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# Inductorless Step-up Piezoelectric Resonator (SUPR) Converter: a Describing Function Analysis

Jack Forrester, Jonathan N. Davidson and Martin P. Foster

**Abstract**—A model for an inductorless step-up piezoelectric-resonator-based DC-DC (SUPR) converter is presented. A describing function approach is used to determine the equivalent circuit of the MOSFET power stage and resonator, allowing the duty cycles of both MOSFETs to be accurately estimated. The model provides a method for estimating the output voltage and resonant current of the converter during operation. The accuracy of the model is verified against experimental and simulation results, with a less than 1.5% mean absolute percentage error in voltage gain estimation, compared to experimental measurements. Design guidelines for optimal operation of the converter are also presented.

**Index Terms**—DC-DC converters, Design Optimization, Piezoelectric Devices, Resonant Converters, Zero Voltage Switching (ZVS).

## I. INTRODUCTION

Piezoelectric-based resonant converters are an alternative to traditional resonant converters, offering several advantages including high power density, high efficiency, low EMI content and the ability to be used in harsh environments [1]–[4]. Piezoelectric devices make use of the direct and converse piezoelectric effects, allowing the transformation of energy from mechanical to electrical and vice versa, respectively. The use of piezoelectric devices in power converters has seen an increase in interest in the past few decades, with improvements to piezoelectric materials leading to performance comparable with traditional magnetic transformer based converters [5].

The most popular piezoelectric device used in resonant power conversion is the piezoelectric transformer (PT). A PT has an input and an output section, with electrodes attached to these sections, which are each formed from one or more layers of piezoelectric material, depending on the PT topology. PTs exhibit high Q factors, have inherent resonant circuits and transformers, making them ideal for use in resonant converters. Several applications of piezoelectric transformers (PTs) have been presented such as cold-cathode fluorescent lamp drivers, battery chargers and circuitry for liquid crystal displays [1], [6]–[8]. PTs suffer from the need for multiple layers, which significantly complicates their manufacture.

Typically, magnetic resonant converters use frequency and duty cycle to control the output voltage (PFM, PWM control)

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[9]. However, voltage regulation is challenging for high Q converters because the ZVS ability is highly sensitive to changes in duty cycle and frequency. In [8], to ensure ZVS with a piezoelectric transformer-based converter, a fixed phase difference between resonant current and switching signals, and a switching arrangement involving 25% on time and 25% off time with two 25% duration deadtime intervals between is used. Control methods [10]–[12] have been developed for PT based converters but are complex to implement and can lead to performance degradation.

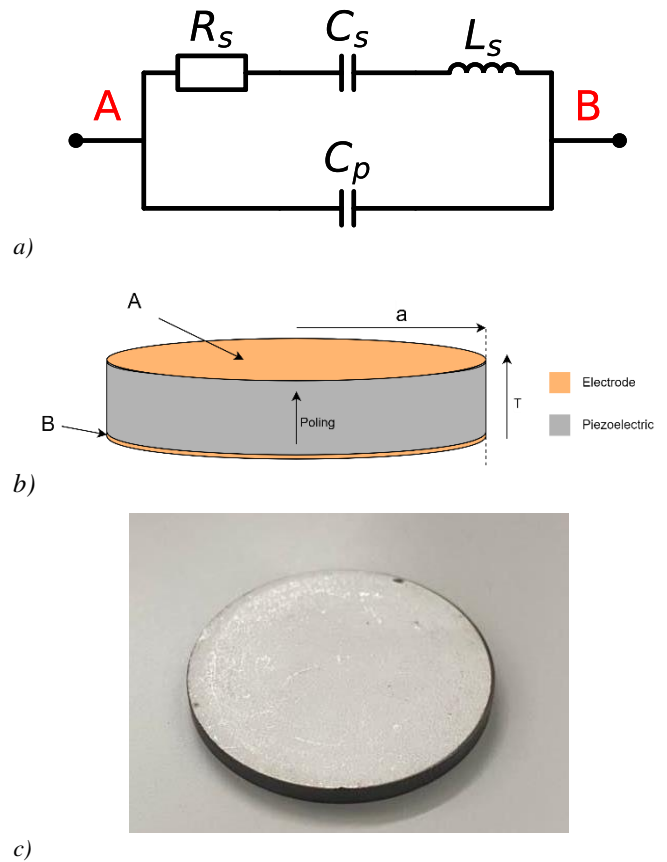


Fig. 1 – a) Piezoelectric resonator equivalent circuit, b) Radial mode PR schematic, c) SMD30T21F1000S Piezo resonator

Piezoelectric resonator-based converters provide an alternative way of utilising piezoelectric devices in resonant converters, whilst avoiding some of the complexities involved in PT based converters. Piezoelectric resonators (PRs) can be formed from a single layer of ceramic and, as such, they are equivalent to the input section of a PT. PRs also have high quality (Q) factors, often higher than PTs owing to their simpler construction. Whilst exhibiting some advantages over PTs, Piezoelectric resonator (PR) devices only possess two

terminals, meaning they do not have an integrated transformer for voltage transformation and as a result they do not provide galvanic isolation. Additionally, similar to PTs, they suffer from additional spurious resonant modes, which can interfere with the operation of a PR based converter [13].

In most PR based converters, the PR acts as a resonating energy store much like the inductor in a boost converter: it accumulates and releases energy according to the instantaneous circuit conditions [14]. The equivalent circuit of the piezoelectric resonator is shown in Fig. 1a. Nodes  $A$  and  $B$  represent the electrodes of the device,  $C_p$  is the electrode capacitance,  $C_s$  and  $L_s$  represent the mechanical resonance and  $R_s$  models the damping.

Current work on PRs has been focused around developing new converter topologies. Pollet et al [14] presented a novel piezoresonator based boost converter, showing significant voltage gain and good efficiency. Boles et al [15] presented several new topologies of PR converters, each offering unique operating voltage ranges, efficiencies and output voltages. Touhami et al [16] presented a new PR converter offering improved performance over those presented by Boles et al, with an output power of 175W at 80% efficiency, significantly higher output power than previous PR converters and higher than PT based converters [17]. These proposed converters do not require magnetic inductors to operate efficiently and can operate over a wider range of voltage conversion ratios than the inductorless half-bridge PT converter.

Other work by Boles et al and Touhami et al focuses on optimal material choice for high power and high efficiency PRs [18], [19]. Authors have subsequently presented control schemes for the proposed PR converters, Piel et al [20] presented a simple feedback controller for a step-down converter topology presented in [14], based on either sensing the PR voltage or by using fixed time durations for each of the switching modes. Both Touhami et al [21] and Forrester et al [22] present control strategies for the PR based converters presented in [16] and [14], respectively. Both papers use comparators and voltage sensing to determine optimal timings for each of the MOSFETs within the converter. In [21] the authors focus on regulating the output voltage using a PI controller, whereas in [22] the authors focused on optimising the operating frequency using a phase-locked loop (PLL).

The converter described in [14] and [23] is named here as the inductorless step-up piezoelectric resonator (SUPR) converter for convenience. It is of particular interest as it can be used for step-up power conversion of DC supplies. Unlike the PT-based half-bridge resonant converter, the SUPR can achieve ZVS under a wide variety of operating conditions if a suitable controller is used.

Modelling of resonant converters is often performed to assist the design process of the converter. Many methods of analysing PT-based resonant converters have been presented in the literature. These methods include a parametric analysis [24], which allow the ZVS characteristics of a PT converter to be estimated. However, this method requires a number of assumptions to be made, limiting its application. Cyclic mode analysis [25] is also commonly used and provides accurate estimation of voltages and current in a PT based converter. One downside to cyclic mode modelling is that only numerical

solutions can be found. Parametric sweeps and curve fitting [26] have also been employed, where PT parameters are varied and the resulting converter simulated (using SPICE), the results are then curve fit against the varied parameters to estimate a simplified equation for the ZVS characteristics, offering an approximate analytical solution. However, this method can get overly complex when assessing many variables and by not considering all variables can make the resulting equation inaccurate. A describing function approach using fundamental mode analysis [27] was also used to analyse the inductorless half-bridge PT based converter offering a good balance of accuracy and complexity to achieve a very simplified equation for designing PTs for achieving ZVS. As a result, the modelling in this paper will use a similar describing function approach.

