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13 **Abstract**

14 Rock, soil and many porous-like materials are often fractured or structured media,
15 which can exhibit dual-porosity behaviour. Studies on solute transport in deformable
16 dual-porosity media remain challenging due to the multi-physics coupled effects and
17 the complex interaction between fracture (or macropore) and porous matrix. Though
18 several studies exist on constitutive modelling of coupled behaviour in deformable
19 dual-porosity, the previously developed models are not systematic in thermodynamical
20 frameworks. This paper proposes a Mixture Coupling Theory approach based on
21 nonequilibrium thermodynamics to develop the solute transport model with
22 consideration of hydro-mechanical coupling in dual-porosity media (referred to as the
23 ST-HM model). This paper derives the constitutive equations of fully hydro-
24 mechanical coupled behaviour in dual-porosity media and considers the pore and
25 fracture porosity evolution influenced by both hydro and mechanical fields. Therefore,
26 the governing equations of ST-HM are capable of predicting non-reactive solute
27 transport with a fully hydro-mechanical coupled effect in dual-porosity media. Then,
28 the model was verified against existing models and validated by relevant experimental
29 results. Further, a numerical example shows that the presented model significantly
30 improves the accuracy of the prediction of porosity, fluid pressure, and solute
31 concentration compared with previous models, which ignore the fully hydro-
32 mechanical coupled effects on solute transport.

33 **Keywords:** Solute transport; Deformable dual-porosity media; Coupled model;

34 Nonequilibrium thermodynamics; Mixture Coupling Theory

35

- 36 **List of notations**
- 37 B material coefficient
- 38 $c_{\alpha c}$ solute mass fraction
- 39 C solute volumetric concentration
- 40 \mathbf{d} displacement vector
- 41 d_i displacement in i axis
- 42 D^α dispersion - diffusion coefficient for solute
- 43 \mathbf{E} Green strain
- 44 \mathbf{F} deformation gradient
- 45 G shear modulus
- 46 H_{ij} material-dependent parameter
- 47 $\mathbf{I}^{\alpha\beta}$ mass flux
- 48 J Jacobian of \mathbf{F}
- 49 $\mathbf{J}^{\alpha\beta}$ diffusion flux
- 50 k_0^α initial permeability
- 51 k^α permeability
- 52 K bulk modulus of total geomaterial
- 53 K_s bulk modulus of solid grain
- 54 K_w bulk modulus of water
- 55 K_{pb} bulk modulus of the porous matrix block
- 56 L fracture spacing

- 57 L_{ijkl} material-dependent parameter
- 58 $m^{\alpha\beta}$ mass density in the reference configuration
- 59 M_{ij} material-dependent parameter
- 60 \mathbf{n} outward unit normal vector
- 61 N dimension of porous matrix block ($N=1,2,3$)
- 62 $p^{\alpha l}$ fluid pressure
- 63 Q material coefficient
- 64 r_{ex}^{β} exchange flux between the pore and fracture
- 65 S_{ij} material-dependent parameter
- 66 S boundary of an arbitrary region
- 67 t Time
- 68 \mathbf{T} second Piola-Kirchhoff stress
- 69 T temperature.
- 70 ν^{α} porosity in the reference configuration
- 71 ν_0^{α} initial porosity in the reference configuration
- 72 $\mathbf{u}^{\alpha l}$ Darcy velocity of fluid
- 73 ν Poisson's ratio
- 74 \mathbf{v}_s velocity of a solid.
- 75 V an arbitrary region in the current configuration
- 76 V_0 an arbitrary region in the reference configuration
- 77 W deformation energy (dual potential)

- 78 \mathbf{x} material point for one phase in an arbitrary current configuration
- 79 \mathbf{X} material point for one phase in an arbitrary reference configuration
- 80 Z material coefficients
- 81 **Greek symbols**
- 82 γ entropy production per unit volume
- 83 ε_{ij} strain tensor
- 84 ζ^α Biot constant
- 85 μ_w dynamic viscosity of water
- 86 $\mu^{\alpha\beta}$ chemical potential
- 87 $\nabla\mu_{ex}^\beta$ gradient of the chemical potential of exchange water or solute
- 88 $\rho^{\alpha\beta}$ mixture density
- 89 $\rho_t^{\alpha\beta}$ In-phase density (true density)
- 90 $\boldsymbol{\sigma}$ Cauchy stress
- 91 ϕ^α porosity in current configuration
- 92 ϕ^S volume fraction of the solid part
- 93 ψ Helmholtz free energy density
- 94 ψ_α Helmholtz free energy density of the bulk fluid
- 95 χ shape factor
- 96 τ diffusion transfer parameter
- 97 Ψ Helmholtz free energy density in the reference configuration
- 98 \mathcal{G}_{al} entropy production of fluid flow

99 \mathcal{G}_{ex} entropy production of fluid exchange between pore and fracture

100 **Subscripts and superscripts**

101 $\alpha = M$ For pore

102 $\alpha = F$ For fracture

103 $\beta = w$ For water

104 $\beta = c$ For solute

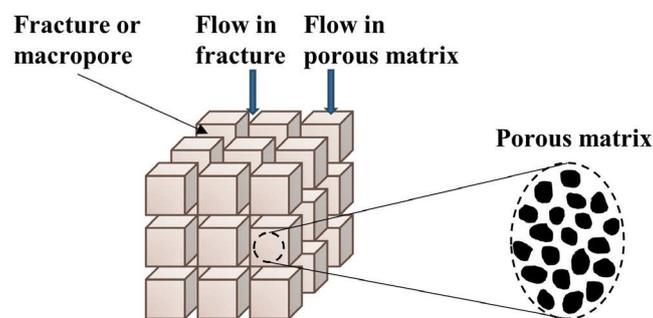
105 $\beta = l$ For liquid (solution)

106 **Introduction**

107 Rock, soil, and many porous-like materials are fractured or structured in reality and
108 exhibit very different characteristics than ideal porous media (Duguid and Lee, 1977,
109 Berre et al., 2019). For fractured and structured porous media, which contains both
110 macropores and porous matrix, the most distinctive properties compared to single
111 porosity media are the very heterogeneous distribution of porosity, permeability, and
112 the interaction between fracture (or macropore) and porous matrix. Here, fractured
113 porous media represents both fractured and structured media. Understanding the fluid
114 flow and solute transport in fractured porous media is very important in groundwater
115 remediation, agriculture engineering, and underground construction (Rutqvist and
116 Stephansson, 2003, Molson et al., 2012, Ray, 1994). Also, the exploitation of naturally
117 fractured reservoirs needs to study the fluid and solute flow in the fractured rock (Liu
118 et al., 2019, Savatorova et al., 2016). Geological carbon storage (GCS) technology often
119 experiences chemo-hydro-mechanical coupled processes in fractured rocks during
120 carbon sequestration (Akono et al., 2019).

121 So far, several methods have been considered to study the fluid and solute transport in
122 fractured porous media, including laboratory experiments (Leij et al., 2012), numerical
123 studies (Kodešova et al., 2006) and mathematical models (Berre et al., 2019). In the
124 previous studies, continuum models (Shu, 1999, Li et al., 2020, Saevik et al., 2013,
125 Dykhuisen, 1990) were used, this includes research on permeability tensor by Oda (Oda,
126 1985), one of the earliest studies, and other damage mechanics studies in fractured

127 media (De Bellis et al., 2017, De Bellis et al., 2016, Arson and Pereira, 2013). Also,
128 many researchers used discrete fracture models (Lee et al., 2001, Jiang and Younis,
129 2017) to represent fractured porous media. The dual-porosity model (also called the
130 double porosity model), one of the multi-continuum models, was firstly presented by
131 Barenblatt in 1960 s (Barenblatt et al., 1960). This theory assumes that fractured porous
132 media can be viewed as two overlapping domains, porous matrix and fracture (shown
133 in **Figure 1**); they are both continua but with different characteristics such as
134 permeability, porosity, bulk modulus and so on. In addition, the water and solute
135 exchange happens between fracture and porous matrix. However, the dual-porosity
136 model cannot present a geometric distinction between porous matrix and fracture, as it
137 employs a continuum to simplify the fracture network.



138

139 **Figure 1.** Conceptual description of the dual-porosity theory

140 A more realistic understanding of solute transport needs a further understanding of the
141 coupled model because, natural rocks always exist in a multi-physics coupled
142 environment. Biot (Biot, 1941) and Terzaghi (Terzaghi, 1943) are the pioneers of hydro-
143 mechanical coupled consolidation theory. Based on these theories, Khalili et al. studied
144 the Hydro-Mechanical and Thermo-Hydro-Mechanical coupling in deformable dual-

145 porosity media (Khalili, 2003, Gelet et al., 2012b, Khalili, 2008a). In addition, several
146 solute transport models (Murad and Moyne, 2008, Revil, 2017) and Hydro-Mechanical-
147 Chemical coupled models (Li et al., 2020, Musso et al., 2013, Vecchia and Musso, 2016)
148 in dual-porosity media were developed recently. However, these studies often ignore
149 the fully hydro-mechanical coupling (especially the porosity evolution induced by fluid
150 pressure), which may strongly influence solute transport. These models can be termed
151 as mechanics approach (most widely used by researchers since Terzhagi developed 1D
152 consolidation theory) models as their derivation is mainly based on mechanics and
153 momentum balance. Mechanics approach can easily develop new coupled models in by
154 introducing equations from other disciplines (e.g. adding chemical reaction equations
155 into solute transport equation), while lacking a systemic framework that covers
156 different fields (Chen et al., 2016).

