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Combined Pre-supernova Alert System with KamLAND and Super-Kamiokande

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Abstract

Preceding a core-collapse supernova (CCSN), various processes produce an increasing amount of neutrinos of all flavors characterized by mounting energies from the interior of massive stars. Among them, the electron antineutrinos are potentially detectable by terrestrial neutrino experiments such as KamLAND and Super-Kamiokande (SK) via inverse beta decay interactions. Once these pre-supernova (pre-SN) neutrinos are observed,



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an early warning of the upcoming CCSN can be provided. In light of this, KamLAND and SK, both located in the Kamioka mine in Japan, have been monitoring pre-SN neutrinos since 2015 and 2021, respectively. Recently, we performed a joint study between KamLAND and SK on pre-SN neutrino detection. A pre-SN alert system combining the KamLAND detector and the SK detector was developed and put into operation, which can provide a supernova alert to the astrophysics community. Fully leveraging the complementary properties of these two detectors, the combined alert is expected to resolve a pre-SN neutrino signal from a $15 M_{\odot}$ star within 510 pc of the Earth at a significance level corresponding to a false alarm rate of no more than 1 per century. For a Betelgeuse-like model with optimistic parameters, it can provide early warnings up to 12 hr in advance.

Unified Astronomy Thesaurus concepts: [Particle astrophysics \(96\)](#); [Neutrino astronomy \(1100\)](#); [Core-collapse supernovae \(304\)](#)

1. Introduction

Neutrinos emitted by a supernova during the first ~ 10 s carry unique information about the physics of supernovae, which hold immense significance in the realm of astrophysics. The first observed supernova neutrinos (Alekseev et al. 1987; Bionta et al. 1987; Hirata et al. 1987) were from SN 1987A in the Large Magellanic Cloud, ~ 50 kpc away from Earth (Pietrzyński et al. 2019). Since then, various neutrino experiments, such as Borexino (Alimonti et al. 2009), IceCube (Kopke 2018), KamLAND (Abe et al. 2022b), LVD (Agafonova et al. 2015), NOvA (Acero et al. 2020), SNO+ (Andringa et al. 2016), and Super-Kamiokande (SK; Abe et al. 2016), equipped with advanced technology and improved capabilities, continued the quest to detect supernova neutrino bursts. Furthermore, a number of next-generation neutrino detectors sensitive to galactic supernova neutrinos are under construction, including DUNE (Abi et al. 2021), Hyper-Kamiokande (Abe et al. 2018), JUNO (Abusleme et al. 2023), and KM3NeT (Aiello et al. 2021). To catch such a fleeting event, it is desirable to be alerted well before the explosion, so that astronomers and particle physicists may prepare for the observation of supernova neutrinos and possible gravitational waves as soon as the explosion happens.

The evolution of a single star, whose initial mass is greater than 8 solar masses (M_{\odot}), to its final stages prior to the core-collapse supernova (CCSN) is characterized by nuclear burning in its core due to its high temperature and density (Woosley et al. 2002).⁷⁸ The change in the chemical composition of a star, forming concentric shells of heavier elements along its volume, is the result of the nuclear fusion of heavier elements in the core. Stars at this stage are called pre-supernova (pre-SN) stars. The main cooling mechanism of a pre-SN star is through neutrino emissions. Starting from the carbon burning stage, neutrinos are produced by pair annihilation $e^+e^- \rightarrow \bar{\nu}\nu$, producing all flavors of neutrino and antineutrino pairs (Odrzywolek et al. 2004a). As the star approaches core collapse, the nuclear beta decay begins to dominate. Nuclear processes, such as beta decay, will eventually contribute more to neutrino emission than thermal processes as the star approaches core collapse (Patton et al. 2017).

These neutrinos, referred to as pre-SN neutrinos, are potentially detectable by terrestrial detectors if the progenitor is close enough to Earth (Odrzywolek et al. 2004b). They not only signal the imminent supernova, but also provide insight into the late stages of stellar evolution of massive stars. A pre-SN neutrino detection can help unravel many uncertainties associated with stellar evolution models: the physical processes

that lead to a CCSN, the shell structure formation, the isotopic composition of stars, etc. It can also provide evidence of neutrino mass ordering (Kato et al. 2020).

The energy of pre-SN neutrinos is of the sub-MeV scale or MeV scale. We focus on inverse beta decay (IBD) $\bar{\nu}_e + p \rightarrow e^+ + n$, which has a relatively large cross section in liquid scintillator (LS) detectors (e.g., KamLAND; Suzuki 2014) and water Cherenkov detectors (e.g., SK; Fukuda et al. 2003) in the energy range of pre-SN neutrinos. Compared to LS detectors, whose energy threshold is typically less than 1 MeV, water Cherenkov detectors are less sensitive to low-energy neutrinos because the higher energy threshold, which is 2.5 MeV in kinetic energy, as shown in Section 4.1, limits the detection of neutron capture signals. However, starting in 2020, the SK detector was loaded with gadolinium (Gd) to improve the neutron detection efficiency (Abe et al. 2022a). Since 2015, KamLAND has been monitoring pre-SN neutrinos and was able to provide pre-SN alerts to the astrophysics community (Asakura et al. 2016). Later in 2021, SK also implemented an online pre-SN alert system (Machado et al. 2022). As of 2024 May, no alert has been issued. Besides IBD in LS detectors and water Cherenkov detectors, distinct detection methods in other detectors could be utilized for pre-SN neutrino detection as well. For example, coherent neutrino–nucleus scattering in future large-scale dark matter direct detection experiments is a method complementary to IBD, because it can detect all flavors of neutrinos (Raj et al. 2020).

In this article, we introduce a joint study between KamLAND and SK on pre-SN neutrino detection. This combination aims at extending the reach to potential CCSN progenitors at a greater distance and reducing the warning time of pre-SN alerts. Additionally, we present new sensitivity results individually for KamLAND and SK. Compared to the previous study (Asakura et al. 2016), KamLAND has now taken more recent pre-SN neutrino models into consideration. The sensitivity of SK to pre-SN neutrinos with 0.01% Gd by mass is presented in Machado et al. (2022). Since 2022, the Gd concentration in SK has increased to 0.03% by mass, further enhancing its capability to identify low-energy electron antineutrinos ($\bar{\nu}_e$) via IBD. Therefore, both experiments have reoptimized their analysis strategies according to these changes.

2. Pre-supernova Neutrino Model

To estimate the expected signal from pre-SN $\bar{\nu}_e$ in SK and KamLAND, two models for pre-SN neutrino emission during the evolution of massive stars were used: Odrzywolek et al. (2004b), Odrzywolek & Heger (2010), and Patton et al. (2017). Both models provide data sets for the calculation of $\bar{\nu}_e$ emission during the pre-SN stage. Odrzywolek et al. (2004b) provides

⁷⁸ With different physical assumptions, such as rotation or the presence of a massive companion, the evolution can be significantly different (Eldridge & Stanway 2022).

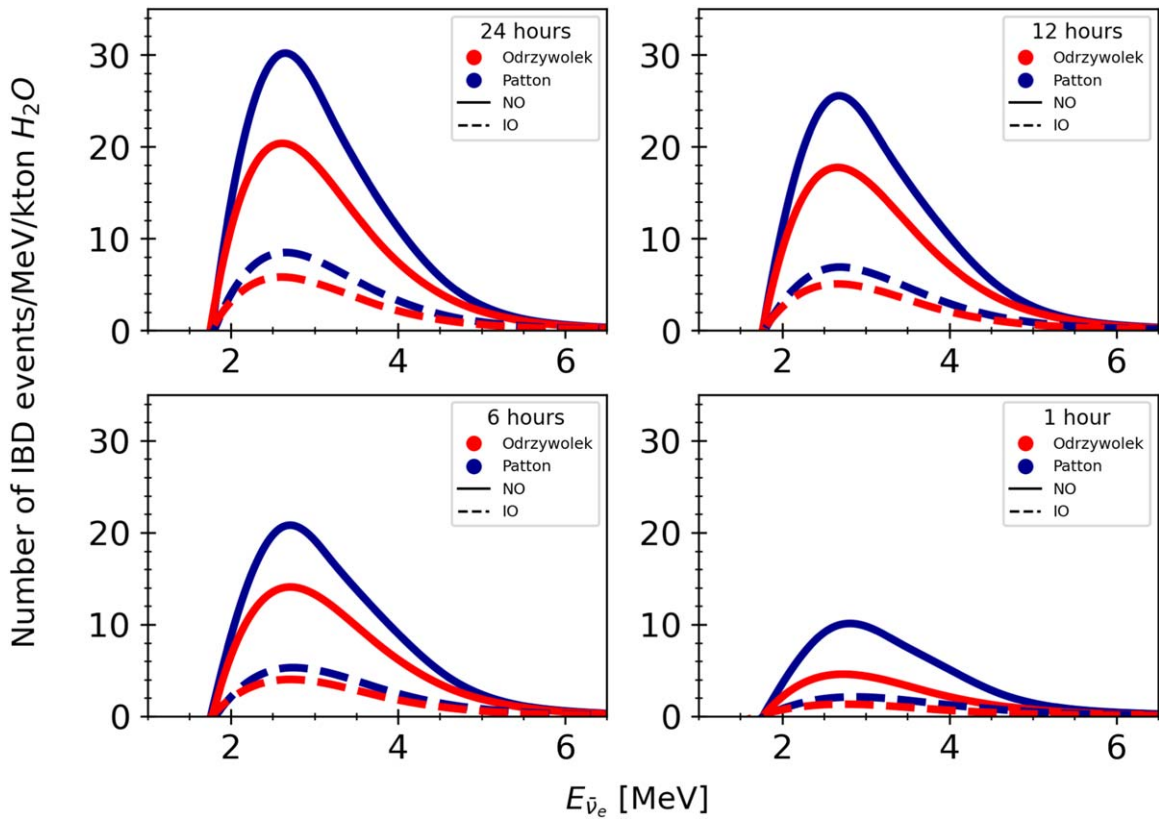


Figure 1. Number of pre-SN IBD interactions per kton of water integrated over the last 24 hr, 12 hr, 6 hr, and 1 hr prior to the CCSN as a function of the $\bar{\nu}_e$ energy, $E_{\bar{\nu}_e}$. The Betelgeuse-like models consider stars with initial masses of $15 M_{\odot}$ located 150 pc away from Earth, for both normal neutrino mass ordering (NO) and inverted neutrino mass ordering (IO).

data sets for stars with $15 M_{\odot}$ and $25 M_{\odot}$ and Patton et al. (2017) for $15 M_{\odot}$ and $30 M_{\odot}$.