In [15], Boles et al applied a modelling approach to several topologies of PR converter. However, obtaining an accurate estimation of the switch timings using this method requires the solution of 15 equations with 18 variables, which is highly computationally complex. In [14], Pollet et al present a model for the SUPR converter, deriving expressions for the voltage gain and piezoelectric current amplitude, which contain only a small number of variables (circuit parameters and control variables). However, whilst offering a simple expression, operating frequency is an input variable and not calculated by the model. Owing to the high Q factor, operating frequency has a significant effect on the operation of PR converters, therefore, additional steps should be taken to determine the optimal operating frequency when using this model. Additionally, both models do not include the effects of semiconductor losses, leading to inaccuracy.

In this paper a derivation and evaluation of an analytical model of the SUPR converter is presented using a describing function approach. The operation of the SUPR converter is discussed and decomposed into six distinct operating modes. A piecewise description of the resonator input voltage,  $v_{C_p}$ , is derived in terms of circuit component values and mode durations. The optimal mode durations for given operating conditions (load and output voltage) are estimated by analysing energy conservation, charge conservation and using a describing function approximation of the resonator voltage. Unlike previous approaches, this analysis then allows the optimum switch conduction times and optimum operating frequency to be found for any set of circuit parameters. In addition, solving for mode durations only requires solving a single equation with only 4 unknowns, a large reduction compared to previous approaches. The model is compared to experimental and Simulink results for particular component values. From this analysis, guidelines for design and control are provided.

## II. OPERATION OF CONVERTER

Fig. 2 shows the circuit diagram for the piezoelectric-resonator boost converter with the resonator being represented by the equivalent circuit within the dashed box. The operating waveforms for the converter are shown in Fig. 3, with the start of a cycle ( $\theta = 0$ ) being the positive zero crossing of resonant current  $i_{L_s}$  (note that this current is equivalent to the particle velocity in the piezoelectric resonator; it cannot be measured directly). A cycle is divided into six modes of operation

depending on the conduction states of the diodes and the control of the MOSFET switches. The circuit is driven at an operating frequency  $f$  such that the zero crossings of the resonant current  $i_{L_s}$  occur at  $\theta = 0, \theta_3 = \pi, \theta_6 = 2\pi$  with a period  $T = 1/f = 2\pi/\omega$ .

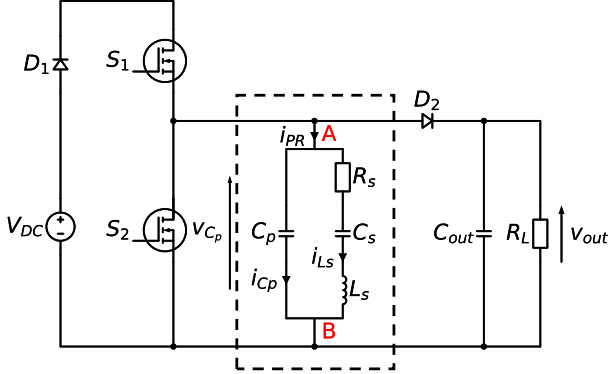


Fig. 2 – Piezoelectric-resonator-based boost converter

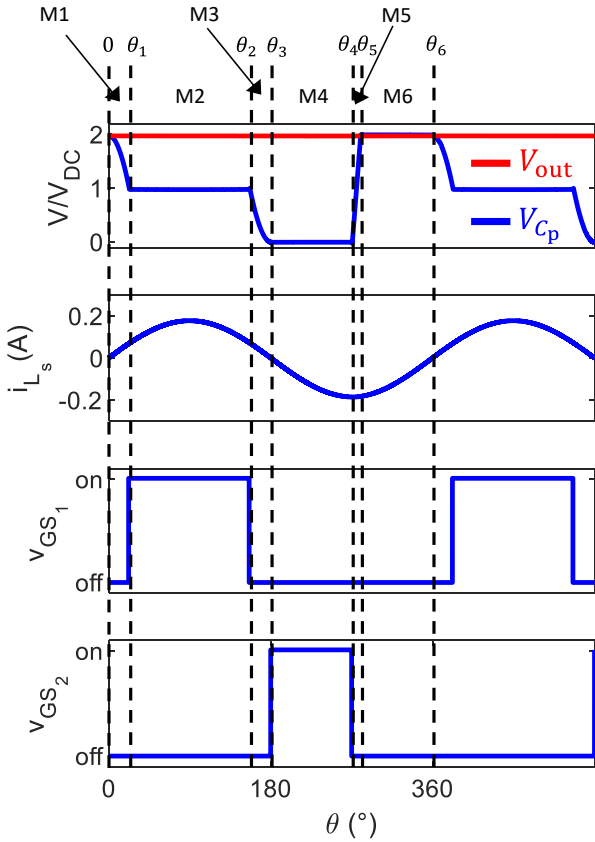


Fig. 3 – Typical operating waveforms for the piezoelectric resonator-based boost converter

The six modes of operation M1  $\rightarrow$  M6 for this converter during ideal operation are:—

M1 ( $0 \rightarrow \theta_1$ ): Prior to start of a cycle,  $D_2$  is conducting such that  $v_{C_p}(0) = V_{out} + V_{df}$ , where  $V_{out}$  is the output voltage and  $V_{df}$  is the forward voltage drops of  $D_1$  and  $D_2$ . Both MOSFETs  $S_1$  and  $S_2$  are turned off at  $\theta = 0$ , as the current  $i_{L_s}$  undergoes a positive zero crossing. During M1,  $i_{L_s}$  circulates within the PR discharging  $C_p$  and causing  $v_{C_p}$  to decrease until it reaches

$V_{DC} - V_{df}$ . Since  $v_{C_p} < V_{out}$ ,  $D_2$  is naturally turned off at  $\theta = 0$  and so  $C_{out}$  provides the energy to the load.

M2 ( $\theta_1 \rightarrow \theta_2$ ): At  $\theta_1$ ,  $v_{C_p} = V_{DC} - V_{df}$  and so  $S_1$  is turned on under ZVS condition. Energy is supplied to the resonator through  $S_1$  during M2.

M3 ( $\theta_2 \rightarrow \theta_3$ ):  $S_1$  is turned off at  $\theta_2$ .  $\theta_3$  is chosen to allow sufficient time for  $i_{L_s}$  to discharge  $v_{C_p}$  to 0.

M4 ( $\theta_3 \rightarrow \theta_4$ ): At  $\theta_3$ ,  $i_{L_s}$  undergoes a negative zero crossing and, if the conditions are correct,  $v_{C_p}(\theta_3) = 0$ . Therefore,  $S_2$  turns-on at  $\theta_3$  under ZVS condition. During M4,  $v_{C_p} \approx 0$  and  $i_{L_s}$  flows through  $S_2$ .

M5 ( $\theta_4 \rightarrow \theta_5$ ): At  $\theta_4$ ,  $S_2$  is turned off, the negative  $i_{L_s}$  charges  $C_p$  thereby increasing  $v_{C_p}$ . At  $\theta_5$ ,  $v_{C_p} = V_{out} + V_{df}$ .

M6 ( $\theta_5 \rightarrow \theta_6$ ):  $v_{C_p} = V_{out} + V_{df}$ ,  $D_2$  is forward-biased, thus energy is provided from the resonator to  $R_L$  and  $C_{out}$ . At  $\theta_6 = 2\pi$ ,  $i_{L_s}$  undergoes another positive zero crossing and the cycle repeats.

### III. MODEL DERIVATION

In this section a model is derived for the SUPR converter, allowing the output voltage, resonant current and resonator voltage ( $v_{C_p}$ ) to be estimated. The derivation begins by developing a piecewise equation for  $v_{C_p}$  based on the resonant current, mode angles and equivalent circuit component values. During operation,  $\theta_4$  (or the  $S_2$ -conduction angle,  $(\theta_4 - \pi)$ ) is used to control the voltage gain of the converter. We will therefore assume  $\theta_4$  is fixed by some external controller and is an input to the model. The unknown mode angles ( $\theta_1, \theta_2$  and  $\theta_5$ ) are found, allowing the converter operation to be modelled across a cycle and allowing metrics such as output voltage and resonant current to be determined.