157 On the other hand, the Mixture Theory, which adopts a different approach, was
158 developed from the study of biological tissues at the early stage (Bowen, 1976) but,
159 gained more and more attention from geological and environmental researchers in
160 recent years. In the Mixture Theory, each point in the mixture is assumed to contain
161 every component (for example, solid and fluid) (Bedford and Drumheller, 1983).
162 Several works have been carried out on coupled fluid transport and porous media
163 deformation (Drumheller, 1978, Bowen, 1984, Rajagopal and K., 2007, Hutter et al.,
164 2015).

165 Mixture Coupling Theory (originally named Modified Mixture Theory) is an emerging

166 approach developed by Heidug and Wong (Heidug and Wong, 1996). This theory is
167 based on nonequilibrium thermodynamics and Mixture Theory, which makes it a
168 energy-based theory. Mixture Coupling Theory has been used to derive the Hydro-
169 Mechanical model (Chen and Hicks, 2011, Chen, 2013), Hydro-Mechanical-Chemical
170 and Thermo-Hydro-Mechanical-Chemical model (Chen and Hicks, 2013, Ma et al.,
171 2021) in porous media. The Mixture Coupling Theory links different physical fields
172 using Helmholtz free energy and obtains the coupling relationships of stress, porosity
173 and other variables by deformation energy analysis.

174 Mixture theory uses the internal frictional interaction force to study information on the
175 interaction between different components, which is hard to obtain (Laloui et al., 2003).
176 In comparison, Mixture Coupling Theory uses entropy production to study the
177 interaction between each component, which makes it easier to consider molecule-scale
178 friction or other chemical-induced dissipative processes.

179 The presented paper used the Mixture Coupling Theory to derive a systematic solute
180 transport model with consideration of hydro-mechanical coupling in dual-porosity
181 media (referred to as the ST-HM model) based on nonequilibrium thermodynamics.

182 This model successfully determines the non-reactive solute transport behaviour with
183 fully hydro-mechanical coupling effects, and considers porosity and permeability
184 evolution influenced by solid deformation, porous/fracture fluid pressure, and the
185 coupled interaction between porous and fracture fluid flow. The new model is compared
186 with previous models derived by the mechanics approach and validated against

187 experimental data. A numerical simulation is further performed to investigate the
188 potential application of the proposed model in groundwater contamination research.

189 **Constitutive equation development**

190 *Balance equation of Helmholtz free energy and entropy production*

191 An arbitrary region V with a boundary S is selected for this to study the dual-porosity
192 medium. Water and solute fluxes can pass through the boundary while the solid phase
193 cannot; and only one solute is considered in this study. The local version balance
194 equation of Helmholtz free energy for the open system can be obtained in a dual-
195 porosity media as

196
$$\dot{\psi} + \psi \nabla \cdot \mathbf{v}^s - \nabla \cdot (\boldsymbol{\sigma} \mathbf{v}^s) + \nabla \cdot (\mu^{Mw} \mathbf{I}^{Mw} + \mu^{Mc} \mathbf{I}^{Mc}) + \nabla \cdot (\mu^{Fw} \mathbf{I}^{Fw} + \mu^{Fc} \mathbf{I}^{Fc}) = -T\gamma \leq 0$$

197 (1)

198 where ψ is the Helmholtz free energy density, $\boldsymbol{\sigma}$ is the Cauchy stress tensor, \mathbf{v}^s is
199 the velocity of a solid, μ^{Mw} , μ^{Fw} , μ^{Mc} , μ^{Fc} are the chemical potential of pore water,
200 fracture water, solute in porous flow, solute in fracture flow, respectively. \mathbf{I}^{Mw} , \mathbf{I}^{Mc} ,
201 \mathbf{I}^{Fw} , \mathbf{I}^{Fc} are the flux of pore water, porous solute, fracture water, and fracture solute,
202 respectively. T is the constant temperature and γ is the entropy production per unit
203 volume.

204 The detailed derivation of equation (1) is presented in Supplementary information.

205 It is assumed that only the friction generated at the solid-fluid (water and solute)
206 interface contributes to the energy dissipation in the mixture. By using the equation of
207 entropy production of water and solute flow (Katachalsky and Curran, 1965) and the
208 equation for fluid exchange between porous and fracture space (Gelet et al., 2012b,

209 Coussy, 2004), the entropy production of the dual-porosity media can be obtained as

$$210 \quad 0 \leq T\gamma = -\mathbf{I}^{Mw} \cdot \nabla \mu^{Mw} - \mathbf{I}^{Mc} \cdot \nabla \mu^{Mc} - \mathbf{I}^{Fw} \cdot \nabla \mu^{Fw} - \mathbf{I}^{Fc} \cdot \nabla \mu^{Fc} + r_{ex}^w (\mu^{Mw} - \mu^{Fw}) + r_{ex}^c (\mu^{Mc} - \mu^{Fc})$$

$$211 \quad (2)$$

212 where r_{ex}^w and r_{ex}^c are the exchange water and solute flux between the pore and

213 fracture. The chemical potential of exchange water changes from μ^{Mw} to μ^{Fw} when

214 the exchange is to fracture, from μ^{Fw} to μ^{Mw} vice versa.

215 *Basic equations of water and solute flow in dual-porosity media*

216 In this manuscript, both mixture density and In-phase (also called true density) density

217 are used. The mixture density can be obtained by the In-phase density of pore water

218 ρ_l^{Mw} , fracture water ρ_l^{Fw} , porous solute ρ_l^{Mc} and fracture solute ρ_l^{Fc} as

$$219 \quad \rho^{Mw} = \phi^M \rho_l^{Mw}, \rho^{Fw} = \phi^F \rho_l^{Fw}, \rho^{Mc} = \phi^M \rho_l^{Mc}, \rho^{Fc} = \phi^F \rho_l^{Fc} \quad (3)$$

220 where ϕ^M , ϕ^F are the porosity of porous matrix and fracture, and they obey

$$221 \quad \phi^M + \phi^F + \phi^S = 1 \quad (4)$$

222 where ϕ^S is the volume fraction of the solid phase.

223 In-phase density refers to the mass of a specified component per unit phase volume (for

224 example, the mass of solute per unit liquid volume). Here, using In-phase density

225 instead of true density avoids confusion since true density can easily be misunderstood

226 with the intrinsic density of one component, such as solid salt.

227 The In-phase density of the solution in the porous matrix and fracture is defined as

$$228 \quad \rho_l^{Ml} = \rho_l^{Mc} + \rho_l^{Mw}, \rho_l^{Fl} = \rho_l^{Fc} + \rho_l^{Fw} \quad (5)$$

229 The balance equation of fluid writes

$$230 \quad (\nu^M \rho_l^{Ml}) \cdot + \nabla \cdot (\rho_l^{Ml} \mathbf{u}^{Ml}) + r_{ex}^l = 0 \quad (6)$$

$$231 \quad (\nu^F \rho_l^{Fl}) \cdot + \nabla \cdot (\rho_l^{Fl} \mathbf{u}^{Fl}) - r_{ex}^l = 0 \quad (7)$$

232 where $\nu^M = J\phi^M$ and $\nu^F = J\phi^F$ are pore volume per unit referential volume and

233 fracture volume per unit referential volume, and $r_{ex}^l = r_{ex}^w + r_{ex}^c$ is the fluid mass

234 exchange between pore and fracture, \mathbf{u}^{Ml} and \mathbf{u}^{Fl} are the Darcy velocity of the porous

235 fluid and fracture fluid, J is the Jacobian of the deformation gradient \mathbf{F} .

236 Also, the balance equation of solute in porous matrix and fracture can be derived as

$$237 \quad (\nu^M \rho_l^{Ml}) \dot{c}_{Mc} + \rho_l^{Ml} \mathbf{u}^{Ml} \nabla c_{Mc} + r_{ex}^c + \nabla \cdot \mathbf{J}^{Mc} = 0 \quad (8)$$

$$238 \quad (\nu^F \rho_l^{Fl}) \dot{c}_{Fc} + \rho_l^{Fl} \mathbf{u}^{Fl} \nabla c_{Fc} - r_{ex}^c + \nabla \cdot \mathbf{J}^{Fc} = 0 \quad (9)$$

239 where $c_{Mc} = \frac{\rho_l^{Mc}}{\rho_l^{Ml}}$, $c_{Fc} = \frac{\rho_l^{Fc}}{\rho_l^{Fl}}$ are the solute mass fraction, \mathbf{J}^{Mc} , and \mathbf{J}^{Fc} are the

240 diffusion flux of porous solute and fracture solute.

241 The detailed derivation of the balance equation is presented in Supplementary

242 information.

243 *Basic equation of state*

244 Based on the assumption that the dual-porosity media maintain mechanical equilibrium

245 and without volume force, so $\nabla \cdot \boldsymbol{\sigma} = \mathbf{0}$. By using equation (1) and (2), it leads to

$$246 \quad \dot{\psi} + \psi \nabla \cdot \mathbf{v}^s - \nabla \cdot (\boldsymbol{\sigma} \mathbf{v}^s) + (\mu^{Mw} \nabla \cdot \mathbf{I}^{Mw} + \mu^{Mc} \nabla \cdot \mathbf{I}^{Mc}) + (\mu^{Fw} \nabla \cdot \mathbf{I}^{Fw} + \mu^{Fc} \nabla \cdot \mathbf{I}^{Fc}) + r_{ex}^w (\mu^{Mw} - \mu^{Fw})$$

$$+ r_{ex}^c (\mu^{Mc} - \mu^{Fc}) = 0$$

$$247 \quad (10)$$

248 Based on the mass balance equation and using equation (10), using finite strain theory
 249 in continuum mechanics, the free energy in the reference configuration can be obtained
 250 as

$$251 \quad \dot{\Psi} = tr(\mathbf{T}\dot{\mathbf{E}}) + \mu^{Mw} \dot{m}^{Mw} + \mu^{Fw} \dot{m}^{Fw} + \mu^{Mc} \dot{m}^{Mc} + \mu^{Fc} \dot{m}^{Fc} \quad (11)$$

252 where $\Psi = J\psi$, \mathbf{E} represents Green strain, \mathbf{T} is the second Piola-Kirchhoff stress,
 253 $m^{Mw} = J\rho^{Mw}$, $m^{Fw} = J\rho^{Fw}$, $m^{Mc} = J\rho^{Mc}$, $m^{Fc} = J\rho^{Fc}$ are the mass density of pore
 254 water, fracture water, porous solute, and fracture solute in the reference configuration.
 255 The detailed derivation of equation (11) is presented in Supplementary information.