The model from Odrzywolek & Heger (2010) assumes that the entire neutrino flux comes from pair annihilation. For the nuclear isotopic composition of the star, the model assumes a nuclear statistical equilibrium (NSE), which is a treatment only dependent on the temperature, density, and electron fraction, making it a simple flux estimated by only postprocessing an already existing stellar model. The model from Patton et al. (2017) includes a more complete evaluation of the neutrino flux from the pre-SN star, including contributions not only from pair annihilation, but also from plasmon decay, photoneutrino process, beta decay, and electron/positron captures. By using the star evolution code MESA (Modules for Experiments in Stellar Astrophysics; Paxton et al. 2011), this model couples the isotopic evolution to the stellar evolution, giving a more robust estimation of the neutrino fluxes from nuclear weak processes.

Neutrinos undergo flavor conversion, i.e., neutrino oscillations, from the point of production to the point of detection. To calculate the expected signal from the considered models, adiabatic neutrino oscillations in the matter of the star and neutrino oscillations in vacuum are taken into account. For the former, the ratio of $\bar{\nu}_e$ is changed at high Mikheyev–Smirnov–Wolfenstein resonance, which depends on the mass ordering of neutrinos (Smirnov 2003). Different transition probabilities are assumed for normal and inverted neutrino mass orderings to account for the change in the ratio of electron flavor neutrinos due to the dense stellar medium and the effects of neutrino oscillations in a vacuum.

We attempt to explore the sensitivities for detecting pre-SN neutrinos from the well-known red supergiant α -Ori (Betelgeuse), which will potentially end up a CCSN.⁷⁹ Current estimation of its mass and distance suggests $16.5\text{--}19 M_{\odot}$ and 168_{-15}^{+27} pc (Joyce et al. 2020). Limited by the available data sets, we choose $15 M_{\odot}$ and 150 pc to simulate a Betelgeuse-like pre-SN star in this work. Although we focus on detecting pre-SN neutrinos from a Betelgeuse-like star, pre-SN neutrinos from other stars could also be observable. A list of candidate pre-SN stars with updated distance and mass estimates can be found in Machado et al. (2022). Figure 1 shows the expected number of IBD candidates per kton of water for different pre-SN models as a function of the $\bar{\nu}_e$ energy, assuming a Betelgeuse-like pre-SN star. The event spectra are obtained by integrating the expected candidates over the last 24 hr, 12 hr, 6 hr, and 1 hr prior to core collapse. These candidates predominantly cluster around $E_{\bar{\nu}_e} \approx 2.6$ MeV, but can be found at higher energies. The expected event rate increases over time, leading to a large fraction of the total IBD candidates concentrated in the last hour.

In Sections 3 and 4, the KamLAND and the SK experiments are introduced, accompanied by the event selection strategies.

3. The KamLAND Experiment

KamLAND is an LS detector located 1000 m underground in the Kamioka mine. KamLAND was originally designed to study reactor neutrinos, geoneutrinos, and low-energy solar

⁷⁹ There are claims that the explosion is imminent (Saio et al. 2023) and opposing views (Molnar et al. 2023).

neutrinos. The primary target volume consists of 1 kton of ultrapure LS contained in a 13 m diameter spherical balloon made of 135 μm -thick transparent nylon ethylene vinyl alcohol copolymer (EVOH) composite film. The components of the KamLAND LS are 80% dodecane and 20% pseudocumene (1,2,4-trimethylbenzene) with 1.36 g l^{-1} of the fluor PPO (2,5-diphenyloxazole). An array of 1325 17 inch photomultiplier tubes (PMTs) and 554 20 inch PMTs mounted on the inner surface of an 18 m diameter stainless steel sphere is used to detect the scintillation light from events occurring within the balloon. Nonscintillating mineral oil fills the space between the balloon and the inner surface of the sphere. This is all surrounded by a 3.2 kton water Cherenkov detector contained in a resin-coated cylindrical rock cavern for a cosmic-ray veto. Detailed information of the detector is given in Suzuki (2014).

KamLAND started its data taking in 2002 March. Pre-SN $\bar{\nu}_e$ are expected to be detected through IBD processes, which is the main interaction channel for these neutrinos in KamLAND. Positrons produced in the process lose their kinetic energy within the LS medium and annihilate with electrons, emitting two 511 keV γ rays (prompt events). Neutrons with a mean lifetime of $207.5 \pm 2.8 \mu\text{s}$ are captured by protons, releasing 2.2 MeV γ rays (delayed events). By using the time and spatial correlation between the prompt and the delayed events, we achieve low-background conditions in the detection of $\bar{\nu}_e$.

In 2011, an inner balloon of 1.54 m radius containing Xe-loaded LS (Xe-LS) was installed in the center of the main balloon as a part of the KamLAND Zero-Neutrino Double-Beta Decay (KamLAND-Zen) experiment (Gando et al. 2012a). The inner balloon was updated to have a 1.92 m radius to house double the amount of Xe-LS in 2018 (Abe et al. 2023b). The center region is not used for the $\bar{\nu}_e$ analysis because of backgrounds from the inner balloon and its support materials.

The energy and vertex of an event can be reconstructed using the timing and charge distributions of scintillation photons recorded by the PMTs. The reconstruction algorithms are calibrated with radioactive sources deployed from the top of the detector (Berger et al. 2009; Banks et al. 2015). Using these calibration sources, the energy resolution is estimated to be $6.4\%/\sqrt{E_{\text{rec}}(\text{MeV})}$ and the vertex resolution is estimated to be $12 \text{ cm}/\sqrt{E_{\text{rec}}(\text{MeV})}$, respectively, where E_{rec} is reconstructed energy. The nonlinear and particle-dependent effects of the conversion between deposited (real) energy and E_{rec} are also calibrated with the Birks formula (Birks 1951) and the contribution of Cherenkov emission.

3.1. Event Selection in KamLAND

KamLAND performs muon vetos prior to selecting prompt-delayed pairs (delayed coincidence method) as neutrino events. Cosmic-ray muons produce events with bright scintillation light and multiple spallation products, including neutrons. This makes it challenging to reconstruct the correct vertex and energy of low-energy events and to select prompt-delayed pairs immediately following the muon event. Thus, all events within 2 ms of the arrival time of muons are vetoed. However, the 2 ms veto is not enough for high-energy muons, which make cascade showers in the detector. Such muons generate a nonnegligible amount of long-lived spallation products such as ${}^9\text{Li}$, which has a lifetime of 257.2 ms. Therefore, KamLAND performs a 2 s whole-volume veto for high-energy muons (Gando et al. 2012b). Alternatively, a cylindrical cut along the trajectory is applied when the reconstruction quality is good.

These three muon vetos are determined by the total observed charge, the residual charge, which means the difference between the observed charge minus the charge that would be expected if the muon simply penetrated the detector, and the quality of muon event reconstruction.

After applying the muon vetos, KamLAND applies the following criteria: (i) reconstructed prompt energy: $0.9 < E_p(\text{MeV}) < 4.0$; (ii) reconstructed delayed energy: $1.8 < E_d(\text{MeV}) < 2.6$ (capture on proton), or $4.4 < E_d(\text{MeV}) < 5.6$ (capture on ${}^{12}\text{C}$); (iii) spatial correlation between the prompt and delayed events: $\Delta R < 200 \text{ cm}$; (iv) time difference between prompt and delayed events: $0.5 < \Delta T(\mu\text{s}) < 1000$; (v) fiducial volume (FV) radii: $R_p, R_d < 600 \text{ cm}$; (vi) inner balloon cut: $R_d < 2.5 \text{ m}$ and $\sqrt{x_d^2 + y_d^2} < 2.5 \text{ m}$ for $z_d > 0 \text{ m}$, where (x_d, y_d, z_d) is the reconstructed delayed vertex. Note that the reconstructed prompt energy (E_p) is the sum of the positron kinetic energy and annihilation γ energies with the quenching effect.

Although the delayed coincidence method strongly suppresses accidental background events, KamLAND performs an additional likelihood-based selection to differentiate $\bar{\nu}_e$ from accidental backgrounds, which become more likely at lower energies and as the vertices are reconstructed closer to the balloon (Gando et al. 2013).

The total selection efficiency is calculated using a Geant4 Monte Carlo simulation. A total of 10^7 $\bar{\nu}_e$ events are generated uniformly in the 750 cm radius volume for each prompt energy bin, and the delayed coincidence selection with the likelihood selection is applied. The selection efficiency is calculated as the ratio of the number of surviving events after the selection to the number of events generated within the 600 cm FV, shown in Figure 2. The efficiency at low energies mirrors the spectrum of the accidental background because of the likelihood-based selection. At high energies, it remains nearly constant. The efficiency loss is dominated by the inner balloon cut. Without this cut, the efficiency is higher, $\sim 90\%$ above 4 MeV. The muon veto has an additional effect of reducing the analysis time. The residual analysis time after this reduction is defined as the *livetime*. The KamLAND data are divided into runs. The average livetime ratio in any single run is approximately 0.903.

All PMT waveforms undergo digitization through front-end electronics and are collected by the Data Acquisition (DAQ) software. Event energy and vertex are reconstructed from these waveforms and are collected as a single file, each covering approximately 6 minutes of data. It takes 300–900 s, on average 700 s, from the time of the last event in the file to the end of the reconstruction. The latency is 800–1200 s, on average 1120 s, when we consider the first event in the file. The latency strongly depends on the status of other processes. The pre-SN monitoring process is scheduled to run at 5 minutes intervals. Upon the identification of a new file, the process applies the selection criteria described above to detect neutrino events. Additionally, the process counts the number of events that passed the selection criteria within the past 24 hr for pre-SN analysis.

3.2. Background Sources in KamLAND

The backgrounds for pre-SN neutrinos through IBD can be categorized into two types. One type includes non-neutrino events, such as ${}^{13}\text{C}(\alpha, n){}^{16}\text{O}$ reactions and accidental prompt and delayed coincidences. The other type is neutrino backgrounds, such as reactor neutrinos and geoneutrinos.

