57'42.91''$	$-59^{\circ}44'58.4''$	7.3	2900	27	24	0.30±0.04
G293.0352-00.7871	IRAS 11233-6143	$11^{h}25'33.22''$	$-61^{\circ}59'51.3''$	3.3	1900	23	17	$0.78 {\pm} 0.05$
G297.1390-01.3510	IRAS 11563-6320	$11^{h}58'54.99''$	$-63^{\circ}37'47.8''$	8.9	12 000	24	17	$1.34{\pm}0.16$
G300.7221+01.2007	IRAS 12300-6119	$12^{h}32'50.15''$	$-61^{\circ}35'27.0''$	4.3	1600	33	17	$0.38{\pm}0.02$
G301.0130+01.1153	-	$12^{h}35'14.47''$	$-61^{\circ}41'46.8''$	4.3	2400	85	48	0.21 ± 0.03
G304.7592-00.6299	IRAS 13046-6310	$13^{h}07'47.42''$	$-63^{\circ}26'37.0''$	11.2	4100	32	19	$0.63 {\pm} 0.02$
G309.9796+00.5496	IRAS 13475-6115	$13^{h}51'02.72''$	$-61^{\circ}30'14.1''$	3.5	7600	64	21	$0.80{\pm}0.03$
G314.3197+00.1125	IRAS 14227-6025	$14^{h}26'26.28''$	$-60^{\circ}38'31.5''$	3.6	38 000	26	19	0.41 ± 0.03
G328.2523-00.5320A	IRAS 15541-5349	$15^{h}57'59.82''$	$-53^{\circ}58'00.4''$	2.9	8000	220	95	$0.92 {\pm} 0.03$
G328.3442-00.4629	IRAS 15543-5342	$15^{h}58'09.62''$	$-53^{\circ}51'18.2''$	2.9	2400	305	57	$0.93 {\pm} 0.03$
G328.9842-00.4361	IRAS 15574-5316	16 ^h 01'19.12"	$-53^{\circ}25'04.2''$	4.7	1800	23	17	$0.63 {\pm} 0.03$
G329.0663-00.3081	IRAS 15573-5307	16 ^h 01'09.93"	$-53^{\circ}16'02.3''$	11.6	36 000	24	16	$0.33 {\pm} 0.02$
G334.7302+00.0052	IRAS 16223-4901	16 ^h 26'04.65"	$-49^\circ08'41.8''$	2.5	3800	27	27	$0.12{\pm}0.01$
G335.0611-00.4261A	IRAS 16256-4905	16 ^h 29'22.98"	$-49^{\circ}12'27.0''$	2.8	4200	37	20	$1.18{\pm}0.16$
G338.2801+00.5419A	IRAS 16344-4605	16 ^h 38'09.11"	$-46^{\circ}11'04.0''$	4.1	5100	26	16	$0.43 {\pm} 0.08$
G338.4712+00.2871	IRAS 16363-4606	16 ^h 39'58.90"	$-46^{\circ}12'36.4''$	13.1	170000	60	84	$0.80{\pm}0.03$
G339.3316+00.0964	IRAS 16404-4535	16 ^h 44'04.39"	$-45^{\circ}41'27.1''$	13.1	17 000	52	27	$2.12{\pm}0.04$
G340.7455-01.0021	-	$16^{h}54'04.05''$	$-45^{\circ}18'50.0''$	2.6	5300	978	877	$1.63 {\pm} 0.05$
G345.2619-00.4188A	IRAS 17033-4119	$17^{h}06'50.54''$	$-41^{\circ}23'46.4''$	2.7	1200	23	30	0.27 ± 0.03
G345.5043+00.3480	IRAS 17008-4040	$17^{h}04'22.87''$	$-40^{\circ}44'23.5''$	2.0	100 000	1923	1937	$1.77 {\pm} 0.07$
G345.7172+00.8166A	IRAS 16596-4012	$17^{h}03'06.40''$	$-40^{\circ}17'08.7''$	1.6	1100	27	21	$0.68 {\pm} 0.03$
G345.9561+00.6123	IRAS 17012-4009	17 ^h 04'43.00"	$-40^{\circ}13'13.4''$	2.5	2400	23	19	< 0.09

Survey (SMGPS; Goedhart et al. 2024), conducted at 1.3 GHz with a sensitivity $\simeq 0.01-0.02$ mJy beam⁻¹, detected a population of radioweak MYSOs in the Southern hemisphere (Obonyo et al. 2024). The low resolution of the observation, however, did not resolve most of the jets, which were identified statistically (Lumsden et al. 2013) using their mid- and near-IR colours and through visual inspection.

While the IR studies of MYSOs have the potential to identify most of these objects, there is a possibility of contamination from HII regions. However, these objects can be resolved using the Australia Telescope Compact Array (ATCA) at higher frequencies, revealing details of their morphologies. For instance, Purser et al. (2016) reported resolved jets in their ATCA observations, where radio emission manifested the separation of some ionized jets' cores from their lobes. In this work, we report a simultaneous observation of the objects with ATCA at both 5.5 and 9.0 GHz to reveal their morphologies and to constrain their spectral indices.

2 OBSERVATION AND DATA

2.1 Sample selection

27 of the MYSOs in Obonyo et al. (2024) had not been studied in previous higher frequency radio studies (e.g. Helfand et al. 2006; Urquhart et al. 2007). This was either due to their low declinations, $\delta \leq 40^{\circ}$, or their low surface brightness at radio wavelength, and were selected for the study. Indeed many of the objects lie within the CORNISH-South survey area but were too weak to detect (Irabor et al. 2023). Additionally, G287.3716+00.6444, which was classified

by Purser et al. (2016) as an H II region, was included in the sample to search for variability. All the selected objects were unresolved in the MeerKAT observation and did not reveal any clear jet morphology even though their IR images display bipolar/monopolar nebulae, indicative of outflow activity. Table 1 shows a list of the selected MYSOs; columns 1–9 show the RMS¹ name, associated IR Astronomical Satellite object, right ascension, declination, distance, bolometric luminosities, rms noise at 5.5 GHz, rms noise at 9.0 GHz, and the integrated fluxes of the radio sources at 1.3 GHz (Obonyo et al. 2024), respectively.

2.2 ATCA observation

The observations were made between 2022 November 21 and 2023 January 16 using the ATCA's 6C array configuration at 5.5 and 9.0 GHz. The observations, done under the project code C3508, made use of all the six antennas of the configuration where the shortest and longest baselines are 153 and 6000 m, corresponding to angular resolutions of 2 and 1 arcsec at 5.5 GHz (C band) and 9.0 GHz (C band), respectively. The bandwidth for each of the frequencies of observation was 2 GHz. The total on-source integration time at each frequency was 72 min, resulting from 12 uv-cuts spread over a 12-h track to provide good uv-coverage. The sources were observed in groups of four sources, for 6 min each, followed by a 1-min observation of a phase calibrator to correct for atmospheric and

instrumental effects on the amplitudes and phases of the data. This was repeated for a further set of three sources.

