







# Lethal Surface Ozone Concentrations Are Possible on Habitable Zone Exoplanets

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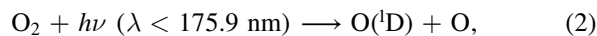
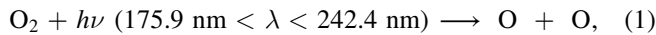
## Abstract

Ozone (O<sub>3</sub>) is important for the survival of life on Earth because it shields the surface from ionizing ultraviolet radiation. However, the existence of O<sub>3</sub> in Earth's atmosphere is not always beneficial. Resulting from anthropogenic activity, O<sub>3</sub> exists as a biologically harmful pollutant at the surface when it forms in the presence of sunlight and other pollutants. As a strong oxidizer, O<sub>3</sub> can be lethal to several different organisms; thus, when assessing the potential habitability of an exoplanet, a key part is determining whether toxic gases could be present at its surface. Using the Whole Atmosphere Community Climate Model version 6 (WACCM6; a three-dimensional chemistry-climate model), 12 atmospheric simulations of the terrestrial exoplanet TRAPPIST-1 e are performed with a variety of O<sub>2</sub> concentrations and assuming two different stellar spectra proposed in the literature. Four atmospheric simulations of the exoplanet Proxima Centauri b are also included. Some scenarios for both exoplanets exhibit time-averaged surface O<sub>3</sub> mixing ratios exceeding harmful levels of 40 ppbv, with 2120 ppbv the maximum concentration found in the cases simulated. These concentrations are toxic and can be fatal to most life on Earth. In other scenarios O<sub>3</sub> remains under harmful limits over a significant fraction of the surface, despite there being present regions that may prove inhospitable. In the case in which O<sub>3</sub> is detected in a terrestrial exoplanet's atmosphere, determining the surface concentration is an important step when evaluating a planet's habitability.

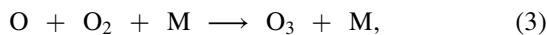
*Unified Astronomy Thesaurus concepts:* [Exoplanets \(498\)](#); [Habitable planets \(695\)](#); [Exoplanet atmospheres \(487\)](#); [Exoplanet atmospheric composition \(2021\)](#)

## 1. Introduction

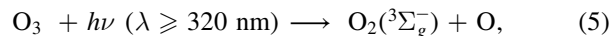
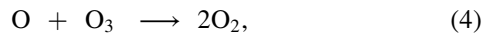
Molecular oxygen (O<sub>2</sub>) makes up 21% by volume of Earth's atmospheric composition and is required for aerobic respiration, provides a fuel for combustion, and gives rise to the "ozone layer." In an atmosphere with O<sub>2</sub> and sufficient ultraviolet (UV) irradiation at wavelengths shortward of 242.4 nm, O<sub>2</sub> is photodissociated into atomic oxygen (O):



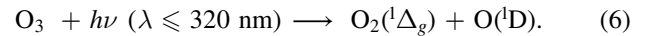
where  $h\nu$  represents a photon of frequency  $\nu$ , and  $h$  is Planck's constant.<sup>3</sup> O<sub>2</sub> and O can form ozone (O<sub>3</sub>) via the following three-body reaction:



where M is any third body (usually N<sub>2</sub> or O<sub>2</sub> on Earth, due to their relatively high abundance). O<sub>3</sub> can also be destroyed through photolysis, or by reacting with atomic oxygen:



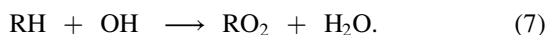
<sup>3</sup> O(<sup>1</sup>D) is the first excited state of atomic oxygen, where 1 represents the spin multiplicity and D is the spectroscopic notation for total orbital angular momentum. Note that the ground state of atomic oxygen can be written as O(<sup>3</sup>P).



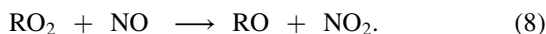
On modern-day Earth, the majority of O<sub>3</sub> resides in the stratosphere, roughly 15–35 km above the surface in the "ozone layer" (Brasseur & Solomon 2005). Here O<sub>3</sub> is beneficial for surface-dwelling life, absorbing biologically harmful UV radiation and providing a partial screen for life exposed to the Sun's radiation. Even though the majority of Earth's O<sub>3</sub> is produced in the equatorial stratosphere, there exists a larger column of O<sub>3</sub> at higher latitudes. This is because O<sub>3</sub> is distributed through a seasonal equator-to-pole circulation driven by atmospheric gravity waves, known as the Brewer–Dobson circulation (Butchart 2014). The Brewer–Dobson circulation has been observed to accelerate and decelerate owing to climate change (Garcia & Randel 2008; Butchart 2014; Fu et al. 2019), consequently affecting regional composition and temperatures near the tropopause and lower stratosphere.

The situation in the troposphere (the lowest atmospheric layer where temperature decreases with altitude) is rather different, because the photolysis rate of O<sub>2</sub> is significantly lower here than in the stratosphere. Near the surface, volatile organic compounds (VOCs) can also contribute to O<sub>3</sub> formation. Hydrocarbon emissions emanate from plants (e.g., isoprene,  $\alpha$ -pinene; Chameides et al. 1988; Sharkey et al. 2008) and, on modern-day Earth, from anthropogenic activity (e.g., naphthalene, acetone, formaldehyde, and many others; Atkinson 2000). When photooxidation of hydrocarbons occurs in the presence of nitrogen oxides, O<sub>3</sub> can eventually be produced through the "smog mechanism" (Haagen-Smit 1952). For example, NO<sub>2</sub> can be produced when OH reacts with a hydrocarbon (Sillman 1999), RH (R is an organic group),

producing RO<sub>2</sub> and H<sub>2</sub>O:



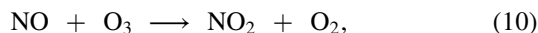
RO<sub>2</sub> leads to NO<sub>2</sub> formation via



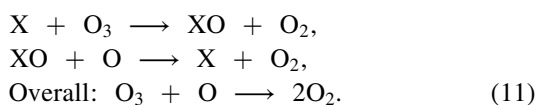
Then, NO<sub>2</sub> is photolyzed in the presence of UV light:



The O produced can lead to reaction 3, making O<sub>3</sub> at low altitudes. On Earth, surface O<sub>3</sub> is low at night when it is removed through reaction with NO,



and O<sub>3</sub> increases during the day owing to photochemistry (Sillman 1999). Reaction 10 is part of a catalytic cycle, where a catalyst (X) leads to the destruction of O<sub>3</sub>, but is ultimately not used up in the overall reaction, such that



At the surface, O<sub>3</sub> is considered a pollutant because it causes oxidative stress to plants, insects, and animals, including humans (Avnery et al. 2011; Silva et al. 2013; Valavanidis et al. 2013; Squire et al. 2014; Démares et al. 2022). Oxidative stress is a chemical imbalance between oxidants and reductants inside an organism that can lead to molecular and biological damage (Lykkesfeldt & Svendsen 2007; Sies et al. 2017). It has been demonstrated in many scenarios that O<sub>3</sub> is an antimicrobial agent, capable of microbial inactivation of fungi, viruses, and bacteria (Kim et al. 1999; Guzel-Seydim et al. 2004; Najafi & Khodaparast 2009; Rojas-Valencia 2011; Fontes et al. 2012; Epelle et al. 2023). For instance, the removal of microbiota was demonstrated using ozonation of the air (Epelle et al. 2022), and aqueous O<sub>3</sub> is effective at inactivating microorganisms (Premjit et al. 2022). Additionally, O<sub>3</sub> has been found to be toxic across a wide range of organisms, including guinea pigs, rats, mice (Giese & Christensen 1954; Stokinger 1965), terrestrial plants (Rich 1964; Bytnerowicz et al. 1993; Sandermann 1996; Rao & Davis 2001; Ramya et al. 2023), aquatic life (Jones et al. 2006), protozoa (Erickson & Ortega 2006), and algae (Hu et al. 2003; Goncalves & Gagnon 2011).

It is useful to consider some quantities to illustrate ozone's danger to life. For example, 40 parts per billion by volume (ppbv) of O<sub>3</sub> is defined by the World Health Organization (WHO) as a critical limit above which crop yield and species biomass may be reduced (World Health Organization et al. 2000). The WHO stated that significant health effects were exhibited by humans at 80 ppbv (World Health Organization et al. 2000), with O<sub>3</sub> damaging lung function at 100 ppbv for 1–8 hr of exposure. Indeed, years of evidence have indicated that long-term exposure to O<sub>3</sub> appears to be related to premature human deaths (Bell et al. 2006; Turner et al. 2016; Sun et al. 2022). For instance, O<sub>3</sub> was attributed to 6000 premature deaths in the EU in 2013 (Nuvolone et al. 2018), and a modeling study by Malashock et al. (2022) calculated a global O<sub>3</sub>-attributable mortality in 2019 of 423,100 deaths (95% confidence interval of 223,200–659,400). The majority (77%) of these were estimated to have occurred in Asia, where ground-level O<sub>3</sub> concentrations were relatively high (Malashock et al. 2022). Furthermore, Feng et al. (2022)

estimated that in East Asia the reduced crop yield to O<sub>3</sub> pollution costs US\$63 billion annually. If O<sub>3</sub> is detrimental to life on Earth, then the same could be possible for extra-terrestrial life. Due to ozone's powerful oxidizing capacity (Menzel 1984; Iriti & Faoro 2007), it is possible that its toxicity to life could be ubiquitous. It is highly reactive, ranking among the highest oxidizers.<sup>4</sup> O<sub>3</sub>, when internal to an organism, causes oxidative stress by releasing reactive oxygen species, which can then cause damage to proteins and DNA and ultimately result in genetic mutations and cell growth that potentially turns into cancer (Klaunig et al. 2010).

O<sub>3</sub>, its spatial distribution on Earth, and its impact on terrestrial organisms have been well studied. Less explored have been the implications of O<sub>3</sub> in exoplanet atmospheres. Hundreds of terrestrial exoplanets, rocky planets orbiting stars other than the Sun, have been detected in our Galaxy. Many of these are in the purported habitable zone (HZ) around their host star (the region in which liquid water could persist on the surface of a rocky exoplanet; Kasting et al. 1993), although the potential for exoplanets and exomoons to be habitable goes beyond the traditional terrestrial-like HZ (see, e.g., Colose et al. 2019; Tjoa et al. 2020; Madhusudhan et al. 2021). If extraterrestrial life exists, then at some point in its evolution it is possible that O<sub>2</sub> could be biologically produced just as it is on Earth, although there are several situations where O<sub>2</sub> could be abiotically produced in high quantities (Des Marais et al. 2002; Domagal-Goldman et al. 2014; Wordsworth & Pierrehumbert 2014; Luger & Barnes 2015; Kleinböhl et al. 2018). These scenarios include major water loss from photolysis (Wordsworth & Pierrehumbert 2014; Luger & Barnes 2015) and high rates of CO<sub>2</sub> photodissociation (Gao et al. 2015; Harman et al. 2015; Schwieterman et al. 2016). Either way, O<sub>3</sub> is a molecule of interest because its detection can indicate the presence of atmospheric O<sub>2</sub> (Leger et al. 1993; Kozakis et al. 2022). Additionally, O<sub>3</sub> has strong spectroscopic signatures in both direct imaging and transmission spectra observations at relatively small volume mixing ratios (e.g., between 10<sup>-7</sup> and 10<sup>-5</sup>; Reinhard et al. 2017; Schwieterman et al. 2018; Kozakis et al. 2022). Due to this property, some work has shown that in particular scenarios O<sub>3</sub> may be easier to detect than O<sub>2</sub> (Reinhard et al. 2017; Kozakis et al. 2022; Cooke et al. 2023b). To date, O<sub>3</sub> has not yet been detected in the atmosphere of a terrestrial exoplanet, so the only estimates of the full O<sub>3</sub> spatial distribution on exoplanets arise from three-dimensional chemistry-climate simulations.