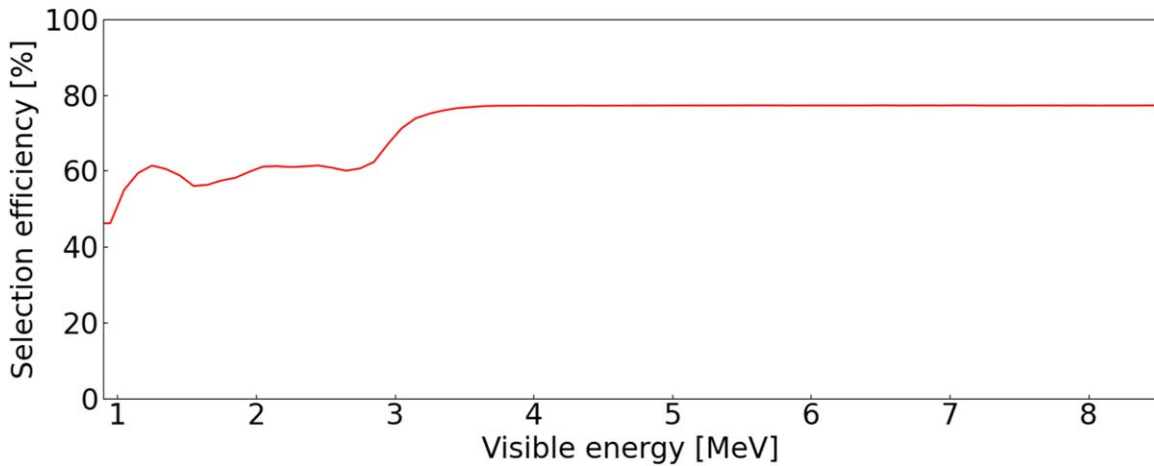


Figure 2. The total IBD selection efficiency in KamLAND.

In the early stages, KamLAND suffered from fake prompt-delayed pairs, which are $^{13}\text{C}(\alpha, n)^{16}\text{O}$ generated from α -decay of ^{210}Po in the KamLAND LS (Abe et al. 2008). However, this $^{13}\text{C}(\alpha, n)^{16}\text{O}$ reaction was substantially reduced during two distillation campaigns in 2007 and 2008. Currently, the rate of $^{13}\text{C}(\alpha, n)^{16}\text{O}$ events is 0.003 day^{-1} . The accidental background is effectively suppressed by the likelihood selection. The accidental event rate is 0.015 day^{-1} .

Reactor $\bar{\nu}_e$ is one of the main backgrounds in this analysis and will be discussed in Section 5.2. Geoneutrinos, generated by beta decays of nuclear isotopes such as ^{238}U and ^{232}Th in the Earth, also constitute a background as their energies can be up to 3.27 MeV. The expected geoneutrino event rate in KamLAND is 0.030 day^{-1} .

4. The Super-Kamiokande Experiment

The SK experiment is a water Cherenkov detector located in the same Kamioka mine as KamLAND. SK is a multipurpose detector that has been operating since 1996 and focuses on nucleon decays (Takenaka et al. 2020) and neutrino properties such as neutrino oscillations by observing atmospheric (Wester et al. 2024), solar (Abe et al. 2024), and accelerator neutrinos (Abe et al. 2023a). Furthermore, SK is a neutrino telescope capable of observing neutrinos from astronomical sources (Abe et al. 2021a; Orii et al. 2021; Mori et al. 2022).

The SK detector is composed of a cylindrical stainless steel tank with a 39.3 m diameter and 41.4 m height (Fukuda et al. 2003). The detector is divided into two regions: the inner detector (ID) and the outer detector (OD). The ID is responsible for the event detection with over 11,000 20 inch PMTs, and it has a volume of 32 kton, although the usual FV used in SK analyses is 22.5 kton. The OD has a thickness of about 2 m and it is composed of 1885 8 inch PMTs, facing the outside of the detector to reduce entering cosmic-ray muon-induced backgrounds.

In 2020, Gd sulfate octahydrate $\text{Gd}_2(\text{SO}_4)_3 \cdot 8\text{H}_2\text{O}$ was dissolved in the water in the detector, starting the SK-Gd phase. The loading of Gd improves SK's sensitivity to low-energy $\bar{\nu}_e$, expanding the physics goals of the experiment. SK now has the potential to reveal neutrinos from the Diffuse Supernova Neutrino Background (DSNB) (Abe et al. 2021b; Harada et al. 2023) and pre-SN stars, which are yet to be observed. In 2022, an additional Gd loading into SK was completed, achieving higher concentrations of Gd in the water (Abe et al. 2022a).

Low-energy $\bar{\nu}_e$ from pre-SN stars are detected in SK via IBD, similarly to KamLAND. However, positrons generated from IBD produce Cherenkov radiation instead of scintillation light, and γ -rays from neutron capture are detected mainly by Compton scattering electrons, producing Cherenkov radiation. In SK-Gd, the majority of thermal neutron captures occur on Gd due to its significantly higher neutron capture cross section. Specifically, while hydrogen has a capture cross section of only 0.3 barns, Gd's effective cross section averages 49,000 barns. The largest contributions for neutron capture come from the isotopes ^{157}Gd and ^{155}Gd (Abe et al. 2022a). The resulting γ -ray cascade from neutron captures on Gd (nGd) releases more energy, approximately 8 MeV compared to captures on hydrogen, leading to a greater photon yield. In the first phase of SK-Gd (2020 July–2022 March), which corresponded to a concentration of 0.01% Gd by mass, approximately 50% of neutron captures were on Gd. For the current phase with 0.03% Gd by mass (since 2022 July), the neutron capture efficiency is approximately 75%.

4.1. Event Selection in SK

The full event selection strategy for pre-SN neutrino detection in SK is described in Machado et al. (2022). Some updates have been made to the selection with the start of the second phase of SK-Gd with 0.03% Gd.

The data used for the pre-SN neutrino analysis come from the Wide-band Intelligent Trigger (WIT; Carminati 2015), a computing farm with approximately 900 hyperthreaded cores dedicated to real-time data processing. Each core handles 23 ms data files sequentially, applying a set of criteria to select good-quality events while ensuring high efficiency, even at energies as low as 2.5 MeV in kinetic energy. After event reconstruction, the processed files are sent to an organizer machine: the files with the reconstructed events arrive time unordered, and are then organized. While organizing the data, the files are grouped into segments lasting about 1.5 minutes each. Subsequently, they are made available to the pre-SN subsystem and then transferred outside the WIT system for low-energy offline analysis. In addition to the pre-SN alert system, the WIT system also hosts an online supernova burst trigger and raw data buffer, which would be preserved in case of a supernova event. Table 1 provides the estimated time between DAQ and a decision by the pre-SN alert system.

Table 1

Estimated Latency Time of Each Step in the pre-SN Alert System and Update Frequency

Process	Estimated Time
Data Fitting (WIT system)	10 s
Data organizing (WIT system)	4 minutes
Process Queue ($\sim 2 \times 10^6$ events)	2 minutes
Alert Decision/Export Results	Performed every 5 minutes

Note. Updated from Machado et al. (2022). Total latency time is the sum of the latency of each step.

The pre-SN alert system receives data from WIT right after the organizer processes sorts the data in time. The system runs the event selection in real time, which is based on the coincidence distance (dR) and coincidence time (dT) of IBD pairs and two Boosted Decision Tree (BDT) methods: one used as preselection for IBD pair identification (BDT_{online}) and another used as final selection based on angular distribution of hits, reconstructed energy and quality (BDT_{offline} ; more details in Machado et al. 2022). For the current SK-Gd phase with 0.03% Gd, BDTs were retrained, and cuts were optimized: $BDT_{\text{online}} > 0.2$, $dR < 300$ cm, $dT < 80$ μ s, and $BDT_{\text{offline}} > -0.1$. Figure 3 shows the signal background separation of the current BDT_{online} used for preselection, and Figure 4 shows the efficiency of applying the selection criteria to 10^7 IBD pairs. The irreducible background rate is approximately 0.5 event hr^{-1} .

4.2. Background Sources in SK

The major backgrounds for the pre-SN neutrino search in SK are reactor neutrinos and accidental coincidences. Other background sources include geoneutrinos, radioactive contaminants, and cosmic-ray muon-induced spallations. The backgrounds from geoneutrinos and accidental coincidences are similar to what has been described for KamLAND in Section 3.2. Reactor neutrino background will be discussed in Section 5.2. Radioactive contaminants that came along with the Gd loading are also a background source: ^{235}U chain isotopes can emit α , contributing to the backgrounds from $^{18}\text{O}(\alpha, n)^{21}\text{Ne}^*$ and $^{17}(\alpha, n)^{20}\text{Ne}^*$ processes. Moreover, the spontaneous fission of ^{238}U can emit neutrons that mimic delayed signals or even IBD candidates. Cosmic-ray muon-induced spallations are expected to have a tiny contribution since its resulting background rate is low and are effectively removed using BDTs.

5. Sensitivity to pre-SN Neutrinos in KamLAND and SK

Previous sensitivity studies for the detection of pre-SN neutrinos in KamLAND and SK are given in Asakura et al. (2016), Simpson et al. (2019), and Machado et al. (2022). In Asakura et al. (2016), the sensitivity for KamLAND's detection of pre-SN neutrinos, assuming the model from Odryzwolek & A. Heger (2010) is presented. Simpson et al. (2019) present a preliminary overall sensitivity for SK doped with 0.1% Gd by mass. In Machado et al. (2022), an improved sensitivity to pre-SN neutrinos in SK is shown, using data from the first phase of SK-Gd (with 0.01% Gd) to predict realistic backgrounds and new event selection methods.

In this study, both analyses in KamLAND and SK have been updated. KamLAND has now taken the additional pre-SN neutrino model from Patton et al. (2017) into consideration and reoptimized the selection parameters, the analysis time

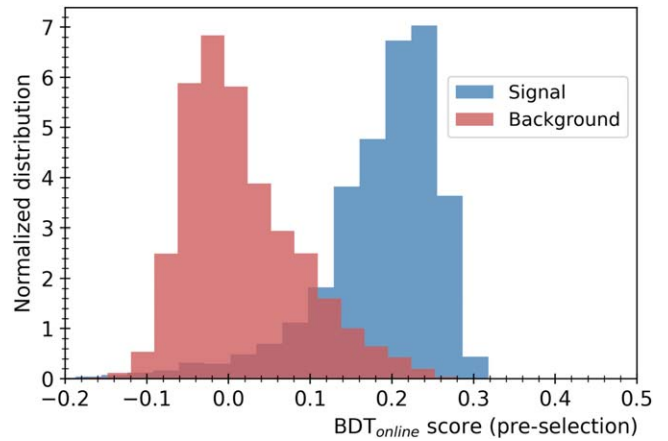


Figure 3. Signal background separation for the Boosted Decision Tree classifier used for preselection (BDT_{online}) using random subsets of SK data with 0.03% Gd as background and a fraction of the simulated IBD coincidence events as signal.

window, and the detector status. As for SK, it has entered a new SK-Gd phase with 0.03% Gd loaded, further enhancing the sensitivity to low-energy $\bar{\nu}_e$. The loading of Gd has brought radioactive contaminants into the FV of the detector, affecting the background rate. Moreover, the reactor neutrino-induced backgrounds in both detectors have changed substantially since the previous studies because many of the nuclear reactors in Japan have been restarted. This section presents a new assessment of the sensitivity to pre-SN neutrinos in KamLAND and SK. Analysis strategies are reoptimized to adapt to the changes.