Calibration and data reduction were performed using the MULTI-CHANNEL IMAGE RECONSTRUCTION, IMAGE ANALYSIS AND DISPLAY astronomy software (Sault, Teuben & Wright 1995). Whenever a radio source with $SNR \ge 10$ was present within the primary beam of the telescope, we used the source to perform self-calibration on the phases of the data. Maps of the fields were made by performing Fourier transform on the visibilities with robust weighting (Briggs 1995) of robustness parameter r = 0.5, obtaining synthesized beams whose typical full width at half-maximum are 2.1 arcsec \times 1.8 arcsec and 1.3 $\operatorname{arcsec} \times 1.1$ arcsec at 5.5 and 9.0 GHz, respectively. The dirty maps obtained were deconvolved using MIRIAD's MFCLEAN task. The average root-mean-square (rms) noise of the deconvolved maps typically lies within the range 22–50 and 15–50 μ Jybeam⁻¹ at 5.5 and 9.0 GHz, respectively. However, in a few of the fields, e.g. in G345.5043+00.3480's field, the rms is almost a factor of a hundred higher than the typical noise of the observation. This is largely due to the presence of bright extended emission that lowers the dynamic range in such a field. The flux density and bandpass calibrator for the observation was 1934-638 (3C84).

The astrometry of the MeerKAT and ATCA data sets were examined as both were used in the analysis. Goedhart et al. (2024) estimated the positional accuracy of MeerKAT observations by comparing the coordinates of point sources to their corresponding positions in the CORNISH and CORNISH-South surveys (Hoare et al. 2012; Irabor et al. 2023). They reported a systematic positional error, with angular offsets ranging from 0 to 2.1 arcsec, as illustrated in fig. 6 in Goedhart et al. (2024). This suggests the potential for positional discrepancies of up to 2.1 arcsec between MeerKAT and ATCA data.

3 RADIO PROPERTIES OF THE MASSIVE PROTOSTARS MANIFESTING LOW SURFACE-BRIGHTNESS IN THE SARAO MEERKAT GPS

We detected centimetre radio continuum emission from 19 sources at 5.5 GHz and 20 sources at 9.0 GHz at 3σ level and above. The positions of each of the detected sources were derived from 2D Gaussian fits at each frequency and found to be consistent at all frequencies within 0.02–0.08 arcsec, with the largest difference registered in sources where structures are resolved at one frequency and not the other. This range of uncertainties is also in agreement with the accuracies stated in the ATCA information manual. The positions of the sources at 5.5 and 9.0 GHz are listed in Table 2, with the sources that were not detected at the frequencies of observation represented by dashes.

The other observed parameters of the sources: flux densities, deconvolved angular sizes, position angles (PAs), and peak fluxes of the MYSOs at 5.5 and 9.0 GHz are presented in Table 3. The fluxes for each of the sources were estimated in two methods: using the MIRIAD's task, imfit, that fits a two-dimensional Gaussian function to the detections and secondly by summing up the emission within a region encompassing the source. For sources with irregular or non-Gaussian morphologies, the fluxes were estimated by calculating the total flux density within a box enclosing the source. In most cases, the source morphologies were well estimated by Gaussian fits, with the major and minor axes listed for each source being the deconvolved dimensions derived from the fits.

The range, average, and standard deviation of the fluxes at 5.5 GHz are 0.07-3.86, 0.67, and 0.84 mJy, respectively, with a majority

of the sources having integrated flux densities $S_{\nu} < 1.5$ mJy. Two of the sources, G287.3716+00.6444-A and G340.7455-01.0021, however, have flux densities $S_{\nu} > 3$ mJy, i.e. 3.86 ± 0.21 and 3.13 ± 0.26 mJy, respectively. A comparable trend is seen at 9.0 GHz, with the same objects manifesting the highest fluxes again at 9.0 GHz.

3.1 Morphologies of the sources

The morphologies of radio jets from massive protostars can provide insights into the formation and evolution of the protostars and the associated jets (Cesaroni et al. 2018). The continuum images of the ATCA observations at 5.5 and 9.0 GHz manifest a higher degree of detail of the unresolved 1.28 GHz MeerKAT images, revealing their morphologies in the majority of cases. Indeed, this is the first high-resolution radio observation of 27 massive protostars in the sample, contributing to the census of massive protostellar jets.

The 20 objects detected at either 5.5 GHz, 9.0 GHz, both, as shown in Fig. 1, exhibit a range of or Seven of these objects display multiple morphologies. components, characteristic of a jet core-lobe(s) system. Examples include G282.8969-01.2727, G287.3716+00.6444, G309.9796+00.5496, G314.3197+00.1125, G328.2523-00.5320, G335.0611-00.4261A, and G338.2801+00.5419A. Another nine objects exhibit single extended elliptical emission, typical of jet-driven structures, such as G271.2225-01.7712, G287.8768-01.3618, G297.1390-01.3510, G328.9842-00.4361, G334.7302+00.0052, G338.4712+00.2871, G339.3316+00.0964, G340.7455-01.0021, and G345.7172+00.8166A. In contrast, objects such as G293.0352-00.7871 and G304.7592-00.6299 exhibit circular morphologies. The observed morphologies were key in identifying and classifying the objects, for example, objects with elliptical morphologies and multiple components are likely to host jet-driving sources.

Typically, collimated jets are manifested by narrow streams of radio emission that extend over large distances with minimal widening, often exhibiting internal shocks and knots along their lengths, e.g. HH 80–81 (Carrasco-González et al. 2010; Rodríguez-Kamenetzky et al. 2017). While the objects in our sample do not manifest clear collimation over large distances, some of the jets show components that resemble core-lobe systems of jets, with others manifesting non-thermal emission suggesting that they host nonthermal lobes. Moreover, some of the sources show elliptical shapes of a few arcseconds, similar to the structures of cores of radio jets (Anglada et al. 2018).

3.2 Spectral indices of the objects

The spectral indices of the sources were calculated using their flux densities at 5.5 and 9.0 GHz. To ensure consistency, the data at both frequencies were re-imaged by restricting the visibilities to baselines sensitive to emission from comparable angular scales. This was achieved by limiting the uv-range to a common range of $3-90 \text{ k}\lambda$, applied to both frequency bands before estimating the indices. The full uv-coverage for the observations at 5.5 and 9.0 GHz were $1.8-89.5 \text{ k}\lambda$ and $3.1-159.5 \text{ k}\lambda$, respectively. Notably, most fluxes from the common uv-range agree with those from the full uv-range within uncertainties. Flux densities at 1.3 GHz were excluded from the calculation of the indices due to their dissimilar uv-range compared to the other observations, and even if re-imaged, would not result in significant detections. However, for sources detected at all three frequencies, trends of their spectral energy distributions

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Table 2.	Coordinates of sources detected at 5.5 GHz, 9.0 GHz, or both frequencies. A - indicates non-detection at a given frequency or cases where the source's
position of	could not be determined due to being resolved. The table also includes the spectral indices of the detected sources.