256 *Helmholtz free energy density of the fluid and solid skeleton*

257 The derivation of the Helmholtz free energy density of the bulk fluid in porous matrix
 258 and fracture requires some of the following assumptions to be made. The fluid flow is
 259 isothermal, the equilibrium is transient (i.e. porosity can be viewed as a constant), and
 260 the chemical is non-reactive. The fluid is non-sorptive, and this is necessitated as
 261 Brochard and Honório points out that the Gibbs-Duhem equation does not hold well, in
 262 general, for an adsorbed fluid ([Brochard and Honório, 2020](#)). Based on the above
 263 assumptions, water/fluid in the pores and fractures is not strongly influenced by
 264 intermolecular and surface forces. So the Helmholtz free energy density of the bulk
 265 fluid in porous matrix and fracture can be derived as

$$266 \quad \psi_M = -p^{Ml} + (\mu^{Mw} \rho_l^{Mw} + \mu^{Mc} \rho_l^{Mc}) \quad (12)$$

$$267 \quad \psi_F = -p^{Fl} + (\mu^{Fw} \rho_l^{Fw} + \mu^{Fc} \rho_l^{Fc}) \quad (13)$$

268 where p^{Ml} is the porous fluid pressure and p^{Fl} is the fracture fluid pressure. Using
 269 Gibbs–Duhem equation, there are

$$270 \quad \dot{\psi}_M = \dot{\rho}_l^{Mw} \mu^{Mw} + \dot{\rho}_l^{Mc} \mu^{Mc} \quad (14)$$

$$271 \quad \dot{\psi}_F = \dot{\rho}_l^{Fw} \mu^{Fw} + \dot{\rho}_l^{Fc} \mu^{Fc} \quad (15)$$

272 Using equation (11), (12), (13), (14), (15) and equation (3), the free energy of the
 273 solid skeleton (including solid phase and water molecules attached to the solid's surface)
 274 can be described as:

$$275 \quad \left(\Psi - J\phi^M \psi_M - J\phi^F \psi_F \right) \dot{} = tr(\mathbf{T}\dot{\mathbf{E}}) + \dot{\nu}^M p^{Ml} + \dot{\nu}^F p^{Fl} \quad (16)$$

276 *Constitutive equations*

277 Subtracting the contribution of porous fluid pressure and fracture fluid pressure from
 278 Helmholtz free energy density of the solid skeleton, the deformation energy (also called
 279 dual potential) can be obtained

$$280 \quad W = \left(\Psi - J\phi^M \psi_M - J\phi^F \psi_F \right) - \nu^M p^{Ml} - \nu^F p^{Fl} \quad (17)$$

281 The time derivative of deformation energy can be obtained using equation (16) and
 282 (17) as

$$283 \quad \dot{W} = tr(\mathbf{T}\dot{\mathbf{E}}) - \nu^M \dot{p}^{Ml} - \nu^F \dot{p}^{Fl} \quad (18)$$

284 Equation (18) shows that W is a function of \mathbf{E} , p^{Ml} and p^{Fl} .

285 From equation (18), the expression of \mathbf{T} , ν^M and ν^F can be obtained as

$$286 \quad T_{ij} = \left(\frac{\partial W}{\partial E_{ij}} \right)_{p^{Ml}, p^{Fl}}, \nu^M = - \left(\frac{\partial W}{\partial p^{Ml}} \right)_{E_{ij}, p^{Fl}}, \nu^F = - \left(\frac{\partial W}{\partial p^{Fl}} \right)_{E_{ij}, p^{Ml}} \quad (19)$$

287 So that

$$288 \quad \dot{W}(\mathbf{E}, p^{Ml}, p^{Fl}) = \left(\frac{\partial W}{\partial E_{ij}} \right)_{p^{Ml}, p^{Fl}} \dot{E}_{ij} + \left(\frac{\partial W}{\partial p^{Ml}} \right)_{E_{ij}, p^{Fl}} \dot{p}^{Ml} + \left(\frac{\partial W}{\partial p^{Fl}} \right)_{E_{ij}, p^{Ml}} \dot{p}^{Fl}$$

289 (20)

290 By differentiating equation (19), and using equation (20), the basic constitutive
 291 equation for the evolution of stress, pore porosity and fracture porosity can be obtained
 292 as

$$293 \quad \dot{T}_{ij} = L_{ijkl} \dot{E}_{kl} - M_{ij} \dot{p}^{Ml} - S_{ij} \dot{p}^{Fl} \quad (21)$$

$$294 \quad \dot{v}^M = M_{ij} \dot{E}_{ij} + Q \dot{p}^{Ml} + B \dot{p}^{Fl} \quad (22)$$

$$295 \quad \dot{v}^F = S_{ij} \dot{E}_{ij} + B \dot{p}^{Ml} + Z \dot{p}^{Fl} \quad (23)$$

296 where the parameters L_{ijkl} , M_{ij} , S_{ij} , H_{ij} , B , Q , Z , defined by the following
 297 equations

298 Coupled field equations

299 *Assumptions in dual-porosity media*

300 Some simplifications and assumptions are made to simplify the further discussion.

- 301 1. Small strain condition is adopted in this paper, so the Green Strain tensor E_{ij} and
 302 Piola-Kirchhoff stress T_{ij} can be replaced by strain tensor ε_{ij} and Cauchy stress σ_{ij} .
- 303 2. The parameters L_{ijkl} , M_{ij} , S_{ij} , Z , B , Q are all material-dependent constants.
- 304 3. The rock is assumed to be isotropic, so that

$$305 \quad M_{ij} = \zeta^M \delta_{ij}, \quad S_{ij} = \zeta^F \delta_{ij}, \quad L_{ijkl} = G(\delta_{ik} \delta_{jl} + \delta_{il} \delta_{jk}) + (K - \frac{2G}{3}) \delta_{ij} \delta_{kl} \quad (24)$$

306 where G and K are the shear modulus and bulk modulus of the rock.

307 Similarly, the pore and fracture porosity equation (22) and (23) become

$$308 \quad \dot{\nu}^M = \zeta^M \dot{\varepsilon}_{ii} + Q\dot{p}^{Ml} + B\dot{p}^{Fl} \quad (25)$$

$$309 \quad \dot{\nu}^F = \zeta^F \dot{\varepsilon}_{ii} + B\dot{p}^{Ml} + Z\dot{p}^{Fl} \quad (26)$$

310 The evolution of pore and fracture permeability can be obtained using the Kozeny-

311 Carman function(Bitzer, 1996)

$$312 \quad k^M = k_0^M \left(\frac{\nu^M}{\nu_0^M} \right)^3 \left(\frac{1-\nu_0^M}{1-\nu^M} \right)^2 \quad (27)$$

$$313 \quad k^F = k_0^F \left(\frac{\nu^F}{\nu_0^F} \right)^3 \left(\frac{1-\nu_0^F}{1-\nu^F} \right)^2 \quad (28)$$

314 where k_0^M and k_0^F are the initial pore and fracture permeability, ν_0^M and ν_0^F are the
315 initial pore and fracture porosity.

316 *Multi-physics governing equations of solute transport in dual-porosity media*

317 Mechanical static equilibrium is assumed here, so $\partial\sigma_{ij}/\partial x_j = 0$, using displacement

318 variables $d_i (i=1,2,3)$ through $\varepsilon_{ij} = \frac{1}{2}(d_{i,j} + d_{j,i})$ and $\dot{\varepsilon}_{ii} = \nabla \cdot \dot{\mathbf{d}}$, the mechanical

319 governing equation can be described as

$$320 \quad G\nabla^2 \dot{\mathbf{d}} + \left(\frac{G}{1-2\nu} \right) \nabla(\nabla \cdot \dot{\mathbf{d}}) - \zeta^M \nabla \dot{p}^{Ml} - \zeta^F \nabla \dot{p}^{Fl} = 0 \quad (29)$$

321 where ν is Poisson's ratio.

322 If the influence of solute on fluid pressure is neglected, the variation of the In-phase

323 density of the porous fluid is

324
$$\dot{\rho}_i^{Ml} = \rho_i^{Ml} \left(\frac{1}{\rho_l^{Ml}} \frac{\partial \rho_l^{Ml}}{\partial p^{Ml}} \right) \frac{\partial p^{Ml}}{\partial t} = \rho_l^{Ml} \frac{1}{K_w} \dot{p}^{Ml} \quad (30)$$

325 where K_w is the bulk modulus of water.

326 Using the pore porosity equation (25), Darcy's law, and displacement \mathbf{d} , substituting
 327 equation (30) into equation (6), the balance equation of porous fluid (equation
 328 (6))becomes

329
$$\zeta^M \nabla \cdot \mathbf{d} + (Q + v^M \frac{1}{K_w}) \dot{p}^{Ml} + B \dot{p}^{Fl} + \frac{r_{ex}^l}{\rho_l^{Ml}} - \frac{k^M}{\mu_w} \nabla^2 p^{Ml} = 0 \quad (31)$$

330 where k^M is the permeability for the porous matrix, μ_w the water dynamic viscosity
 331 (normally related to water temperature), p^{Ml} the fluid pressure in the porous matrix.