5.1. Analysis Strategies

The general analysis strategy for each experiment is as follows. A rapid increase in the candidate event rate is sought without explicit reference to any of the pre-SN neutrino models. In each experiment, the background rate is measured over a relatively long period (30 days or more) using recent data, in order to reduce the effects of random fluctuations in the data. A sliding analysis window of a few hours is used to measure the observed event rate for the purpose of searching for signal events. The detection significance is calculated by comparing the observed event rate to the expected background rate. In other words, a test of significance is performed, with the null hypothesis being that the observed event rate is consistent with the expected background rate within the sampling error.

Although KamLAND and SK are at nearly identical locations, they are in rather different experimental conditions, such as target mass, detection energy threshold, background rates, duty cycle, etc. These factors affect the choices of background time window and analysis time window. SK chooses a 30 days background time window, while KamLAND's choice is a longer 90 days background time window due to its lower background rate. The analysis time windows are chosen based on the principle of achieving the longest warning time. As a result, KamLAND has chosen an optimal time window of 24 hr. For SK, the time window was optimized to 12 hr, maximizing the warning time for Betelgeuse-like models and reducing the impact that potential interruptions in DAQ and calibration work in the detector may have in the pre-SN alert system.

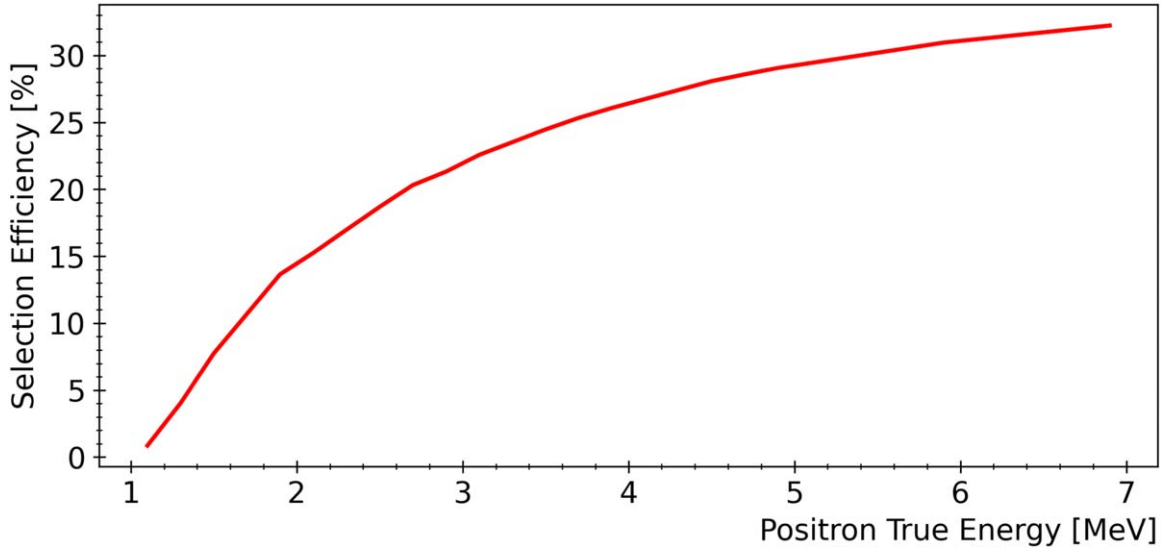


Figure 4. Evaluation of the efficiency of selection after the application of cut criteria to 10^7 IBD pairs as a function of positron true total energy for SK.

In this analysis, both experiments are considered as Poisson counting experiments, and their Poisson likelihoods \mathcal{L}_{SK} and \mathcal{L}_{KL} are constructed. The subscripts SK and KL denote Super-Kamiokande and KamLAND, respectively. The Poisson likelihood for each experiment can be written as

$$\mathcal{L}_x = \frac{(\lambda_x)^{N_x} \exp^{-\lambda_x}}{N_x!}, \quad (1)$$

where the subscript x can be SK or KL. N_x is the observed number of events within the sliding analysis time window. The term λ_x , being the expected number of events, is given by

$$\lambda_x = S_x + B_x, \quad (2)$$

where S_x is the parameter for the number of signals and B_x is the expected number of backgrounds. The test statistic based on the likelihood ratio is given by

$$\Lambda_x = -2 \ln \frac{\text{Max}(\mathcal{L}_x |_{S_x=0})}{\text{Max}(\mathcal{L}_x)}, \quad (3)$$

where the numerator and denominator are the maximum likelihoods with and without imposing a background-only scenario $S_x = 0$ (Cowan et al. 2011). The more the observation disagrees with the background-only hypothesis, the larger Λ_x is. As the test statistic Λ_x asymptotically approaches χ^2 distribution, we consider $\sqrt{\Lambda_x}$ as the detection significance for each experiment.

5.2. Background Assumptions

The background rates in KamLAND and SK can be largely affected by the nuclear reactors near the Kamioka mine. Reactor $\bar{\nu}_e$ are mainly generated through the beta decays of the fission products of ^{235}U , ^{238}U , ^{239}Pu , and ^{241}Pu in nuclear reactors. The Great East Japan Earthquake of 2011 caused numerous reactors in Japan to be shut down, and they have been gradually restarting since 2015. A constant monitoring of the situation in Japan's nuclear power plants is being conducted, and backgrounds are being updated. Nonetheless, it is difficult to predict when and which reactors will start operation in the future. Therefore, we assume different reactor

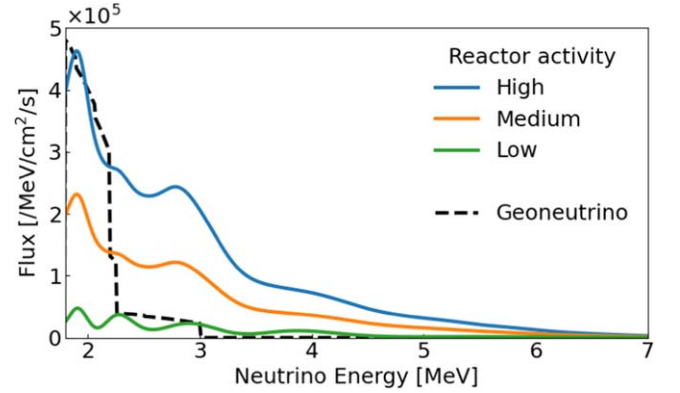


Figure 5. Reactor $\bar{\nu}_e$ fluxes at the Kamioka mine assuming low, medium, and high reactor activities. The relative fission yields (^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu) are assumed to be (0.570, 0.078, 0.295, 0.057) (Eguchi et al. 2003). The $\bar{\nu}_e$ spectra per fission are from Huber (2011), Mueller et al. (2011), and Vogel et al. (1981). The values of neutrino oscillation parameters are $\Delta m_{21}^2 = 7.53 \times 10^{-5} \text{ eV}^2$, $\tan^2 \theta_{12} = 0.436$, and $\sin^2 \theta_{13} = 0.023$. Geoneutrino flux (black) is calculated using the parameters from Enomoto et al. (2007).

background conditions according to three scenarios: low, medium, and high reactor activities. The low reactor activity scenario assumes that all reactors in Japan are not in operation. For the medium reactor activity scenario, reactors near the Kamioka mine (Mihama 3 unit, Ohi 3, 4 units, and Takahama 1, 2, 3, 4 units) are assumed to be operating with a 100% load factor. We note this is close to the situation as of the winter of 2023–2024. The high reactor activity scenario assumes the amount of reactor neutrinos is doubled compared to the medium reactor activity scenario.

Figure 5 shows the expected reactor fluxes at the Kamioka mine considering these three reactor scenarios, as well as the expected geoneutrino fluxes. The fluxes are calculated based on the $\bar{\nu}_e$ spectra per fission from Huber (2011), Mueller et al. (2011), and Vogel et al. (1981), with relative fission yields (^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu) assumed to be (0.570, 0.078, 0.295, 0.057) (Eguchi et al. 2003). The values of neutrino oscillation parameters are $\Delta m_{21}^2 = 7.53 \times 10^{-5} \text{ eV}^2$, $\tan^2 \theta_{12} = 0.436$, and $\sin^2 \theta_{13} = 0.023$. The thermal power of each reactor is

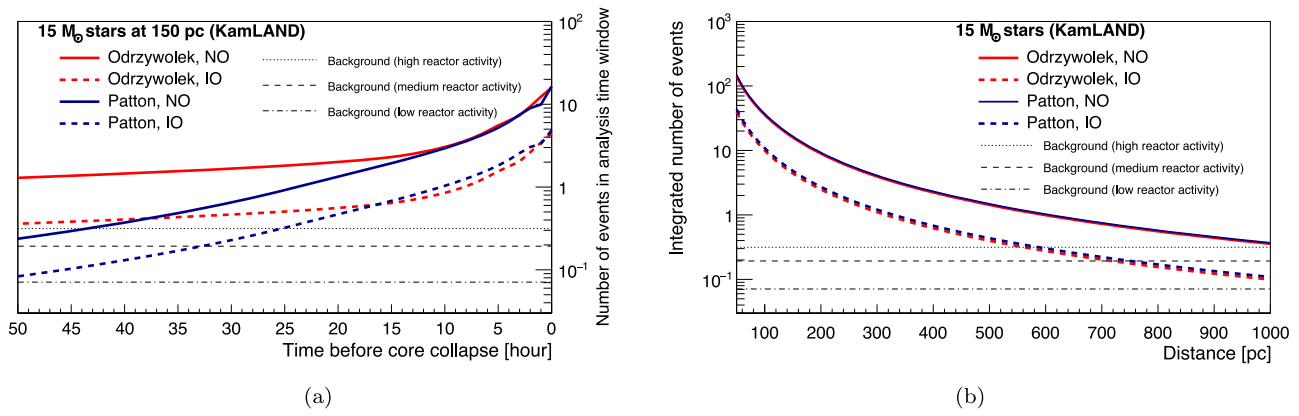


Figure 6. Expected number of signal events in KamLAND as a function of (a) time to core collapse, and (b) distance. Pre-SN $\bar{\nu}_e$ fluxes from a star with $15 M_{\odot}$ are considered, following the Odrzywolek model (red) and the Patton model (blue). For (b), the signal rates are integrated over the last 24 hr before the CCSN. Solid curves show normal neutrino mass ordering and dashed curves show inverted neutrino mass ordering.

estimated from publicly available data on reactor electric power. In the low reactor activity scenario, where all reactors in Japan are assumed to be off, the only contribution we consider comes from reactors in Korea. In addition to the reactor neutrino fluxes, we also plot the geoneutrino flux at the Kamioka mine for comparison. The geoneutrino flux is calculated based on the model in Enomoto et al. (2007). For $\bar{\nu}_e$ energy below 2.2 MeV, the geoneutrino flux is comparable to the reactor neutrino flux under the high reactor activity assumption. It decreases rapidly at ~ 2.2 MeV, and is roughly the same as the reactor neutrino flux assuming low reactor activity between 2.2 and 3 MeV. Above 3 MeV, the geoneutrino flux becomes negligible.