Tidally locked exoplanets are exoplanets that have a rotational period equal to their orbital period (*P*), such that they rotate synchronously (Joshi et al. 1997; Showman & Polvani 2011; Pierrehumbert & Hammond 2019). Carone et al. (2018) simulated tidally locked terrestrial exoplanets with orbital periods of 1–100 days, finding that their atmospheric circulation depends in part on rotation rate. For *P* < 25 days, it was established that stratospheric transport could occur from the pole to the equator (described as an “Anti-Brewer Dobson circulation”), or vice versa, depending on stratospheric wind breaking and the location of the planetary-scale Rossby waves (e.g., tropical or extratropical). At rotational periods greater

<sup>4</sup> F<sub>2</sub> is the strongest oxidizer with a standard electrode potential of 2.87 eV, while O<sub>3</sub> usually ranks second and has a standard electrode potential of 2.075 eV (Kishimoto & Arai 2022). The standard electrode potential is “the value of the standard emf of a cell in which molecular hydrogen under standard pressure is oxidized to solvated protons at the left-hand electrode” (McNaught et al. 1997).

than 25 days, the results from Carone et al. (2018) showed that a thermally driven circulation between the dayside and nightside could widely distribute air parcels. Yates et al. (2020) used the Unified Model (UM) to simulate Proxima Centauri b (assuming a terrestrial exoplanet with a 11.18-day rotation period) in a slab ocean aquaplanet configuration and found that the nightside  $O_3$  lifetime is much higher than it is on the dayside. The same conclusion was reached by Proedrou & Hocke (2016), who simulated a tidally locked Earth with a rotational period of 365 days (no Brewer–Dobson circulation was present on this simulated exoplanet). Chen et al. (2019) used WACCM4 and reported that the pole-to-equator transport predicted by Carone et al. (2018) was present in two of their chemistry-climate simulations for terrestrial exoplanets with periods of 4.11 and 7.91 days and total irradiation of  $1.0S_0$  and  $1.1S_0$ , respectively. Recently, Braam et al. (2023) used the UM and found that  $O_3$  is produced on the dayside and transported to the nightside, with downwelling motions causing  $O_3$  to move into the troposphere at the positions of the nightside gyres. The use of a slab ocean aquaplanet configuration results in highly symmetric winds and chemical transport.

Only a few studies have commented on surface  $O_3$  in paleo atmospheres and exoplanet atmospheres. Grenfell et al. (2013) used a one-dimensional radiative–convective–photochemical model to investigate the atmospheric properties of super-Earths around M0–M7 stars and with surface gravity of  $1g$  and  $3g$  (where  $g = 9.81 \text{ m s}^{-2}$ ). While the smog mechanism was important for  $O_3$  production around later spectral types, the surface  $O_3$  concentrations did not exceed harmful levels. Grenfell et al. (2006) used a box model and showed how the smog mechanism could produce ground-level  $O_3$  up to 3500 ppbv during the Proterozoic (2.4–0.541 Gyr ago) on Earth at 1% the present atmospheric level (PAL) of  $O_2$ . During the Proterozoic,  $O_2$  concentrations could have ranged between  $10^{-5}$  and 1 times the PAL of  $O_2$  (Large et al. 2019; Catling & Zahnle 2020; Steadman et al. 2020; Lyons et al. 2021).

The study by Grenfell et al. (2006) is the only example of a simulated atmosphere that differs from modern Earth where harmful levels of  $O_3$  have been discussed, although the narrative focused on how  $O_3$  would have shielded the early Earth from UV radiation. No previous work has discussed the hypothetical dangers from  $O_3$  for extraterrestrial life on exoplanets and also used a 3D chemistry-climate model that accounts for horizontal transport. This work presents simulations of the exoplanets TRAPPIST-1 e and Proxima Centauri b using WACCM6, a 3D chemistry-climate model. Both exoplanets are located in the supposed HZ of their host stars, and TRAPPIST-1 e is a target for JWST transmission spectra observations. Proxima Centauri b orbits the star Proxima Centauri (M5.5V spectral type, with a stellar effective temperature of 2992 K; Pineda et al. 2021), which is the closest star to the Sun (1.3 pc; Gaia Collaboration et al. 2016), making it an exciting target for future observations (Fowler et al. 2023). TRAPPIST-1 e is a roughly Earth-sized exoplanet orbiting in the HZ around its ultracool M8V (stellar effective temperature of 2566 K; Agol et al. 2021) dwarf host star, TRAPPIST-1. While faint, TRAPPIST-1 is relatively close at a distance of 12.4 pc (40.5 lt-yr). As host of six other terrestrial exoplanets, the TRAPPIST-1 system is a prime target to test theories of planetary system formation and evolution, by confirming whether atmospheres exist on any of the exoplanets and characterizing their properties if they do. To date, analysis

of observations of the exoplanetary thermal emission with JWST suggests that the two innermost exoplanets, TRAPPIST-1 b and c, have either thin atmospheres or no atmosphere at all (Greene et al. 2023; Zieba et al. 2023). Assuming that Earth-like atmospheres exist on both TRAPPIST-1 e and Proxima Centauri b, we investigate the abundance and distribution of  $O_3$  concentrations in different simulated scenarios and discuss the implications for the habitability of oxygenated worlds.

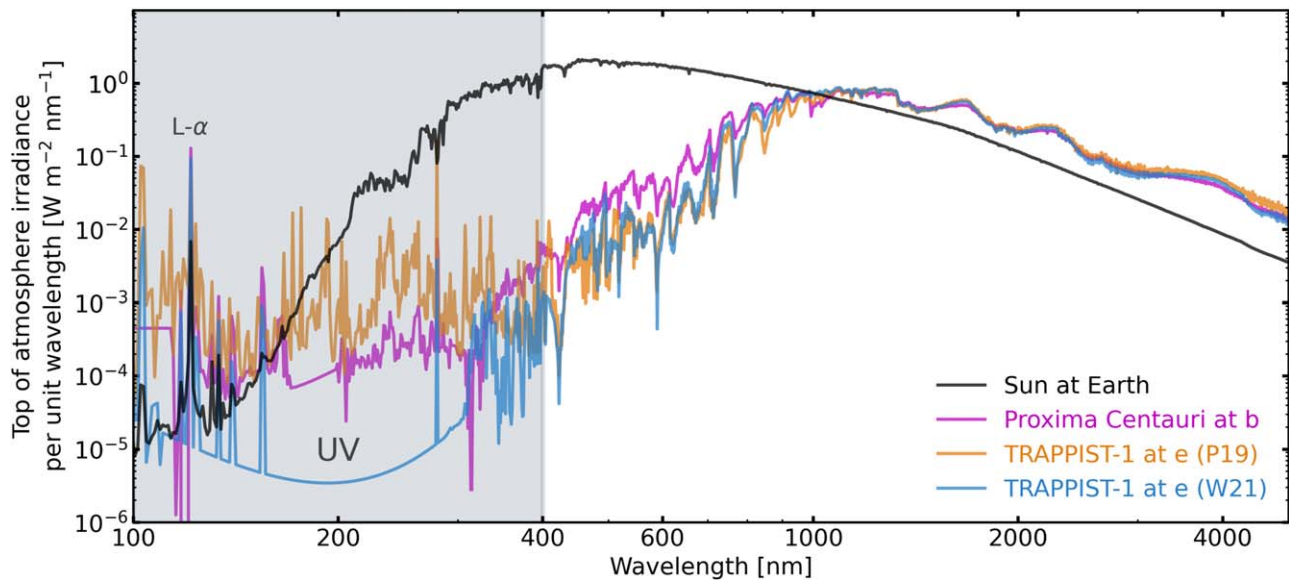
## 2. Simulations

We use the three-dimensional Earth System Model WACCM6 (Gettelman et al. 2019) to perform 12 simulations of possible climates of TRAPPIST-1 e and 4 possible climate simulations of Proxima Centauri b. WACCM6 is a configuration of the Coupled Earth System Model (CESM), and we use version 2.1.3 (CESM2.1.3). In each simulation the initial conditions represent the approximate conditions of Earth’s preindustrial (PI) atmosphere for the year 1850. The simulations have the modern ocean and land configuration, a horizontal resolution of  $1^\circ 875 \times 2^\circ 5$  (96 latitude points and 144 longitude points), and 70 vertical atmospheric levels from 1000 hPa to  $4.5 \times 10^{-6}$  hPa. Both the atmosphere and ocean models are set up to be fully interactive so that they respond to physical changes such as temperature. Because it is likely that Proxima Centauri b and TRAPPIST-1 e rotate synchronously (they may be tidally locked to their host star, although spin-orbit resonance states are plausible; Ribas et al. 2016; Renaud et al. 2021), the substellar point is fixed. This is done by fixing the solar zenith angle in each grid cell. The exoplanet’s obliquity and orbital eccentricity are set to zero. We run WACCM6 with middle atmosphere chemistry, which is described in Emmons et al. (2020), where further details can be found. This chemical mechanism in the WACCM6 simulations has 98 chemical species and 298 chemical reactions, including both the photochemical and heterogeneous reactions that are necessary to simulate the atmospheric conditions of 1850 and, crucially,  $O_3$  chemistry.  $O_3$  pollution due to VOCs is not simulated. The atmospheric time step,  $\Delta t$ , is 30 minutes. The concentrations of 75 species are computed using the implicit method, which considers the chemical system at time  $t$  and  $t + \Delta t$  to evaluate the system at the future time step  $t + \Delta t$  (Sandu et al. 1997). A total of 22 long-lived species are computed with the explicit method, which calculates the chemical system at a later time  $t + \Delta t$  by considering the current system at time  $t$  (Brasseur & Solomon 2005).  $N_2$  is invariant in each simulation, and its mixing ratio in each simulation is adapted to ensure that the atmosphere maintains a surface pressure of 1000 hPa. Following the work by Ji et al. (2023), we include absorption by  $O_3$ ,  $O_2$ ,  $CO_2$ , and  $H_2O$  in the Schumann–Runge bands (175–192 nm).

“Dry deposition” is the process through which atmospheric trace gases and particulate matter are deposited on Earth’s surface and are removed from the atmosphere, and it is an atmospheric sink of  $O_3$ . Dry deposition in WACCM6 (Emmons et al. 2020) was updated following Val Martin et al. (2014) and was originally based on a parameterization from Wesely (1989). The parameterization accounts for variables such as the aerodynamic resistance and the surface resistance and is influenced by vegetation, if present.

We assume that TRAPPIST-1 e receives  $900 \text{ W m}^{-2}$  of irradiation and that Proxima Centauri b receives  $884 \text{ W m}^{-2}$





**Figure 1.** The three stellar input spectra used for the WACCM6 simulations are the PC MUSCLES spectrum for Proxima Centauri at b (magenta) and the P19 (orange) and W21 (blue) spectra for TRAPPIST-1 at e. The top-of-atmosphere irradiance per unit wavelength is shown against the wavelength in nm. In the simulations, TRAPPIST-1 e receives  $900 \text{ W m}^{-2}$  of irradiation ( $0.66 S_{\oplus}$ , where  $S_{\oplus}$  is the total insolation received by Earth), and Proxima Centauri b receives  $884 \text{ W m}^{-2}$  ( $0.65 S_{\oplus}$ ). The UV range is highlighted in gray between 100 and 400 nm, and the Ly $\alpha$  line is labeled. The average difference between the TRAPPIST-1 spectra in the UV range is a factor of 500, with a difference of up to 5000 in some wavelength bins.

( $0.66S_0$  and  $0.65S_0$ , respectively, where  $S_0$  is the total insolation that Earth receives). This is consistent with previous work on Proxima Centauri b (Boutle et al. 2017; Yates et al. 2020; Braam et al. 2022; Ridgway et al. 2023) and with the TRAPPIST-1 Habitable Atmosphere Intercomparison (THAI) project (Fauchez et al. 2020; Turbet et al. 2022). Proxima Centauri b was detected using the radial velocity method and has a minimum mass measured of  $M_p \sin i = 1.07 M_{\oplus}$  (Faria et al. 2022) only, where  $M_p$  is the mass of the exoplanet and  $i$  is the inclination angle of the planetary orbit. Therefore, a recently estimated mass–radius relationship from Otegi et al. (2020), given as  $R_p = 1.03M_p^{0.29}$ , was used to estimate the planetary radius. Assuming an optimistic mass of  $M_p = 1.07 M_{\oplus}$ , this places the radius of Proxima Centauri b at  $1.05 R_{\oplus}$  and the surface gravity of Proxima Centauri b at  $12.2 \text{ m s}^{-2}$ . In our simulations, TRAPPIST-1 e is set to have a mass of  $0.772 M_{\oplus}$  and a radius of  $0.91 R_{\oplus}$ , consistent with the THAI project and transit timing variations from Grimm et al. (2018). The surface gravity of TRAPPIST-1 e is therefore set to  $9.14 \text{ m s}^{-2}$ .