332 Similar to the above steps of porous fluid, the governing equation for fracture fluid
 333 transport can be obtained as

334
$$\zeta^F \nabla \cdot \mathbf{d} + (Z + v^F \frac{1}{K_w}) \dot{p}^{Fl} + B \dot{p}^{Ml} - \frac{r_{ex}^l}{\rho_l^{Fl}} - \frac{k^F}{\mu_w} \nabla^2 p^{Fl} = 0 \quad (32)$$

335 where k^F is the permeability for fracture, p^{Fl} is the fluid pressure in fracture.

336 Using Darcy's law and Fick's law in the balance equation of porous solute transport,
 337 the governing equation for porous solute transport is

338
$$(v^M \rho_l^{Ml}) \dot{c}_{Mc} - \rho_l^{Ml} \frac{k^M}{\mu_w} \nabla p^{Ml} \nabla c_{Mc} + r_{ex}^c + \nabla \cdot (\rho_l^{Ml} D^M \nabla c_{Mc}) = 0 \quad (33)$$

339 Similarly, the governing equation for fracture solute transport can be obtained as

340
$$(v^F \rho_l^{Fl}) \dot{c}_{Fc} - \rho_l^{Fl} \frac{k^F}{\mu_w} \nabla p^{Fl} \nabla c_{Fc} - r_{ex}^c + \nabla \cdot (\rho_l^{Fl} D^F \nabla c_{Fc}) = 0 \quad (34)$$

341 where D^M , D^F are the dispersion -diffusion coefficient for the solute in porous matrix

342 and fracture.

343 The transfer term has been studied for a long time since 1960 s (Barenblatt et al., 1960,
344 Khalili, 2003). Here a linear transfer equation is adopted, which can be described as:

$$345 \quad \frac{r_{ex}^l}{\rho_l^{\alpha l}} = \frac{\chi k^M}{\mu_w} (p^{Ml} - p^{Fl}) \quad (35)$$

346 where the χ means shape factor, $\rho_l^{\alpha l}$ is the In-phase density of porous or fracture
347 fluid. χ can be calculated as (Warren and Root, 1963)

$$348 \quad \chi = \frac{4N(N+2)}{L^2} \quad (36)$$

349 where N denotes the dimension of porous matrix block ($N=1,2,3$), L is the fracture
350 spacing.

351 The transfer term of the solute can be divided into two parts, convection and diffusion
352 (Dykhuizen, 1987)

$$353 \quad \frac{r_{ex}^c}{\rho_l^{\alpha l}} = \frac{\chi k^M}{\mu_w} (p^{Ml} - p^{Fl}) (c_{Mc} - c_{Fc}) + \tau (c_{Mc} - c_{Fc}) \quad (37)$$

354 where τ is the diffusion transfer parameter.

355 *Parameter identification*

356 The parameters in the governing equations were identified by MA et al. (2023). Below
357 are the expressions of these parameters

$$358 \quad \zeta^F = 1 - \frac{K}{K_{pb}} \quad (38)$$

359 where K_{pb} denotes the bulk modulus of the porous matrix block (Valliappan and

360 Khalili-Naghadeh, 1990).

$$361 \quad \zeta^M = \frac{K}{K_{pb}} - \frac{K}{K_s} \quad (39)$$

362 where K_s is the bulk modulus of solid grain.

$$363 \quad Z = \frac{1}{K_{pb}} \left(1 - \phi^F - \frac{K}{K_{pb}} \right) \quad (40)$$

$$364 \quad B = \left(1 - \phi^F - \frac{K}{K_{pb}} \right) \left(\frac{1}{K_s} - \frac{1}{K_{pb}} \right) \quad (41)$$

$$365 \quad Q = \left(\frac{1}{K_{pb}} - \frac{1}{K_s} \right) \left(1 - \phi^F - \frac{K}{K_{pb}} + \frac{K}{K_s} \right) - \frac{\phi^M}{K_s} \quad (42)$$

366 These parameters can be determined directly in the laboratory test.

367 **Theoretical verification**

368 The mathematical model in the present paper (from equation (29) to (34)), obtains the
369 effects of mechanical deformation and fluid flow on solute transport in dual-porosity
370 media. The constitutive equation of hydro-mechanical coupling is the same as research
371 from the mechanics approach (Khalili, 2003, Khalili, 2008b), which mainly focus on
372 the coupling between fluid flow (including multiphase flow) and solid deformation,
373 without exploration of the influence of hydro-mechanical coupling on solute transport.
374 Equation (25) and (26) obtain the porosity evolution in dual-porosity media; which is
375 determined by solid deformation, both pore and fracture pressure change. If Equation
376 (25) and (26) are ignored, as well as the mechanical part (equation (29)), the ST-HM
377 model can degenerate to the classical Dual-Permeability Model proposed by (Simunek

378 and van Genuchten, 2008).

379 The fluid transfer term is consistent with previous research (Khalili, 2003, Zhao and
380 Chen, 2006). The solute transfer term is consistent with solute transport research in
381 dual-porosity media (Dykhuizen, 1987, Li et al., 2020, Gerke and Vangenuchten, 1993).

382 If it ignores the solute diffusion between pore and fracture, it converts to a pure
383 convection term (Sharma et al., 2021). It converts to a pure concentration-depended
384 term (Leij et al., 2012) if it ignores the convection between pore and fracture.

385 Musso and Li (Musso et al., 2013, Li et al., 2020) studied HMC coupling in swelling
386 clay. Compared with the solute transport part in their model (i.e. ignoring swelling), the
387 ST-HM model considers the fully coupled effects between pore and fracture fluid flow
388 (equation (31) and (32)).

389 **Experimental validation of the coupled model**

390 *Validation strategy*

391 The presented model is validated by two existing laboratory tests. The first simulation
392 is used to validate the solute transport behaviour in dual-porosity media, while the
393 second one is used to validate the dual porosity Hydro-Mechanical coupling part.

394 The numerical models based on the equations proposed in this paper were used to
395 simulate the experiments. The initial parameter's value was adopted by the experimental
396 research, the simulated results were compared with experimental data for model
397 validation. During the optimization, the parameters were changed to different values,

398 and the final optimized parameter values were obtained when the modelling results
399 fitted well with the experimental data.

400 *Validation of the solute transport behaviour*

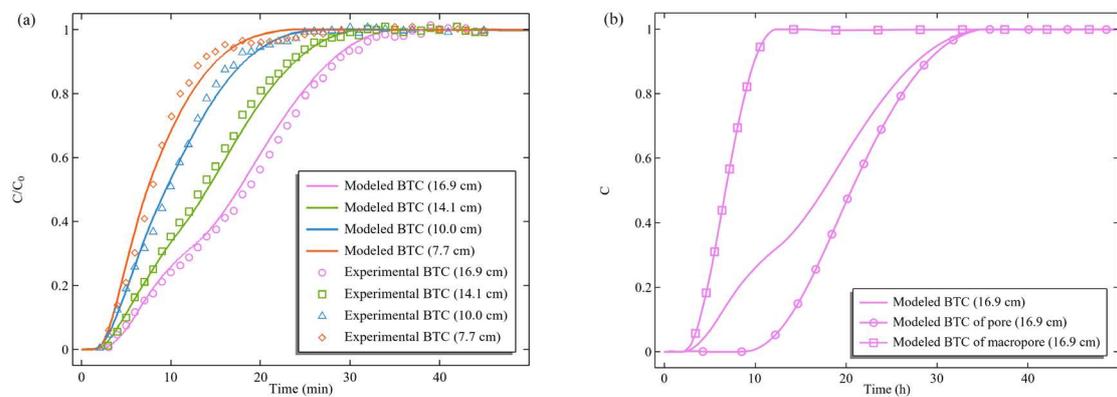
401 The experimental information is detailed in ([Leij et al., 2012](#)). The details of this
402 experiment and how the model was calibrated are provided in the supplementary
403 material. At first, the soil column which contains both macropore and pore was
404 saturated, and a steady flow rate of 1397.5 cm/d was established. The solute
405 concentration (CaCl₂) of the inflow fluid C_i was 60 mol/m³. Then the concentration
406 (CaCl₂) of inflow fluid C_i was changed from 60 mol/m³ to 90 mol/m³, and the
407 concentration at each probe was measured. The relative concentration ratio $\frac{C}{C_0}$ was
408 calculated, C means the relative concentration (measured concentration minus the
409 initial concentration 60 mol/m³) and C_0 the concentration step (30 mol/m³).

410 Note that the soil is dual-porosity media. So, the measured concentration was the
411 average volumetric concentration, which was calculated as

$$412 \quad C_{av} = \frac{C_{Mc} \times v^M + C_{Fc} \times v^F}{v^M + v^F} \quad (43)$$

413 A numerical model was set up in COMSOL MultiphysicsTM software ([AB](#)) based on
414 the experiment. Mechanical deformation was ignored in this simulation, and no
415 adsorption was considered (details of the numerical model setting are provided in the
416 supplementary materials). All simulation parameters are listed in **Table 1**.

417 The modelled and experimental BTC data are plotted together in **Figure 2**. The
 418 calculation error between the model results and experimental data is presented in
 419 Supplementary information.



420
 421 **Figure 2.** Experimental and modelled BTC at depths of 7.7, 10.0, 14.1, and 16.9 cm
 422 in (a). Modelled BTC for volume average, pore and macropore concentration at a
 423 depth of 16.9 cm in (b).