RMS name	Position	at C band	Position	at X band
	RA(J2000)	Dec.(J2000)	RA(J2000)	Dec.(J2000)
G270.8247-01.1112	09 ^h 10'30.77"	-049°41′27.06″	_	_
G271.2225-01.7712	$09^{h}09'11.17''$	$-050^{\circ}25'55.04''$	$09^{h}09'11.17''$	$-050^{\circ}25'55.08''$
G281.2206-01.2556	_	_	_	-
G282.8969-01.2727-A	$10^{h}11'31.56''$	$-057^{\circ}47'03.95''$	_	_
G282.8969-01.2727-В	$10^{h}11'31.71''$	$-057^{\circ}47'04.99''$	_	_
G287.3716+00.6444-A	$10^{h}48'04.66''$	$-058^{\circ}27'02.10''$	$10^{h}48'04.65''$	-058°27′02.21″
G287.3716+00.6444-B	$10^{h}48'05.19''$	$-058^{\circ}26'57.17''$	$10^{h}48'05.32''$	$-058^{\circ}26'57.78''$
G287.8768-01.3618	$10^{h}44'17.91''$	$-060^{\circ}27'45.08''$	$10^{h}44'17.94''$	-060°27′45.33″
G289.0543+00.0223	_	_	$10^{h}57'42.89''$	$-059^{\circ}44'57.65''$
G293.0352-00.7871	_	_	$11^{h}25'33.14''$	-061°59′51.55″
G297.1390-01.3510	$11^{h}58'54.95''$	$-063^{\circ}37'47.67''$	_	_
G297.1390-01.3510A	_	_	$11^{h}58'54.86''$	$-063^{\circ}37'47.90''$
G297.1390-01.3510B	_	_	$11^{h}58'54.90''$	-063°37′46.41″
G300.7221+01.2007	_	_	_	_
G301.0130+01.1153	_	_	_	_
G304.7592-00.6299	$13^{h}07'47.42''$	-063°26′36.76″	$13^{h}07'47.40''$	$-063^{\circ}26'37.08''$
G309.9796+00.5496	_	_	_	_
G309.9796+00.5496A	$13^{h}51'02.70''$	-061°30′13.43″	$13^{h}51'02.66''$	-061°30′13.63″
G309 9796+00 5496B	$13^{h}51'03.00''$	-061°30′13.15″	$13^{h}51'03.06''$	$-061^{\circ}30'12.71''$
G314.3197+00.1125-A	$14^{h}26'26.26''$	-060°38′31.29″	$14^{h}26'26.30''$	$-060^{\circ}38'31.31''$
G314 3197+00 1125-B	$14^{h}26'26.12''$	$-060^{\circ}38'29\;44''$	_	_
G314 3197+00 1125-C	$14^{h}26'26.66''$	-060°38′39 88″	$14^{h}26'26~64''$	-060°38′39.91″
G328 2523-00 5320A	_	_	_	-
G328 2523-00 5320A-A	15 ^h 57'59 46''	-053°58′01 56″	$15^{h}57'59~44''$	-053°58′01 70″
G328 2523-00 5320A-B	$15^{h}57'59~82''$	$-053^{\circ}57'59.99''$	$15^{h}57'59.83''$	-053°58′00.23″
G328 2523-00 5320A-C	$15^{h}57'59.05''$	-053°58′03 39″	$15^{h}57'59.05''$	-053°58′03.12″
G328.3442-00.4629	-		-	-
$G_{328} 9842 - 004361$	$16^{h}01'19.09''$	-053°25′03 76″	$16^{h}01'19\ 12''$	-053°25′04.06″
G320.0663_00.3081		-		-
$G_{334} 7_{302} \pm 00.0052$	16 ^h 26'04 65"	-049°08'42 03"	16 ^h 26'04 66"	-040008/42 22/
$G_{335} = 0.0000000000000000000000000000000000$	$16^{h}20'23 16''$	$-049^{\circ}0842.03^{\circ}$		-049 00 42.22
$G_{335,0011=00.4201A=(A + B)}$	10 29 25.10	-049 12 28:25	16 ^h 29'23 19"	_040°12′28 84″
G335.0611_00.4261A_R	_	_	$16^{h}20'23.11''$	-049' 12' 20.04' $040^{\circ} 12' 27' 30''$
G335.0611-00.4261A-C	$\frac{-}{16^{h}20'23}$ 13"		16 ^h 29'23.09"	$-049^{\circ}12^{\circ}27.30^{\circ}$ $-049^{\circ}12'31^{\circ}40''$
G338 2801 + 00 5/10 A	16 ^h 38'08 07''	046°10′59.00″	16 ^h 38'08 06''	-049 12 51.40 046°10′58 81″
G228 2801 + 00 5410 A P	16 ^h 28'00.06"	$-040^{\circ}10^{\circ}39.00^{\circ}$	16 ^h 28'00.07"	$-040\ 10\ 30.01$ $046^{\circ}11'02\ 27''$
G338 4712 + 00 2871	16 ^h 30'58 50"	-040 11 03.43 $046^{\circ}12'37$ $00''$	16 ^h 30'58 50"	$-040\ 11\ 03.27$ $046^{\circ}12'37\ 06''$
$G_{220}^{2216+00.2671}$	$16^{h}44'04~26''$	-040 12 37.90 $045^{\circ}41'27$ 51″	16 ^h 44'04 26''	$-040^{\circ}12^{\circ}37.90^{\circ}$
$G_{240,7455} = 01,0021$	16h54/04.02″	-045 41 27.51 $045^{\circ}18'40 64''$	$16^{+}44^{-}04^{-}.50^{-}$	$-045 \ 41 \ 27.59$ $045^{\circ}18'50 \ 02''$
C245 2610 00 4188A	$10^{-54} 04.05$	-043 18 49.04 $041^{\circ}22'44$ 20''	10 54 04.05	-045 18 50.02
G345 5043 + 00 3480	17 00 30.70	-041 23 44.39	_	_
C_{245} 7172 + 00 8166 *	17/02/06 27/	 040°17′0° 91″	17/02/06 20/1	
G245.0561+00.6122	17 03 00.57	-040 17 08.81	17 05 00.59	-040 1/ 08.81
0343.9301+00.0123	—	—	—	_

(SEDs) were inspected to help characterize the nature of their radio sources.

For radio sources detected at only one frequency, their spectral indices were constrained by setting upper or lower limits. These limits were determined using the measured flux at the detected frequency and three times the noise level (3σ) at the undetected frequency. Additionally, if a source was resolved into multiple components at a particular frequency, its spectral index was estimated by summing the fluxes of all components at that frequency.

Fig. 2 presents a histogram illustrating the distribution of spectral indices for the detected sources and their resolved components, while Fig. 3 displays the SEDs of the sources. Additionally, Table 3 lists the spectral indices, which range from approximately -1.3

to 1.7. This distribution indicates the presence of both thermal $(\alpha \ge -0.1)$ and non-thermal $(\alpha < -0.1)$ objects within the sample. The histogram shows two prominent peaks at $\alpha \simeq -0.57$ and 0.34, suggesting a mix of thermal and non-thermal radio emission in the jets. Notably, two sources exhibit highly negative spectral indices, e.g. $\alpha < -1.0$, atypical of synchrotron emission. Despite this, these sources, G297.1390-01.3510 and G338.4712+00.2871, were still classified as non-thermal jets or lobes.

Overall, the classification of the objects and their components based on their morphology, nature of emission, and association with IR and molecular emission are shown in Table 3. Out of the 28 objects observed, 22 were detected. Of these, 13 exhibited either non-thermal emission or were associated with non-thermal lobes, while some potential lobes were identified as thermal. Five radio sources were **Table 3.** A table of massive protostars and their components detected in the ATCA observations. Their deconvolved sizes, integrated fluxes, and peak fluxes at both 5.5 and 9.0 GHz are shown. – indicates a non-detection, combined emission, or when a source is split into components at a given frequency, making it impossible to determine its parameters at that frequency.