Two semiempirical stellar spectra were used in the TRAPPIST-1 e simulations. Peacock et al. (2019, hereafter P19) modeled the spectral energy distribution (SED) of TRAPPIST-1 and produced models 1A, 2A, and 2B, of which we use model 1A (ver. 1; Peacock 2020). More recently, Wilson et al. (2021b, hereafter W21) used further HST observations to produce a semiempirical SED of TRAPPIST-1 (ver. 1; Wilson et al. 2021a) as part of the Mega-MUSCLES series (Froning et al. 2019; Wilson et al. 2021b). Details of their spectra can be found in the aforementioned references. Both the stellar spectra are included here to illustrate how different strengths and shapes of the incoming UV radiation environment can affect the abundance and distribution of surface  $\text{O}_3$ . For Proxima Centauri b, we use the GJ 551 MUSCLES (ver. 2.2; France et al. 2016; Loyd et al. 2016; Youngblood et al. 2016) spectrum as input.<sup>5</sup> GJ 551 is the

Gliese–Jahreij catalog name for Proxima Centauri. We present the model spectra used in this study in Figure 1.

Note that the TRAPPIST-1 e simulations were started in the year 2020, before Agol et al. (2021) published updates to planetary parameters in the TRAPPIST-1 system. For TRAPPIST-1 e, Agol et al. (2021) gave mass and radius values of  $0.69 M_{\oplus}$  and  $0.92 R_{\oplus}$ , respectively, meaning that the surface gravity would be  $8.015 \text{ m s}^{-2}$ , instead of  $9.14 \text{ m s}^{-2}$  as used here. Using these updated values, the scale height of the atmosphere would increase, but we expect that simulations with the parameters from Agol et al. (2021) would produce similar surface  $\text{O}_3$  mixing ratios to the ones we present here. Only the minimum mass has been measured for Proxima Centauri b, so it is conceivable that it may have a larger mass and radius than the values used here. Brugger et al. (2016) estimated the radius to be in the range  $0.94\text{--}1.40 R_{\oplus}$ , placing it somewhere between a Mercury-like exoplanet and an ocean-like world. Regardless, with M dwarf stars being so numerous, it is plausible that somewhere there exists an exoplanet with similar size and instellation, such that these simulations remain useful should Proxima Centauri b eventually be confirmed to have a mass or radius that is significantly larger.

A summary of the simulations is given in Table 1. For TRAPPIST-1 e, six simulations use the P19 (stronger UV) spectrum, and six simulations use the W21 (weaker UV) spectrum. For both exoplanets, atmospheric concentrations of  $\text{O}_2$  at the PAL (PI; 0.21 by volume), 10% PAL, 1% PAL, and 0.1% PAL are simulated. For TRAPPIST-1 e, we move the substellar point for the 100% PAL simulation between the Pacific Ocean ( $180^\circ$  longitude; PI case) and Africa ( $30^\circ$  longitude; PI SPL case). We run two simulations that are not tidally locked and have a rotational period of 1 day. Whether a slab or dynamic ocean is implemented, and whether the land or ocean is at the substellar point, can modulate the climatology of exoplanets (Hu & Yang 2014; Lewis et al. 2018; Del Genio et al. 2019; Salazar et al. 2020; Zhao et al. 2021; Macdonald et al. 2022; Olson et al. 2022). Salazar et al. (2020) found that

<sup>5</sup> GJ 551 found at <https://archive.stsci.edu/prepds/muscles/>

**Table 1**  
The Sixteen Simulations Used in This Study

Simulation	Planet	Spectrum	O <sub>2</sub> Mixing Ratio (PAL)	Orbital Parameters
W21 PI	TRAPPIST-1 e	W21	1.000	$P = 6.1$ days, SP = 180° lon
W21 PI noTL	TRAPPIST-1 e	W21	1.000	$P = 1$ day
W21 PI SPL	TRAPPIST-1 e	W21	1.000	$P = 6.1$ days, SP = 30° lon
W21 10% PAL	TRAPPIST-1 e	W21	0.100	$P = 6.1$ days, SP = 180° lon
W21 1% PAL	TRAPPIST-1 e	W21	0.010	$P = 6.1$ days, SP = 180° lon
W21 0.1% PAL	TRAPPIST-1 e	W21	0.001	$P = 6.1$ days, SP = 180° lon
P19 PI	TRAPPIST-1 e	P19	1.000	$P = 6.1$ days, SP = 180° lon
P19 PI noTL	TRAPPIST-1 e	P19	1.000	$P = 1$ day
P19 PI SPL	TRAPPIST-1 e	P19	1.000	$P = 6.1$ days, SP = 30° lon
P19 10% PAL	TRAPPIST-1 e	P19	0.100	$P = 6.1$ days, SP = 180° lon
P19 1% PAL	TRAPPIST-1 e	P19	0.010	$P = 6.1$ days, SP = 180° lon
P19 0.1% PAL	TRAPPIST-1 e	P19	0.001	$P = 6.1$ days, SP = 180° lon
PCb PI	Proxima Centauri b	PC MUSCLES	1.000	$P = 11.18$ days, SP = 180° lon
PCb 10% PAL	Proxima Centauri b	PC MUSCLES	0.100	$P = 11.18$ days, SP = 180° lon
PCb 1% PAL	Proxima Centauri b	PC MUSCLES	0.010	$P = 11.18$ days, SP = 180° lon
PCb 0.1% PAL	Proxima Centauri b	PC MUSCLES	0.001	$P = 11.18$ days, SP = 180° lon

**Note.** Twelve for TRAPPIST-1 e: six with the P19 spectrum, and six with the W21 spectrum. There are four simulations of Proxima Centauri b, where the MUSCLES spectrum of Proxima Centauri (see text for details) is used as stellar input. Each simulation started with the PI WACCM6 simulation composition. Each set of six TRAPPIST-1 e simulations includes three with the PAL of O<sub>2</sub> (0.21 by volume, all denoted as “PI”), one where the substellar point is placed over the Pacific Ocean, one where it is placed over Africa (SPL), and one where it is not tidally locked and the rotation rate is 1 Earth day (noTL). Then, the 10% PAL, 1% PAL, and 0.1% PAL simulations have reduced O<sub>2</sub> mixing ratios from the PI simulation by 10, 100, and 1000 times, respectively. Each of the TRAPPIST-1 e simulations receive a total instellation of 900 W m<sup>-2</sup>, and the Proxima Centauri b simulations receive 884 W m<sup>-2</sup> of irradiation. The Proxima Centauri b simulations include the PI, 10% PAL, 1% PAL, and 0.1% PAL cases. The simulated radius and mass of TRAPPIST-1 e are 0.91 R<sub>⊕</sub> and 0.772 M<sub>⊕</sub>, respectively. For Proxima Centauri b, the radius and mass are 1.05 R<sub>⊕</sub> and 1.07 M<sub>⊕</sub>, respectively. The ⊕ subscript denotes values relative to Earth. The orbital parameters assume zero eccentricity and 0° obliquity, and the table lists the period  $P$  and the longitude of the substellar point (SP) relative to Earth’s coordinates (the latitude of the SP is always 0°). Each simulation has been run out for at least 250 model Earth years.

broad climate differences between models (ROCKE-3D and the UM) were larger than the difference between a slab ocean and a dynamic ocean. Because of this previous work and the fact that the PI and PI SPL simulations show only small globally averaged chemical differences in O<sub>3</sub>, and due to computational expense (WACCM6 takes 1332 core-hours per simulated year to run), we do not simulate the substellar point over land in any of the reduced O<sub>2</sub> cases. Each simulation has been run out for at least 250 model Earth years, and then we present the last year of data (365 Earth days). The full details of the model setup, alongside simulation scripts, are available via GitHub.<sup>6</sup>

### 3. Results

#### 3.1. Surface Ozone Concentrations

Figure 2 shows the time-averaged distribution of O<sub>3</sub> at the atmospheric level closest to surface for the cases that have time-averaged surface O<sub>3</sub> mixing ratios of 40 ppbv or greater. We use 40 ppbv as a lower cutoff for harmful levels of surface O<sub>3</sub> (World Health Organization et al. 2000). Gray indicates regions below 40 ppbv, while the yellow/orange/purple color map indicates regions where O<sub>3</sub> exceeds 40 ppbv.

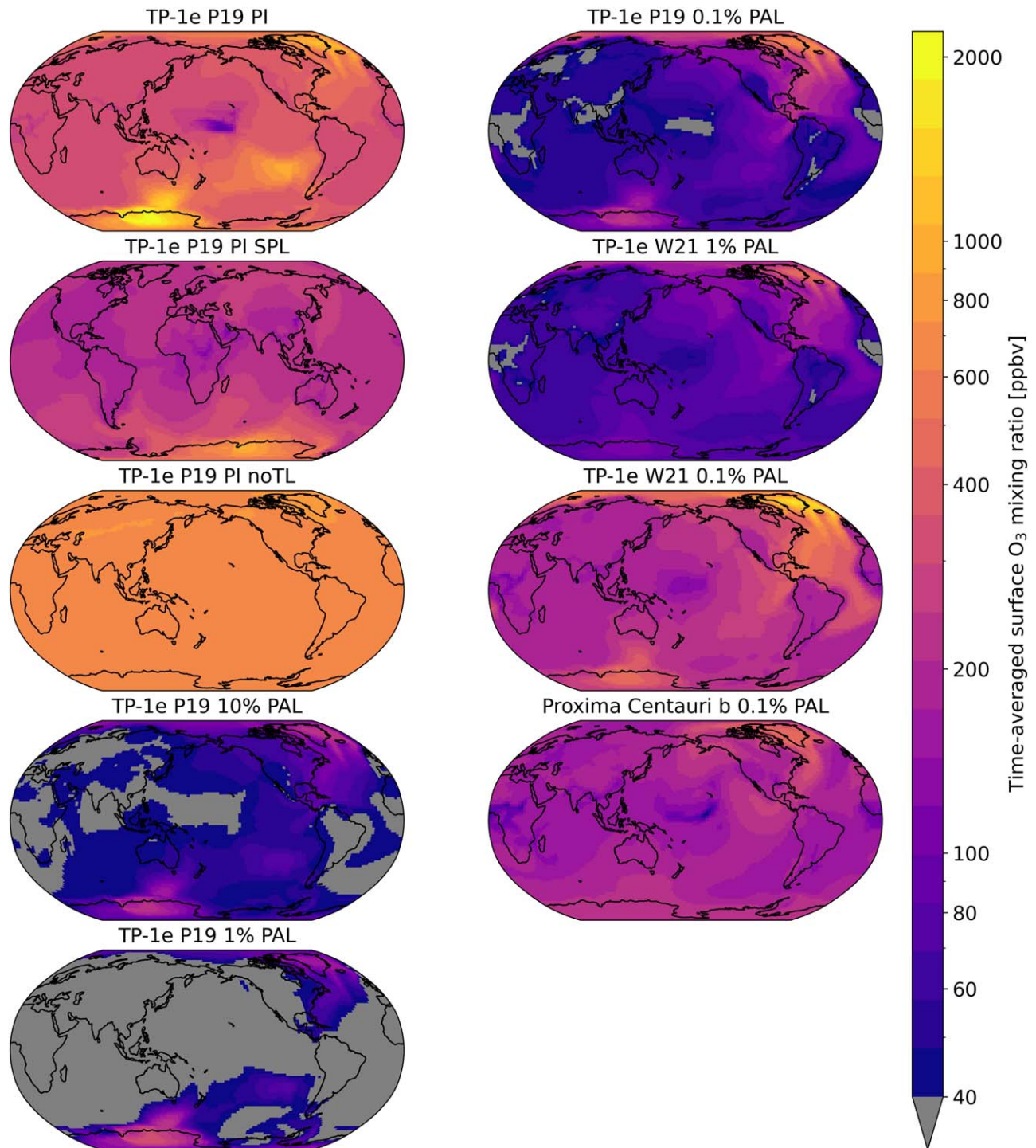
The TRAPPIST-1 e (TP-1 e) P19 PI, PI SPL, noTL, and W21 0.1% PAL simulations everywhere exceed 40 ppbv for surface O<sub>3</sub>. The P19 PI simulation has a maximum mixing ratio of 2120 ppbv, which is the largest surface O<sub>3</sub> mixing ratio in all of the simulations presented. In the P19 PI and 0.1% PAL simulations and the W21 0.1% PAL simulation, specific

locations (e.g., Antarctica or Greenland) have extremely high mixing ratios at certain times, exceeding 1000 ppbv, which is deadly to some organisms on Earth. On the other hand, the W21 PI, PI SPL, and noTL simulations everywhere have O<sub>3</sub> mixing ratios below 40 ppbv and are not shown. The low O<sub>3</sub> surface concentrations are a consequence of the upper atmosphere efficiently absorbing UV such that insufficient UV reaches altitudes closer to the planetary surface to synthesize enough O<sub>3</sub>. On time average, the P19 10%, 1%, and 0.1% PAL simulations and the W21 1% simulation have some areas where O<sub>3</sub> exceeds 40 ppbv, while maintaining regions below this limit. For the Proxima Centauri b (PCb) cases, the PI, 10% PAL, and 1% PAL cases have surface O<sub>3</sub> levels below 40 ppbv everywhere. For the 0.1% PAL PCb scenario, surface O<sub>3</sub> everywhere exceeds 40 ppbv and has a global mean mixing ratio of 200 ppbv.