424 **Figure 2a** shows that the modelled results are in very close agreement with the
 425 experimental data. **Figure 2b** shows the modelled BTC of volume average, pore and
 426 macropore concentration at 16.9 cm. The peak of macropore BTC appears at about 12
 427 h, much ahead of the peak of pore BTC (about 35 h), which leads to bimodal transport.

428 *Validation of hydro-mechanical coupling behaviour*

429 The experimental information is detailed in Callari's research (Callari and Federico,
 430 2000). A artificial dual porosity media experienced a 1D consolidation process.
 431 Drainage of water during the consolidation was allowed through the upper boundary.
 432 The preliminary compaction was conducted under an initial effective pressure of 50
 433 kPa; three successive load increments were applied. Here the third load increment (from

434 200 KPa to 400 KPa) is considered.

435 The details of the numerical model and how the model was calibrated are provided in

436 the supplementary materials. All simulation parameters are listed in **Table 2**. The

437 modelled and experimental excess fluid pressure data are plotted together in **Figure 3**.

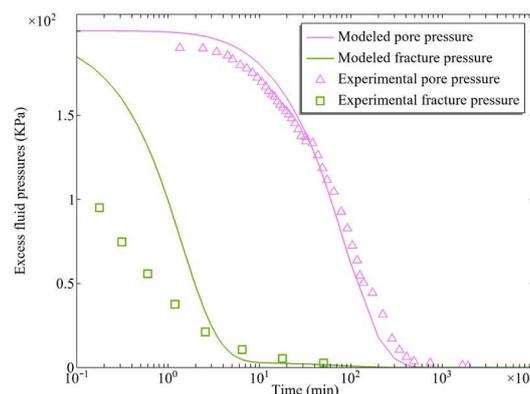
438 The modelled results fit well with the experimental one, except in the very early stage

439 (less than 2 minutes). The measurement of pressure in the geotextile is quite difficult

440 due to the high rate of pressure dissipation in the geotextile ([Callari and Federico, 2000](#)).

441 Both modelled and experimental results show that the fluid pressures in clayey block

442 (porous space) and geotextile (fracture space) are quite different.



443

444 **Figure 3.** Experimental and modelled excess fluid pressure at the base of the cell

445

versus time

446 **Numerical case setup**

447 Two numerical models were set up to further illustrate the solute transport behaviour in

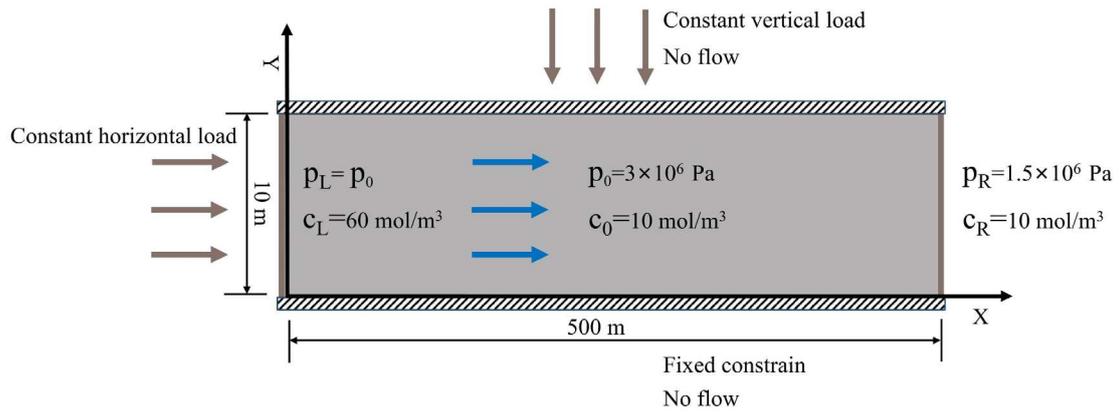
448 deformable dual-porosity. The first model used the ST-HM equations, and the second

449 used the equations in the research of Zhenze Li ([Li et al., 2020](#)), hereinafter referred to

450 as the Zhenze model. Zhenze model deals with coupled hydro-mechanical-chemical

451 (HMC) behaviour in dual-porosity media, considering osmosis flow and swelling,
452 which can interpret the swelling behaviour of MX-80 when infiltrated with brine.
453 Osmosis flow and swelling were ignored in the presented paper, so only the non-
454 reactive solute transport part in Zhenze model is compared with ST-HM. This
455 comparison investigates the fully coupling effects between micropore and macropore
456 flow on the solute transport and porosity evolution, which is not considered in Zhenze
457 model.

458 As shown in **Figure 4**, the initial fluid pressure was 3×10^6 Pa, and the initial solute
459 concentration was 10 mol/m^3 . The bottom boundary was a fixed constraint and no flow
460 boundary. The overburden applies a constant vertical load of 4×10^6 Pa on the rock,
461 with no flow crossing this boundary. The left boundary was constant horizontal load
462 2×10^6 Pa, and the right boundary was fixed in the X direction. In the simulation, the
463 fluid pressure on the left boundary remained 3×10^6 Pa while the fluid pressure on the
464 right boundary decreased to 1.5×10^6 Pa due to external human activity. The solute
465 concentration on the left boundary changed to 60 mol/m^3 while the concentration on
466 the right remained at 10 mol/m^3 . The boundary conditions are referenced from
467 (Nguyen, 2010).



468

469 **Figure 4.** Model description and boundary conditions for an assumed aquifer

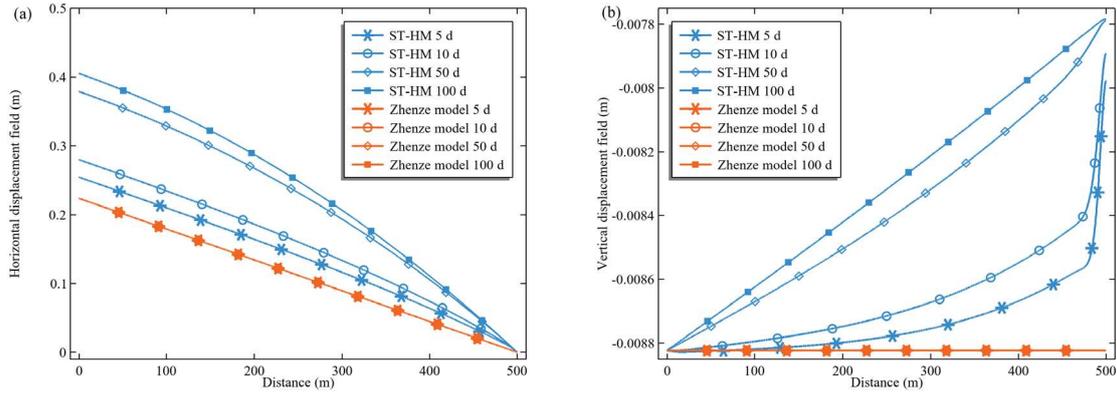
470 The parameters adopted in these two simulations are the same and are listed in **Table 3.**

471 Parameters are from the research of ([Gelet et al., 2012a](#)) and ([Abousleiman and Nguyen,](#)
472 [2005](#)).

473 **Numerical case results**

474 *Mechanical deformation in dual-porosity media*

475 The effects of fluid pressure change on solid deformation are well predicted by the ST-
476 HM model. As shown in **Figure 5**, the rock consolidates in the simulation. The
477 displacement of the ST-HM model increases with time, finally reaching almost a steady
478 value at about 100 d. In contrast, the displacement of Zhenze model does not change
479 during the simulation. The reason is that fluid pressure does not influence solid
480 deformation in Zhenze model. As the fluid pressure decreases in the rock, it leads to
481 further consolidation until the fluid pressure reaches an equilibrium stage. Deformation
482 is related to the evolution of fluid pressure (equation (29)), so the increased
483 deformation rates slow with time (as the decreased rates of fluid pressure become slow).



484

485

Figure 5. Horizontal and vertical displacement in X direction of rock

486

Fluid flow coupling in dual-porosity media

487

Figure 6 shows the pressure evolution in the pore and fracture of the rock. The fluid

488

pressure in both pores and fractures decreases gradually from the right side. Fluid

489

pressure in fracture changes much faster than in pore due to the much higher

490

permeability of fracture. Comparing the results of these two models, the fluid pressure

491

of the Zhenze model in both pore and fracture decreases faster than the fluid pressure

492

of the ST-HM model. It can be attributed to the coupled effect between pore and fracture

493

flow which is not considered in the Zhenze model. In equation (31) and (32), the

494

coupling term $Q\dot{p}^{Ml} + B\dot{p}^{Fl}$ in the porous fluid equation and $Z\dot{p}^{Fl} + B\dot{p}^{Ml}$ in fracture

495

fluid equation is one of the novelties of ST-HM model. Since pore and fracture pressure

496

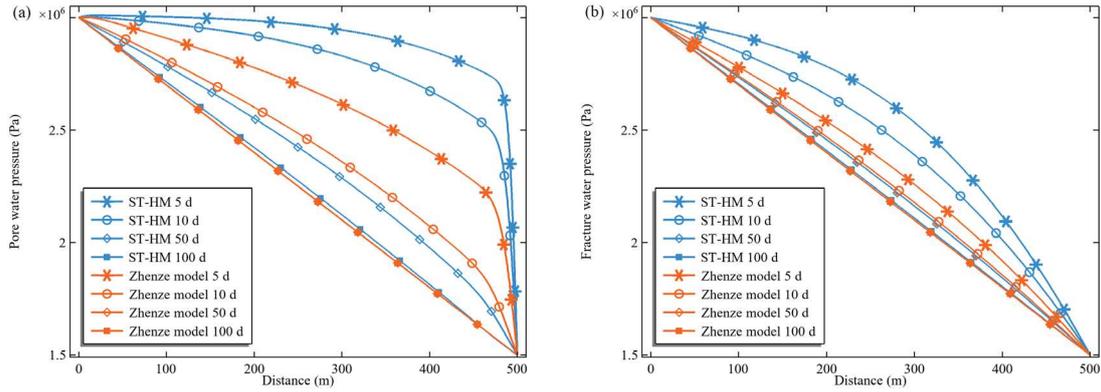
decrease in the simulation domain, these two coupling terms contribute to a lower value

497

of fluid pressure gradient (means lower water flow rates and slower decreased rates of

498

fluid pressure) than the Zhenze model.