5.3. False Alarm Rate

It is a common practice to quantify the false positive rate of a statistical test using the p -value of the test. However, in this search, we find it misleading to report the results using the p -value or the significance of a single test. The online search of pre-SN neutrinos is performed continuously, while the time when the pre-SN neutrino signal appears is not known in advance. An appropriate way to estimate the p -value is to determine the probability, assuming background-only, of finding a signal at any time that is at least as extreme as the one observed. Such a p -value can be substantially higher than the p -value of a single test. This is the so-called “look-elsewhere effect” (Lyons 2008).

To resolve this issue, we introduce the quantity “false alarm rate” to report the result of the search. The false alarm rate is the expected frequency that a false alarm may happen per century. A false alarm is caused by a significant signal due to background fluctuations. The false alarm rate can be derived from toy Monte Carlo simulations, assuming a background-only scenario. In practice, we generate a time series of Poisson random events with the expected value being the background rate. The statistical test described in Section 5.2 is performed. Then we evaluate the frequency with which the random events cause a significant signal. Thereby, a false alarm rate is found. Considering the frequency of CCSN in the galaxy is approximately once every few decades (Adams et al. 2013), we set a false alarm rate ≤ 1 per century as the ultimate pre-SN alert criteria.

5.4. Sensitivity to pre-SN Neutrinos at KamLAND

We evaluate the expected numbers of signal events in KamLAND for the two pre-SN neutrino models with neutrino oscillation effects, assuming a Betelgeuse-like pre-SN star. Figure 6(a) shows the expected number of signal events, integrated over a sliding 24 hr time window, as a function of time to CCSN. Figure 6(b) shows the integrated number of signals for the last 24 hr before core collapse at different distances. The expected background counts integrated over 24 hr are 0.07, 0.19, and 0.32 events for low, medium, and high reactor activities, respectively. The background rate in KamLAND is sufficiently low, such that a few candidate events can cause a significant signal.

Figure 7 presents the time evolution of the expected detection significance assuming medium reactor activity. The significance corresponding to false alarm rates of 1, 10, and 100 per century are also plotted as horizontal dashed-dotted lines. The remaining time to core collapse, when KamLAND observes an excess of pre-SN neutrino candidates as extreme as a false alarm rate ≤ 1 per century, is defined as the warning time. Note that the latency due to data processing is not taken into account when calculating the warning time. For a Betelgeuse-like pre-SN star, KamLAND is capable of issuing a pre-SN alert 6.5 hr before the CCSN, assuming the Odrzywolek model and normal mass ordering. In the case of inverted mass ordering, the warning time is largely shortened as the pre-SN neutrino fluxes become lower. The worst case is that of the Odrzywolek model and inverted mass ordering; the expected detection significance cannot reach the alert criteria. The discussions above are based on the medium reactor activity scenario. Table 2 summarizes the warning time for all three reactor activity assumptions. If the reactor activity assumption shifts from medium to high, for the normal ordering cases, the warning time will be shortened by roughly 1 hr. For the inverted ordering cases, with the high reactor activity assumption, KamLAND is unable to issue an alert with a false alarm rate of 1 per century, noted as “N/A (Not Applicable)” in Table 2.

Figure 8 shows the warning time as a function of distance. The lines are estimations assuming medium reactor activity. The upper edges of the bands are for low reactor activity, and the lower edges are for high reactor activity. These results indicate that, for nearby pre-SN candidates, KamLAND can send alerts tens of hours before the explosion. For the medium

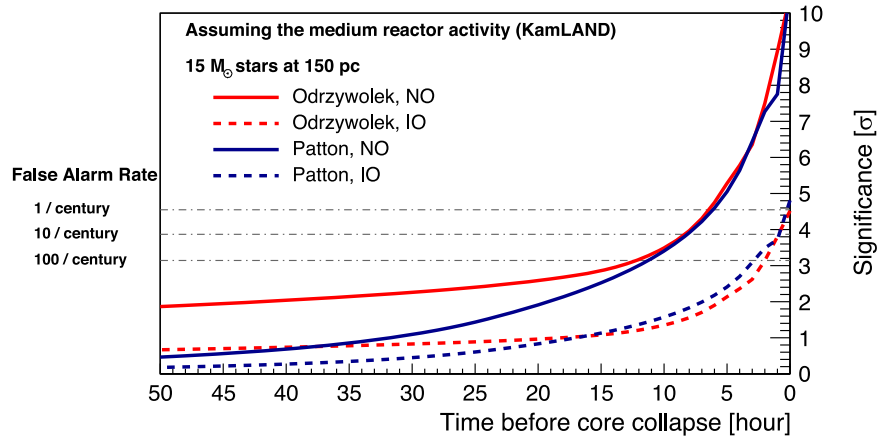


Figure 7. Time evolution of the sensitivity to pre-SN neutrinos in KamLAND, assuming medium reactor activity, following the Odrzywolek model (red) and the Patton model (blue). Solid (dashed) lines are for normal (inverted) neutrino mass ordering. Horizontal dashed-dotted lines indicate false alarm rates of 1, 10, and 100 per century.

Table 2
Warning Times of the Individual and Combined Alert Systems

Alert System	Pre-SN Model	Mass Ordering	Warning Time		
			(hr)		
			Low Reactor Activity	Medium Reactor Activity	High Reactor Activity
KamLAND	Odrzywolek	NO	8.3	6.5	5.5
		IO	0.9	N/A	N/A
	Patton	NO	8.1	6.1	5.0
		IO	0.8	0.2	N/A
SK	Odrzywolek	NO	6.7	6.3	5.9
		IO	2.4	2.1	1.9
	Patton	NO	12.0	10.9	9.8
		IO	4.7	4.3	3.9
Combined	Odrzywolek	NO	9.8	8.0	7.3
		IO	3.0	2.5	2.2
	Patton	NO	14.2	12.4	11.2
		IO	5.4	4.6	4.2

Note. The warning times are estimated assuming a Betelgeuse-like pre-SN star of $15 M_{\odot}$. The latency due to data processing is not taken into account. N/A denotes not applicable, meaning the expected significance does not reach the alert criteria.

reactor activity case, KamLAND is sensitive to pre-SN candidates within an optimistic distance of 280 pc away from Earth.

5.5. Sensitivity to pre-SN Neutrinos at SK

The fluxes of pre-SN neutrinos are taken from the two pre-SN models, with neutrino oscillation effects assuming normal and inverted mass orderings. The expected signal rates are aggregated over a sliding 12 hr time window, resulting in the expected number of signal events as a function of time, as pictured in Figure 9(a). Figure 9(b) presents the number of signals integrated over the last 12 hr for different distances. The expected background counts for low, medium, and high reactor activities in SK are 4.6, 6.2, and 8.1 events. Although the background rate is much higher than that in KamLAND, the large target volume allows SK to collect signal events an order of magnitude larger than KamLAND. Thus the statistical significance in SK can increase rapidly when approaching core collapse.

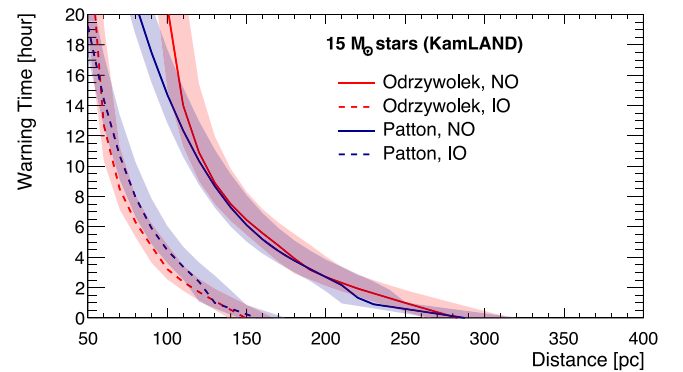


Figure 8. Expected warning time in KamLAND as a function of distance. The lines are estimations assuming medium reactor activity. The upper (lower) edges of the bands are for the low (high) reactor activity case.