In terms of time variability, the surface O<sub>3</sub> concentrations are not static. Taking the last year of simulated data and averaging each calendar month, the fraction of land for each simulation where surface O<sub>3</sub> concentrations are under 40 ppbv is given in Table 2. For example, the P19 10%, 1%, and 0.1% PAL simulations have monthly O<sub>3</sub> surface concentrations under harmful levels varying between 12%–44%, 75%–83%, and 4%–9% of the total surface area. Considering all these scenarios, the prospect is raised for safe areas on exoplanets that are sheltered from hazardous O<sub>3</sub> concentrations found at other locations. Meanwhile, some locations will fluctuate between toxic and safe levels. Only the P19 PI noTL simulation has surface O<sub>3</sub> mixing ratios that everywhere exceed 40 ppbv throughout the final year of data.

The bars in Figure 3(a) indicate the full range of surface O<sub>3</sub> mixing ratios in each of the tidally locked simulations where

<sup>6</sup> [https://github.com/exo-cesm/CESM2.1.3/tree/main/Tidally\\_locked\\_exoplanets](https://github.com/exo-cesm/CESM2.1.3/tree/main/Tidally_locked_exoplanets)



**Figure 2.** The surface mixing ratio of  $O_3$ , in ppbv, is displayed for the simulations in this work that have time-averaged surface  $O_3$  mixing ratios exceeding 40 ppbv. This includes all TP-1 e **P19** cases, the TP-1 e **W21** 1% PAL and 0.1% PAL simulations, and only a single Proxima Centauri b case (0.1% PAL). The PI cases start with an initial PI atmospheric composition. The PI case has the substellar point placed over the Pacific Ocean, the SPL case has it placed over Africa, and the noTL case is not tidally locked, so that the substellar point moves with time. PAL is the PAL of  $O_2$ , which is a mixing ratio of 21% by volume. The TP-1 e **P19** simulations have stronger incident ultraviolet radiation than the TP-1 e **W21** simulations. See Table 1 for a more detailed description of the simulations. Gray indicates where the  $O_3$  mixing ratio is below 40 ppbv and thus at “safe” levels, while the different shades of yellow/orange/purple indicate places that exceed 40 ppbv, i.e., these concentrations are known to be harmful to life on Earth. The color map has a log scale that extends from 40 to 2120 ppbv.

the substellar point is placed over ocean, with the global mean surface concentrations indicated by the points. Simulations of Earth (using WACCM6) at atmospheric  $O_2$  mixing ratios between 0.1% and 150% PAL from data from Cooke et al. (2022) are given for comparison. Figure 3(b) presents the same data for the TP-1 e PI scenarios. The TP-1 e **W21** PI and 10%

PAL cases have a large range in surface  $O_3$  concentrations, spanning 7 and 8 orders of magnitude, respectively. All other TP-1 e simulations span 5 orders of magnitude or less, with the Earth simulations spanning approximately 1 order of magnitude. The PCb simulations span between 9 and 2 orders of magnitude. While both the PI and PI SPL simulations with



**Table 2**  
Dry Deposition and Surface Mixing Ratios of O<sub>3</sub>

Simulation	O <sub>3</sub> Dry Deposition Flux (kg m <sup>-2</sup> s <sup>-1</sup> )	O <sub>3</sub> Surface Mixing Ratio (ppbv)	Fraction of Surface with O <sub>3</sub> Mixing Ratios <40 ppbv	O <sub>3</sub> Column (DU)
Earth PI	$2.1 \times 10^{-11}$	12	N/A	297
W21 PI	$1.0 \times 10^{-12}$	3	100%–100%	53
W21 PI noTL	$4.3 \times 10^{-15}$	5	100%–100%	46
W21 PI SPL	$1.3 \times 10^{-14}$	6	100%–100%	54
W21 10% PAL	$6.2 \times 10^{-13}$	2	100%–100%	154
W21 1% PAL	$1.8 \times 10^{-11}$	82	1%–4%	901
W21 0.1% PAL	$4.9 \times 10^{-11}$	246	0%–0%	1227
P19 PI	$9.1 \times 10^{-11}$	404	0%–0%	1289
P19 PI noTL	$4.5 \times 10^{-12}$	692	0%–0%	1245
P19 PI SPL	$4.3 \times 10^{-11}$	243	0%–0%	1098
P19 10% PAL	$1.2 \times 10^{-11}$	51	12%–44%	498
P19 1% PAL	$6.3 \times 10^{-12}$	31	75%–83%	260
P19 0.1% PAL	$1.5 \times 10^{-11}$	81	4%–9%	407
PCb PI	$1.7 \times 10^{-12}$	5	99%–100%	179
PCb 10% PAL	$4.0 \times 10^{-12}$	11	100%–100%	134
PCb 1% PAL	$2.2 \times 10^{-12}$	7	100%–100%	256
PCb 0.1% PAL	$4.6 \times 10^{-11}$	203	0%–1%	790

**Note.** The time-averaged and global mean dry deposition flux of O<sub>3</sub> is given in terms of kg m<sup>-2</sup> s<sup>-1</sup>, for all the exoplanet simulations used in this work, as well as the Earth PI Simulation. Dry deposition is a sink for atmospheric surface O<sub>3</sub>. The time-averaged and global mean O<sub>3</sub> surface mixing ratios are given. The fraction of the surface where O<sub>3</sub> concentrations are under 40 ppbv and thus considered “safe” is given for each simulated exoplanet scenario. The O<sub>3</sub> concentrations vary every calendar month, so the fraction is given as a range over a 1 yr period and as a percentage. Additionally, the global mean O<sub>3</sub> column is given in Dobson units (DU), where 1 DU is equal to  $2.69 \times 10^{20}$  molecules m<sup>-2</sup>.

the **P19** spectrum have harmful O<sub>3</sub> mixing ratios at the surface, the mean surface O<sub>3</sub> mixing ratio is reduced by 1.7 times when the substellar point is placed over land (PI SPL case). In contrast, the **W21** PI SPL case has 1.7 times the mean surface O<sub>3</sub> mixing ratio of the **W21** PI case. Additionally, the noTL cases have a smaller range than both of the PI tidally locked cases, similar to the Earth simulations. These results imply that surface topography, the rotation rate, whether or not a diurnal cycle exists, and the position of the substellar point will be important for modulating surface concentrations of biologically toxic gases such as O<sub>3</sub>. For the PCb scenarios, the PCb 0.1% PAL case reaches the largest O<sub>3</sub> mixing ratio of 466 ppbv, and everywhere has mixing ratios exceeding 40 ppbv. None of the WACCM6 Earth simulations from Cooke et al. (2022) have time-averaged O<sub>3</sub> mixing ratios at dangerous levels, which is to be expected because industrial pollutants were not included in the simulations. If pollutants were included near urban areas, for instance, then harmful O<sub>3</sub> levels would be localized, rather than on a planetary scale. All simulations with O<sub>2</sub> concentrations at 0.1% PAL have toxic mixing ratios exceeding 400 ppbv, and the simulations that have mean surface O<sub>3</sub> concentrations the most toxic to life are the PI, PI SPL, and PI noTL simulations with the **P19** assumed spectrum. Quantitatively, we find harmful O<sub>3</sub> mixing ratios at the surface when global mean stratospheric O<sub>3</sub> number density exceeds  $2 \times 10^{18}$  molecules m<sup>-3</sup>.

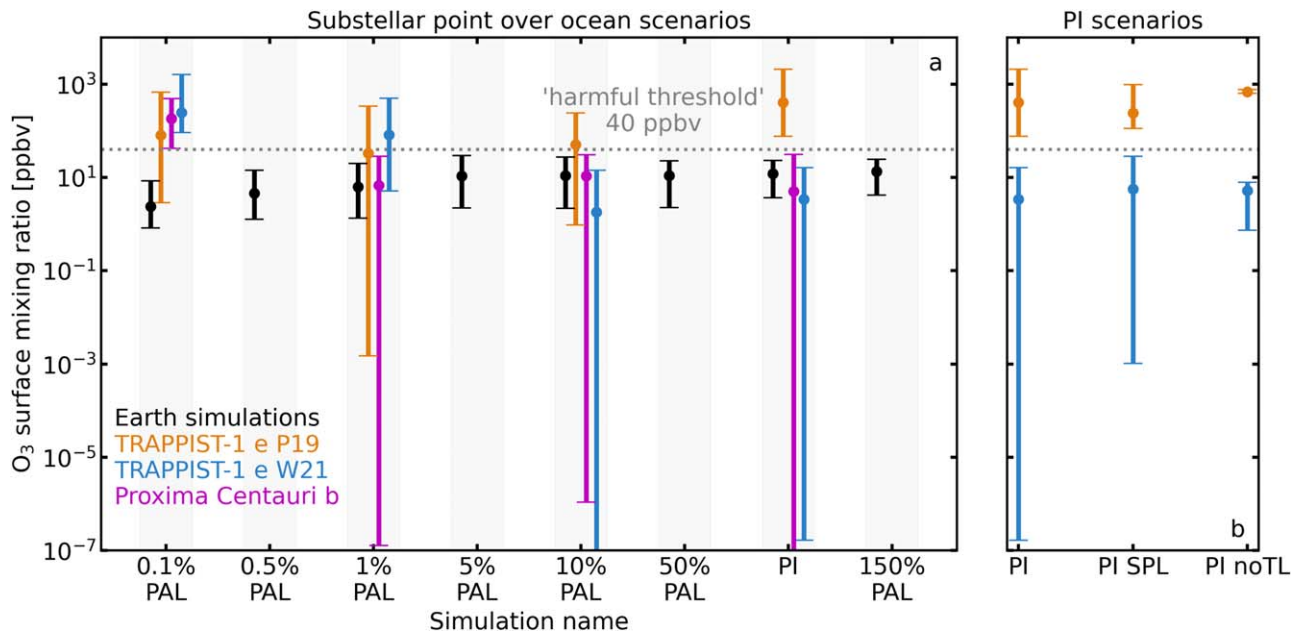
### 3.2. Cause of Ozone Production at the Surface

It is important to state here that initially in the simulations O<sub>3</sub> was present in Earth-like quantities throughout the troposphere, and no toxic concentrations existed at the start of the tidally locked simulations. The O<sub>3</sub> profile depends on UV radiation, O<sub>2</sub> number density, O<sub>3</sub> production rates, O<sub>3</sub> loss, and the transport of O<sub>3</sub>. Figure 4 shows the photolysis rates of O<sub>2</sub> leading to O production (reactions 1 and 2), the photolysis rates

of O<sub>3</sub> leading to O production (reactions 5 and 6), the production rates of O<sub>3</sub> (reaction 3), and the O<sub>3</sub> number density in each simulation. These quantities are important for understanding where O<sub>3</sub> is produced and its resulting number density. O<sub>2</sub> has an approximately constant mixing ratio up until the homopause, where gases start to diffusively separate, but in contrast, the mixing ratio of O increases with altitude until the homopause.

In the **W21** simulations, as O<sub>2</sub> decreases, the total amount of O<sub>3</sub> in the atmosphere and at the surface increases. The opposite is generally true in the **P19** simulations, although there is an increase in surface O<sub>3</sub> between 1% PAL and 0.1% PAL. This difference between the **P19** and **W21** scenarios occurs as a result of the weak UV radiation in the **W21** simulations and the pressure dependency on the reaction that produces O<sub>3</sub> (reaction 3). When the peak of O<sub>2</sub> photolysis occurs at higher altitudes and thus lower pressure, O does not react with O<sub>2</sub> as quickly to produce O<sub>3</sub> compared to the rate lower in the atmosphere where the density of the third body, M, is higher. In the **P19** cases, while the UV can penetrate deeper into the atmosphere when the concentration of O<sub>2</sub> is reduced in the simulations, the availability of O<sub>2</sub> becomes the limiting factor for the production of O<sub>3</sub>, instead of UV radiation. In the PCb cases, there is an intermediate amount of UV radiation compared to the **P19** and **W21** TP-1 e cases. These PCb cases follow the same trend as the TP-1 e **W21** scenarios, with the surface O<sub>3</sub> increasing as O<sub>2</sub> is reduced.