The unknown mode angles are determined by solving simultaneous equations (named  $J_1, J_2, J_3$  and  $J_4$ ) which are derived based on diode commutation conditions, energy balance, output charge conservation and by equating the signal-based fundamental mode impedance of the resonator to the circuit-based impedance using a describing-function approximation of the  $v_{C_p}$ . Thus, the resulting four simultaneous equations can be solved numerically to find the three unknown mode angles and operating frequency. Finally, equations are derived for the output voltage and resonant current amplitude based on the mode angles and equivalent circuit component values thereby completing the model to allow  $V_{out}$ ,  $v_{C_p}$  and  $i_{L_s}$  to be estimated for any PR based converter.

#### A. Developing a piecewise model for $v_{C_p}$

In the derivation that follows it is assumed that the PR is constructed using high quality factor piezoelectric material ( $Q \gg 1$ ) thus the PR is considered to be a highly selective band-pass filter. Under these assumptions, the PR current  $i_{L_s}$  is sinusoidal and is defined as

$$i_{L_s} = I_{L_s} \sin(\theta) \quad (1)$$

where  $I_{L_s}$  is the amplitude of the resonant current,  $i_{L_s}$ . During modes M1, M3 and M5, the power devices are off and so  $i_{L_s}$

circulates through  $C_p$  charging or discharging it. It should be noted that the MOSFET drain-source capacitance also appears in parallel with  $C_p$ .  $v_{C_p}$  during these modes is given by

$$v_{C_p}(\theta) = \frac{1}{\omega C_p} \int -i_{L_s} d\theta = \frac{I_{L_s} \cos(\theta)}{\omega C_p} + V_i \quad (2)$$

where  $i$  is the mode number,  $V_i$  is the constant of integration associated with the mode and  $\omega$  is the operating frequency (rad s<sup>-1</sup>). We assume that the resonant current crosses through zero at  $0, \theta_3$  and  $\theta_6$  and  $C_p \gg C_{DS}$ , where  $C_{DS}$  is the MOSFET drain-source capacitance.

In M1,  $V_1$  can be found by solving (2) with  $v_{C_p}(0) = V_{out} + V_{df}$ , thus giving

$$V_1 = V_{out} + V_{df} - \frac{I_{L_s}}{\omega C_p} \quad (3)$$

where we assume the output voltage  $V_{out}$  is constant. Thus, for M1,

$$v_{C_p} = \frac{I_{L_s}(\cos \theta - 1)}{\omega C_p} + V_{out} + V_{df} \quad 0 \leq \theta < \theta_1 \quad (4)$$

In M2,  $S_1$  is on, allowing energy to be injected into the resonator through  $S_1$  and  $D_1$  and, thus,  $v_{C_p}$  is clamped to the supply voltage, less the diode forward voltage drop and the voltage drop across the  $R_{DS_{on}}$  of the switch,

$$v_{C_p} = V_{DC} - V_{df} - I_{L_s} R_{DS_{on}} \sin \theta \quad \theta_1 \leq \theta < \theta_2 \quad (5)$$

In M3,  $S_1$  is off and  $i_{L_s}$  circulates through  $C_p$  discharging it to 0 at  $\theta_3 = \pi$ .  $V_3$  can be found by considering  $v_{C_p}(\theta_2) = V_{DC} - V_{df} - I_{L_s} \sin \theta_2 R_{DS_{on}}$

$$V_3 = V_{DC} - V_{df} - I_{L_s} \left( R_{DS_{on}} \sin \theta_2 + \frac{\cos(\theta_2)}{\omega C_p} \right) \quad (6)$$

Thus,

$$v_{C_p} = V_{DC} - V_{df} + I_{L_s} \left( \frac{\cos \theta - \cos \theta_2}{\omega C_p} - R_{DS_{on}} \sin \theta_2 \right) \quad \theta_2 \leq \theta < \pi \quad (7)$$

In M4,  $S_2$  is on and the  $v_{C_p}$  voltage is clamped at 0, minus the voltage drop across the  $R_{DS_{on}}$  of the switch, therefore

$$v_{C_p} = -I_{L_s} R_{DS_{on}} \sin \theta \quad \pi \leq \theta < \theta_4 \quad (8)$$

In M5,  $S_2$  has turned off and  $i_{L_s}$  charges  $C_p$ .  $V_5$  can be found solving  $v_{C_p}(\theta_4) = -I_{L_s} \sin \theta_4 R_{DS_{on}}$

$$V_5 = -I_{L_s} \left( \frac{\cos \theta_4}{\omega C_p} + R_{DS_{on}} \sin \theta_4 \right) \quad (9)$$

Thus

$$v_{C_p} = I_{L_s} \left( \frac{\cos \theta - \cos \theta_4}{\omega C_p} - R_{DS_{on}} \sin \theta_4 \right) \quad \theta_4 \leq \theta < \theta_5 \quad (10)$$

Finally, during M6,  $v_{C_p}$  is given by

$$v_{C_p} = V_{out} + V_{df} \quad \theta_5 \leq \theta < 2\pi \quad (11)$$

A piecewise description of the  $v_{C_p}$  voltage can be written

$$v_{C_p} = \begin{cases} V_{out} + V_{df} + \frac{I_{L_s}}{\omega C_p} (\cos \theta - 1) & 0 \leq \theta < \theta_1 \\ V_{DC} - V_{df} - I_{L_s} R_{DS_{on}} \sin \theta & \theta_1 \leq \theta < \theta_2 \\ I_{L_s} \left( \frac{\cos \theta - \cos \theta_2}{\omega C_p} - R_{DS_{on}} \sin \theta_2 \right) + V_{DC} - V_{df} & \theta_2 \leq \theta < \pi \\ -I_{L_s} R_{DS_{on}} \sin \theta & \pi \leq \theta < \theta_4 \\ I_{L_s} \left( \frac{\cos \theta - \cos \theta_4}{\omega C_p} - R_{DS_{on}} \sin \theta_4 \right) & \theta_4 \leq \theta < \theta_5 \\ V_{out} + V_{df} & \theta_5 \leq \theta < 2\pi \end{cases} \quad (12)$$

## B. Finding mode angles

In order to generate the  $v_{C_p}$  waveform and calculate the output voltage and resonant current amplitude, the unknown mode angles ( $\theta_1, \theta_2$  and  $\theta_5$ ) and the operating frequency ( $\omega$ ) need to be determined.

### 1) Output capacitance charge conservation

In steady state there is no change in the output capacitor voltage from the end of one cycle to the end of the next (i.e.  $V_{out}(t) = V_{out}(t + T)$ ). Thus, the charge transferred through  $D_2$ , denoted by  $Q_{D_2}$ , must balance with the charge that is delivered to the load,  $Q_{load}$  over a whole cycle. Charge is transferred through  $D_2$  only during M6. The current  $i_{D_2}$  during conduction is given by,

$$i_{D_2} = i_{D_1} - i_{L_s} - i_{C_p} \quad (13)$$

During M6,  $S_1$  is off, therefore  $i_{D_1} = 0$ . Similarly, since  $v_{C_p}$  is constant  $\frac{dV_{C_p}}{dt} = 0 \Rightarrow i_{C_p} = 0$ . Therefore,  $i_{D_2} = -i_{L_s}$  during M6. Thus, an equation for the charge through  $D_2$  is given by,

$$Q_{D_2} = \frac{1}{\omega} \int_{\theta_5}^{2\pi} -i_{L_s} d\theta = \frac{I_{L_s} (1 - \cos \theta_5)}{\omega} \quad (14)$$

Since  $I_{out}$  is consumed by the load, the charge delivered to the load over the course of one cycle is given by

$$Q_{load} = I_{out} T = 2\pi \frac{I_{out}}{\omega} = 2\pi \frac{V_{out}}{R_L \omega} \quad (15)$$