Based on the estimations of signal and background, we assessed the sensitivity of pre-SN neutrino detection in SK-Gd with 0.03% Gd loading. Figure 10 presents the time evolution

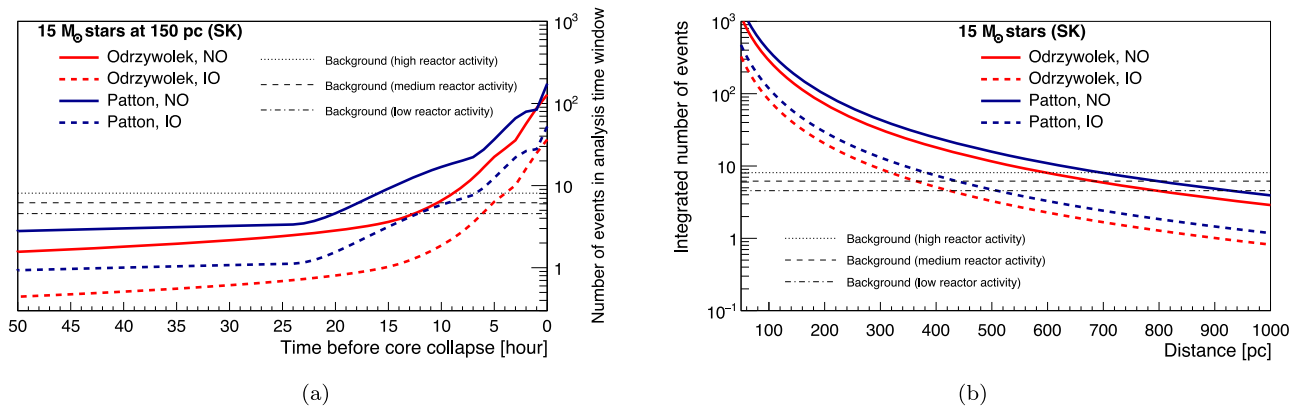


Figure 9. Expected number of signal events in SK with 0.03% Gd concentration as a function of (a) time to core collapse, and (b) distance. Pre-SN $\bar{\nu}_e$ fluxes from a star with $15 M_{\odot}$ is considered, following the Odrzywolek model (red) and the Patton model (blue). For (b), the signal rates are integrated over the last 12 hr before the CCSN. Solid curves show normal neutrino mass ordering and dashed curves show inverted neutrino mass ordering.

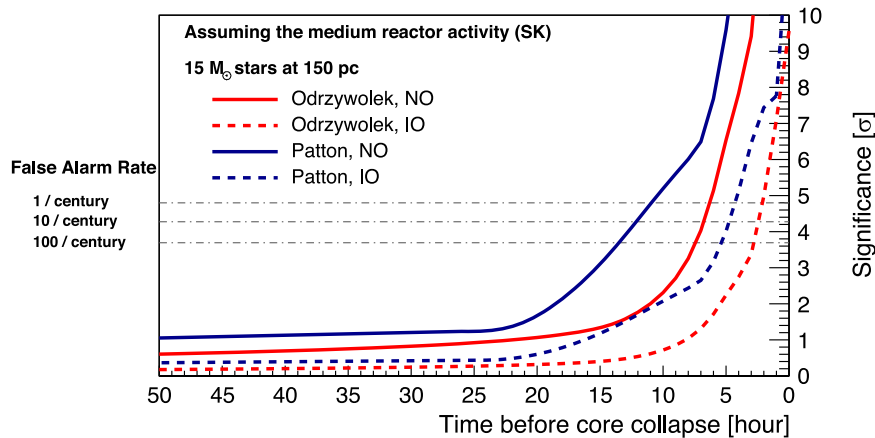


Figure 10. Time evolution of the sensitivity to pre-SN neutrinos in SK with 0.03% Gd concentration, assuming medium reactor activity, following the Odrzywolek model (red) and the Patton model (blue). Solid (dashed) lines are for normal (inverted) neutrino mass ordering. Horizontal dashed-dotted lines indicate false alarm rates of 1, 10, and 100 per century.

of the expected detection significance in SK-Gd, assuming medium reactor activity. The results show that SK-Gd is capable of providing an early warning before the CCSN, at most 10.9 hr assuming the Patton model and normal ordering, and at least 2.1 hr for the Odrzywolek model and inverted ordering. The warning time for all of the simulated scenarios is summarized in Table 2. If the reactor neutrino fluxes around the Kamioka mine are doubled, the warning time can be shortened by 0.2–1.1 hr.

We plot the warning time as a function of distance in Figure 11. The upper (lower) edges of the bands are for low (high) reactor activity, and the lines in between are for medium reactor activity. Under neutrino flux assumptions of the Patton model and normal ordering, the SK alert can cover $15 M_{\odot}$ stars with a distance of 500 pc from Earth, for the medium reactor activity case.

5.6. Discussion

The results shown above suggest that the two detectors, KamLAND and SK, have different advantages in pre-SN neutrino detection. For KamLAND, the low-background rate is an advantage in terms of resolving a small signal. As shown in Figure 6 (a), for a Betelgeuse-like pre-SN candidate, the expected number of signal events exceeds the background count even tens of hours prior to the CCSN. As a result, KamLAND

can provide a warning for nearby supernovae earlier than SK. On the other hand, the number of signal events in KamLAND is limited by the target mass, making it hard to reach stars at far distances. SK is sensitive to CCSN candidates further away from Earth compared to KamLAND. The large target mass of SK can increase the significance rapidly when approaching the CCSN, resulting in a much higher $\bar{\nu}_e$ event rate. But due to its relatively high background rate, SK is less sensitive to small signals. By noting that these two detectors are complementary in pre-SN neutrino searches, a joint search combining measurements from these two detectors should improve the current detection sensitivity. We show in Section 6 that the combined alert benefits from the advantages of these two detectors. We expect the complementary properties of the two detectors to create a synergistic bond, enhancing the sensitivity to pre-SN neutrino signals.

6. Combined Sensitivity to pre-SN Neutrinos

6.1. Statistical Approach for the Combined Search

The purpose of the combined pre-SN alert system is to provide a semi-real-time result of an online search of pre-SN neutrino signals. The chosen strategy is to perform a test of significance every 5 minutes upon the observed numbers of candidates and the expected numbers of background events in KamLAND and SK. The following likelihood function is

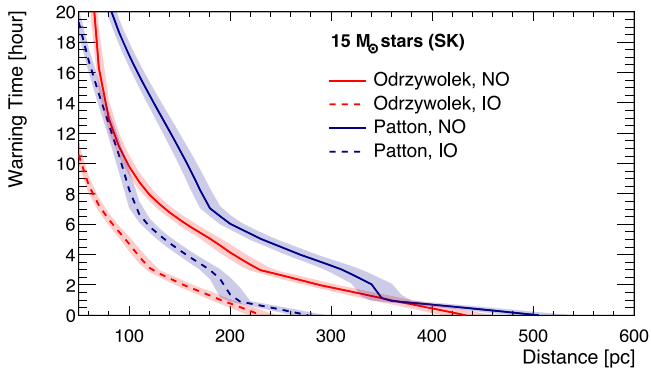


Figure 11. Expected warning time in SK as a function of distance. The lines are estimations assuming medium reactor activity. The upper (lower) edges of the bands are for the low (high) reactor activity case.

constructed, which is a product of the Poisson likelihood of each experiment,

$$\mathcal{L}_{\text{combine}} = \mathcal{L}_{\text{SK}} \times \mathcal{L}_{\text{KamLAND}}. \quad (4)$$

The test statistic based on likelihood ratio Λ_{combine} can be calculated by substituting $\mathcal{L}_{\text{combine}}$ for \mathcal{L}_x in Equation (3). The corresponding significance is found by calculating the chi-square quantile for two degrees of freedom.

For any pre-SN neutrino model, the neutrino fluxes in both detectors should be the same, and thus there are correlations between the numbers of signals in KamLAND and SK. However, we note that the test statistic $\mathcal{L}_{\text{combine}}$ merely reflects the level of agreement between data and the background-only hypothesis. Whether the data is consistent with a pre-SN neutrino model is not tested. Therefore, the parameters of the signal in SK and KamLAND (S_{SK} and S_{KamLAND}) are treated as two independent parameters, and the degrees of freedom are considered as two.

Likewise, the reactor neutrino background in the two detectors is correlated. This correlation does not affect the statistical test, because the expected number of background B is estimated from the data taken in the background time window and normalized to the analysis time window, based on the assumption that the background rate is consistent with the background data taken before. However, such an assumption may not hold, since the reactor $\bar{\nu}_e$ background can change significantly within a week if nuclear reactors are turned on or off. An unexpected increase of reactor $\bar{\nu}_e$ flux, for example, when several nearby reactors are turned on, may cause an excess of $\bar{\nu}_e$ events in the detectors. Although the background rate is still far below the alert criteria even in an extreme case, we can imagine, i.e., the high reactor activity scenario, a higher baseline can increase the risk of sending a false alarm. The reactor neutrino background is irreducible, as it consists of true $\bar{\nu}_e$ events, and its energy range overlaps that of the signal. Concerning this issue, we perform frequent background measurements in both detectors.

As mentioned in Section 5.3, the frequency to find a single combined search at least as extreme as the observation, i.e., the false alarm rate, is utilized to report the results of the pre-SN neutrino search. Figure 12 shows an example contour of the false alarm rate resulting from a toy Monte Carlo simulation. The x -axis and y -axis are the numbers of candidates observed in KamLAND (in 24 hr) and SK (in 12 hr), respectively. The color of each box indicates the false alarm rate level of the corresponding numbers of observed events. If the

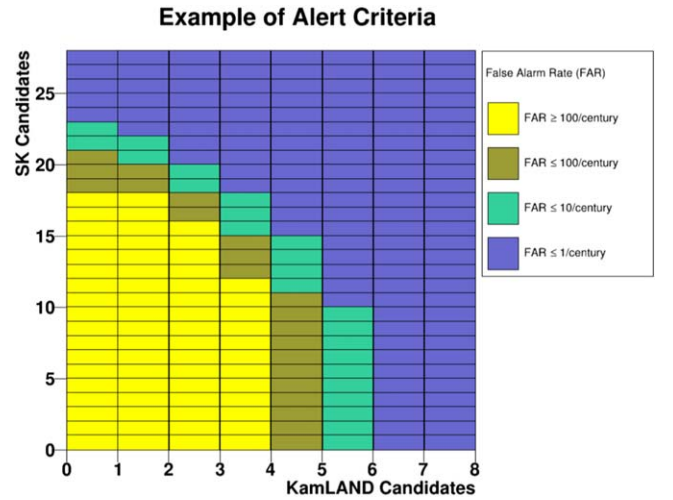


Figure 12. Example contour of the false alarm rates extracted from toy Monte Carlo simulation assuming the background-only hypothesis. The x -axis and y -axis are the numbers of candidates observed in KamLAND (in 24 hr) and SK (in 12 hr). The assumed background rates are 12.4 events per day in SK and 0.19 events per day in KamLAND, the same as those in the medium reactor activity scenario.

observation drops in the yellow regions, which indicates a false alarm rate ≥ 100 per century, there is no hint of a pre-SN neutrino emission. If the observation is in the blue regions, where the false alarm rate ≤ 1 per century, it is considered as a significant excursion of the observed event rate, which could be caused by pre-SN $\bar{\nu}_e$.