The destruction of O<sub>3</sub> plays an important role in these atmospheres too. Photolysis of O<sub>3</sub> is not counted as a loss of O<sub>3</sub> because the O produced quickly cycles back to produce O<sub>3</sub>. The peak in O<sub>2</sub> photolysis, O<sub>3</sub> production, and O<sub>3</sub> number density occurs in each simulation at pressures less than 100 hPa (above the troposphere). O<sub>3</sub> formation also takes place in the troposphere because O<sub>3</sub> photolysis there is fast as a result of O<sub>3</sub> being present in relatively high quantities. O<sub>3</sub> is being destroyed (by HO<sub>x</sub> and NO<sub>x</sub> catalytic cycles) and remade in



**Figure 3.** (a) Surface  $O_3$  mixing ratios are presented for the tidally locked **P19** (orange) and **W21** (blue) TP-1 e simulations, and the PCb simulations (magenta). All of these simulations have the substellar point placed over ocean. The circles show the mean mixing ratio, while the top bar shows the maximum and the bottom bar shows the minimum surface  $O_3$  mixing ratios. Also shown in black are the time-averaged  $O_3$  mixing ratios of data taken from Cooke et al. (2022). The horizontal axis indicates the simulations with a fixed  $O_2$  mixing ratio at the lower boundary, such that the horizontal ticks are categories rather than absolute values (the values are offset from each other for clarity). The gray dotted line indicates the 40 ppbv “harmful threshold,” above which  $O_3$  surface mixing ratios are considered dangerous to some forms of life on Earth. (b) The PI scenarios are compared in the TP-1 e simulations. These include the substellar point over ocean (PI), over land (PI SPL), and the nontidally locked cases (PI noTL).

the troposphere, but compared to above the troposphere, its production via  $O_2$ ,  $CO_2$ ,  $NO_2$ , or  $H_2O$  photolysis is significantly slower. In the simulations that have toxic quantities of  $O_3$ , tropospheric destruction of  $O_3$  is dominated by  $HO_x$  catalytic cycles rather than  $NO_x$  catalytic cycles. When  $O_3$  is below harmful levels,  $NO_x$  catalytic cycles (predominantly  $NO$  and  $NO_2$ ) dominate over  $HO_x$  catalytic cycle destruction of  $O_3$ . In other words, when  $NO$  and  $NO_2$  are significantly depleted in the troposphere,  $O_3$  is able to accumulate to harmful levels. The tropospheric column amount of  $NO_x$  exhibits negative correlation with the mean surface  $O_3$  mixing ratio, as shown in Figure 5. Therefore, the smog mechanism is not the reason for harmful levels of  $O_3$ , in contrast to modern-day Earth.

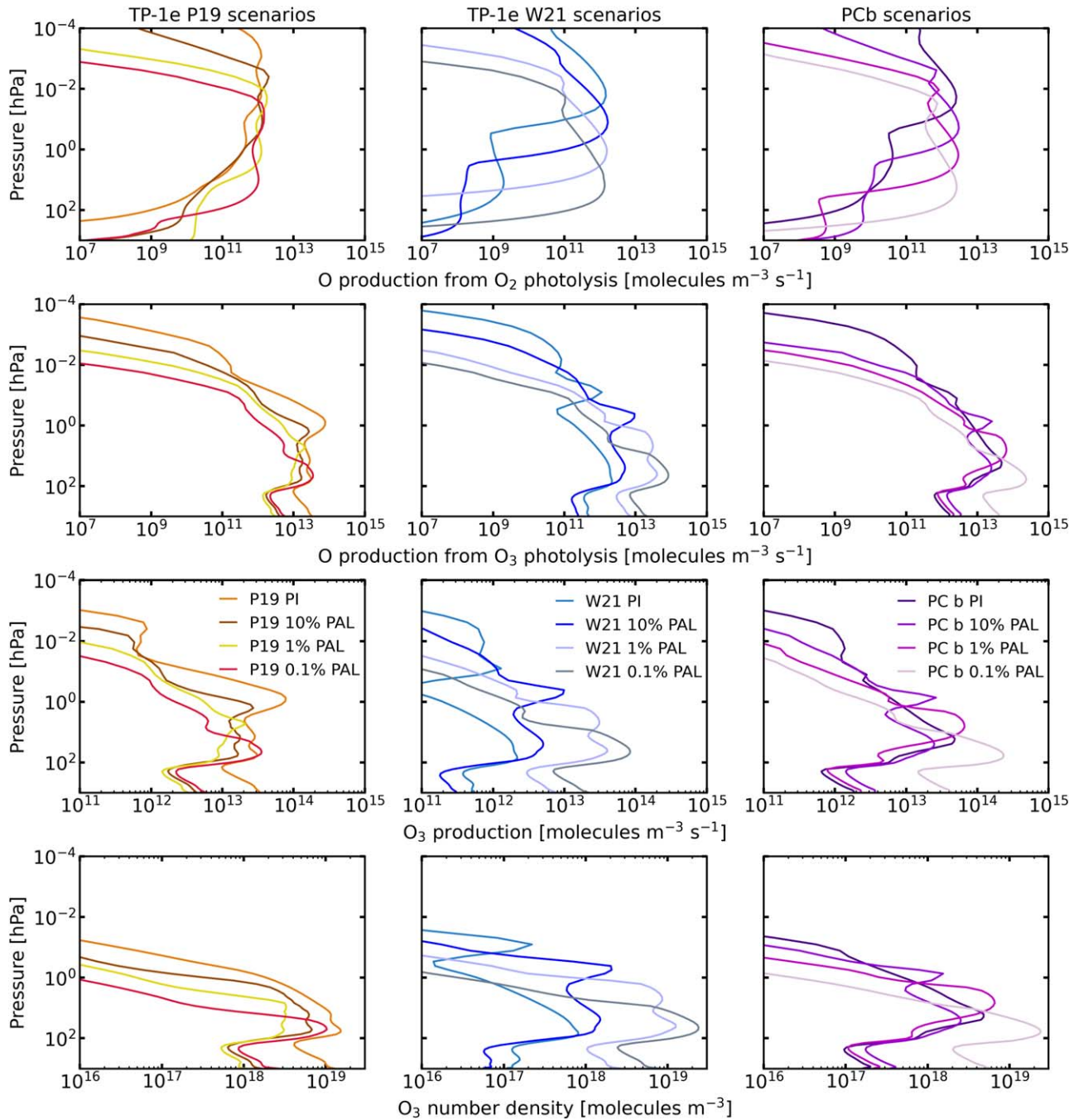
The atmospheric temperatures and dynamics influence the abundance and distribution of  $O_3$ . Figure 6 shows the surface temperatures (color map) and surface winds (arrows) in the TP-1 e models, with Figure 7 showing the same in the PCb models. Several TP-1 e tidally locked simulations have substellar points with temperatures above 273 K and surface winds converging toward this point (associated with the dayside upwelling at the substellar point), with some of the lowest surface mixing ratios also found near the substellar point. Away from the substellar point, temperatures drop below freezing and can be as low as 170 K on the nightside. The PCb cases have similar surface temperatures and winds because the total irradiance is quantitatively similar to the TP-1 e scenarios ( $0.65S_0$  vs.  $0.66S_0$ ). The cold temperatures result in reduced destruction from catalytic cycles (see reaction 11), which proceed slower at lower temperatures (e.g., from  $HO_x$  and  $NO_x$  families), allowing  $O_3$  to persist in relatively high quantities.

Figure 8 shows the dry deposition flux of  $O_3$  for some of the TRAPPIST-1 e simulations and the Earth PI simulation, and Table 2 shows the global mean dry deposition flux for the Earth

PI, TRAPPIST-1 e, and Proxima Centauri b simulations. Dry deposition over snow and ice is slow compared to that over other surfaces (Wesely & Hicks 2000; Helmig et al. 2007; Barten et al. 2023). Additionally, marine surface deposition is slower than land deposition when plant stomata are available to take up  $O_3$  (Ainsworth et al. 2012). If surface plants do not exist or have died from prolonged  $O_3$  exposure (Rich 1964; Bytnerowicz et al. 1993; Sandermann 1996; Rao & Davis 2001; Ramya et al. 2023), then dry deposition would occur more slowly, which results in even more  $O_3$  buildup. Even if there are no surface or gaseous molecules for  $O_3$  to interact with, the surface thermal decomposition of  $O_3$  can take place, and its importance may vary depending on the surface type, although this process is currently not sufficiently understood (Fowler et al. 2009; Clifton et al. 2020). The Earth PI simulation has a global mean dry deposition flux of  $2.1 \times 10^{-11} \text{ kg m}^{-2} \text{ s}^{-1}$ . All TRAPPIST-1 e and Proxima Centauri b simulations have reduced loss rates when compared to this, apart from the **P19** PI, **W21** 0.1% PAL, and PCb 0.1% PAL simulations. Despite these cases having relatively high rates of dry deposition, all retain harmful concentrations of surface  $O_3$ . For the Earth PI case, most  $O_3$  is deposited over land (see Figure 8). For the exoplanet simulations, the majority of  $O_3$  is deposited near the substellar point, regardless of whether it is placed over land or ocean. The dry deposition flux around the substellar point contributes to the relatively reduced  $O_3$  concentrations at the substellar point (see, e.g., the **P19** 10% and 0.1% PAL simulations and **W21** 1% PAL and 0.1% PAL simulations in Figure 2). The noTL cases have globally averaged dry deposition rates that are slower than their PI tidally locked counterparts by a factor of 220 and 20 for the **W21** and **P19** scenarios, respectively.

To summarize the  $O_3$  distribution in the simulations,  $O_3$  is made primarily in the middle atmosphere on the modeled



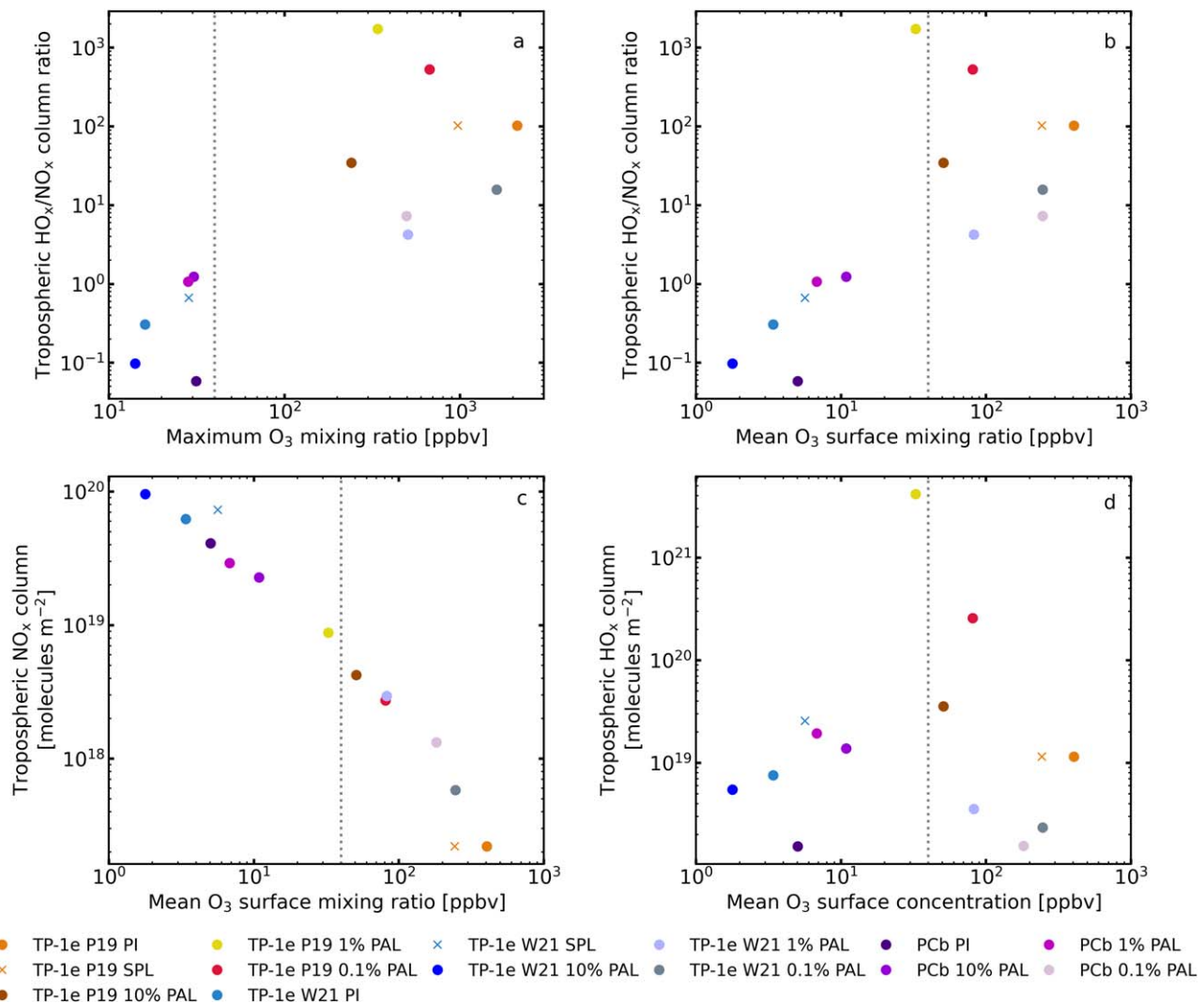


**Figure 4.** The production of O from O<sub>2</sub> photolysis (top row) and O<sub>3</sub> photolysis (second row) is shown against pressure, with O<sub>3</sub> production (third row) and O<sub>3</sub> number density also shown (bottom row). All profiles are time-averaged global means. The TP-1 e P19 (orange, brown, yellow, and red, respectively) and TP-1 e W21 (light blue, blue, lilac, and gray, respectively) stellar spectra in the left and middle columns, respectively. The right column show the PCb PI (indigo), 10% PAL (violet), 1% PAL (magenta), and 0.1% PAL (light pink) simulations.