PMS name			hand naramete	2			X	hand naramata	9		5	Obiact
DIAMIN HAILIN	C-bane	d size	balance PA	Int. flux	Peak flux	X-bane	d size	uanu paramete PA	s Int. flux	Peak flux	3	Cujat
	$\theta_{maj}(\operatorname{arcsec})$	$\theta_{min}(\operatorname{arcsec})$	(_)	(mJy)	$(mJy beam^{-1})$	$\theta_{maj}(\operatorname{arcsec})$	$\theta_{min}(arcsec)$	(_)	(mJy)	$(mJy beam^{-1})$		
G270.8247-01.1112	2.38±0.92	0.65±0.91	54土43	0.13 ± 0.04	0.08 ± 0.02	1	I	I	I	<0.06	<-1.57	Lobe
G271.2225-01.7712	2.86 ± 0.32	1.64 ± 0.11	177±5	0.35 ± 0.05	0.24 ± 0.02	<1.19	<1.11	<169	0.28 ± 0.04	0.25 ± 0.02	-0.45 ± 0.09	Jet
G281.2206-01.2556	I	I	I	I	< 0.08	I	I	I	I	< 0.08	I	I
G282.8969-01.2727-A	I	I	I	0.09 ± 0.03	0.07 ± 0.01	I	I	I	I	< 0.08	<-0.24	Core
G282.8969-01.2727-B	I	I	Ι	0.07 ± 0.02	0.06 ± 0.01	I	I	I	I	< 0.08	< 0.27	Lobe
G287.3716+00.6444-A	1.83 ± 0.20	1.31 ± 0.22	137±19	3.86 ± 0.21	2.28 ± 0.09	1.42 ± 0.13	1.30 ± 0.15	112±79	4.43土0.28	1.90 ± 0.09	0.28 ± 0.04	Core
G287.3716+00.6444-B	3.94 ± 0.64	2.04±0.47	66±11	0.67 ± 0.10	0.20 ± 0.02	1.99 ± 0.22	1.25 ± 0.17	165±11	0.48 ± 0.05	0.14 ± 0.02	-0.68 ± 0.08	Lobe
G287.8768-01.3618	1.82 ± 0.32	0.19 ± 0.53	148 ± 10	0.58 ± 0.06	0.42 ± 0.03	2.28 ± 0.25	0.45 ± 0.24	161 ± 3	0.93 ± 0.09	0.40 ± 0.03	0.96 ± 0.06	Jet
$G289.0543 \pm 00.0223$	I	I	I	Ι	<0.09	1.32 ± 0.41	1.08 ± 0.38	83土91	0.22 ± 0.04	0.11 ± 0.02	>1.81	H II-region
G293.0352-00.7871	I	I	I	I	< 0.08	<0.81	<0.3	<166	0.14 ± 0.02	0.12 ± 0.02	>1.14	Unknown
G297.1390-01.3510	4.76 ± 0.54	1.06 ± 0.37	77±3	0.63 ± 0.07	0.19 ± 0.02	I	I	T	I	I	-1.25 ± 0.08	Jet
G297.1390-01.3510A	I	I	I	I	Ι	1.65 ± 0.23	1.21 ± 0.34	109±35	0.23 ± 0.04	0.09 ± 0.01	I	Lobe
G297.1390-01.3510B	I	I	I	Ι	Ι	1.47 ± 0.31	0.13 ± 0.46	78±10	0.12 ± 0.02	0.07 ± 0.01	I	Unknown
G300.7221+01.2007	I	I	I	Ι	<0.09	I	I	Ι	I	<0.05	I	I
G301.0130+01.1153	I	I	I	I	<0.26	I	I	I	I	<0.14	I	I
G304.7592-00.6299	2.51 ± 0.38	1.00 ± 0.48	6 ∓ 86	0.61 ± 0.07	0.31 ± 0.03	1.53 ± 0.33	0.82 ± 0.31	140土21	$0.38 {\pm} 0.06$	0.19 ± 0.02	-0.93 ± 0.09	H II-region
G309.9796+00.5496	I	I	I	0.46 ± 0.10	I	I	I	I	0.59 ± 0.09	I	0.51 ± 0.12	Jet
G309.9796+00.5496A	I	I	I	0.20 ± 0.08	0.16 ± 0.04	2.70 ± 0.67	1.00 ± 0.47	100 ± 10	0.30 ± 0.07	0.09 ± 0.02	0.82 ± 0.20	Core
$G309.9796 \pm 00.5496B$	1.74 ± 0.59	0.55 ± 0.64	30土47	0.26 ± 0.06	0.19 ± 0.03	1.83 ± 0.42	0.94 ± 0.39	53 ± 18	0.29 ± 0.05	0.11 ± 0.02	0.22 ± 0.13	Lobe
G314.3197+00.1125-A	I	I	I	0.08 ± 0.02	0.06 ± 0.01	1.28 ± 0.13	0.26 ± 0.19	116 ± 6	0.17 ± 0.02	0.11 ± 0.01	1.53 ± 0.12	Core
G314.3197+00.1125-B	I	I	I	0.07 ± 0.02	0.04 ± 0.01	I	I	I	I	<0.05	<-0.68	Lobe
G314.3197+00.1125-C	I	I	I	0.12 ± 0.02	0.05 ± 0.01	I	I	I	0.09 ± 0.03	0.08 ± 0.02	-0.58 ± 0.16	Unknown
G328.2523-00.5320A-A	I	I	I	0.48 ± 0.08	0.45 ± 0.04	I	I	I	0.42 ± 0.07	$0.34{\pm}0.05$	-0.27 ± 0.10	Lobe
G328.2523-00.5320A-B	I	I	I	0.33 ± 0.04	0.29 ± 0.03	I	I	I	0.54 ± 0.09	0.46 ± 0.06	1.00 ± 0.09	Core
G328.2523-00.5320A-C	I	I	I	0.23 ± 0.07	0.21 ± 0.02	1.53 ± 0.32	0.79 ± 0.33	34土12	$0.61 {\pm} 0.08$	$0.31 {\pm} 0.03$	1.98 ± 0.14	Unknown
G328.3442-00.4629	I	Ι	Ι	Ι	<0.63	Ι	Ι	I	< 0.17	I	Ι	Ι
G328.9842-00.4361	1.93 ± 0.57	0.24 ± 0.69	143土21	0.22 ± 0.04	0.16 ± 0.02	2.24 ± 0.70	0.32 ± 0.49	141 ± 13	$0.24 {\pm} 0.06$	0.12 ± 0.02	0.18 ± 0.13	Jet
G329.0663-00.3081	I	I	I	I	< 0.07	I	I	I	I	<0.05	I	I
G334.7302+00.0052	2.17 ± 0.15	1.69 ± 0.09	1 ± 9	0.31 ± 0.04	0.30 ± 0.02	1.12 ± 0.28	0.75 ± 0.33	150 ± 51	0.55 ± 0.05	$0.38 {\pm} 0.02$	1.20 ± 0.07	Jet
G335.0611-00.4261A-	3.13 ± 0.44	2.15 ± 0.35	160 ± 23	0.76 ± 0.08	0.30 ± 0.02	I	I	I	I	I	-0.38 ± 0.06	Jet
(A + B)												
G335.0611-00.4261A-A	I	I	I	I	I	0.73 ± 0.26	0.71 ± 0.38	32±83	0.25 ± 0.03	0.18 ± 0.02	I	Lobe
G335.0611-00.4261A-B	I	I	I	I	I	2.04 ± 0.24	0.86 ± 0.26	66 ± 8	0.38 ± 0.04	0.15 ± 0.02	I	Core
G335.0611-00.4261A-C	I	I	I	0.11 ± 0.02	0.09 ± 0.01	I	I	I	0.24 ± 0.04	0.14 ± 0.02	1.58	Unknown
$G338.2801 \pm 00.5419 A -$	I	I	I	0.99 ± 0.08	I	I	I	I	$0.74{\pm}0.04$	I	I	Jet
(A + B)												
G338.2801+00.5419A-A	1.87 ± 0.70	0.97 ± 0.50	$144{\pm}45$	0.41 ± 0.06	0.31 ± 0.03	1.18 ± 0.28	0.43 ± 0.20	167 ± 17	$0.28 {\pm} 0.03$	0.20 ± 0.01	-0.77 ± 0.08	Lobe
G338.2801+00.5419A-B	2.12 ± 0.47	1.46 ± 0.43	165±87	0.58 ± 0.06	0.35 ± 0.03	0.71 ± 0.16	$0.46 {\pm} 0.16$	177±55	$0.46 {\pm} 0.03$	$0.38 {\pm} 0.01$	-0.46 ± 0.05	Core
G338.4712+00.2871	3.82 ± 0.43	1.90 ± 0.12	162 ± 4	0.80 ± 0.11	0.52 ± 0.04	2.32 ± 0.43	1.12 ± 0.12	152±5	0.45 ± 0.11	0.32 ± 0.05	-1.17 ± 0.12	Jet
G339.3316+00.0964	2.54 ± 0.11	1.71 ± 0.05	16 ± 6	1.31 ± 0.08	1.28 ± 0.04	0.97 ± 0.19	0.69 ± 0.20	21 ± 46	1.60 ± 0.11	1.14 ± 0.05	0.41 ± 0.04	Jet
G340.7455-01.0021	1.95 ± 0.18	1.15 ± 0.06	171 ± 4	3.13 ± 0.26	3.54 ± 0.23	2.04 ± 0.25	0.89 ± 0.05	157±3	4.06 ± 0.58	2.91 ± 0.25	0.53 ± 0.08	Jet
G345.2619-00.4188A-A	5.87±1.16	1.83 ± 0.90	59±8	0.19 ± 0.06	0.09 ± 0.02	I	I	I	I	<0.09	<-1.52	Lobe
G345.2619-00.4188A-B	I	I	I	0.09 ± 0.03	0.08 ± 0.02	I	I	I	I	<0.09	<0.00	Lobe
$G345.5043 \pm 00.3480$	I	I	I	Ι	<16.80	I	I	I	I	<4.20	I	I
G345.7172+00.8166A	2.56 ± 0.21	2.32 ± 0.18	139±32	0.26 ± 0.04	0.20 ± 0.02	4.20 ± 1.22	1.01 ± 0.11	158±3	0.19 ± 0.05	0.13 ± 0.02	$-0.64{\pm}0.13$	Jet
G345.9561 + 00.6123	I	I	I	I	<7.10	I	I	I	I	<5.82	I	I