6.2. Combined Sensitivity

Following the above-mentioned statistical approach, we perform a joint sensitivity study based on the estimations of S_{KamLAND} , S_{SK} , B_{KamLAND} , and B_{SK} presented in Section 5. Figure 13 shows the time evolution of the combined sensitivity of pre-SN neutrinos from a Betelgeuse-like star of $15 M_{\odot}$, assuming medium reactor activity. For the Patton model and normal mass ordering, the warning time is extended to 12.4 hr prior to the CCSN. Similar to the discussions in Section 5, we assessed the sensitivities for the two pre-SN models, the three reactor activity cases, and the two neutrino mass orderings. Table 2 summarizes the warning time of the combined alert and the individual alerts. These results indicate that the combined alert presents an improved performance, because the warning times are longer compared to either of the individual alerts. It is important to note that, even in the high reactor activity case, the warning time is at least 2.2 hr before the CCSN.

Figure 14 presents the expected warning time and the star distance coverage of the combined alert. Variations due to changes in the reactor neutrino flux are shown as shaded, enveloped by the upper edges resulting from low reactor activity, and the lower edges resulting from high reactor activity. Significant improvements in star distance coverage are observed when comparing to the individual alerts shown in Figures 8 and 11. Assuming $15 M_{\odot}$ stars, the combined alert is able to cover 510 pc for the medium reactor activity case.

These results demonstrate the complementarity of the KamLAND and the SK-Gd detectors. Taking advantage of the low-background rate of KamLAND and the large target mass of SK, the combined alert presents improvement in extending the warning time as well as the distance coverage. In

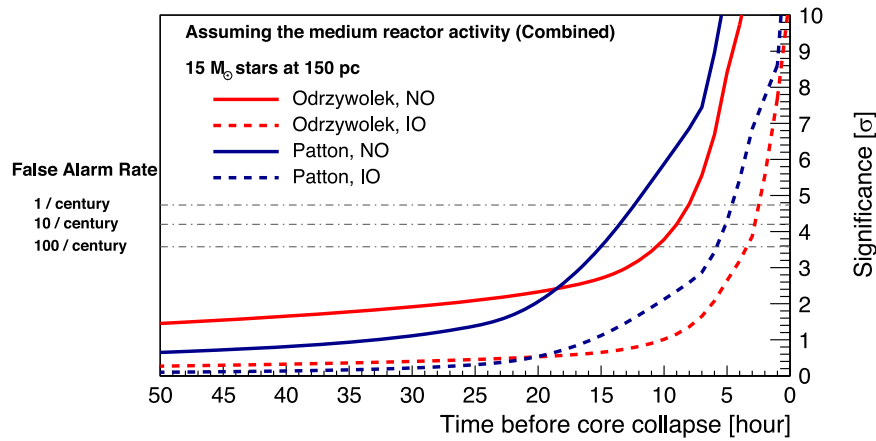


Figure 13. Combined sensitivity to pre-SN neutrinos as a function of time based on the detection capability of KamLAND and SK-Gd with 0.03% Gd concentration, assuming medium reactor activity, following the Odrzywolek model (red) and the Patton model (blue). Solid (dashed) lines are for normal (inverted) neutrino mass ordering. Horizontal dashed–dotted lines indicate false alarm rates of 1, 10, and 100 per century.

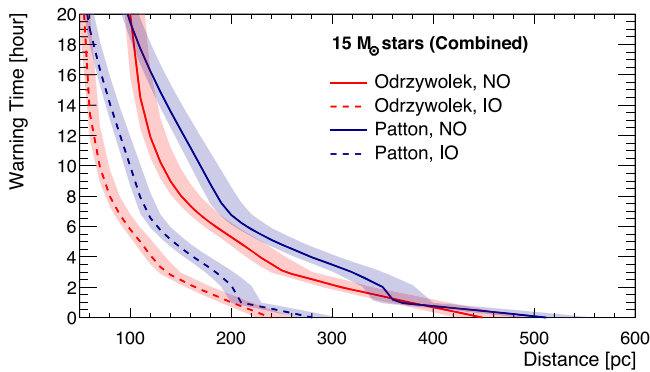


Figure 14. Expected warning time of the combined search as a function of distance. The lines are estimations assuming medium reactor activity. The upper (lower) edges of the bands are for the low (high) reactor activity case.

light of this, a combined pre-SN neutrino alert system was developed, and will be discussed in Section 7.

7. Combined Online Search for pre-SN Neutrinos

The combined pre-SN alert system aims to provide early warning of a potential CCSN upon the detection of pre-SN neutrinos in the KamLAND and SK detectors. It is now operational and ready to issue alarms for CCSNs.

The workflow of the system is introduced in the following. The system receives from both detectors the number of pre-SN neutrino candidates and the expected number of backgrounds. Processed by the DAQ systems, events in the two detectors are selected by their own selection processes following the descriptions in Sections 3.1 and 4.1. The individual pre-SN alert software of KamLAND (SK) then counts the number of observed candidates N_{KL} (N_{SK}) within a 24 hr (12 hr) time window. The expected number of background B_{KL} (B_{SK}) is estimated using data from a background time window of ~ 90 (~ 30) days, as described in Section 5.1. The validity of data is also taken into consideration. Detector status is monitored and marked by a status code. When a detector undergoes calibration work, a test run, or shutdown, the status of the detector is marked as “abnormal” and the data from this period will be invalidated. In addition to the above situations, there may also be network connection problems that can delay the data transfer. Therefore, the differences between the current time

and the time when data are processed will also be checked. The individual KamLAND and SK pre-SN alert systems gather the above information, and exchange them between the servers of KamLAND and SK, as illustrated in Figure 15. The update frequency of the input is once every 5 minutes for each of the experiments. These inputs will be processed by the combined pre-SN alert software, yielding a result of the combined pre-SN search. The result will be exported to users, and the alert decision will be made based on the result.

On the two servers, identically functioning software for the combined pre-SN alert system is installed. If either one of the software pipelines fails, the other one can still output the search result and issue alerts. Figure 16 illustrates the workflow of this software. The software runs on a precise repeating schedule once every 5 minutes. The validity of inputs is first checked by examining the status of each detector and the timestamps of the inputs. If a detector has an abnormal status, or if data from it is delayed for over 15 minutes, input from this detector is invalidated. In this case, instead of exporting the combined search result, the system will output a result based only on the valid input. If all of the inputs are invalid, the result is not applicable. Only when both KamLAND and SK are in normal status and the data are up to date, the system exports the result of the combined search.

In order to determine the corresponding false alarm rate, the software loads three precalculated false alarm rate tables, for KamLAND-only, SK-only, and their combination. Once the software finds a false alarm rate ≤ 1 per century, an alarm will be sent to the Gamma-ray Coordinates Network via an email-based circular. In addition, a text file containing the false alarm rate, along with a timestamp, and a code that denotes whether the result is for KamLAND-only, SK-only, or the combined search, is available to users who have registered on the official website of the combined pre-SN alert system.⁸⁰

The above processes, called main processes, are identical in both servers, as pictured in the blue box with a solid border in Figure 16. An additional process as shown in the orange box with a dashed border is uniquely installed on the SK server. In this process, the false alarm rate tables are updated automatically upon any changes $>5\%$ in the expected numbers of background. This process typically takes ~ 40 minutes, and is

⁸⁰ <https://www.lowbg.org/presnalarm/>

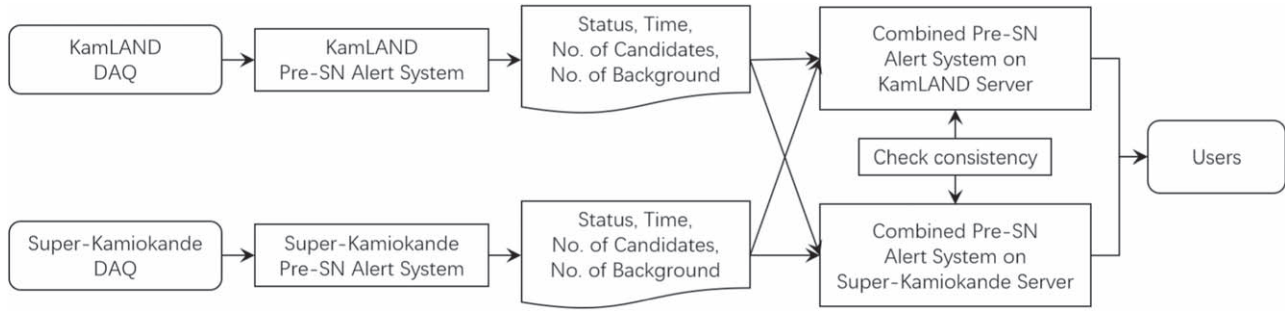


Figure 15. Illustration of the input of the combined pre-SN alert system.

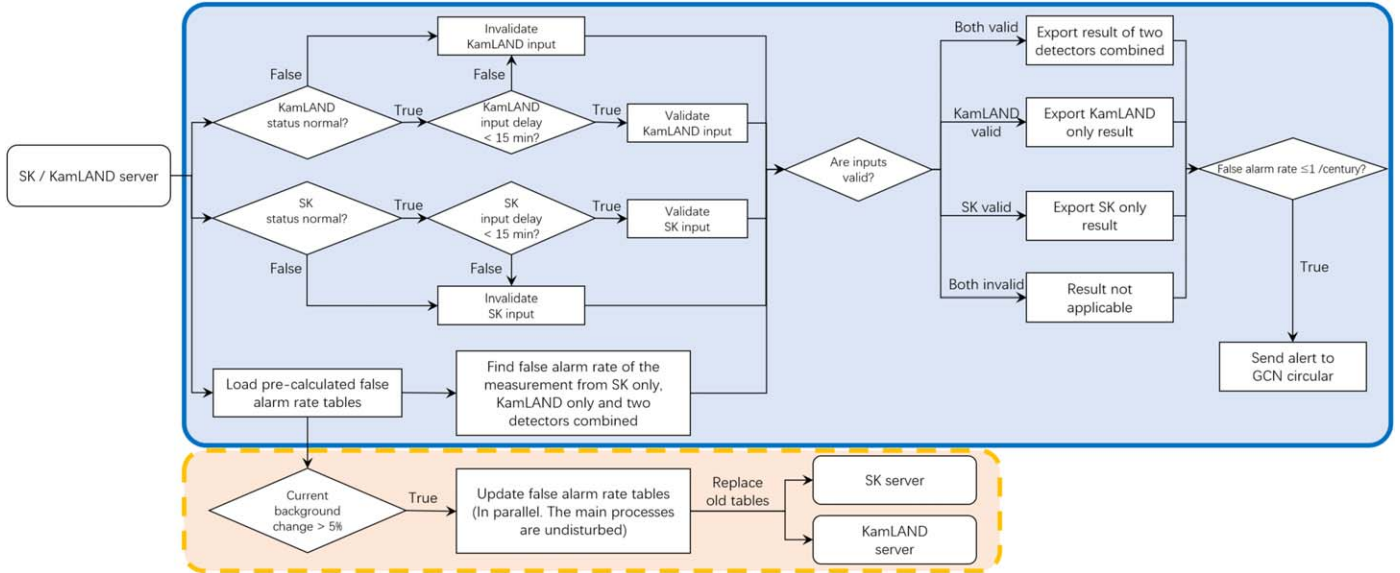


Figure 16. Illustration of the workflow of the combined pre-SN alert system. The main processes to produce results of statistical tests as pictured in the blue box with a solid border are identical for the systems on both servers. The process to update false alarm rate tables as plotted in the orange box with a dashed border is installed exclusively on the SK server.

in parallel with the main processes, in order not to delay the output of the results. Therefore, the precalculated false alarm rate tables do not always correspond to the current background values. However, we note that a significant change in the background rates within an hour is unusual, because the background rates are obtained from measurements of a specific time window long enough to mitigate the effects of statistical fluctuations.

