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Systematic analysis of inner crust composition using the extended Thomas-Fermi approximation with pairing correlations

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We perform a systematic investigation of the chemical composition of the inner crust of a neutron star, using the extended Thomas-Fermi approximation, the Strutinsky integral correction for shell effects, and the BCS approximation for pairing. 15 Skyrme functionals were selected, which cover the range of values of important bulk properties of infinite nuclear matter, while also having pure neutron matter (PNM) equation of states (EoS) with varying degrees of stiffness. We find that a functional’s low-density PNM EoS is correlated with the number of protons found in the inner crust’s nuclear clusters and, in the lower-density region of the inner crust, with the pressure.

I. INTRODUCTION

Due to a strong pressure gradient, the matter within a Neutron Star (NS) organises itself into different layers [1]. Near the surface of the star lies the *crust*, where neutrons and protons are arranged in nuclear clusters in a crystalline structure, that transitions gradually into a liquid phase [2] approaching the core of the star [3].

Although the crust only accounts for a small fraction of the mass of the NS, it plays a major role in a variety of astrophysical phenomena, including the r-process [4], short gamma-ray burst precursors caused by resonant shattering [5], soft gamma-ray repeater giant flares [6, 7], and thermal relaxation in soft x-ray transients. The equation of state (EoS) of the crust is also believed to influence many properties of NS, such as the moment of inertia [8] which influences pulsar glitches, transport phenomena within the star [9–11], the relation between the radius and tidal deformability in low-mass NS [12], and the value of the second Love number, k_2 [13, 14]. Therefore it is crucial to have a reliable model of the crust of a NS, in order to achieve accurate descriptions of these phenomena.

As discussed in Ref.[1], current models predict that the crust has a crystalline structure, comprising neutron-rich nuclei surrounded by ultra-relativistic electrons [15–17]. The crust can be further separated into the *inner* and *outer* parts. The *outer* crust is the lower-density part of the crust, where neutrons are found only in bound states. The *inner* crust begins above baryonic densities of $n_b \approx 10^{11} \text{ g cm}^{-3}$, where neutrons begin to drip out of the nuclei, forming a gas.

Since the pioneering work of Negele and Vautherin [18], many theoretical models have been developed to study the properties of the inner crust. The presence of a neutron gas dramatically changes the properties of the clusters [19], in contrast to the outer crust, whose nuclei are isolated. Consequently, determining the composition of the inner crust requires a simultaneous treatment of the

clusters and the neutron gas. To simplify this task, it is customary to adopt the Wigner-Seitz (WS) approximation, in which the inner crust is decomposed into independent spheres, named *WS cells*, centered around each cluster, with a radius R_{WS} . Each cell is at β -equilibrium, with a certain number of protons, Z , and the same number of electrons (under the condition of charge-neutrality) spread through the cell. Using this approximation, the values of Z (the *chemical composition*) and R_{WS} can be determined at a given baryonic density, n_b , by minimising the total energy per particle, e , of the cell [20]. This procedure is valid only in the zero-temperature limit, which is applicable to the case of a non-accreting NS.

Since the density of electrons in the crust is essentially uniform [21], it is possible to calculate their contribution to e analytically [1]. The nuclear contribution is more complex and requires the use of a model. In the literature, several models are used to determine the nuclear binding energy of the system, such as the compressible liquid drop model [22], semi-classical models using the Thomas-Fermi approximation [20, 23–28], or the Hartree-Fock(-Bogoliubov) equations [18, 29].

Ideally, in order to model the EoS of a NS in a unified way, one should use the same model for all of its layers. The ideal choice is a fully microscopic method based on solving the Hartree-Fock-Bogoliubov (HFB) equations [29–31], using an effective nucleon-nucleon interaction [32, 33] adjusted to reproduce selected nuclear observables [34]. A big disadvantage of this approach is how the HFB equations treat continuum states in the inner crust using Dirichlet-Neumann boundary conditions [31, 35, 36], which leads to spurious shell effects. New methods to overcome such a problem have been suggested [37, 38], but no systematic calculations of the WS cells have been performed yet.