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Figure 1. IR images of the sources together with their L-, C-, and X-band contours of levels 3σ , 5σ , 7σ , 11σ , 15σ , 25σ , and 50σ ,... shown in green, grey, and blue colours, respectively. The synthesized beams for L, C, and X bands are shown in the lower left corner of each frame. The locations of mm-cores and methanol masers within the vicinity of the massive protostars are marked with asterisks (*) and plus (+) signs, respectively. The sources of the IR images are indicated in the lower right corners of the plots. GLIMPSE images are made from composite images of 3.6, 4.5, and 8.0 μ m.



Figure 1. (Continued.)



Figure 1. (Continued.)



Figure 1. (Continued.)



Figure 2. A histogram of the spectral indices of the radio sources detected in the ATCA observations. Indices of both the components of the objects, if resolved, as well as the objects are shown. The dashed vertical line at $\alpha = -0.1$ represents the typical dividing line between thermal and nonthermal sources.

classified as unknowns, two as H II-regions, and the remaining were categorized as jets. These results suggest that non-thermal emission, and by extension magnetic fields, plays a significant role in the formation of massive stars.

3.3 Nature of the objects

3.3.1 G270.8247-01.1112

The radio source associated with G270.8247–01.1112 is detected at both 1.28 and 5.5 GHz but not 9.0 GHz, where the rms noise level is 0.021 mJy beam⁻¹. Although its spectral index is highly negative, <-1.57, its flux density at 1.28 GHz is lower than expected based on the SED generated from 5.5 and 9.0 GHz. This discrepancy suggests that the source may be a non-thermal lobe with a typical index of -0.6. Additionally, its position is offset to the north-west of IRS A, the potential jet driver, further supporting the interpretation that it is a potential lobe of a jet driven by IRS A.

G270.8247–01.1112 exhibits strong NIR CO emission (Ilee et al. 2013). Ilee et al. (2013) modelled the CO band-head emission using a disc model, establishing that it hosts a disc of size 0.1-1.6 au, inclined at 89° , i.e. edge-on. Using data from the



Figure 3. SEDs of detected radio sources (continued on subsequent pages).

ALMA Evolutionary Study of High Mass Protocluster Formation in the Galaxy (ALMAGAL) survey (Wells et al. 2024), we extracted the 13 CO (2 - 1) transition emission line at mm wavelength - 220.39868 GHz, revealing blue and red wings aligned in a SE-NW direction (see the red and blue contours on Fig. 4). Additionally, the millimetre continuum emission shows a single core coincident with IRS A. The similarity between the alignment of the radio source with IR source A and the direction of the molecular outflow suggests that the radio emission originates from a lobe of an ionized jet aligned with this outflow and driven by IRS A.



Figure 3. (Continued.)

3.3.2 G271.2225-01.7712

G271.2225-01.7712 has a spectral index of -0.45 ± 0.09 , indicating non-thermal emission characteristic of radiation from non-thermal lobes of protostellar jets. The morphology of the object at both 5.5 and 9.0 GHz, generated from data of similar uv-range, depicts a jet-like structure that is aligned in a N-S direction. However, no clear elongation is observed in the 2MASS and VVV (Minniti et al. 2010) IR emissions. Similarly, Navarete et al. (2015) did not detect both molecular hydrogen emission at 2.12 µm, which is typically associated with shocks, or extended continuum IR emission.



Figure 3. (Continued.)

3.3.3 G282.8969-01.2727

G282.8969–01.2727 was detected at 5.5 GHz but not 9.0 GHz, where the field rms is $0.027 \text{ mJy beam}^{-1}$. The emission at 5.5 GHz manifests an extended source consisting of two components, A and B, with flux densities 0.09 ± 0.03 and 0.07 ± 0.02 mJy, respectively.

The spectral index of the extended emission, calculated using the flux at 5.5 GHz, and the upper limit at 9.0 GHz, is <-1.41, indicating non-thermal emission. The steep negative index might result from the resolution of flux at higher frequencies, as evidenced by significant size differences between emissions at 1.3 and 5.5 GHz. The upper limits of the spectral indices of components A and B, estimated using

a compact H II-region, suggesting that it may be a jet-driving H II region, similar to G345.4938+01.4677 (Guzmán et al. 2016).