499

500

Figure 6. Pore and fracture fluid pressure evolution on X direction of rock

501

Porosity evolution in dual-porosity media

502

Figure 7 shows the porosity evolution in pore and fracture in the X direction of rock.

503

The pore porosity and fracture porosity all decrease in the consolidation process.

504

However, the ST-HM model and the Zhenze model give different predictions for the

505

evolution trend. In the the Zhenze model, porosity decreases from the initial value

506

instantaneously and remains constant throughout the simulation because this model

507

does not consider the influence of fluid pressure on porosity. According to equation (25)

508

and (26), the ST-HM model considers the influence of both pore pressure and fracture

509

pressure on porosity, so the porosity decreases gradually from the right side to the left

510

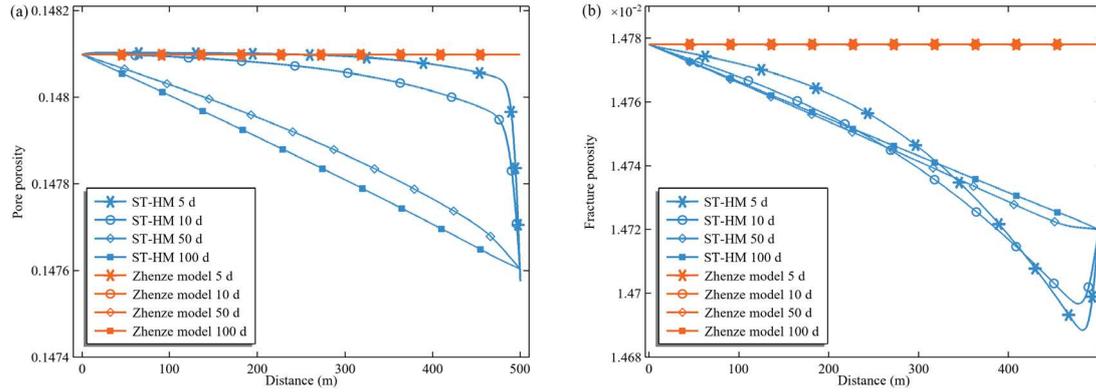
side (as the pressure change from the right side to the left side). Compared with the

511

Zhenze model, ST-HM model can better predict the long-term porosity change induced

512

by fluid pressure change.



513

514

Figure 7. Pore and fracture porosity evolution on X direction of rock

515

Note that the fracture porosity at the right part reaches the lowest value in the early

516

stage (less than 100 d) and subsequently increases to a higher value, while the left part

517

shows the opposite trend. This can be explained based on the effects of pore pressure

518

and fracture pressure on fracture porosity. According to equation (26), the fracture

519

porosity is determined by solid strain, both pore and fracture pressure. In this simulation,

520

$B = -6.65 \times 10^{-11}$ and $Z = 6.95 \times 10^{-11}$, which means that pore pressure change and

521

fracture pressure change affect fracture porosity oppositely in this case. Between 5 to

522

100 days \dot{p}^{Ml} is much lower than \dot{p}^{Fl} at the right part of the simulation domain

523

(because the response of fracture pressure is much quicker than pore pressure), which

524

leads to an increase of fracture porosity at the right. While at the left part of the

525

simulation domain, the difference of \dot{p}^{Ml} and \dot{p}^{Fl} is not so high to overcome the

526

total decreased trend of fracture porosity.

527

The porosity change is induced by solid deformation, pore and fracture pressure. As

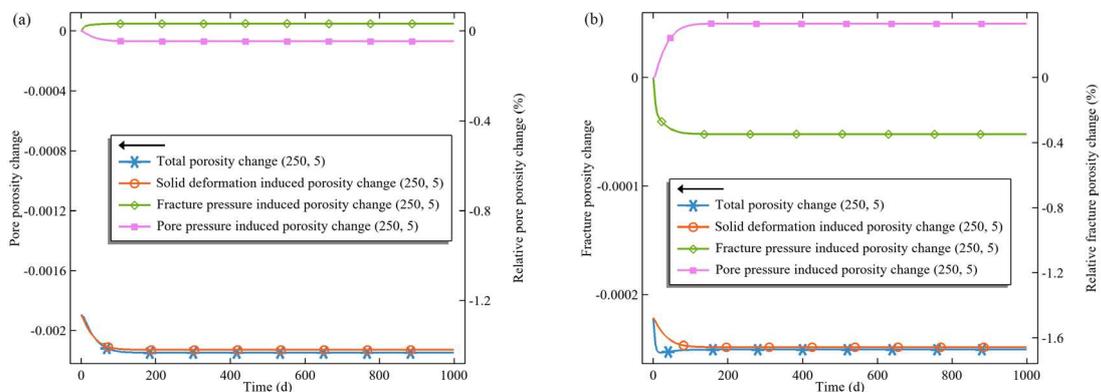
528

shown in **Figure 8**, the maximum pore porosity change is about -0.0022 (accounts for

529

1.46% of the initial pore porosity). In comparison, the maximum fracture porosity

530 change is about -0.00025 (accounts for 1.67% of the initial fracture porosity). The main
 531 contribution of porosity change comes from solid deformation. Remarkably, in dual-
 532 porosity media, with the decrease of fluid pressure in both pore and fracture, the
 533 porosity responds to both pore and fracture pressure. For pore porosity, the decrease of
 534 pore pressure induces consolidation and the decrease of pore porosity. At the same time,
 535 the decrease of fracture pressure induces the expansion of pore water which increase
 536 the pore porosity. Similarly, the decrease of pore pressure induces fracture water
 537 expansion for fracture porosity, while the decrease of fracture pressure induces
 538 consolidation. As the permeability of fracture is much higher than pore, the fracture
 539 pressure reaches its equilibrium much earlier, so the influence of the pore pressure on
 540 pore and fracture porosity last longer than fracture pressure.



541

542 **Figure 8.** Pore and fracture porosity evolution at the centre of the rock (X=250,

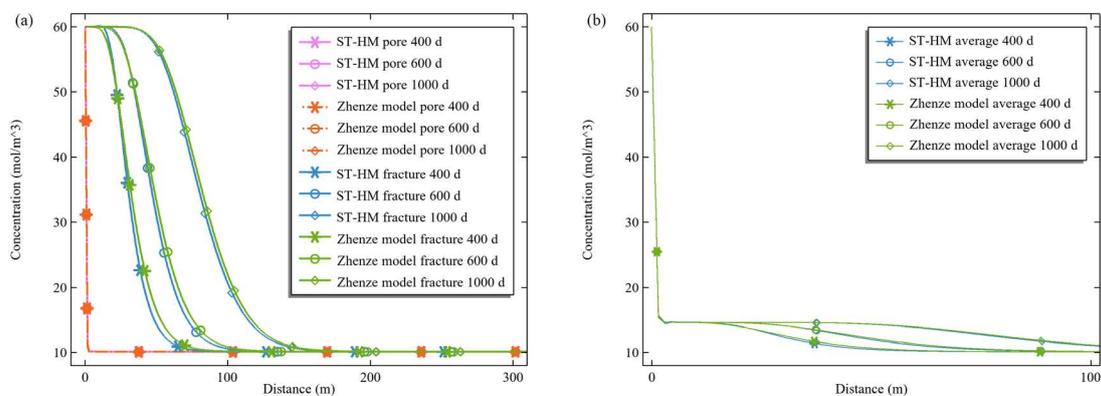
543 Y=5)

544 *Effects of hydro-mechanical coupling on solute transport*

545 The solute concentration evolution in the X direction of rock is illustrated in **Figure 9**.

546 The solute transport is quite slow due to the low permeability of the rock. The transport

547 rates in the fracture are much higher than in the pore. That is, the transport in fracture
548 dominates the transport phenomenon in this simulation. For the ST-HM model, it gives
549 a prediction that the solute transport front reached approx. 180 m after 1000 days, while
550 for Zhenze model this value is approx. 210 m. This difference is due to the
551 overestimation of porosity in the Zhenze model, which also overestimates permeability
552 and diffusivity. Hence, in most parts of the simulation domain, the concentration values
553 predicted by Zhenze model are higher than the ST-HM model. The ST-HM model
554 reduces about 16.67 % error for the prediction of solute transport front after 1000 days,
555 considering the influence of fluid pressure on porosity and the coupled effect between
556 pore and fracture flow. After 1000 days, the overestimation of fracture solute
557 consideration is 0.41 mol/m³ (accounting for 0.73 %) at 50 m, 2.85 mol/m³ (accounting
558 for 9.71 %) at 85 m, 2.26 mol/m³ (accounting for 12.39 %) at 103 m. The increased rate
559 of the difference in solute transport font location between ST-HM and Zhenze model
560 becomes slow with time (because the pressure gradient and concentration gradient
561 decrease throughout the simulation).