dayside and the O<sub>3</sub> column maximizes at the poles (Cooke et al. 2023a). O<sub>3</sub> loss on the nightside and at the poles is slower than on the dayside, due to relatively low temperatures and a lack of photolysis producing the molecules that become involved in O<sub>3</sub> destroying catalytic cycles. O<sub>3</sub> is lost at the surface owing to dry deposition, but the flux is not large enough to mitigate for the dangerous concentrations of O<sub>3</sub>. As an example, in the P19 PI case, the global mean dry deposition flux increases by  $\approx 4$  times greater than the Earth PI case, with a corresponding increase in surface O<sub>3</sub> concentrations of  $\approx 34$ . In the tidally locked cases, the O<sub>3</sub> chemical loss rate at the surface

is approximately 2–29 times less than the peak stratospheric loss rate. Surface winds (of order  $10 \text{ m s}^{-1}$ ), which are much stronger than vertical winds (of order  $0.1 \text{ m s}^{-1}$ ), transport O<sub>3</sub> across the surface.

We hypothesize that O<sub>3</sub> is transported from where it is produced in the dayside stratosphere to the nightside and toward the poles as discussed in Braam et al. (2023), who simulated Proxima Centauri b assuming an initial condition of a modern-Earth atmosphere. The difference between our work and Braam et al. (2023) is that, in several of our simulations, O<sub>3</sub> accumulates to harmful and lethal quantities. With loss



**Figure 5.** The tropospheric HO<sub>x</sub>/NO<sub>x</sub> column ratio is shown against the maximum and mean surface O<sub>3</sub> mixing ratio in each tidally locked simulation in panels (a) and (b), respectively. The tropospheric NO<sub>x</sub> and HO<sub>x</sub> columns are shown against the mean O<sub>3</sub> surface mixing ratio in panels (c) and (d), respectively. The gray vertical dotted lines show the 40 ppbv cutoff for harmful O<sub>3</sub> levels. The tropospheric column is calculated as the column abundance of molecules between 120 hPa and the surface. Its units are molecules m<sup>-2</sup>. The circles indicate simulations where the substellar point is placed over ocean, and the crosses indicate simulations where it is placed over land.

processes in the troposphere less effective than in the stratosphere, the lifetime of surface O<sub>3</sub> increases and enables a buildup of toxic O<sub>3</sub> levels. The atmospheric transport throughout the entire atmosphere will be explored in a follow-up paper to confirm whether such a scenario is occurring in the simulations.

## 4. Discussion

### 4.1. Prior Results for Toxic Levels of Ozone

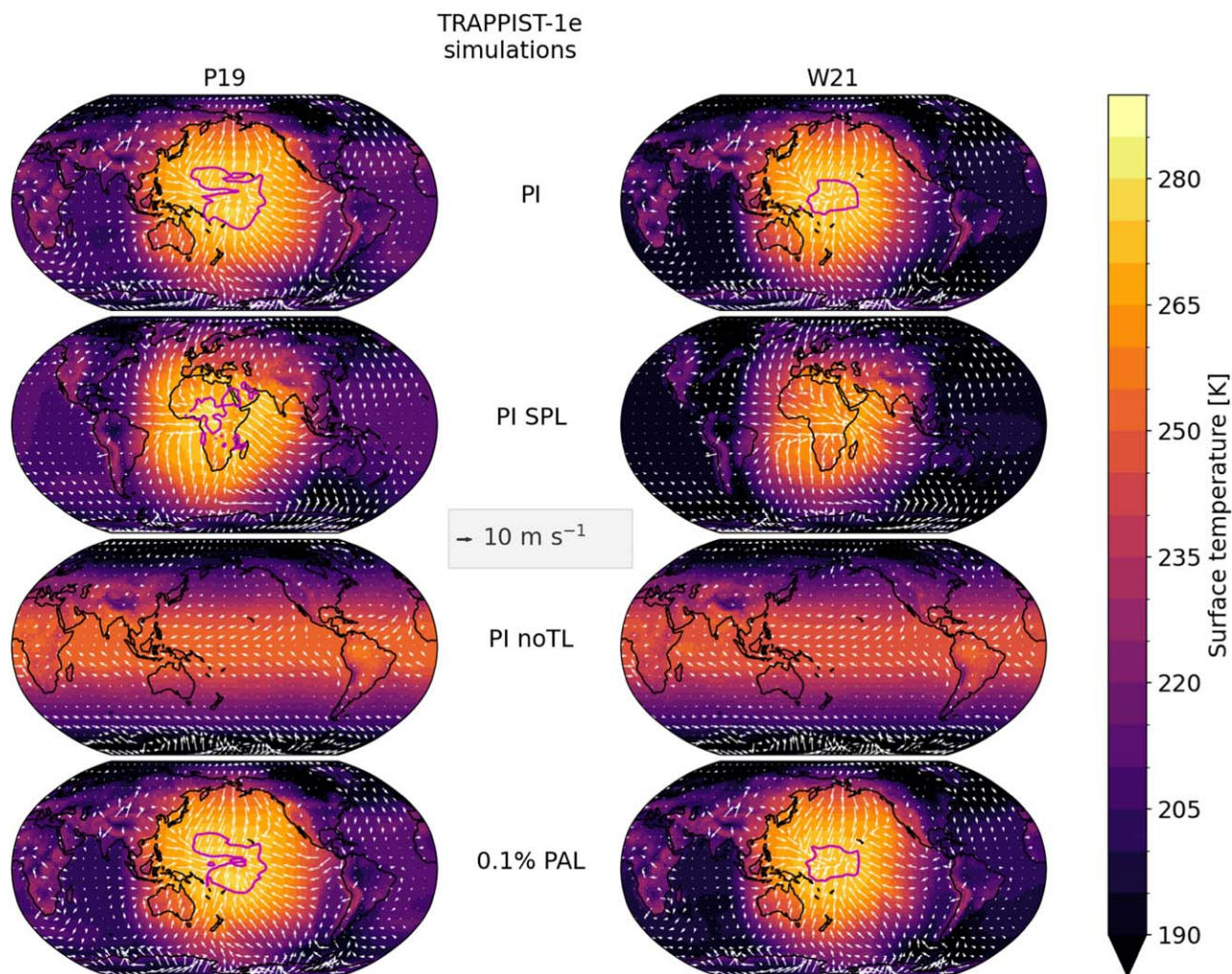
As far as can be discerned from the presented data (e.g., globally averaged vertical profiles), other 3D simulations of oxygenated exoplanets (Proedrou & Hocke 2016; Way et al. 2017; Chen et al. 2019; Yates et al. 2020), excluding Braam et al. (2022), have not produced surface O<sub>3</sub> mixing ratios above 40 ppbv. This is likely due to the investigated scenarios that differ between each work, although model differences will be important (Ji et al. 2023, 2024). Figure 5 in Braam et al. (2022) shows dayside mixing ratios of O<sub>3</sub> reaching  $\approx 45$  ppbv at the surface. It is worth noting here that all of these studies,

including our simulations, assume a surface pressure of 1 bar (1000 hPa), so that a mixing ratio of 40 ppbv at 1000 hPa surface pressure corresponds to a number density of  $1.0 \times 10^{18}$  molecules m<sup>-3</sup> at 288 K. Chen et al. (2021) used WACCM4 and Ridgway et al. (2023) used the UM (two 3D chemistry-climate models) to investigate the impact of flares and coronal mass ejections on terrestrial exoplanets. The flares significantly perturbed O<sub>3</sub> concentrations,<sup>7</sup> but the changes were in the middle and upper atmosphere and the flares did not cause surface concentrations to exceed 40 ppbv. However, there could be specific cases where flares act to increase O<sub>3</sub> surface concentrations to harmful levels, likely depending on the atmospheric properties, incoming flare strength, and flare frequency.

1D photochemical models have also simulated the impact of flares, coronal mass ejections, and cosmic rays on atmospheric chemistry. The calculations in Segura et al. (2010), Grenfell

<sup>7</sup> The amount of atmospheric O<sub>3</sub> decreased by a factor of 3 in Chen et al. (2021) for active M dwarf stars, whereas total atmospheric O<sub>3</sub> increased by up to a factor of 20 in Ridgway et al. (2023).



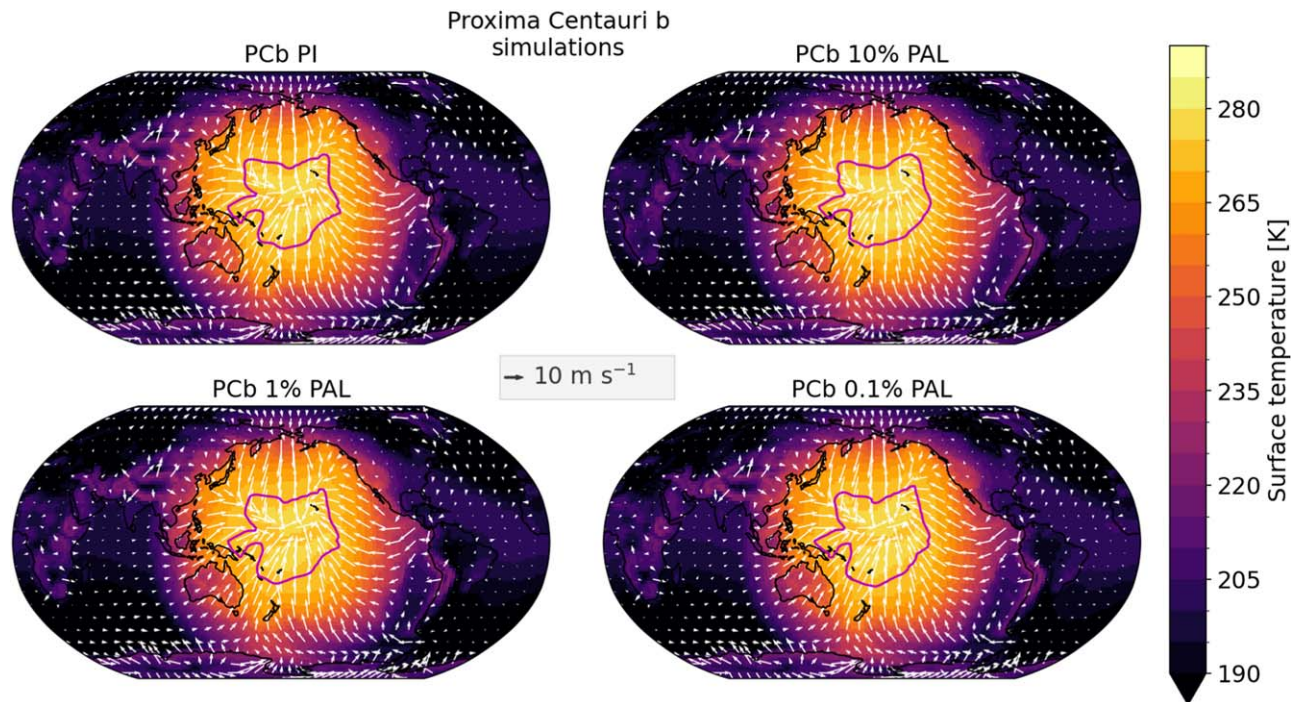


**Figure 6.** The surface temperature is displayed for the 12 TP-1 e simulations used in this work. The left column shows the **P19** simulations, and the right column shows the **W21** simulations. From top to bottom: PI, PI SPL, PI noTL, and the 0.1% PAL simulations are displayed. White arrows indicate the magnitude and direction of the surface winds. The magenta contours show surface temperatures of 273 K. The 10% PAL and 1% PAL simulations are not shown for brevity, but their surface temperatures are very similar to the PI and 0.1% PAL cases. For scale, a  $10 \text{ m s}^{-1}$  arrow is shown in the middle of the figure.