Therefore, equating  $Q_{D_2}$  (14) and  $Q_{load}$  (15) and rearranging gives an equation for  $V_{out}$ ,

$$V_{out} = \frac{I_{L_s} R_L (1 - \cos \theta_5)}{2\pi} \quad (16)$$

Another equation for  $V_{out}$  can be generated by noting that  $v_{C_p}(\theta_5) = V_{out} + V_{df}$ , thus, by substituting  $v_{C_p}(\theta_5) = V_{out} + V_{df}$  into (10) and rearranging,

$$V_{out} = -V_{df} + I_{L_s} \left( \frac{\cos \theta_5 - \cos \theta_4}{\omega C_p} - R_{DS_{on}} \sin \theta_4 \right) \quad (17)$$

Both (16) and (17) are dependent on  $I_{L_s}$ . An expression for  $I_{L_s}$  can be found by substituting  $v_{C_p}(\pi) = 0$  into (7),

$$V_{DC} = V_{df} + I_{L_s} \left( \frac{\cos \theta_2 + 1}{\omega C_p} + R_{DS_{on}} \sin \theta_2 \right) \quad (18)$$

Then, (18) can be rearranged for  $I_{L_s}$  giving,

$$I_{L_s} = \frac{V_{DC} - V_{df}}{\frac{\cos \theta_2 + 1}{\omega C_p} + R_{DS_{on}} \sin \theta_2} \quad (19)$$

Finally, by subtracting (16) from (17), and using (19) gives the first simultaneous equation  $J_1$  in terms of circuit components and mode angles

$$J_1 = 0 \quad (20)$$

where

$$\begin{aligned} J_1 = & -(\omega C_p R_L + 2\pi) \cos \theta_5 + 2\pi \cos \theta_4 \\ & + 2\pi \omega C_p R_{DS_{on}} \sin \theta_4 \\ & + \frac{1}{V_{dc} - V_{df}} \left( 2\pi V_{df} \left[ \cos \theta_2 \right. \right. \\ & + \omega C_p R_{DS_{on}} \sin \theta_2 + \left. \left. \left\{ 1 - \frac{R_L \omega C_p}{2\pi} \right\} \right] \right. \\ & \left. + \omega C_p R_L V_{dc} \right) \end{aligned} \quad (21)$$

## 2) Energy conservation

The second simultaneous equation  $J_2$  is developed based on energy conservation. The energy into the converter will equal the energy supplied to the load plus the losses in the circuit.

Energy is supplied to the resonator during M2, therefore an equation for the input energy is given by

$$E_{in} = \frac{1}{\omega} \int_{\theta_1}^{\theta_2} i_{L_s}(\theta) V_{DC} d\theta \quad (22)$$

Evaluating this gives,

$$E_{in} = \frac{I_{L_s} V_{DC} (\cos(\theta_1) - \cos(\theta_2))}{\omega} \quad (23)$$

Similarly, an equation for the energy delivered to the load can be derived by noting that energy is transferred to the load throughout the whole cycle, thus

$$E_{out} = \frac{1}{\omega} \int_0^{2\pi} \frac{V_{out}^2}{R_L} d\theta = \frac{2\pi V_{out}^2}{\omega R_L} \quad (24)$$

An equation for the energy dissipated in the resonator is given by

$$E_{R_{sLoss}} = \frac{R_s}{\omega} \int_0^{2\pi} I_{L_s}^2 \sin^2(\theta) d\theta = \frac{I_{L_s}^2 R_s \pi}{\omega} \quad (25)$$

Additionally, we also consider the conduction losses in the two MOSFETs (we assume ZVS thus no switching losses). Conduction losses occur when the MOSFET is conducting, thus equations for these losses are given by

$$\begin{aligned} E_{S_{1Loss}} &= \frac{R_{DS_{on}}}{\omega} \int_{\theta_1}^{\theta_2} I_{L_s}^2 \sin^2(\theta) d\theta \\ &= \frac{I_{L_s}^2 R_{DS_{on}}}{2\omega} \left( \theta_2 - \theta_1 \right. \\ &\quad \left. + \frac{1}{2} (\sin 2\theta_2 - \sin 2\theta_1) \right) \end{aligned} \quad (26)$$

$$\begin{aligned} E_{S_{2Loss}} &= \frac{R_{DS_{on}}}{\omega} \int_{\pi}^{\theta_4} [I_{L_s} \sin(\theta)]^2 d\theta \\ &= \frac{I_{L_s}^2 R_{DS_{on}}}{2\omega} \left( \theta_4 - \pi - \frac{\sin 2\theta_4}{2} \right) \end{aligned} \quad (27)$$

Finally, we also consider the conduction losses in the two diodes. Conduction losses occur when the diode is conducting thus equations for these losses are given by

$$\begin{aligned} E_{D_{1Loss}} &= \frac{V_{df}}{\omega} \int_{\theta_1}^{\theta_2} I_{L_s} \sin(\theta) d\theta \\ &= \frac{V_{df} I_{L_s} (\cos \theta_1 - \cos \theta_2)}{\omega} \end{aligned} \quad (28)$$

$$E_{D_{2Loss}} = \frac{V_{df}}{\omega} \int_{\theta_5}^{2\pi} I_{L_s} \sin(\theta) d\theta = \frac{V_{df} I_{L_s} (\cos \theta_5 - 1)}{\omega} \quad (29)$$

where  $V_{df}$  is the diode forward voltage drop. Therefore, due to energy conservation, the input energy is equated to the output energy plus losses

$$E_{in} = E_{out} + E_{R_{sLoss}} + E_{S_{1Loss}} + E_{S_{2Loss}} + E_{D_{1Loss}} + E_{D_{2Loss}} \quad (30)$$

Substituting (16), (23)-(29) into (30) gives (35). (19) is used to eliminate  $I_{L_s}$  in (35), then by rearranging to move all terms on to the right-hand side, gives the second simultaneous equation  $J_2$

$$J_2 = 0 \quad (31)$$

Where  $J_2$  is defined in (36).

### 3) Describing function for $v_{C_p}$

A third equation for the mode angles can be generated by equating the signal-based equivalent impedance of the resonant circuit to the component-based impedance of the resonant circuit. The signal-based impedance is generated by first performing fundamental mode analysis to generate a describing function approximation of the  $v_{C_p}$  voltage,  $v_{C_p<1>}$ .

We define the fundamental component (FC) of a waveform as its first term of the Fourier series which in general can be written as

$$v_x(\theta) \xrightarrow{\text{FC}} V_{x<1>} = \frac{1}{\pi} \int_0^{2\pi} v_x(\theta) e^{-j\theta} d\theta \quad (32)$$

$v_{C_p<1>}$  is evaluated using the piecewise description of  $v_{C_p}$  given in (12) and the result simplified using (16) and (18). The resulting equation for  $v_{C_p<1>}$  has been omitted due to length and instead is written as a function of its dependent variables

$$v_{C_p<1>} = f(I_{L_s}, V_{df}, C_p, R_L, \omega, \theta_1, \theta_2, \theta_4, \theta_5) \quad (33)$$

A Python script for generating  $v_{C_p<1>}$  using the SymPy library [28] is provided in the Appendix. The signal-based impedance can then be calculated using (19) and Ohm's law, with the result given in (34).

$$Z_{SB} = \frac{v_{C_p<1>}}{I_{L_s}} = g(V_{df}, C_p, R_L, \omega, \theta_1, \theta_2, \theta_4, \theta_5) \quad (34)$$

We assume this impedance provides a good model of the converter under its particular operating conditions. The component-based impedance of the piezoelectric resonator is given by

$$Z_{CB} = j\omega L_s + \frac{1}{j\omega C_s} + R_s \quad (37)$$

Since  $Z_{SB}$  and  $Z_{CB}$  will be equal under our assumption (and close in practice),  $0 = Z_{SB} - Z_{CB}$ . Substituting (37), into this gives

$$0 = Z_{SB} - j\omega L_s - \frac{1}{j\omega C_s} - R_s \quad (38)$$

An equation for  $J_3$  is generated by subtracting the imaginary part from the real part of (38), which must both equal 0.

$$J_3 = 0 \quad (39)$$

$$J_3 = \Re\left(Z_{SB} - j\omega L_s - \frac{1}{j\omega C_s} - R_s\right) - \Im\left(Z_{SB} - j\omega L_s - \frac{1}{j\omega C_s} - R_s\right) \quad (40)$$