562

563

Figure 9. Solute concentration evolution on the X direction of rock (fracture

564 concentration, pore concentration and average volumetric concentration)

565 **Discussion of advantages and limitations**

566 This study presents a novel method based on nonequilibrium thermodynamics rather
567 than the traditional mechanics approach. Since it is an energy-based approach, it is easy
568 to bridge different behaviour and extendable for other processes such as thermo-hydro-
569 mechanical-chemical and bio-hydro-mechanical-chemical coupling.

570 The governing equation in this paper is based on the assumption of small deformation
571 and isotropic, these equations for hydro and mechanical fields are similar to research
572 from the mechanics approach ([Khalili, 2003](#), [Khalili, 2008b](#)). If these assumptions are
573 released, more general governing equations are likely to be obtained, which shows an
574 advantage of mixture-coupling theory over the classical mechanics approach on
575 deriving complicated multi-physics in a systematic way.

576 The consideration of hydro-mechanical coupling on solute transport in dual-porosity
577 media is an advancement to the classical Dual-Permeability Model proposed by
578 [Simunek and van Genuchten \(2008\)](#). This model captures the fully coupling effects of
579 pore and fracture flow with mechanical deformation. By engaging this effect, the solid
580 stress and the porosity change can be better predicted than the classical model. Since
581 ignoring the fully coupling effects of pore and fracture fluid can induce significant error,
582 especially for long-term prediction, the ST-HM model can significantly reduce the error
583 for the prediction of the solute transport front compared with the equations by [Li et al.](#)
584 [\(2020\)](#).

585 There are also some limitations existed in this model. The potential effects of exchange
586 on the transport phenomenon in pore and fracture were not studied. In addition, high-
587 velocity flow would affect the dissipation process in porous media (Takhanov, 2011).
588 So, the modification of general Darcy's law in dual-porosity media is needed to
589 implement these limitations.

590 Finally, the basic assumption of the dual-porosity theory is that the different spaces in
591 the media and the exchange phenomenon between those spaces can be captured by the
592 equivalent continuous parameters and variables (Berre et al., 2019). The proposed
593 model is not capable to solve solute transport in some fractured or structured media
594 with insignificant dual porosity characteristics.

595 **Conclusion**

596 This research extends the Mixture Coupling Theory and derives the constitutive
597 equations for non-reactive solute transport in deformable dual-porosity media (ST-HM),
598 based on nonequilibrium thermodynamics. The final governing equations of ST-HM
599 enable the prediction of non-reactive solute transport considering fully hydro-
600 mechanical coupling, which was often ignored by existing research.

601 A conceptual case was simulated to illustrate the major improvement of the presented
602 model. Simulated results show that the presented model provides a better prediction
603 than a previous model on the porosity evolution, fluid pressure and solute concentration
604 due to the consideration of fully coupling effects between pore and fracture flow. The
605 ST-HM model reduces about 16.67 % error for the prediction of solute transport front

606 after 1000 days and reduces up to 12.39 % error for the prediction of fracture
607 concentration.

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- 758

759 **Table 1** *Parameters used in the numerical model against solute transport experiment*

Parameter	Value	Meaning
ρ	$8.5 \times 10^2 \text{ kg} \cdot \text{m}^{-3}$	bulk density of Andisol
μ_w	$8 \times 10^{-4} \text{ Pa} \cdot \text{s}$	Viscosity of water
ν^M	0.41	Initial pore porosity
ν^F	0.16	Initial macro pore porosity
k^M	$2.75 \times 10^{-9} \text{ m}^2$	Initial pore permeability
k^F	$8.5 \times 10^{-9} \text{ m}^2$	Initial macro pore permeability
K_w	4.3 GPa	Compressibility of water
L	$1 \times 10^{-2} \text{ m}$	Macro pore spacing
D^M	$7 \times 10^{-5} \text{ m}^2 \cdot \text{s}^{-1}$	Pore Hydrodynamic dispersion coefficient
D^F	$2.5 \times 10^{-4} \text{ m}^2 \cdot \text{s}^{-1}$	Macro pore Hydrodynamic dispersion coefficient
τ	$1 \times 10^{-8} \text{ s}^{-1}$	Coefficient of diffusive transfer
T	70 min	Calculation period
Step	0.1 min	Timestep

760 **Table 2** *Parameters used in the numerical model against hydro-mechanical coupling*
 761 *experiment*

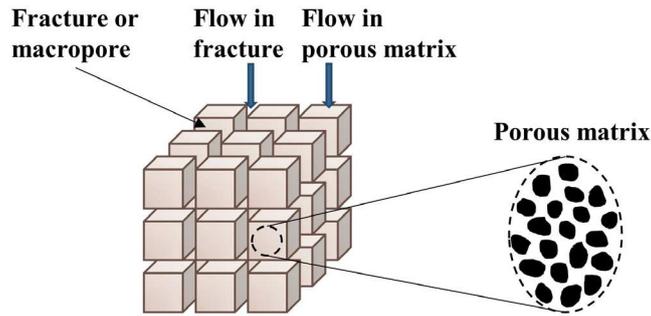
Parameter	Value	Meaning
ρ	$2.5 \times 10^3 \text{ kg} \cdot \text{m}^{-3}$	bulk density of soil
μ_w	$8 \times 10^{-4} \text{ Pa} \cdot \text{s}$	Viscosity of water
ν	0.3	Poisson's ratio
K	3.2 MPa	Bulk modulus of total rock
K_{pb}	6.6 MPa	Bulk modulus of porous block
K_s	27.5 GPa	Bulk modulus of solid grain
ν^M	0.55	Initial pore porosity
ν^F	0.008	Initial fracture porosity

k^M	$4.8 \times 10^{-18} \text{ m}^2$	Initial pore permeability
k^F	$4 \times 10^{-11} \text{ m}^2$	Initial fracture permeability
K_w	4.3 GPa	Compressibility of water
L	$6 \times 10^{-3} \text{ m}$	Fracture spacing
T	10^4 min	Calculation period
Step	0.01 min	Timestep

762 **Table 3** Parameters used in the numerical simulation for discussion of solute

763 *transport behaviour in deformable dual-porosity*

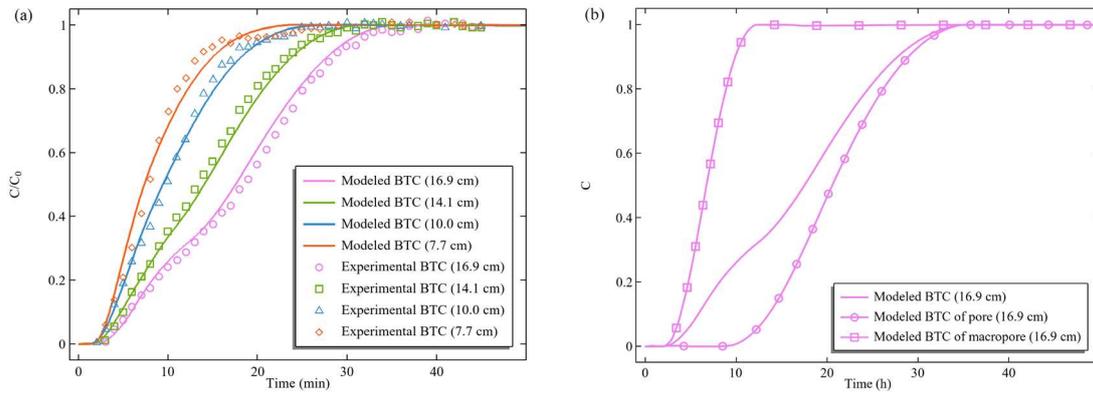
Parameter	Value	Meaning
K	1.103 GPa	Bulk modulus of rock
K_{pb}	1.226 GPa	Bulk modulus of porous block
K_s	27.5 GPa	Bulk modulus of solid grain
ν	0.22	Poisson's ratio
μ_w	$8 \times 10^{-4} \text{ Pa}\cdot\text{s}$	Viscosity of water
ν^M	0.15	Initial pore porosity
ν^F	0.015	Initial fracture porosity
k^M	$5 \times 10^{-18} \text{ m}^2$	Initial pore permeability
k^F	$5 \times 10^{-15} \text{ m}^2$	Initial fracture permeability
K_w	4.3 GPa	Compressibility of water
D^M	$9 \times 10^{-12} \text{ m}^2\cdot\text{s}^{-1}$	Pore Hydrodynamic dispersion coefficient
D^F	$7 \times 10^{-11} \text{ m}^2\cdot\text{s}^{-1}$	Fracture Hydrodynamic dispersion coefficient
τ	$1 \times 10^{-11} \text{ s}^{-1}$	Coefficient of diffusive transfer
L	2 m	Fracture spacing
T	1000 d	Calculation period
Step	0.1 d	Timestep



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Figure 1. Conceptual description of the dual-porosity theory



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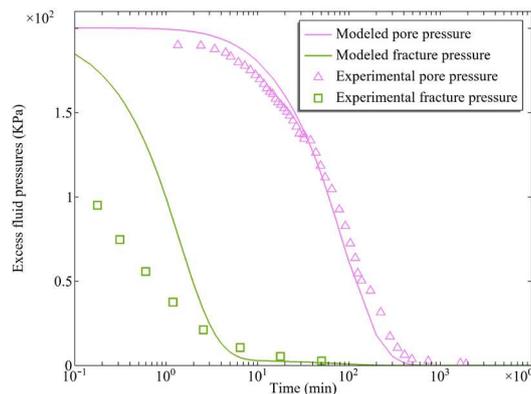
Figure 2. Experimental and modelled BTC at depths of 7.7, 10.0, 14.1, and 16.9 cm

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in (a). Modelled BTC for volume average, pore and macropore concentration at depth

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16.9 cm in (b).