To avoid this difficulty, several authors have opted to instead use the extended Thomas-Fermi (ETF) method [24, 26, 39]. Due to its semi-classical nature, the ETF method is not affected by the spurious shell effects of the neutron gas encountered in the standard HFB approach. In Ref. [25], the ETF method was extended to use the Strutinsky integral (SI) correction to recover the important shell effects for the protons in the clus-

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ters. This method, named ETFSI, calculates the nuclear energy contribution using parameterised nuclear density profiles, while still using the same energy density functional to generate the fields as in the HFB method. In Ref. [40], the ETFSI method was further developed to take into account the effects of proton pairing correlations.

In a recent series of articles [41, 42], we have performed a systematic comparison between the HFB and ETFSI methods, which have typically yielded different results for the structure of the inner crust. In particular, we observed that by adding to ETFSI the effects of neutron pairing correlations under a simple local density approximation, the resulting energy per particle agrees well with that obtained using the HFB method, apart from the very low density region where the outer-inner crust transition takes place.

Having quantified the agreement of this ETFSI+pairing method with HFB, we can now address an interesting question, namely, why Zirconium isotopes (i.e. clusters with $Z = 40$) are consistently predicted throughout the inner crust [20, 25, 39, 43–45]. To answer this, we perform systematic calculations of the inner crust EoS, selecting several Skyrme functionals whose various infinite nuclear matter (INM) properties [46] cover reasonable ranges. By investigating possible relationships between the INM properties of Skyrme functionals and the proton content and pressure of the WS cells, we aim at providing a better understanding of previous investigations of the structure of the inner crust.

The article is organised as follows: in Sec. II, we outline the main features of the ETFSI+pairing method, and in Sec. III we discuss the Skyrme functionals selected for this work. In Sec. IV, we present results for the chemical composition and EoS of the inner crust, and we give our conclusions in Sec. V.

II. THE ETFSI+PAIRING METHOD

We now briefly review the ETFSI+pairing method, used to study the structure of the inner crust. For a more detailed discussion on the ETFSI method we refer the reader to Refs. [20, 25].

In this article, we consider only cold non-accreting NS, and so we neglect the effects of temperature [47, 48]. Under the WS approximation, for a fixed baryonic density n_b in the inner crust, we minimise the total energy per particle of the WS cell

$$e = e_{Sky} + e_e - Y_p Q_{n,\beta} . \quad (1)$$

e_{Sky} is the contribution to e from the interaction of the baryons via the strong force and from the Coulomb interaction between the protons. e_e is contribution from the kinetic and potential energies of ultra-relativistic electrons [49] and from the proton-electron interaction [29].

The last term accounts for the mass difference between neutrons and protons, $Q_{n,\beta} = 0.782$ MeV. $Y_p = Z/A$ is the proton fraction in the cell, where A is the total number of baryons in the cell.

The nuclear part of the energy per particle is expressed as functional of local densities as [33]

$$e_{Sky} = \frac{1}{A} \int \mathcal{E} \left(n_q(\mathbf{r}), \tau_q(\mathbf{r}), \vec{J}(\mathbf{r}) \right) d\mathbf{r} , \quad (2)$$

where $q = n, p$ stands for neutrons (n) and protons (p). Considering only time-reversal invariant systems [50], the *standard* Skyrme functional [51] depends only on a linear combination of the matter densities $n_q(r)$, kinetic densities $\tau_q(r)$, and spin current densities $\vec{J}(r)$, and on their derivatives.

To calculate e_{Sky} , we use the ETFSI+pairing method [25, 40, 42]. In this approach, the densities $n_q(r)$ are not calculated using the wave-functions, as in the HFB method [52], but are instead parameterised using Fermi-Dirac density profiles [25] as

$$n_q(r) = \frac{\rho_q^{liq} - \delta_{q,n} \rho^{gas}}{1 + \exp\left(\frac{r-r_q}{a_q}\right)} + \delta_{q,n} \rho^{gas} . \quad (3)$$

$\rho_{q=n,p}^{liq}$ are the densities of the neutrons and protons at the centre of the WS cell, $r = 0$, while ρ^{gas} is the neutron density at the cell edge, $r = R_{WS}$. $r_{q=n,p}$ are the radii of the neutrons and protons in the cluster, and $a_{q=n,p}$ are their cluster surface diffusivities. These 7 adjustable parameters are determined by the minimisation of the energy per particle, given in Eq. 1, under the constraints of charge neutrality and β -equilibrium. We also use the relations presented in Ref. [25], reducing the number of free parameters to 5. In Ref. [26] the authors introduced an exponent for the denominator in Eq. 3, but they concluded that the results are largely unchanged, so we do not consider such a parameterisation.