### 4) Turn-on angle of $S_1$

As the  $v_{C_p}$  waveform is continuous, then at the end of M1  $v_{C_p}(\theta_1^-) = v_{C_p}(\theta_1^+)$ . Substituting (12) into this and rearranging for  $V_{DC}$  gives

$$V_{DC} = 2V_{df} + V_{out} + I_{L_s} \left( \frac{\cos \theta_1 - 1}{\omega C_p} + R_{DS_{on}} \sin \theta_1 \right) \quad (41)$$

$V_{out}$  is eliminated by substituting (16) into (41)

$$V_{DC} = 2V_{df} + I_{L_s} \left( \frac{\cos \theta_1 - 1}{\omega C_p} + R_{DS_{on}} \sin \theta_1 - \frac{R_L(\cos \theta_5 - 1)}{2\pi} \right) \quad (42)$$

Finally, substituting (18) into (42) and using (19) to eliminate  $I_{L_s}$ , the result is rearranged to move all terms on to the right-hand side, giving an equation for the fourth and final simultaneous equation  $J_4$

$$\begin{aligned} & 2\pi(I_{L_s} + \omega C_p I_{L_s} R_{DS_{on}} \sin \theta_2 + \omega C_p V_{df} + I_{L_s} \cos \theta_2) \frac{\cos \theta_1 - \cos \theta_2}{\omega C_p} \\ & = I_{L_s} R_L \cos^2 \theta_5 + (2\pi V_{df} - 2I_{L_s} R_L) \cos \theta_5 + \pi(2V_{df} + I_{L_s} R_{DS_{on}} \sin \theta_1) \cos \theta_1 \\ & - \pi(2V_{df} + I_{L_s} \sin \theta_2 R_{DS_{on}}) \cos \theta_2 - \pi I_{L_s} R_{DS_{on}} \cos \theta_4 \sin \theta_4 + 2\pi^2 \left( R_s - \frac{R_{DS_{on}}}{2} \right) I_{L_s} \\ & - \pi(2V_{df} - R_{DS_{on}}(\theta_2 + \theta_4 - \theta_1) I_{L_s}) + I_{L_s} R_L \end{aligned} \quad (35)$$

$$\begin{aligned} J_2 = \frac{1}{V_{DC} - V_{df}} & \left( -\omega C_p (V_{DC} - V_{df}) R_L \cos^2 \theta_5 - 2\pi (V_{DC} - V_{df}) \cos^2 \theta_2 \right. \\ & - \left( (-2V_{DC} + 2V_{df}) \cos \theta_1 + 2V_{df} \cos \theta_5 + \omega C_p (V_{DC} - V_{df}) R_{DS_{on}} \sin \theta_2 + 2V_{DC} - 4V_{df} \right) \pi \cos(\theta_2) \\ & - \pi (V_{DC} - V_{df}) (\omega C_p R_{DS_{on}} \sin \theta_1 - 2\omega C_p R_{DS_{on}} \sin \theta_2 - 2) \cos \theta_1 \\ & + (-2\pi \omega C_p V_{df} R_{DS_{on}} \sin \theta_2 - 2\pi V_{df} + 2\omega C_p (V_{DC} - V_{df}) R_L) \cos \theta_5 + 2\pi \omega C_p V_{df} R_{DS_{on}} \sin \theta_2 \\ & + \pi \omega C_p (V_{DC} - V_{df}) R_{DS_{on}} \cos \theta_4 \sin \theta_4 - 2\pi^2 \omega C_p (V_{DC} - V_{df}) \left( R_s - \frac{R_{DS_{on}}}{2} \right) \\ & \left. + \pi (-\omega C_p R_{DS_{on}} (V_{DC} - V_{df}) (\theta_2 + \theta_4 - \theta_1) + 2V_{df}) - \omega C_p (V_{DC} - V_{df}) R_L \right) \end{aligned} \quad (36)$$

$$J_4 = 0 \quad (43)$$

where  $J_4$  is defined as,

$$J_4 = \frac{1}{V_{DC} - V_{df}} \left( 2\pi(V_{DC} - 2V_{df})(\cos \theta_2 + \omega C_p R_{DS_{on}} \sin \theta_2) + 4\pi V_{DC} - 6\pi V_{df} \right) + \omega C_p (R_L \cos \theta_5 - 2\pi R_{DS_{on}} \sin \theta_1) - 2\pi \cos \theta_1 - C_p R_L \omega \quad (44)$$

#### 5) Calculating mode angles

A non-linear minimisation is performed on a cost function comprising of equations (21), (36), (40) and (44) as below,

$$J_{sum} = J_1^2 + J_2^2 + J_3^2 + J_4^2 \quad (45)$$

where  $J_{sum} = 0$  is achieved with ideal mode angles and at the ideal operating frequency. Since model precision and component tolerances provide limits to any model, the minimum value of  $J_{sum}$  is the closest approximation. The sequential quadratic programming (SQP) algorithm is used within MATLAB (the *FMINCON* function) to find the unknowns, including optimal mode angles ( $\theta_1$ ,  $\theta_2$ ,  $\theta_5$ ) and operating frequency ( $\omega$ ). The SQP algorithm is a state of the art non-linear solving method, outperforming other algorithms in terms of efficiency and accuracy [29]. A series of logical constraints are used:  $0 \leq \theta_1 \leq \theta_2 \leq \pi$ ,  $\theta_4 \leq \theta_5 \leq 2\pi$  and  $\omega \geq 1/\sqrt{L_s C_s}$ . Initial estimates are used as follows:  $\omega = 1/\sqrt{L_s C_s}$ ,  $\theta_1 = \pi/10$ ,  $\theta_2 = 3\pi/10$ ,  $\theta_5 = \theta_4 + \pi/20$  where  $\theta_4$  is an input to the cost function. Although this process is involved, it is still significantly less computationally intensive than previously presented methods.

#### C. Finding $I_{L_s}$ , $V_{out}$ and $V_{C_p}$

Based on the calculated mode angles and operating frequency, the output voltage and resonant current amplitude can be found, after which the  $v_{C_p}$  waveform can be plotted. To find  $V_{out}$ , (19) is substituted in to (17) and simplified giving

$$V_{out} = \frac{1}{\sin \theta_2 R_{DS_{on}} \omega C_p + \cos(\theta_2) + 1} \left( (V_{DC} - V_{df})(\cos \theta_4 - \cos \theta_5 + \omega C_p R_{DS_{on}} \sin \theta_4) - V_{df}(1 + \omega C_p \sin \theta_2 R_{DS_{on}} + \cos \theta_2) \right) \quad (46)$$

Finally, (19) and (46) can be substituted into (12) leaving an expression for  $v_{C_p}$  exclusively in terms of circuit component values, mode angles and frequency, thus allowing the whole  $v_{C_p}$  waveform to be evaluated.

## IV. EXPERIMENTAL VERIFICATION

To verify the proposed model, a practical piezo resonator-based converter is constructed using the Steminc SMD30T21F1000S piezo resonator shown in Fig. 1b. The equivalent circuit properties of this resonator, given in Table 1, were extracted from impedance measurements taken using an

Omicron Bode 100 vector network analyser, using the characterisation technique described in [30].

A SUPR converter was constructed using IRF510 MOSFETs (with an IR2110 gate drive IC), 1N5819 diodes, a 12V DC supply is used and  $C_{out} = 10\mu F$ . The converter is driven by a PLL-based controller which synchronises the MOSFET switching events to the resonant current  $i_{L_s}$ . Since  $i_{L_s}$  is internal to the PR, a current estimator is required, as presented in [22]. The current estimator measures the current through the PR ( $i_{L_s} + i_{C_p}$ ) and subtracts the estimated current through  $C_p$  ( $i_{C_p}$ ), giving an estimate of the resonant current ( $i_{L_s}$ ). The zero crossings of the estimated resonant current are extracted using a comparator to give a square wave, which is in-phase with the resonant current.

A Xilinx PYNQ FPGA is used to control the converter. This FPGA contains both programmable logic (PL) and a microprocessor (PS) running Linux. The PS runs a web server allowing communication with both the PS and PL using a python script. First, a simple all digital PLL is implemented on the FPGA, which subsequently locks onto the estimated resonant current of the PR, producing a square wave to drive the rest of the logic.