et al. (2012), Tabataba-Vakili et al. (2016), and Tilley et al. (2019) did not demonstrate surface  $\text{O}_3$  mixing ratios reaching biologically toxic quantities. The same is true for abiotic  $\text{O}_2$  production simulations (Segura et al. 2007; Harman et al. 2015, 2018), although it is unclear in the 100-bar atmosphere from Schwietzman et al. (2016) because surface  $\text{O}_3$  mixing ratios are not shown (their Figure 1), but  $\text{O}_3$  is at 1 ppbv by 10 km and its mixing ratio is strongly decreasing with decreasing height. Other 1D photochemical modeling results for HZ exoplanets have shown that  $\text{O}_3$  mixing ratios may exceed harmful levels (Kozakis et al. 2018; Fauchez et al. 2019; Kaltenecker et al. 2020; Kopparapu et al. 2021). For example, Kozakis et al. (2018) used a 1D photochemical model (EXO-Prime; Kaltenecker & Sasselov 2010) to simulate Earth-like exoplanets with various surface pressures of 0.3, 1.0, 1.5, and 2.0 bars, orbiting white dwarf stars (stellar effective temperatures of 4000, 5000, and 6000 K). Almost all simulations had surface  $\text{O}_3$  below the 40 ppbv threshold, with only a single simulation (0.3-bar surface pressure and 4000 K stellar effective temperature) exceeding it. However, if it is the concentration (number of molecules per unit volume), rather than the mixing ratio (fractional concentration), that is important, one must take into account the surface density of

the atmosphere. Using this criterion, a few more simulations from Kozakis et al. (2018) would be close to  $1.0 \times 10^{18} \text{ O}_3$  molecules  $\text{m}^{-3}$  and therefore potentially dangerous to any surface life present for simulated atmospheres around white dwarf stars. However, the 1-bar atmospheres in Kozakis & Kaltenecker (2020) that were simulated around red giant stars did not surpass dangerous surface  $\text{O}_3$  concentrations. Regarding smog (see Section 1 and reactions 7–9), Kopparapu et al. (2021) used the 1D photochemical model that is part of “Atmos” (see Arney et al. 2016; Arney 2019) to simulate varying surface-to-atmosphere fluxes of  $\text{NO}_2$  to test whether it could be used as a signature that extraterrestrial technology existed on an exoplanet. In two simulations (using a Sun-like star and a K6V star with  $20 \times$  the present Earth flux of  $\text{NO}_2$ ), the surface  $\text{O}_3$  mixing ratios were  $\approx 90$  and  $\approx 100$  ppbv, respectively. Alongside modern-day pollution on Earth due to  $\text{NO}_x$  emissions from vehicles, this study by Kopparapu et al. (2021) demonstrates that the surface fluxes of molecules will be important for determining whether ground-level  $\text{O}_3$  concentrations reach concerning levels for life. Fauchez et al. (2019) simulated the HZ exoplanets TRAPPIST-1 e, f, and g using the 3D model LMD-G and then performed terminator photochemistry simulations with Atmos. The surface  $\text{O}_3$  concentrations





**Figure 7.** The surface temperature is displayed for the four PCb simulations used in this work: PCb PI (top left), PCb 10% PAL (top right), PCb 1% PAL (bottom left), and PCb 0.1% PAL (bottom right). White arrows indicate the magnitude and direction of the surface winds. The magenta contours show surface temperatures of 273 K. For scale, a  $10 \text{ m s}^{-1}$  arrow is shown in the middle of the figure.

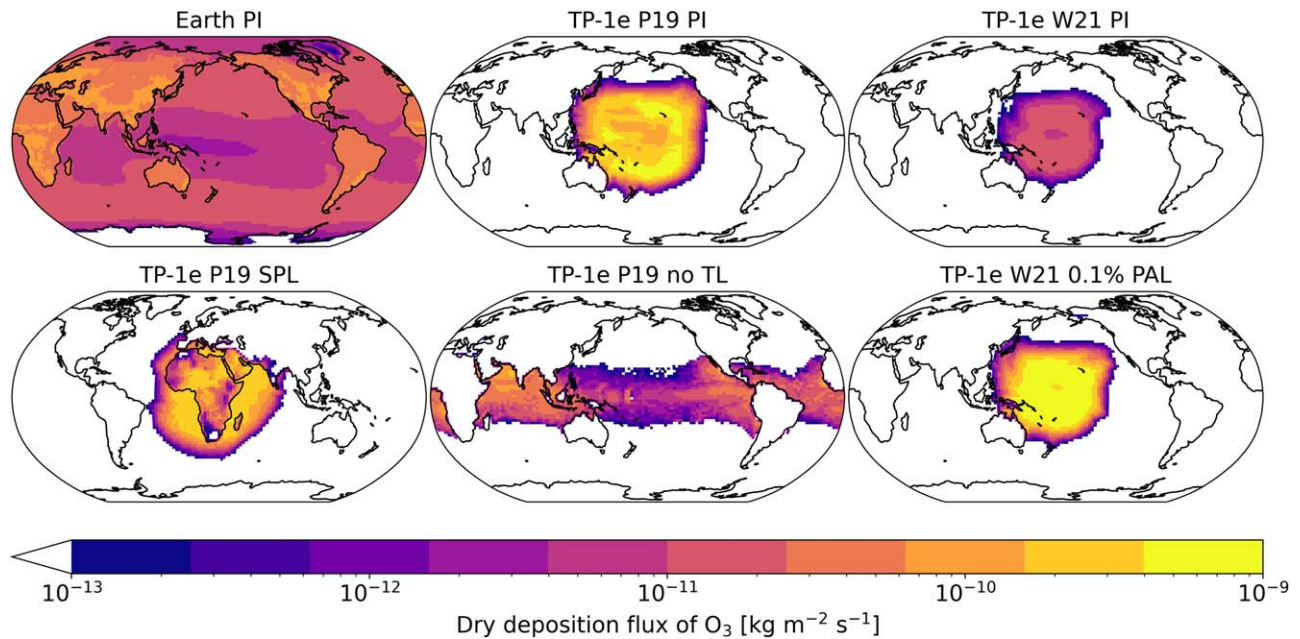
are  $<40$  ppbv for TRAPPIST-1 e but are  $\sim 60$  and  $\sim 120$  ppbv for TRAPPIST-1 f and g, respectively. In their simulations, surface  $\text{O}_3$  increases with decreasing illumination. The simulations presented here predict, for the P19 PI case, the highest ozone surface concentrations ( $1.4 \times 10^{20}$  molecules  $\text{m}^{-3}$ ) compared with other results for exoplanets simulated in the literature.

WACCM6 predicts lower concentrations of  $\text{O}_3$  for various oxygenation states (0.1%–150% PAL) for Earth when compared to 1D models (Cooke et al. 2022; Ji et al. 2023), so there is the question of why it predicts higher concentrations of  $\text{O}_3$  when compared to 1D models for M dwarf stars. There will likely be many reasons that the results differ (see, e.g., Ji et al. 2023), but we suspect that the main causes are differences in atmospheric transport and temperatures. The atmospheric temperatures in Kozakis et al. (2022) with the M5V star are warmer below 30 km when compared to the WACCM6 simulations. The same is true in the M3V and M8V simulations of Rugheimer & Kaltenecker (2018), the GJ 436 and GJ 876 simulations of Rugheimer et al. (2015), the Gebauer et al. (2018) simulations around AD Leo (M3.5V), the TRAPPIST-1 e simulations of Lin et al. (2021), the Proxima Centauri b simulations of Scheucher et al. (2020), and the 1-bar habitable Proxima Centauri b simulations of Meadows et al. (2018). Note that the desiccated Proxima Centauri b atmospheres in Meadows et al. (2018) do have harmful levels of  $\text{O}_3$ , but these would not be considered habitable owing to the lack of  $\text{H}_2\text{O}$ .

As an example of the temperature differences, global mean surface temperatures are approximately 220–230 K in the WACCM6 tidally locked simulations, compared to a surface temperature of  $\approx 310$  K in Kozakis et al. (2022), 260–280 K in Scheucher et al. (2020), and 273 K in Meadows et al. (2018). This means that below 30 km, where most of the atmospheric  $\text{O}_3$  resides, loss rates will be slower and formation will be faster

when compared to the 1D models. Furthermore, in 1D models there is constant illumination, but in 3D models  $\text{O}_3$  is transported to the nightside (where there is no direct starlight) and down to the surface (Braam et al. 2023). The TRAPPIST-1 e 1D photochemical simulations in Pidhorodetska et al. (2020) used temperature and pressure profiles from 3D simulations of TRAPPIST-1 e (Fauchez et al. 2019). Here the deviations from our results may be due to the synthetic BT-Settl spectrum (a model of stellar atmospheres; Rajpurohit et al. 2013) used by Pidhorodetska et al. (2020) and the fact that photochemistry was performed at the terminator. Between all of these simulations, there will be differences in photochemical cross sections and chemical schemes, as well as the UV spectra used and the total instellation, but temperatures and transport may largely explain the substantial discrepancies in both predicted surface  $\text{O}_3$  and  $\text{O}_3$  columns, when compared to similar 1D model simulations.

A detailed model intercomparison will be needed to determine why WACCM6 predicts higher surface concentrations of  $\text{O}_3$  compared to the UM. For now, we speculate that at least one difference is the use of a slab ocean with no ice formation in the UM, such that  $\text{O}_3$  will be lost in greater numbers on the nightside than would occur in reality. When simulating Proxima Centauri b with a slab ocean, Yates et al. (2020) and Braam et al. (2022) assumed a dry deposition velocity of  $0.05 \text{ cm s}^{-1}$  based on previous work (Ganzeveld & Lelieveld 1995; Giannakopoulos et al. 1999). The global mean dry deposition velocities for  $\text{O}_3$  in our tidally locked ocean substellar point cases are approximately 4 times slower, and slower still when considering the “noTL” and “SPL” cases.  $\text{O}_3$  dry deposition depends on multiple interlinked parameters that are poorly known (El-Madany et al. 2017), so whether existing dry deposition parameterizations can be used for exoplanet simulations is not well known.



**Figure 8.** The dry deposition flux of  $O_3$ , given in terms of  $kg\ m^{-2}\ s^{-1}$ , is plotted for six simulations: the Earth PI simulation; the TRAPPIST-1 e **P19** PI, SPL, and noTL simulations; and the **W21** PI and 0.1% PAL simulations. Yellow indicates relatively large amounts of dry deposition, while red indicates relatively low amounts. The white areas indicate regions where the dry deposition flux of  $O_3$  is below  $10^{-13}\ kg\ m^{-2}\ s^{-1}$ .

#### 4.2. Modeling Limitations

The model used for the atmospheric simulations is an important factor in these predictions because varying parameterizations and chemical schemes will impact the results. WACCM6 is a model that is tuned to Earth’s atmosphere, land, ice, and ocean. WACCM6 accounts for scattering longward of 200 nm, but it does not account for scattering in the Schumann–Runge bands (175–192 nm; these wavelengths photolyze  $O_2$  above  $\sim 80$  km in Earth’s atmosphere), becoming pertinent for Earth-like simulations at  $O_2$  mixing ratios of 1% PAL or less (Ji et al. 2023). In our simulations, WACCM6 accounts for absorption in the Schumann–Runge Bands from  $O_3$ ,  $O_2$ ,  $CO_2$ , and  $H_2O$ . The integrated flux in the Schumann–Runge Bands is 1.15, 656, and 26.6 times lower than Earth for the TP-1 e **P19**, TP-1 e **W21**, and PCb cases, respectively. Even with the effects of scattering included, it seems still possible that harmful  $O_3$  surface concentrations could form on the surface of terrestrial exoplanets because all of the **P19** simulations have dangerous concentrations of surface  $O_3$ . Moreover, the **W21** simulations have a relatively low amount of radiation in the Schumann–Runge bands, yet toxic  $O_3$  concentrations are found in the **W21** 1%, and 0.1% PAL cases. We tested these assumptions in the 0.1% PAL cases by reducing the incoming photon flux in the **W21** and **P19** spectra by a factor of  $10^5$ . The total integrated  $O_2$  photolysis in the **P19** 0.1% PAL case reduced by a factor of 1.6, and in the **W21** case the decrease was negligible. The decrease in  $O_2$  photolysis in the **P19** 0.1% PAL case primarily takes place between 10 and 0.01 hPa. As seen from the results presented in Figure 7 in Ji et al. (2023), scattering becomes important in the 0.1%–10% PAL simulations at altitudes below 10 hPa. Therefore, we expect our conclusions to be unaffected by including scattering in the Schumann–Runge band parameterization in WACCM6; however, the test should still be done when that parameterization is eventually updated.