The ETF method expresses $\tau_q(r)$ and $\vec{J}(r)$ as a function of the matter densities and their derivatives, up to 4th order in the Wigner-Kirkwood expansion [53]. The full expressions of these quantities are given in Ref. [54].

In this *semiclassical* approach, shell effects are lost; they are recovered with the Strutinsky integral (SI) correction suggested in Ref. [25]. The main advantages of this ETFSI method over the HFB method are the reliable treatment of the neutron gas [35, 55], and the dramatic reduction in computational cost needed for the energy minimisation.

In the original ETFSI method [20, 25], pairing correlations were not taken into account, despite their fundamental role in describing several important phenomena related to the physics of the inner crust [30, 56–58]. Recent improvements of the model presented in Refs. [40, 42] now allow the inclusion of pairing correlations for both protons and neutrons.

The resulting ETFSI+pairing model leads to results that agree well with the most accurate HFB results currently available [31]. At the interface region between the outer and the inner crust, a non-negligible discrepancy was observed [42], but this concerns a very limited density region and it will not modify the conclusions of this article.

Since Skyrme functionals are often fitted on doubly-magic nuclei, they are not usually equipped with a consistent pairing interaction. Therefore, for each functional used in this work, we add a simple density-dependent pairing interaction of the form [59]

$$v_q^{pair}(\mathbf{r}_1, \mathbf{r}_2) = v_{0q} \left[1 - \eta \left(\frac{n_q(r)}{n_0} \right)^\alpha \right] \delta(\mathbf{r}_1 - \mathbf{r}_2). \quad (4)$$

We choose the parameters $\eta = 0.7$ and $\alpha = 0.45$. n_0 is the saturation density of the functional. We assume that the pairing strength v_{0q} is the same for neutrons and protons, and we fix it to obtain a maximum pairing gap in pure neutron matter (PNM) of ≈ 3 MeV, as done in Ref. [29]. These choices may appear arbitrary [60], but it was shown in Ref. [42] that a variation of the pairing strength does not impact the resulting chemical composition of the inner crust. To avoid the ultraviolet divergence of the interaction given in Eq.4 [61], we adopt a smooth cut-off in quasi-particle space as defined in Ref. [29].

For functionals of the BSk family [62], we keep the pairing interaction used and developed by the authors.

III. CHOICE OF FUNCTIONALS

According to Ref. [46], more than 240 Skyrme functionals have been published in the literature. Typically, their parameters are adjusted to reproduce the binding energy of some (or all) atomic nuclei, and also some INM properties. There is no standard parameter fitting protocol; various (pseudo)-observables are used in fitting, with widely varying uncertainties. As a consequence, it is difficult to assess the quality of functionals.

The neutron gas in the WS cell is similar to an infinite nuclear system, and comprises the majority of the baryonic matter in the inner crust. We therefore focus on 4 INM properties of the functionals: the saturation density n_0 , the energy per particle E_0 , the symmetry energy J , and its slope L , all evaluated at n_0 . These quantities depend exactly on the EoS of INM [63], and so one can expect that variations in the EoS lead to modifications in the energy per particle of the neutron gas, thus changing the relative energy contributions from baryons and electrons in the WS cell, and ultimately affecting the chemical composition and pressure in the inner crust. In Fig. 1, we show as histograms the distributions of n_0 , E_0 , J and L , for all the functionals in Ref. [46].