3.3.5 G287.8768-01.3618

G287.8768-01.3618 exhibits a spectral index of 0.96±0.06, consistent with a thermal radio source. The morphology of its radio emission at 5.5 and 9.0 GHz is aligned in a SE-NW direction, with PAs of $148^{\circ}\pm10$ and $161^{\circ}\pm3$, respectively. The consistent orientation of the radio emission and PAs is characteristic of protostellar jets, suggesting that G287.8768-01.3618 is a thermal jet. However, Navarete et al. (2015) detected bipolar emission at 2.122 µm, a tracer of shock-excited molecular hydrogen at a PA of 65°. This suggests the presence of another jet propagating along the NE-SW axis, which differs from the orientation of the radio emission. Furthermore, IR observations by Mottram et al. (2007) detected two point sources aligned with the radio emission in a SE-NW direction at 10.4 µm. These two IR sources indicate the potential presence of two independent jet-driving sources in the field. Nevertheless, it remains unclear whether the H₂ jet is associated with either of the individual IR sources as both are aligned with the radio emission.

3.3.6 G289.0543+00.0223

G289.0543+00.0223 was detected at 9.0 GHz but not 5.5 GHz, where the rms noise is 0.27 mJy beam⁻¹. The spectral index of the source is >1.81, suggesting that it is a thermal radio source. However, its flux at 1.3 GHz is higher than the expected based on the SED derived from 5.5 and 9.0 GHz. This discrepancy suggests either the presence of non-thremal emission or the loss of flux due to resolvedout emission at higher frequencies. The IR and radio morphologies of the source appear generally spherical, suggesting that it could be a H II region.

3.3.7 G293.0352-00.7871

Two radio sources denoted as A and B in Fig. 1 were detected within the field of G293.0352–00.7871 at 9.0 GHz but not 5.5 GHz, where the rms noise level is 0.02 mJy beam⁻¹. Their flux densities are 0.14 ± 0.02 and 0.08 ± 0.02 mJy, respectively. The positions of G293.0352–00.7871 at 1.3 and 9.0 GHz exhibit an offset of ~1.4 arcsec, likely due to the positional accuracy limitations of the SMGPS (Goedhart et al. 2024). Additionally, the radio source shows a spatial offset relative to its associated IR counterpart, located to the NE. The 2MASS IR emission from the source displays elongation in the NE-SW direction, comparable to the orientation of a line joining A to B, potentially indicating the presence of an outflow.

The nature of components A and B remains uncertain. Their spectral indices, >1.14 for A and >0.3 for B, suggest that they are thermal sources. However, the flux of source A at 1.3 GHz is higher, suggesting that it might be non-thermal. The discrepancy, however, could be attributed to the low resolution of the 1.3 GHz data.

3.3.8 G297.1390-01.3510

G297.1390–01.3510 was detected at both 5.5 and 9.0 GHz, exhibiting distinct morphologies at each frequency. At 5.5 GHz, its emission aligns in an E-W direction, whereas at 9.0 GHz, the orientation is in a



Figure 4. Red and blue wings of ${}^{13}CO(2 \rightarrow 1)$ emission line at 220.39868 GHz in the field of G270.8247-01.1112. The cross (x) and asterisk (*) represent the locations of the IR and mm sources.

the fluxes at 5.5 GHz and $3 \times \sigma$ at 9.0 GHz, are < -0.24 and < 0.27, respectively.

Navarete et al. (2015) classified the 2.12 μ m molecular hydrogen emission associated with the source as monopolar, with a PA of 350°, indicating the presence of a shock. Furthermore, the alignment of the morphology of the source at 5.5 GHz with that at 1.3 GHz suggests that it is a jet. The PA of the radio source at 1.3 GHz, 118°±3, is consistent with the PA of the alignment of components A and B, i.e. 124°. However, the difference between the PA of the emission reported by Navarete et al. (2015) and the PA of the radio emission suggests the presence of multiple outflows in the field. While component A exhibits a negative spectral index, its proximity to the methanol maser source (Green et al. 2012) suggests that it could be one of the cores or a lobe that is very close to the core.

3.3.4 G287.3716+00.6444

G287.3716+00.6444 shows two radio sources, A and B, aligned in a NE-SW direction at 5.5 and 9.0 GHz (see Fig. 1). The emission observed at 1.3 GHz displays a similar orientation, confirming the alignment of components A and B. The two components were also detected in Purser et al. (2016), further supporting the presence of a jet. The alignment of A and B at a PA $\simeq 48^{\circ}$ is comparable to the observed PA of the radio source, which is $35^{\circ}\pm13$ (Obonyo et al. 2024).

The spectral index of component A is 0.28 ± 0.04 , consistent with ionized winds, while component B exhibits a spectral index of -0.68 ± 0.08 , indicative of non-thermal emission. The indices suggest that A is a thermal core, while B is a non-thermal lobe, implying a core-lobe system. Moreover, the proximity of a methanol maser source identified by Green et al. (2012) to component A reinforces its classification as a core, while B is a likely lobe.

Both Koumpia et al. (2021) and Ilee et al. (2018) detected Br_{γ} hydrogen recombination line emission at 2.167 µm, which they interpreted as evidence for the presence of an accretion disc. This suggests that G287.3716+00.6444 is an MYSO that hosts a disc and a jet. However, Purser et al. (2016) classified component A as

N-S direction. These orientations are also seen in images with similar uv-range.

The spectral index of the source is -1.25 ± 0.08 , indicating that it is a non-thermal source. Additionally, the flux of the source at 1.3 GHz is lower than the expected based on the SED generated using 5.5 and 9.0 GHz observations. This suggests that the index could be closer to -0.6, which is characteristic of synchrotron emission from jets. This non-thermal property of extended emission from massive protostars, as observed in Obonyo et al. (2024), suggests the dominance of the jet emission by unresolved non-thermal lobes within the jet.

Mid-IR observations by Mottram et al. (2007) detected two sources, IRS1 and IRS2 (see Fig. 1). IRS1 is spatially coincident with an MYSO that could be driving the jet responsible for the 5 GHz emission. However, the alignment of radio sources A and B suggests the presence of another outflow oriented in the N-S direction. This indicates the presence of multiple outflows in the field.

3.3.9 G304.7592-00.6299

G304.7592–00.6299 has a spectral index of -0.93 ± 0.09 , indicating that it is a non-thermal emitter. Its radio emission at 1.28, 5.5, and 9.0 GHz does not exhibit a clear or consistent orientation of its morphology. However, it manifests a compact source (Mottram et al. 2007) with nebulosity at 8 μ m, which is aligned in the NE-SW direction, comparable to the orientation of 1.28 GHz emission. Furthermore, Navarete et al. (2015) did not detect either IR continuum emission or molecular hydrogen emission at 2.12 μ m. This suggests the absence of an outflow, implying that G304.7592–00.6299 may be an H II region.

3.3.10 G309.9796+00.5496

G309.9796+00.5496 exhibits two radio components, A and B, aligned in a NE-SW direction at both 5.5 and 9.0 GHz. Their spectral indices are 0.82 ± 0.20 and 0.22 ± 0.13 , respectively, indicating that both components are thermal emitters. The combined spectral index of G309.9796+00.5496 is 0.51 ± 0.12 , which is typical of thermal sources. Cyganowski et al. (2008) classified the source as a possible outflow candidate of an Extended Green Object (EGO). The properties of G309.9796+00.5496 suggest that it is a jet, with component A being the potential jet driver due to its proximity to the location of the IR counterpart (Mottram et al. 2007).