$S_2$  is comparatively easy to drive, the MOSFET turns on at the falling edge of the PLL output signal and operates with a fixed duty cycle, which is controlled by a python script. Two comparators are used to generate a signal when the PR voltage drops is  $\leq V_{DC}$  and  $\leq 0$ .  $S_1$  turns on when the PR voltage equals  $V_{DC}$ , then the duty cycle of  $S_1$  is adjusted by a control loop on the FPGA, to cause the PR voltage to cross equal 0 just before  $S_2$  turns on. Additional elements are also included to account for propagation delays, noisy signals and to stop shoot-through. Full details of the operation of the controller is described in [22].

Table 1 - SMD30T21F1000S equivalent circuit properties

$R_s$	$L_s$	$C_s$	$C_p$
2.22 $\Omega$	4.47mH	1.02nF	2.54nF

Fig. 4 shows experimental waveforms for this converter when driving a 1 k $\Omega$  load, measured using a Rigol DS1054Z oscilloscope. PLL<sub>out</sub> in Fig. 4 is the output from the FPGAs PLL and when locked, represents the polarity of the estimated resonant current  $i_{L_s}$  (estimated using signals  $V_A$  and  $V_B$ ). PLL<sub>out</sub> is then used to derive the gate signals,  $V_{gs1}$  and  $V_{gs2}$  for MOSFETs  $S_1$  and  $S_2$  respectively.

A similar converter is constructed within Simulink using the Simscape electrical package. A state machine-based controller is constructed to drive the two MOSFETs with parameters set to match the IRF510. After initialization (with fixed mode time durations calculated using the proposed model), the state machine determines the ideal mode durations based on the resonant current and the  $v_{C_p}$  voltage, as described in section II.



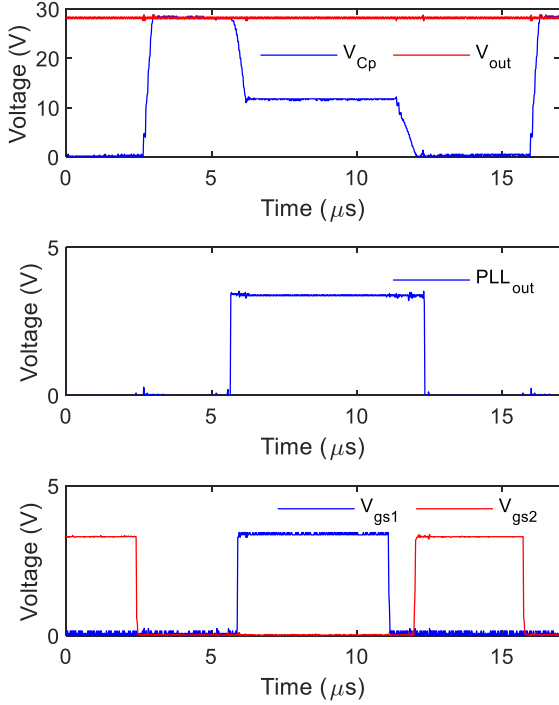


Fig. 4 – Experimental waveforms for the SUPR converter when driving a  $1k\Omega$  load with a 26%  $S_2$  duty cycle [22]. The piezo resonator and output voltages (top), the PLL output signal (middle) and MOSFET gate signal (bottom) are shown

To provide validation, two voltage gain  $M_{out} = V_{out}/V_{DC}$  measurement experiments are performed. First, the duty cycle of the low-side switch,  $S_2 \left( \frac{(\theta_4 - \pi)}{2\pi} \right)$  is varied through a range of 20-32%, and  $M_{out}$  is measured using an oscilloscope when driving a  $1k\Omega$  load. Second, the converter is operated with the duty cycle fixed at 26% and the load varied between 500 and  $10k\Omega$ . Similar analysis is performed using the Simulink model for comparison. Finally, the proposed describing function model is used to estimate  $M_{out}$  under the same load and duty cycle conditions. The results of these experiments are shown in Fig. 5 and Fig. 7.

Fig. 5 shows the voltage gain predicted by the proposed model ('Descrb. Func'), the model in [14] (using frequencies calculated using the proposed model), Simulink simulations and experimental results with changes in  $S_2$  duty cycle  $\left( \frac{(\theta_4 - \pi)}{2\pi} \right)$ , when driving a  $1k\Omega$  load. The voltage gain estimated by the proposed model shows good agreement with both the experimental and Simulink measurements, with a slight overestimation at high  $S_2$  duty cycles. Overall, the model shows close agreement with the experimental results, exhibiting a mean absolute percentage error (MAPE) of 1.44% and a standard deviation in the percentage error of 1.12%, where MAPE is defined as

$$MAPE = \frac{1}{n} \sum_{i=0}^n \left| \frac{V_{out,Model}(i) - V_{out,Exp}(i)}{V_{out,Exp}(i)} \right| \quad (47)$$

This compares to a 4.2 – 39.5% MAPE and a standard deviation in the percentage error of 2.1 - 2.9% from the model in [14], depending on frequencies used, either those estimated by the proposed model or those from the experimental results.

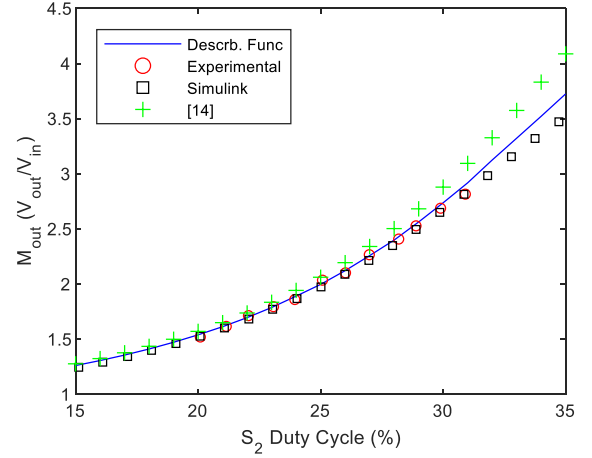


Fig. 5 – Voltage gain  $M_{out}$  vs  $S_2$  duty cycle, with a load  $R_L=1k\Omega$

The overestimation that is observed is due mainly to inaccuracies in the modelling of the diode voltage drop, inaccuracies in the resonator properties (again mainly  $R_s$ ) and unmodelled parasitic effects.

The mode durations and operating frequency estimated by the describing function model when generating the results presented in Fig. 5 are shown in Fig. 6 For convenience, the following mode duration angles are defined

$$\alpha = \theta_1 - 0 \quad (48)$$

$$\beta = \theta_2 - \theta_1 \quad (49)$$

$$\rho = \theta_5 - \theta_4 \quad (50)$$

Fig. 7 shows the voltage gain predicted by the proposed model ('Descrb. Func'), Simulink simulations and experimental results with variations in load, when driven with a duty cycle  $\left( \frac{(\theta_4 - \pi)}{2\pi} \right) = 26\%$ . The estimated voltage gain agrees well with the experimental results obtained from the practical converter, achieving greater accuracy estimations than the Simulink model, which overestimates the gain of the converter at larger value loads. However, at the maximum load tested, the experimental gain is lower than estimated by both the proposed model and Simulink. As a result, the model shows a MAPE of 2.03% compared to experimental measurements, with a standard deviation in the percentage error of 1.95%.