By inspecting the lower panels of Fig.1, we observe that not all functionals given in Ref. [46] are shown. In

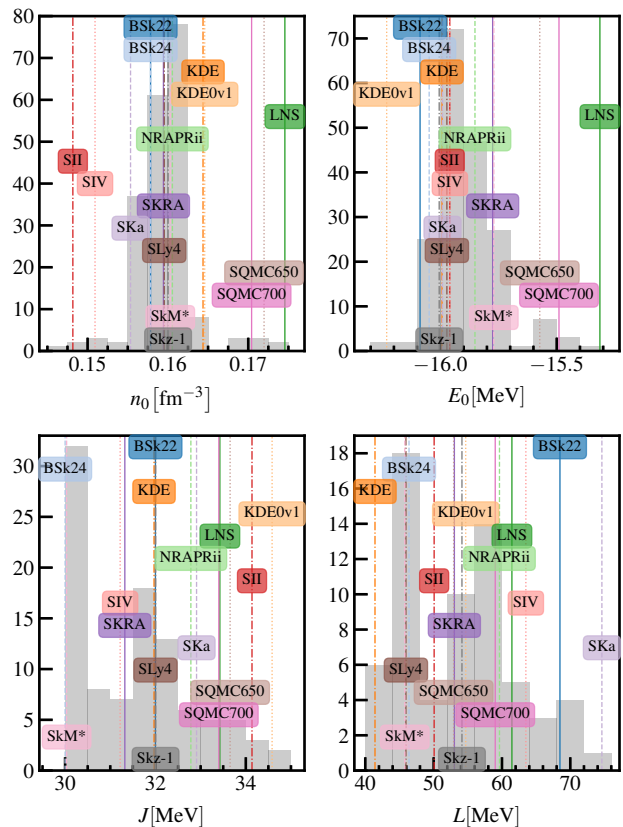


FIG. 1. (Colors online) Each panel represents the distribution of the Skyrme functionals given in Ref. [46] for 4 INM properties. Starting from the top left panel and going clockwise we show n_0 , E_0 , L and J . The values of each quantity for the 15 functionals selected in this work are shown as labelled vertical lines. See text for details.

our selection process, we imposed two additional constraints on the symmetry energy and its slope, namely $30 \leq J \leq 35$ MeV and $L = 58 \pm 18$ MeV. These ranges of values are the ones suggested in Ref. [46]. These constraints are to some extent arbitrary, as there is no consensus on what constitutes a reasonable range of values for J and L , but they are consistent with the findings of recent analyses [64–68], and with results and associated uncertainties from new chiral EFT calculations [69, 70].

From this group, we finally selected 15 functionals which have a wide range of values for each of n_0 , E_0 , J , and L : BSk22 and BSk24 [62], KDE [71], KDE0v1 [71], LNS [72], NRAPRii [73] (NRAPR [74] with the spin-orbit strength modified), SII [75], SIV [76], SKRA [77], SKA [78], SLy4 [79], SQMC650 and SQMC700 [80], SkM* [81], and Skz-1 [82]. Their INM properties are labelled in Fig. 1. Five of them — **KDE0v1**, **LNS**, **NRAPR**, **SKRA** and **SQMC700** — are consistent with all the INM constraints presented in Ref. [46], and are named in that reference and hereafter as the CSkP* set.

For our functional selection, we also show the energy

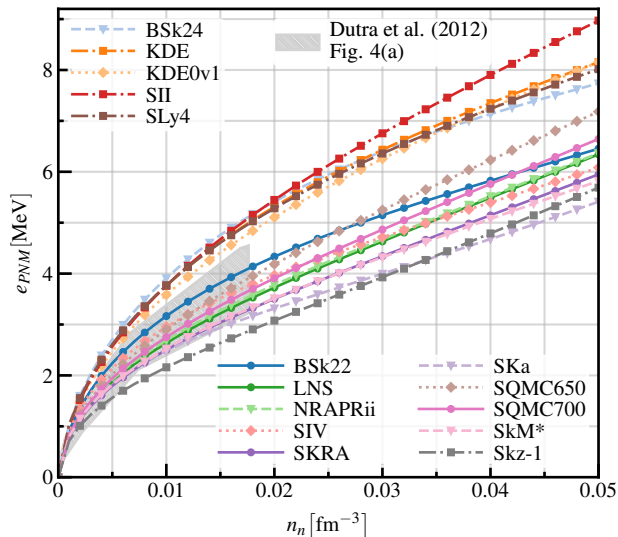


FIG. 2. (Colors online) The EoS of PNM as a function of the neutron density n_n for the 15 functionals selected in this work. The shaded area represents the constraint given in Fig. 4(a) of Ref. [46]. The functionals labelled in the upper left are the Stiff set, and those in the lower right the Soft set. See text for details.