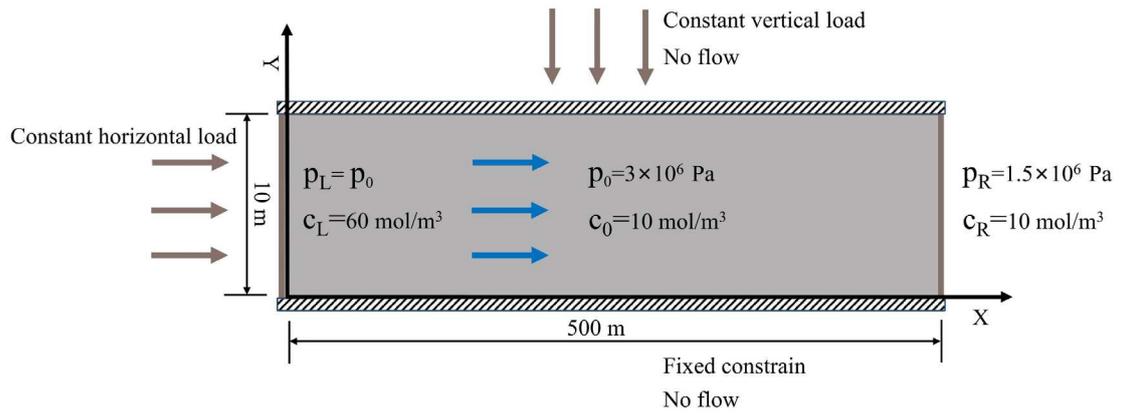
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Figure 3. Experimental and modelled excess fluid pressure at the base of the cell

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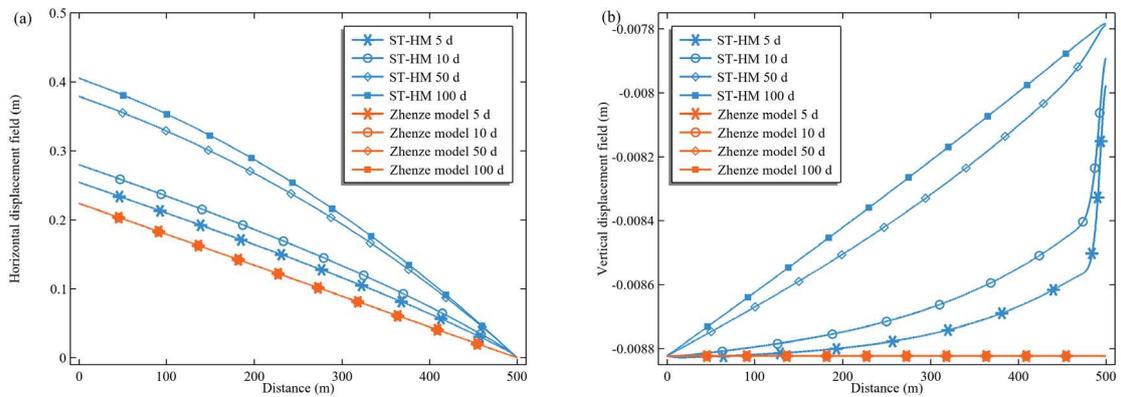
versus time



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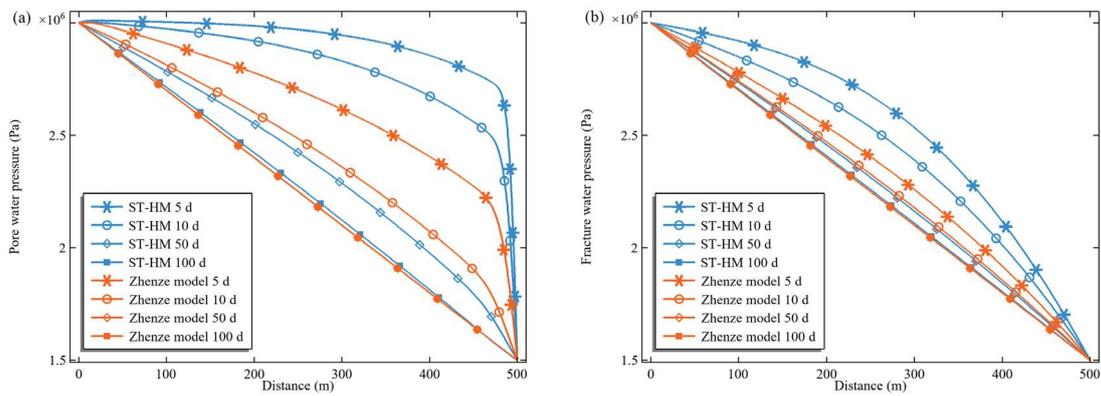
Figure 4. Model description and boundary conditions for an assumed aquifer



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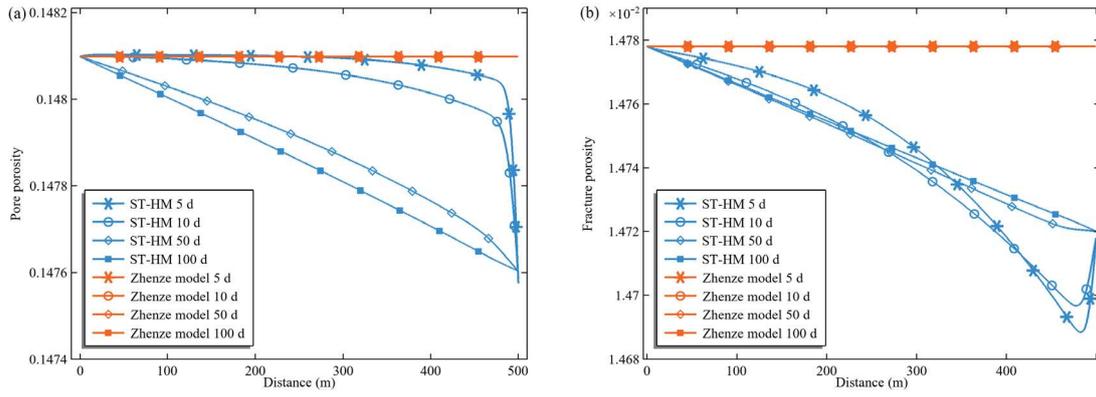
Figure 5. Horizontal and vertical displacement on X direction of rock



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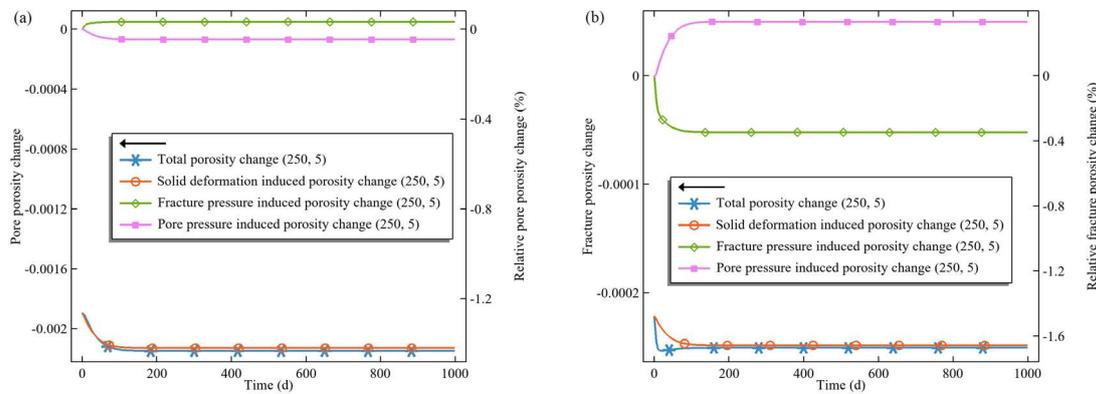
Figure 6. Pore and fracture fluid pressure evolution on X direction of rock



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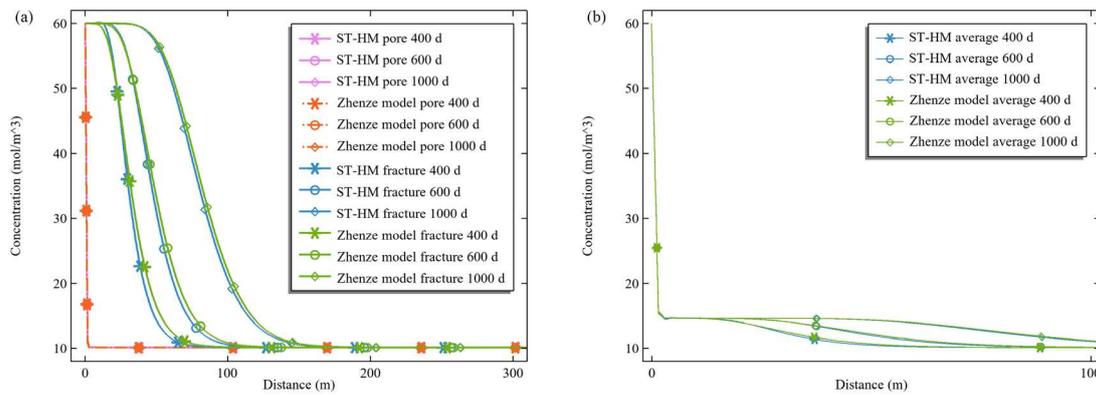
Figure 7. Pore and fracture porosity evolution on X direction of rock



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Figure 8. Pore and fracture porosity evolution at centre of the rock (X=250, Y=5)



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Figure 9. Solute concentration evolution on the X direction of rock (fracture

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concentration, pore concentration and average volumetric concentration)