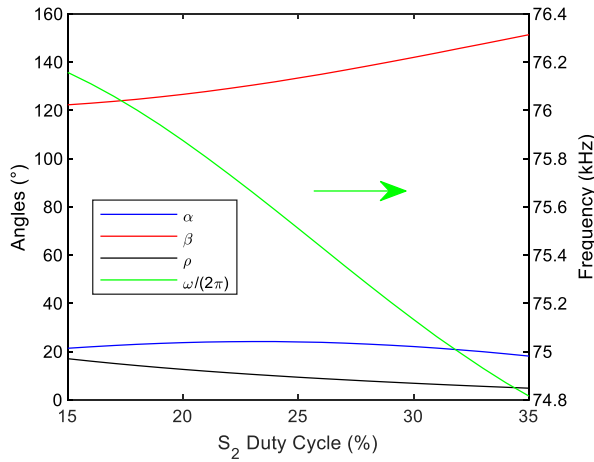


Fig. 6 – Mode angles vs  $S_2$  duty cycle, with a  $1k\Omega$  load

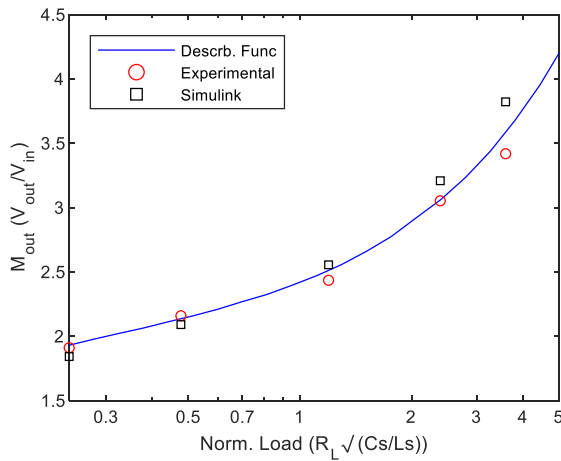


Fig. 7 -  $M_{out}$  vs normalised load (to the resonator characteristic impedance  $\sqrt{L_s/C_s}$ ) with 26%  $S_2$  duty cycle

## V. RESONATOR SENSITIVITY ANALYSIS

### A. Accuracy of the proposed model against LTSpice

A sensitivity analysis is performed by varying the resonator equivalent circuit properties ( $R_s$ ,  $L_s$ ,  $C_s$  and  $C_p$ ) and observing changes in gain. This analysis is performed using both the model derived here and LTSpice. This allows the error of the model with changes in resonator properties to be evaluated.

The converter simulated in LTSpice was operated with an  $R_L = 2k\Omega$  ( $\cong \sqrt{L_s/C_s}$ ) and a 12V DC supply. The converter was operated near its series resonant frequency, such that the negative going zero crossing of  $i_{L_s}$  occurs when  $S_2$  turns on.  $S_2$  was operated at a duty cycle  $\left(\frac{\theta_4 - \pi}{2\pi}\right) = 25\%$  for all variations. Given the complexity associated with determining an appropriate value for the  $S_1$  duty cycle  $\left(\frac{\theta_2 - \theta_1}{2\pi}\right)$ , the mode timing predicted by the proposed model was manually tuned in LTSpice until the response was deemed appropriate (i.e. where the  $v_{C_p}$  waveform crosses zero closest to the resonant current zero crossing at  $\theta_3$ ).

For this analysis, the characteristic impedance  $Z_0 = \sqrt{L_s/C_s}$ , resonant frequency  $f_0$ , damping resistance  $R_s$  and parallel capacitance  $C_p$  of the SMD30T21F1000S resonator were each individually varied. The ranges of variation were selected by assuming radial-mode PRs typically have radii,  $a$ , of 7.5mm to 25mm and thickness,  $h$ , of 0.5mm to 5mm. Based on the ranges of  $a$  and  $h$ , the ranges of  $C_p$ ,  $Z_0$  and resonant frequency can be calculated using the model presented in [31] for the extreme values of  $a$  and  $h$ . Therefore, the range of variation is chosen to be 0.1 to 10 times the nominal value given in Table 1 for  $C_p$  and  $Z_0$  and 0.6 to 2 times the nominal value for  $f_0$ . It is difficult to select a good range for  $R_s$  variation because its value depends on the physical construction of the device and its mounting. A range of 0.1 to 10 times the nominal value was chosen based on a study of existing resonators. It should be noted this investigation assumes that behaviour of the PR is dominated by radial mode vibration and other vibration modes (inc. spurious modes) was omitted from this analysis.

For each step in the parameter sweep, the gain of the converter  $V_{out,avg}/V_{DC}$ , where  $V_{out,avg}$  is the cycle average, was estimated by the model and measured using LTSpice and the percentage error calculated to determine the accuracy of the model.

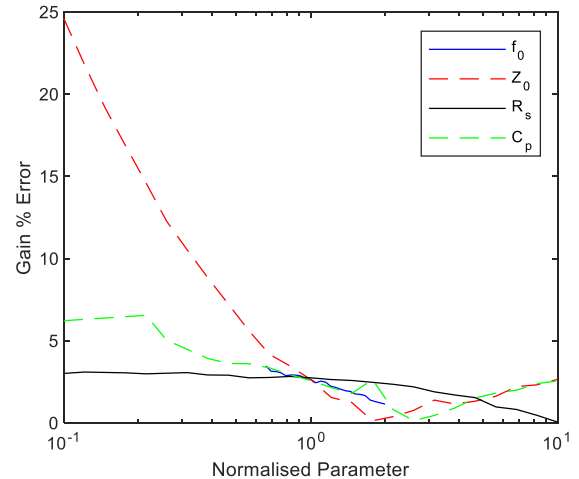


Fig. 8 – Sensitivity analysis evaluating the accuracy of the model when estimating gain with changes in resonator properties

Observing Fig. 8, the model shows acceptable accuracy across the range of variation, with only a  $\sim 3\%$  error at the nominal point and a maximum error of 25%. It is noticeable that the error is largest when each of the parameters are at values below the nominal value, this is especially clear with  $C_p$  and  $Z_0$ . At parameter values greater than nominal, error in the gain predicted by the model was  $<3\%$ . Errors are likely due to the assumptions and simplifications that are used within the model, this is especially true of  $Z_0$ , where we assume that the resonant current only contains a fundamental component, which is increasingly invalid at lower  $Z_0$ .

### B. Effect of the Resonator on converter performance

The previously performed ‘one-at-a-time’ sensitivity analysis can subsequently be used to analyse the impact of changing resonator properties on the gain and additionally the

efficiency and the operating frequency of the converter. This gives insight into optimal resonator designs for this converter.

As can be seen in Fig. 9, increases in damping resistance  $R_s$  cause an approximately quadratic decrease in voltage gain. This is as expected as more power is lost in the resonator, less power is supplied to the load and thus lower output voltages (gains) are observed. Increases in  $C_p$  cause corresponding increases in voltage gain and thus cause increased resonant current to flow through the resonator. Changes in  $f_0$  cause an almost identical effect to changes in  $C_p$  (as  $|Z_{C_p}| = 1/(\omega C_p)$ ), with increases in resonant frequency (or  $C_p$ ), increasing the gain of the converter.  $Z_0$  has a minimal impact on the voltage gain, with only a slight decrease in gain with increases in  $Z_0$ .

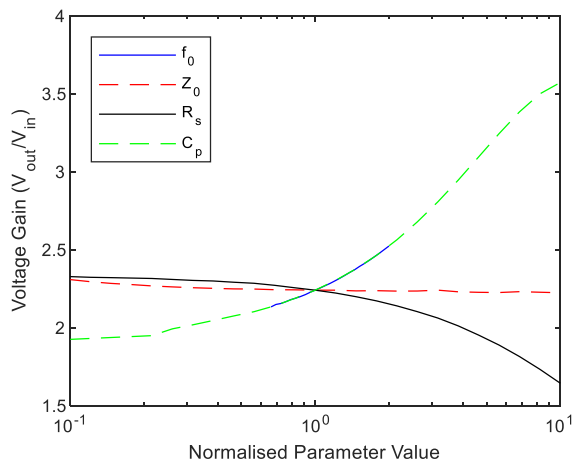


Fig. 9 – Voltage gain sensitivity analysis

The efficiency of the SUPR converter is given by

$$\eta = \frac{E_{out}}{E_{in}} = \frac{V_{out}^2 2\pi}{V_{DC} I_{L_s} R_L [\cos(\theta_1) - \cos(\theta_2)]} \quad (51)$$

Fig. 10 shows the efficiency (calculated using (51)) with changes in each of the parameters.  $Z_0$  has minimal impact on the efficiency of the converter, with only a slight increase with increase in  $Z_0$ . Increases in  $C_p$ ,  $f_0$  and  $R_s$  cause a similar decrease in efficiency, with an approximately quadratic trend. Again, changes in  $f_0$  and  $C_p$  cause an almost identical change in efficiency whereas changes in  $R_s$  cause a slightly shallower drop in efficiency, compared to  $C_p$ .