This alert system, operational since 2023 May and accessible to the public, is designed to promptly notify astronomers and particle physicists to maintain operational readiness at their observatories, ensuring they do not miss any impending supernova events. Users can acquire identical results contained in the above-mentioned text file from either of the servers and are encouraged to check the consistency of the results from the two servers before putting them into scientific use. Directional information on pre-SN neutrinos is not available from either SK or KamLAND. More information can be found on the official website.⁸⁰

8. Conclusion

In this study, we present updates on the sensitivity to pre-supernova neutrinos from a Betelgeuse-like star of the individual alert systems of KamLAND and SK, and introduce a combined pre-supernova alert system with the two detectors.

Pre-supernova neutrino fluxes are calculated based on the models from Odrzywolek & A. Heger (2010) and Patton et al. (2017), with neutrino oscillation effects. Due to their similar energy range, reactor neutrinos originating from Japanese nuclear reactors constitute an important source of background for pre-supernova neutrinos. Different reactor activity conditions in Japan are considered in this study, where reactor fluxes vary from low to high. The following results are estimated assuming that nuclear reactors near the Kamioka mine (Mihama 3 unit, Ohi 3, 4 units, and Takahama 1, 2, 3, 4 units) operate with a 100% load factor. The corresponding background rates are 0.19 events per day and 12.4 events per day in KamLAND and SK, respectively.

The best warning times are attained by each detector under different neutrino flux assumptions. In the ideal case, for the Odrzywolek model and normal ordering, KamLAND can provide an early warning 6.5 hr prior to core collapse, and a pre-supernova neutrino emission can be observed up to 280 pc from Earth. The SK pre-supernova alert has an optimistic warning time of 10.9 hr, and is able to cover a distance of 500 pc, assuming the Patton model and normal ordering.

The combined pre-supernova alert system performs a joint statistics test based on the data from the KamLAND and the SK detectors. It has been operational and accessible to the public since 2023 May. Integrating the complementary properties of

the two detectors, the combined alert shows improved sensitivity to pre-supernova neutrinos. An optimistic warning time of 12.4 hr is obtained, for the Patton model and normal ordering, 1.5 hr longer than the SK alert and ~ 6.3 hr longer than the KamLAND alert, with the medium reactor activity assumption. At the same background level, its distance coverage for $15 M_{\odot}$ progenitors is 510 pc, which exceeds those of the individual alerts. While doubling the neutrino fluxes from nearby reactors significantly increases the backgrounds and affects the sensitivity of both detectors individually, the combined alert remains sensitive to pre-supernova neutrino emission with an expected warning time of no less than 2.2 hr for a Betelgeuse-like pre-supernova star, sufficiently long to cover the latency due to data processing. In addition, the combined alert system reduces the dead time for pre-supernova neutrino detection, promoting continuous monitoring even if one of the detectors is temporarily offline. All of these demonstrate the benefits of having a combined search for pre-supernova neutrinos.

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References

- Abe, K., Haga, Y., Hayato, Y., et al. 2016, *Aph*, **81**, 39
Abe, K., Abe, Ke., Aihara, H., et al. 2018, arXiv:1805.04163
Abe, K., Bronner, C., Hayato, Y., et al. 2021a, *ApJ*, **918**, 78
Abe, K., Bronner, C., Hayato, Y., et al. 2021b, *PhRvD*, **104**, 122002
Abe, K., Bronner, C., Hayato, Y., et al. 2022a, *NIMPA*, **1027**, 166248
Abe, K., Akhlaq, N., Akutsu, R., et al. 2023a, *EPJC*, **83**, 782
Abe, K., Bronner, C., Hayato, Y., et al. 2024, *PhRvD*, **109**, 092001
Abe, S., Ebihara, T., Enomoto, S., et al. 2008, *PhRvL*, **100**, 221803
Abe, S., Asami, S., Eizuka, M., et al. 2022b, *ApJ*, **934**, 85
Abe, S., Asami, S., Eizuka, M., et al. 2023b, *PhRvL*, **130**, 051801
Abi, B., Acciarri, R., Acero, M. A., et al. 2021, *EPJC*, **81**, 423
Abusleme, A., Adam, T., Ahmad, S., et al. 2024, *JCAP*, **01**, 057
Acero, M. A., Adamson, P., Agam, G., et al. 2020, *JCAP*, **10**, 014
Adams, S. M., Kochanek, C. S., Beacom, J. F., Vagins, M. R., & Stanek, K. Z. 2013, *ApJ*, **778**, 164
Agafonova, N. Y., Aglietta, M., Antonoli, P., et al. 2015, *ApJ*, **802**, 47
Aiello, S., Albert, A., Garre, S. A., et al. 2021, *EPJC*, **81**, 445
Alekseev, E. N., Alekseeva, L. N., Volchenko, V. I., & Krivosheina, I. V. 1987, *JETPL*, **45**, 589
Alimonti, G., Arpesella, C., Back, H., et al. 2009, *NIMPA*, **600**, 568
Andringa, S., Arushanova, E., Asahi, S., et al. 2016, *AdHEP*, 2016, 6194250
Asakura, K., Gando, A., Gando, Y., et al. 2016, *ApJ*, **818**, 91
Banks, T. I., Freedman, S. J., Wallig, J., et al. 2015, *NIMPA*, **769**, 88
Berger, B. E., Busenitz, J., Classen, T., et al. 2009, *JInst*, **4**, P04017
Bionta, R. M., Blewitt, G., Bratton, C. B., et al. 1987, *PhRvL*, **58**, 1494
Birks, J. B. 1951, *PPSA*, **64**, 874
Carminati, G. 2015, *PhPro*, **61**, 666
Cowan, G., Cranmer, K., Gross, E., & Vitells, O. 2011, *EPJC*, **71**, 1554
Eguchi, K., Enomoto, S., Furuno, K., et al. 2003, *PhRvL*, **90**, 021802
Eldridge, J. J., & Stanway, E. R. 2022, *ARA&A*, **60**, 455
Enomoto, S., Ohtani, E., Inoue, K., & Suzuki, A. 2007, *E&PSL*, **258**, 147
Fukuda, S., Fukuda, Y., Hayakawa, T., et al. 2003, *NIMPA*, **501**, 418
Gando, A., Gando, Y., Hanakago, H., et al. 2012a, *PhRvC*, **85**, 045504
Gando, A., Gando, Y., Ichimura, K., et al. 2012b, *ApJ*, **745**, 193
Gando, A., Gando, Y., Hanakago, H., et al. 2013, *PhRvD*, **88**, 033001
Harada, M., Abe, K., Bronner, C., et al. 2023, *ApJL*, **951**, L27
Hirata, K., Kajita, T., Koshihara, M., et al. 1987, *PhRvL*, **58**, 1490
Huber, P. 2011, *PhRvC*, **84**, 024617
Joyce, M., Leung, S.-C., Molnar, L., et al. 2020, *ApJ*, **902**, 63
Kato, C., Ishidoshiro, K., & Yoshida, T. 2020, *ARNPS*, **70**, 121
Kopke, L. 2018, *JPhCS*, **1029**, 012001
Lyons, L. 2008, *AnApS*, **2**, 887
Machado, L. N., Abe, K., Hayato, Y., et al. 2022, *ApJ*, **935**, 40
Molnar, L., Joyce, M., & Leung, S.-C. 2023, *RNAAS*, **7**, 119
Mori, M., Abe, K., Hayato, Y., et al. 2022, *ApJ*, **938**, 35
Mueller, T., Lhuillier, D., Fallot, M., et al. 2011, *PhRvC*, **83**, 054615
Odrzywolek, A., & Heger, A. 2010, *AcPPB*, **41**, 1611
Odrzywolek, A., Misiaszek, M., & Kutschera, M. 2004a, *AcPPB*, **35**, 1981
Odrzywolek, A., Misiaszek, M., & Kutschera, M. 2004b, *Aph*, **21**, 303
Orii, A., Abe, K., Bronner, C., et al. 2021, *PTEP*, **2021**, 103F01
Patton, K. M., Lunardini, C., Farmer, R. J., & Timmes, F. X. 2017, *ApJ*, **851**, 6
Paxton, B., Bildsten, L., Dotter, A., et al. 2011, *ApJS*, **192**, 3
Pietrzynski, G., Graczyk, D., Gallenne, A., et al. 2019, *Natur*, **567**, 200
Raj, N., Takhistov, V., & Witte, S. J. 2020, *PhRvD*, **101**, 043008
Saio, H., Nandal, D., Meynet, G., & Ekstrom, S. 2023, *MNRAS*, **526**, 2765
Simpson, C., Abe, K., Bronner, C., et al. 2019, *ApJ*, **885**, 133
Smirnov, A. Y. 2003, arXiv:hep-ph/0305106
Suzuki, A. 2014, *EPJC*, **74**, 3094
Takenaka, A., Abe, K., Bronner, C., et al. 2020, *PhRvD*, **102**, 112011
Vogel, P., Schenter, G. K., Mann, F. M., & Schenter, R. E. 1981, *PhRvC*, **24**, 1543
Wester, T., Abe, K., Bronner, C., et al. 2024, *PhRvD*, **109**, 072014
Woolsey, S. E., Heger, A., & Weaver, T. A. 2002, *RvMP*, **74**, 1015