3.3.11 G314.3197+00.1125

G314.3197+00.1125 is resolved into two radio components at 1.3 GHz, aligned in a SE-NW direction. At 5.5 GHz, the northwestern component further splits into two components, A and B, maintaining a comparable alignment with PA \simeq 155°. However, only component A was detected at 9.0 GHz. Additionally, a third radio source, denoted as C, was detected south of both A and B at 1.3, 5.5, and 9.0 GHz, with a BC PA \simeq 166°. The spectral indices of components A, B, and C are 1.53±0.12, <-0.68, and -0.58±0.16, respectively, indicating a bipolar jet morphology with component A as the core and B and C as non-thermal lobes. Moreover, the comparable PAs of AB and BC support the interpretation of a bipolar outflow. The presence of 6.7 GHz methanol maser emission, detected by Green et al. (2012) at the position marked with a plus(+) sign in Fig. 1, supports the identification of A as the jet driver. While the morphology suggests a bipolar jet, the absence of 2.12 μ m emission in the field (Navarete et al. 2015), the large angular separation of component C from A, and the lack of IR emission associated with C imply that it could potentially be a non-thermal extragalactic source.

3.3.12 G328.2523-00.5320A

G328.2523–00.5320A splits into three components, labelled A, B, and C in Fig. 1, which are aligned in a NE-SW direction. These components, detected at both 5.5 and 9.0 GHz, have spectral indices of -0.27 ± 0.10 , 1.00 ± 0.09 , and 1.98 ± 0.14 , indicating that they are non-thermal, thermal, and thermal, respectively. Source B is associated with methanol maser (Green et al. 2012), a dense mm-core (Csengeri et al. 2017), and IR emission, suggesting that it is the driving source. Components A and C, on the other hand, may be its non-thermal and thermal lobes, respectively.

3.3.13 G328.9842-00.4361

G328.9842–00.4361 was detected at both 5.5 and 9.0 GHz, with a spectral index of 0.18±0.13, indicating that it is a thermal jet. However, its SED between 1.3 and 5.5 GHz shows a negative index, $\alpha = -0.73\pm0.04$, suggesting that while the core may be thermal, it might be hosting non-thermal lobes that dominate the emission at lower frequencies. Morphologically, the source exhibits an elongated structure in a SE-NW direction at 1.3, 5.5, and 9.0 GHz, a characteristic feature of jets. The PAs of its morphology at these frequencies are $164\pm8^{\circ}$, $143\pm21^{\circ}$, and $141\pm13^{\circ}$, respectively. The consistency in the PAs further supports the interpretation of the source as a jet.

3.3.14 G334.7302+00.0052

G334.7302+00.0052 was detected at both 5.5 and 9.0 GHz, with a spectral index $\alpha = 1.20\pm0.07$, consistent with ionized winds, a feature commonly associated with jets driven by MYSOs. While the radio morphology of the source at 1.3, 5.5, and 9.0 GHz does not exhibit a clear or consistent elongation, the presence of extended 4.5 μ m green emission, similar to the EGOs described by Cyganowski et al. (2008), suggests a potential EGO SW of the source. The detection of this EGO indicates that G334.7302+00.0052 likely hosts a jet.

3.3.15 G335.0611-00.4261A

G335.0611–00.4261A exhibits a complex structure, splitting into two components at 5.5 GHz and three components, A, B, and C, at 9.0 GHz. These components are aligned in a N-S direction, which is consistent with the orientation of the blue wing of $^{13}CO(2-1)$ emission observed in the ALMAGAL survey data (see Fig. 5). The red wing of the molecular emission is located NE of the associated IR source, aligned along a NE-SW direction.

The orientation of the CO molecular emission are comparable to the morphology of the extended 4.5 μ m green emission, which traces shock-excited gas (Cyganowski et al. 2008). One of the green emission components extends towards the West, suggesting the presence of multiple outflows in the region. The classification of G335.0611–00.4261A as an EGO by Cyganowski et al. (2008) further supports the interpretation that this source drives a jet.



Figure 5. Red and blue wings of ${}^{13}\text{CO}(2 \rightarrow 1)$ emission line at 220.39868 GHz in the field of G335.0611–00.4261A. The plus (+) sign represents the locations of methanol maser emission in the field.

The spectral indices of the radio components reveal distinct emission mechanisms. Component (A+B) has a spectral index of $\alpha = -0.38 \pm 0.06$, indicative of non-thermal emission, while component C has a spectral index of $\alpha = 1.58 \pm 0.11$, suggesting thermal emission. Only B is associated with 6.7 GHz methanol maser emission (Green et al. 2012), suggesting that it may be the core driving the jet, with component A as its lobe. Component C, which has a thermal spectral index, is another potential core.

3.3.16 G338.2801+00.5419A

G338.2801+00.5419A is resolved into two radio components, A and B, at both 5.5 and 9.0 GHz in ATCA observations. The morphology observed with both MeerKAT and ATCA aligns in a SE-NW direction, depicting a jet-like core-lobe structure, consistent with the characteristics of a collimated outflow. The spectral index analysis reveals a non-thermal nature for the overall system, with an index of -0.59 ± 0.04 . The individual components also exhibit non-thermal characteristics, with component A having an index of -0.77 ± 0.08 and component B -0.47 ± 0.05 , indicating that both A and B are non-thermal lobes. Interestingly, component B is associated with 6.7 GHz methanol maser emission (Green et al. 2012), (sub-)mmcore (Elia et al. 2017), and an IR source, implying that it may either be the driving source of the jet or is located closer to the jet driver. Additionally, a 4.5 µm green point source is observed to the SE of B, potentially an EGO, further supporting the presence of a bipolar outflow with B as the potential jet driver.

3.3.17 G338.4712+00.2871

G338.4712+00.2871 exhibits a jet-like morphology aligned in a SE-NW direction at both 5.5 and 9.0 GHz, with a steep negative spectral index of -1.17 ± 0.12 . While this spectral index suggests non-thermal radio emission, its value is atypical. However, the lower 1.3 GHz flux, compared to the SED derived from 5.5 and 9.0 GHz data, implies that the index could be closer to typical value of -0.6.

The morphology of the radio emission is characteristic of jets from massive protostars. Additionally, the PAs of the radio emission at 5.5 and 9.0 GHz are comparable, i.e. 171 ± 4 and 157 ± 3 , respectively,

further supporting the interpretation that the source is a jet. Despite the jet-like characteristics, the radio source does not show association with the 6.7 GHz methanol maser (Green et al. 2012) or IR sources, which are marked by plus(+) and cross signs in Fig. 1, respectively. The maser and IR sources are located approximately 3.5 arcsec east of the radio source, suggesting that the radio emission is not linked to the IR source. Thus, the radio emission could be from a highly obscured protostellar object, comparable to a class 0 type protostar, which is undetectable at near-IR wavelengths (Kern et al. 2016).

The 8 μ m emission associated with G338.4712+00.2871 appears as a point source adjacent to a bright ridge. This source, however, was not detected in 2MASS and TIMMI2 (Mottram et al. 2007) observations, potentially due to strong absorption seen on the 8 μ m map, especially towards the west where the radio source is located. This absorption could be obscuring the driving source of the radio jet.