#### 4.3. Atmospheric Evolution

We have assumed planetary conditions based on the last 2.4 billion years of Earth’s history, where the atmospheric pressure has been  $\sim 1$  bar and atmospheric  $O_2$  has varied from  $1000\times$  less than the PAL to  $\sim 1.5\times$  greater (Large et al. 2019; Catling & Zahnle 2020; Steadman et al. 2020; Lyons et al. 2021). However, for an HZ tidally locked exoplanet around an M dwarf star, is such a scenario realistic?

First, atmospheres with high  $O_2$ ,  $CO_2$ , or  $H_2O$  mixing ratios and sufficient UV may develop an  $O_3$  layer close to the surface because photodissociation could produce  $O_2$  and O, which then lead to  $O_3$  formation. In previous work simulating a 1-bar atmosphere around the host star GJ 876,  $O_2$  amounts comparable to the Proterozoic Earth were produced in abiotic scenarios (Domagal-Goldman et al. 2014; Tian et al. 2014). When  $CO_2$  is photolyzed, it produces CO and O, but lightning flashes can produce NO, which can catalyze the recombination of CO and O (instead of O going on to produce  $O_2$ ), meaning that abiotic  $O_2$  may not increase to detectable levels (Harman et al. 2018). However, on some exoplanets lightning may not be sufficient to prevent  $O_2$  buildup (Barth et al. 2024).

Oxygenic photosynthesis on Earth, utilizing chlorophyll as a pigment, requires photosynthetically active radiation (PAR) between 400 and 700 nm (Alados et al. 1996). Due to the luminosity of M dwarf stars peaking at longer wavelengths when compared to the Sun, the availability of PAR may be lacking on habitable M dwarf exoplanets. Nevertheless, oxygenic photosynthesis is thought to be possible on exoplanets orbiting M dwarf stars (Gale & Wandel 2017; Claudi et al. 2020; Duffy et al. 2023), although some exoplanets may be limited by a lack of photons reaching the surface (Lehmer et al. 2018), and this could mean that anoxygenic photosynthesis (where oxygen is not a product of photosynthesis) is preferentially selected (Duffy et al. 2023). In particular, the results from Lingam & Loeb (2019) showed how



planets orbiting stars with masses below  $0.13 M_{\odot}$  may have difficulties in producing  $O_2$  through oxygenic photosynthesis, with Proxima Centauri and TRAPPIST-1 falling under this limit. However, the results of Covone et al. (2021) contrasted with those of Lingam & Loeb (2019), and while TRAPPIST-1 f and g might be energy limited (Lehmer et al. 2018; Covone et al. 2021), TRAPPIST-1 e may be able to host an Earth-like biosphere. Furthermore, cyanobacteria can use radiation longward of 700 nm to photosynthesize (Gan & Bryant 2015; Claudi et al. 2020), and a recent study has also shown this in more complex vegetation (Zhen et al. 2022). A new metric quantifying the photoabsorption rate of photosynthetic pigments around different stars suggests that M8V stars could activate photosynthesis (Marcos-Arenal et al. 2022). Moreover, cyanobacteria can survive and photosynthesize in caves and under low-light conditions (Hanelt et al. 1997; Behrendt et al. 2020; Jung et al. 2023). Regardless of the possibility of photosynthesis on M dwarfs,  $O_2$  production needs to surpass  $O_2$  sinks in order to produce an oxygenated atmosphere (Lehmer et al. 2018).

It is hypothesized that the oxidation of Earth's lithosphere (the solid outer layer of a planet) occurred as a result of two primary factors: oxygenic photosynthesis from cyanobacteria over millions of years, and hydrogen escape to space resulting from water photolysis (Catling et al. 2001; Zahnle et al. 2013). Modeling an Earth-like planet orbiting the star AD Leo (M3.5V), Gebauer et al. (2018) found that it may be possible for a Great Oxidation Event (see, e.g., Lyons et al. 2014; Gumsley et al. 2017; Poulton et al. 2021) to occur earlier in the history of the modeled planet when compared to Earth. If  $O_2$  increases and results in large amounts of  $O_3$  at the surface, because  $O_3$  is a strong oxidizer, it could speed up the oxidation of Earth's lithosphere, reducing the sinks of  $O_2$  and further enabling the generation of an oxygenated atmosphere.

#### 4.4. Habitability Feedback

Continuing with this scenario in mind, there is the potential for biological feedback. Surface  $O_3$  may build up to toxic concentrations and kill organisms that produce  $O_2$  on the surface, subsequently reducing the production of  $O_2$  until  $O_3$  reaches safe levels again. Even so, we argue here that if  $O_3$  exceeded lethal surface concentrations, a photosynthetic biosphere could still be safely present in the ocean and under ice.

First, refuge from the dangerous  $O_3$  concentrations may be found in any liquid water ocean present, because  $O_3$  has low solubility in water (Egorova et al. 2015). Although it is possible for  $O_3$  to form in water that is irradiated by the Sun (Lushchak 2011),  $O_3$  then quickly decomposes. Additionally, for disinfection applications,  $O_3$  has to be artificially inserted into water (e.g., via bubble diffusion; Wert et al. 2017), such that it seems unlikely that marine life would be adversely affected by high surface concentrations of  $O_3$  in the air. The known effects of  $O_3$  exposure on life are limited to observations of terrestrial organisms, and extraterrestrial organisms may adapt to survive in an atmosphere with high surface  $O_3$  mixing ratios.

On Earth, phytoplankton blooms have been found under Arctic (Suzuki et al. 1997; Arrigo et al. 2012; Clement Kinney et al. 2020) and Antarctic sea ice (McMinn et al. 2007), algae are found at varying depths (Norris & Olsen 1991; Pritchard et al. 2013; Borlongan et al. 2017), and lichens (Kappen 1993), aquatic plants (Adams et al. 1974; Campbell et al. 2007), and

seaweed (Ramus et al. 1976; Huovinen & Gómez 2013) can photosynthesize below the surface. The long day length during the summers at each pole on Earth can increase the rate of primary productivity (Henshaw & Laybourn-Parry 2002). Given phytoplankton's contribution to roughly half of Earth's primary productivity (Field et al. 1998) and the potential for subsurface photosynthesis on tidally locked exoplanets with constant dayside illumination, significant oxygen accumulation may occur without surface life. The amount of ice coverage and available area of open ocean may limit the  $O_2$  flux to the atmosphere, although gaseous diffusion through sea ice is possible (Delille et al. 2007; Loose et al. 2011, 2011; Bortkovskii 2012), just slower when compared to open water (Bortkovskii 2012).

In summary, whether due to biological or abiotic production,  $O_2$  could be produced in quantities large enough to induce toxic concentrations of  $O_3$  at the surface. If life is present on such a planet, there is then the potential for biogeochemical feedback as organisms are hindered or destroyed by  $O_3$ . We encourage simulations of these scenarios using biogeochemical models to quantitatively determine the possible outcomes we have discussed.

#### 4.5. Future Work

Future work should aim to determine the parameter space (in UV irradiation, composition, and atmospheric pressure) for which detrimental levels of surface level  $O_3$  may occur. Ideally, 1D photochemical models would be used for this, as they are less computationally expensive than 3D models, and the use of 1D photochemical models is indeed a viable investigation for nontidally locked exoplanets. However, for tidally locked exoplanets, where  $O_3$  is transported to the nightside and shielded from destruction (Proedrou & Hocke 2016; Yates et al. 2020; Braam et al. 2023), 3D chemistry-climate models will be needed to predict where surface  $O_3$  concentrations maximize owing to atmospheric dynamics and climatology from varied ocean and land configurations (Zhao et al. 2021; Macdonald et al. 2022). The ocean salinity could be varied to determine the effect on the sea ice distribution (Olson et al. 2022), which is important for habitability estimates and  $O_3$  dry deposition calculations. Simulations where only the ocean model is changed between dynamic ocean, slab ocean with ice formation (e.g., Olson et al. 2022), and slab ocean without ice formation (e.g., Braam et al. 2022) would be useful to determine how important these factors are when considering toxic  $O_3$  at the surface. Additionally, one could explore various topographies, gravitational accelerations, and several different atmospheric compositions (e.g.,  $H_2O$  or  $CO_2$  dominated). As the exoplanets simulated here have limited areas where surface temperatures are above 273 K (and some have no surface temperatures above 273 K), subsequent research could investigate the parameter space with warmer exoplanets to determine how surface  $O_3$  concentrations are affected. Examples of other 3D models that could investigate this chemical phenomenon are LMD-G (Yassin Jaziri et al. 2022), ROCKE-3D (Way et al. 2017), and the UM (Boutle et al. 2017). Once a chemical scheme is implemented, the LFRic-Atmosphere model could also be used (Sergeev et al. 2023).

The classical HZ depends on whether a planet can sustain surface liquid water, but the notion of an HZ for complex life (Schwieterman et al. 2019) is influenced by the presence of toxic gases. So far, the toxic gases that have been suggested are



carbon monoxide (CO) and relatively high concentrations of CO<sub>2</sub> (Schwieterman et al. 2019), as well as N<sub>2</sub> at high pressure (e.g., >2 bar; Ramirez 2020). These molecules impede biochemical processes that may be unique to organisms on Earth. Resulting from its reactivity as a strong oxidant, O<sub>3</sub> could pose a more significant threat to extraterrestrial life. Therefore, based on our modeling work and the properties of O<sub>3</sub>, we recommend that O<sub>3</sub> should now be added to the list of molecules that can influence the HZ for complex life.

## 5. Conclusions

This work used WACCM6 to simulate the climate of two exoplanets: TRAPPIST-1 e and Proxima Centauri b. For each exoplanet, we considered O<sub>2</sub> mixing ratios between 0.1% PAL and 100% PAL. Additionally, two different stellar spectra were used for the TRAPPIST-1 e cases to investigate the effect on surface O<sub>3</sub> due to their large differences in the strength of incoming UV radiation. In multiple simulations, surface concentrations of O<sub>3</sub> exceed 40 ppbv, with maximum time-averaged concentrations reaching up to 2120 ppbv in the TP-1 e P19 PI case. Such concentrations are harmful to life on Earth and may be potentially fatal through oxidative stress. In these simulated atmospheres, O<sub>3</sub> exists not as a pollutant but as a consequence of the planetary atmospheric conditions, such as the 1000 hPa surface pressure, the incoming UV strength and shape, and the O<sub>2</sub> number density vertical profile. Our work suggests that the potential presence of toxic O<sub>3</sub> concentrations should be included when evaluating the habitability of an exoplanet.

The simulations examined in this exploratory work represent a small proportion of the parameter space in which atmospheres may form relatively high O<sub>3</sub> concentrations at the surface. Different planetary rotation rates, topography, atmospheric pressures, total irradiation and UV irradiation environments, and various chemical fluxes from the surface to the atmosphere should all be explored. Upcoming work should consider the potential presence of high surface concentrations of O<sub>3</sub> when simulating oxygenated atmospheres. If O<sub>3</sub> is detected in any future observations of terrestrial exoplanet atmospheres, ascertaining the O<sub>3</sub> surface concentration should be incorporated into frameworks that aim to determine planetary habitability and decide on the most promising targets for follow-up observations (see, e.g., Truitt et al. 2020; Méndez et al. 2021; Safonova et al. 2021). In practice, this will require a combination of planetary modeling, transmission, and direct imaging spectra, as well as precise knowledge of the UV irradiation environment of the atmosphere. 3D chemistry-climate models are essential for understanding how transport can create areas with comparatively lower and thus safer O<sub>3</sub> concentrations. Just as on Earth, the entire surface does not need to be hospitable for life to flourish.





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