As can be seen in Fig. 11, the greatest change in normalised operating frequency occurs with changes in  $Z_0$ ; this is common property of resonant converters. As  $Z_0$  decreases from the nominal value, a sharp increase in normalised frequency is required for optimal operation. Similarly, increases in  $Z_0$  cause a decrease in operating frequency.  $Z_0$  has a larger effect on operating frequency than the other 3 parameters in this study. Both  $f_0$  and  $C_p$  have a similar effect, agreeing with the results shown in Fig. 8 and Fig. 9, with decreases in value causing slight increases in operating frequency, and with increases in value causing a drop in operating frequency. Finally,  $R_s$  has negligible impact on operating frequency, as one would expect

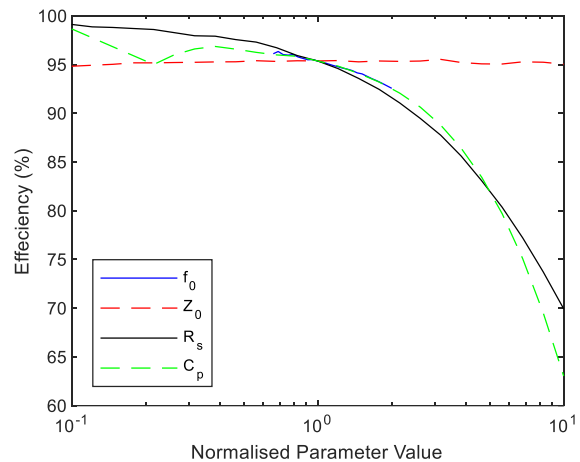


Fig. 10 – Efficiency sensitivity analysis

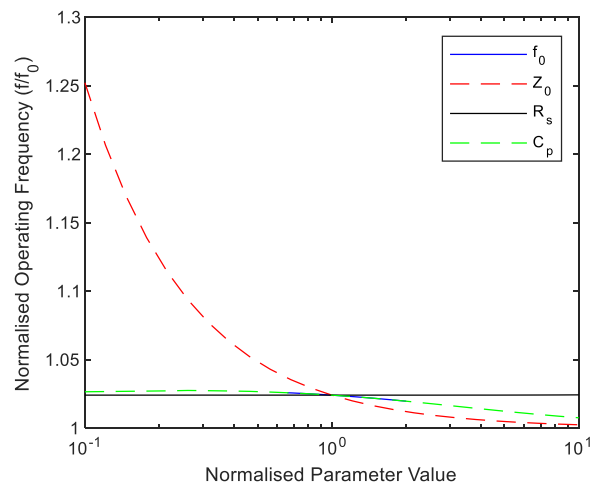


Fig. 11 – Operating frequency sensitivity analysis

## VI. DESIGN CONSIDERATIONS

Based on Fig. 9 and Fig. 10, care should be taken to choose an appropriate resonator for this converter. Ideally, the resonator should have no losses ( $R_s = 0$ ), large  $C_p$  and high resonant frequency, to maximise the output voltage gain that can be achieved. Whilst minimising  $R_s$  has no negative effects on the performance of the converter, choosing a large  $C_p$  and high resonant frequency can cause adverse effects. As Fig. 10 shows, increasing  $C_p$  or  $f_0$  causes the efficiency of the converter to decrease, as shown in (21) from [14]. For a practical resonator, a larger radius could be used to increase  $C_p$  for higher efficiency, however, as a result of the larger radius, the resonant frequency decreases (reducing the efficiency of the converter), resulting in an overall minimal change in efficiency. Therefore,  $C_p$  and  $f_0$  should be designed to be a balance between achieving higher output voltages and reducing efficiency.

In contrast,  $Z_0$  causes minimal impact of the performance of the converter, other than to change the operating frequency of the converter. Whilst this is not necessarily a negative, an increase in normalised operating frequency increases the chances of interacting with spurious modes, which could cause

a decrease in efficiency and reliability, and therefore should be avoided [13]. Additionally, low  $Z_0$  invalidates some of the analysis presented here and in common analysis techniques, therefore the design process becomes more complex as the effect of additional harmonics need to be accounted for.

Determining  $R_s$  of a resonator design is complex as it is a combination of several loss mechanisms occurring in piezoelectrics. However, some steps can be taken at the design stage to minimise  $R_s$ , such as using a high Q ('hard') piezoelectric material, achieving a high-quality construction with minimal defects, and carefully designing mounting mechanisms as to not damp the vibrational characteristics of the resonator.

The other elements of the converter should also be carefully chosen as not to impair performance. First, the output capacitance should be chosen to minimise ripple on the output voltage, with larger capacitances causing lower ripple but with the trade-off of larger size and longer initial charging times. The MOSFETs should be chosen to have very low ( $<1\Omega$ ) on state resistance (lowering losses) and with low  $C_{DS}$  capacitance, as this capacitance is in parallel with  $C_p$ , increasing the effective parallel capacitance of the resonator. Finally, diodes should ideally be Schottky diodes with low forward voltage drops and sufficiency high PIV rating. Ideally, diodes would be replaced with additional MOSFETs as in [14] but this increases control and drive complexity.

## VII. CONCLUSION

A model of the SUPR converter is presented, based on the describing function method. A full derivation of the model is shown, and a method for the solution of the operational mode durations is proposed. The model is validated against experimental and simulated results, with good accuracy. Further analysis is performed using this model to determine how the operational conditions (load and desired output voltage) affect the performance of the converter. Finally, a sensitivity analysis of the resonator parameters is presented, allowing an informed discussion on the design of the resonator for optimal performance.

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## IX. APPENDIX

The following is Python code for the generation of the  $v_{Cp}<1>$  equation given in (33) using the Sympy library:

```
from sympy import pi, cos, sin, symbols, integrate,
simplify, I

omega, Cp, Ils, theta, Vdf, Vout, Vdc, Rds = symbols('omega
Cp Ils theta Vdf Vout Vdc Rds')
theta1, theta2, theta3, theta4, theta5, theta6 =
symbols('theta1 theta2 theta3 theta4 theta5 theta6')

#define vcp
M1 = 1 / omega / Cp * Ils * cos(theta) + (Vdf * omega * Cp
+ Cp * Vout * omega - Ils) / omega / Cp
M2 = Vdc - Vdf - Ils * sin(theta) * Rds
M3 = 1 / omega / Cp * Ils * cos(theta) - (Ils *
sin(theta2) * Rds * omega * Cp - Vdc * omega * Cp + Vdf *
omega * Cp + Ils * cos(theta2)) / omega / Cp
M4 = -Ils * sin(theta) * Rds
M5 = 1 / omega / Cp * Ils * cos(theta) - Ils * (Cp *
sin(theta4) * Rds * omega + cos(theta4)) / omega / Cp
M6 = Vout+Vdf

#Define imaginary elements
vcp_1c = 1/pi*integrate(cos(theta)*M1, (theta, 0, theta1))
vcp_2c = 1/pi*integrate(cos(theta)*M2, (theta, theta1,
theta2))
vcp_3c = 1/pi*integrate(cos(theta)*M3, (theta, theta2, pi))
vcp_4c = 1/pi*integrate(cos(theta)*M4, (theta, pi, theta4))
vcp_5c = 1/pi*integrate(cos(theta)*M5, (theta, theta4,
theta5))
vcp_6c = 1/pi*integrate(cos(theta)*M6, (theta, theta5,
2*pi))
vcp_c = simplify(vcp_1c+vcp_2c+vcp_3c+vcp_4c+vcp_5c+vcp_6c)

#Define real elements
vcp_1s = 1/pi*integrate(sin(theta)*M1, (theta, 0, theta1))
vcp_2s = 1/pi*integrate(sin(theta)*M2, (theta, theta1,
theta2))
vcp_3s = 1/pi*integrate(sin(theta)*M3, (theta, theta2, pi))
vcp_4s = 1/pi*integrate(sin(theta)*M4, (theta, pi, theta4))
vcp_5s = 1/pi*integrate(sin(theta)*M5, (theta, theta4,
theta5))
vcp_6s = 1/pi*integrate(sin(theta)*M6, (theta, theta5,
2*pi))
vcp_s = simplify(vcp_1s+vcp_2s+vcp_3s+vcp_4s+vcp_5s+vcp_6s)

#Define vcp<1>
vcp_a=vcp_s+I*vcp_c
print(vcp_a)
```



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