3.3.18 G339.3316+00.0964

G339.3316+00.0964 was detected at both 5.5 and 9.0 GHz. The morphologies at these frequencies are elliptical and aligned in a N-S direction, with PAs of $16\pm6^{\circ}$ at 5.5 GHz and $21\pm46^{\circ}$ at 9.0 GHz. The spectral index of the source between 5.5 and 9.0 GHz is 0.41 ± 0.04 , indicating that it is a thermal radio source. This value lies between those expected for spherical and collimated ionized winds (Reynolds 1986). However, 1.3 GHz flux suggests a negative spectral index, which is the characteristic of non-thermal emitters, suggesting a combination of thermal and non-thermal emission mechanisms.

The association of the radio source with IR emission, along with its morphologies at 5.5 and 9.0 GHz and the spectral indices, suggests that G339.3316+00.0964 is a thermal jet. However, the orientation of the IR emission is misaligned with that of the radio source towards the SW–NE direction. Additionally, the IR emission exhibits a 4.5 μ m green radiation feature to the NE, suggesting the presence of another jet in that direction or precession.

3.3.19 G340.7455-01.0021

The 5.5 and 9.0 GHz emission from G340.7455-01.0021 displays an extended morphology oriented in a NE-SW direction, similar to alignment of the blue- and red-shifted wings of CO emission in the field (see Fig. 6). The spectral index of the radio emission is 0.53 \pm 0.08, consistent with thermal jets, and the 1.3 GHz emission aligns well with the SED of the jet.

G340.7455–01.0021, classified by Cyganowski et al. (2008) as an EGO outflow candidate, displays 4.5 μ m green emission at a PA $\simeq 225^{\circ}$. Additionally, Navarete et al. (2015) detected a single knot at a PA of 290° which may be associated with the radio jet. The difference in PAs of the knots, and radio emission suggests the presence of multiple outflows in the field.

3.3.20 G345.2619-00.4188A

G345.2619–00.4188A exhibits two radio components, labelled A and B, detected at 5.5 GHz but not at 9.0 GHz where rms noise is 0.03 mJy. The flux densities of A and B at 5.5 GHz are 0.19 ± 0.06 and 0.09 ± 0.03 mJy, respectively. The spectral indices of components A and B are <-1.52 and <0.00, respectively, suggesting that they are non-thermal lobes of a jet.

The alignment of the IR source with radio components A and B in a NE-SW direction suggests the presence of a bipolar outflow, with



Figure 6. Red and blue wings of ${}^{13}CO(2 \rightarrow 1)$ emission line at 220.39868 GHz in the field of G340.7455-01.0021, shown in red and blue colours, respectively. The plus (X) sign represents the locations of IR source.

component B also exhibiting characteristics of a bow-shock. This jet is likely driven by a core associated with the mid-IR source, whose location is marked with a cross(x) in Fig. 1.

3.3.21 G345.7172+00.8166A

G345.7172+00.8166A exhibits an extended morphology oriented in a N-S direction, particularly evident at 1.3 and 9.0 GHz, where it displays a characteristic jet-like structure. Its spectral index of -0.64 ± 0.13 is consistent with non-thermal jets. Moreover, the PA of its emission at 5.5 and 9.0 GHz are comparable, i.e. $139\pm32^{\circ}$ and $158\pm3^\circ$, respectively, further supporting the jet-like nature. Cyganowski et al. (2008) classified it as an EGO and an outflow candidate that is associated with an IR dark cloud. This association suggests that G345.7172+00.8166A may be a very young massive protostar.

3.4 Radio luminosity versus bolometric luminosity

The radio luminosity of the cores of the objects at 5.5 GHz was calculated and plotted against their bolometric luminosities, taken from Lumsden et al. (2013). For sources with unresolved cores, the fluxes of the jets, which host the cores, were used. This plot is instrumental in distinguishing between free-free emission from collisionally ionizing jets and photoionizing UCHII regions (Anglada et al. 2018). Shock-ionized jets typically exhibit a shallower slope in the radio versus bolometric luminosity plot, whereas photoionized UCHII regions show a steeper dependence due to their strong correlation with the Lyman continuum flux.

Fig. 7 shows the plot of radio luminosity versus bolometric luminosity for the objects classified as jets in the sample. Data from previous observations, such as Anglada (1995) and Anglada et al. (2018), are included to illustrate the expected trends. While some sources lie closer to the empirical jet relation from Purser et al. (2021), others are closer to the Lyman continuum line, particularly those with negative spectral indices. Notably, the sources with higher radio luminosities compared to previous studies are predominantly non-thermal. This trend was also observed by Kavak et al. (2021) and Obonyo et al. (2024), who attributed it to differences in the resolution



unresolved in the SMGPS (Goedhart et al. 2024), at both 5.5 and 9.0 GHz, resulted in 22 significant detections. Among these, 19 were classified as jets, two as H II-regions, and one as an unknown source. Notably, of the 19 jets, 13 exhibited non-thermal emission or were associated with non-thermal lobes. This indicates that approximately 68 per cent of the jets are associated with non-thermal emission, while the remaining 32 per cent are thermal. These findings suggest that magnetic fields play a crucial role in driving massive protostellar jets (Blandford & Payne 1982, Shu et al. 1994a), supporting the concept of disc-fed accretion similar to that observed in low-mass protostars. However, the mechanism driving thermal jets, especially for creating thermal lobes, remains unclear. The morphology of most of the sources manifests jet-like structures with resolved components aligned linearly, i.e. a jet-lobe(s) system. Moreover, molecular outflows also manifested the presence of jets in some cases. The objects were further plotted on a radio versus bolometric luminosity plot to eliminate potential UCHII regions. The plot revealed that most of the radio-luminous sources are also non-thermal and are potential jets, implying that their radio luminosity is perhaps due to additional emission from non-thermal radiation.



Figure 7. Radio continuum luminosity at 5.5 GHz versus bolometric luminosity of the jets in the sample. The asterisks represent the jets in our sample with the face colours showing their spectral indices. The arrows show the upper limits. The solid line represents the radio luminosities expected from Lyman continuum radiation for a zero-age main-sequence star of a given bolometric luminosity (Thompson 1984). The dash-dotted line is the leastsquares fit to the radio jets reported in Purser et al. (2021).

of observations used in estimating the indices. However, the higher radio luminosities exhibited by non-thermal jets compared to thermal counterparts may reflect an intrinsic property of the sources, perhaps an indicator of non-thermal contribution to the free-free emission. Around a bolometric luminosity of $10^3 L_{\odot}$, there is significant overlap between jets and UCHII regions within the observed scatter. This region also corresponds to where zero-age main-sequence protostars and massive protostellar jets are expected to lie. In such cases, additional diagnostics such as source morphology or radio recombination line widths are essential to distinguish between the two types.

The higher resolution ATCA observations of 28 MYSOs, previously

4 CONCLUSIONS

MNRAS 539, 3347-3363 (2025)

ACKNOWLEDGEMENTS

Special thanks to the DARA project for sponsoring the observer's trip to CSRIO. Part of the data used here were retrieved from the Japan Virtual Observatory (JVO) portal (http://jvo.nao.ac.jp/portal/), which is operated by ADC/NAOJ.

DATA AVAILABILITY

The data underlying this paper are available at the SARAO Data Archive: https://meerkatgps.org/data_products/Mosaic_Momen t0/, ATCA archive: https://atoa.atnf.csiro.au, and JVO portal: http://jvo.nao.ac.jp/portal/.

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This paper has been typeset from a TEX/LATEX file prepared by the author.

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