per particle e_{PNM} in PNM in Fig. 2, for the density range relevant for spherical inner crust calculations. The functionals clearly fall into two distinct families: the set BSk24, KDE, KDE0v1, SII, and SLy4, with a stiff EoS at these low densities (hereafter the *Stiff* set), and the set BSk22, LNS, NRAPRii, SIV, SKRA, SKa, SQMC650, SQMC700, SkM*, and Skz-1, with a very soft EoS (hereafter the *Soft* set). On the same figure, we have added a shaded area which corresponds to the range spanned by several *ab-initio* calculations used to derive the EoS in PNM, and which is discussed in Fig. 4(a) of Ref. [46].

The Stiff set, including SLy4 and BSk24 which have been widely used for NS calculations, are in disagreement with the *ab-initio* calculations at these densities. In contrast, the majority of the Soft set are in reasonable agreement, apart from Skz-1 which is very soft. However, the *ab-initio* calculations selected in Ref. [46] do not represent the entire range available in the literature. As shown in Ref.[62], BSk24 and SLy4 are compatible with the error bars provided by chiral effective field theory calculations from Ref [83].

The inner crust does not only comprise neutrons, but also has a significant proton fraction, and so it is interesting to observe the behaviour of the symmetry energy S in the relevant density region and not only at n_0 . This is shown in Fig. 3, where we again observe a similar grouping into the Stiff and Soft sets as shown in Fig. 2. The reason is quite simple and can be understood from Fig. 1: the vast majority of the functionals have values of n_0 and E_0 in a very narrow range. This means that all these functionals have a very similar EoS in symmetric nuclear

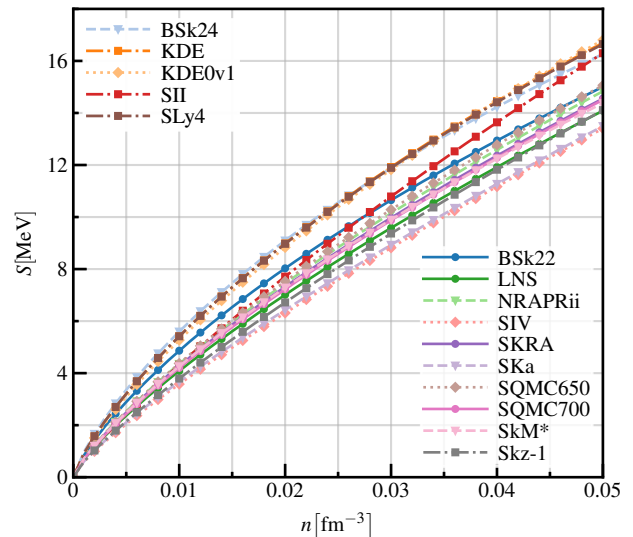


FIG. 3. (Colors online) The symmetry energy S as a function of density for selected functionals. The functionals labelled in the upper left are the Stiff set, and those in the lower right the Soft set. See text for details.

matter (SNM) in the low density region. Since S is just the difference between the EoS in PNM and SNM (within the parabolic approximation [84]), it follows that the pattern observed in PNM repeats here in a very similar way. The exception to this simple rule are SII and SIV. They have extremely low values of n_0 , shown in Fig. 1, and so a significantly different EoS in SNM compared to the other functionals.

IV. INNER CRUST COMPOSITION

To calculate the EoS and determine the chemical composition of the inner crust, we minimise the total energy per particle e in the WS cell, given in Eq.1, using the method explained in Sec. II and Ref [42]. We cover the range of baryonic densities $n_b \in [0.00025, 0.05] \text{ fm}^{-3}$, above which non-spherical pasta phases are expected to appear [45], and the range of proton numbers $Z \in [16, 60]$.

We first show results for a few selected baryonic densities $n_b = 0.01, 0.02$ and 0.03 fm^{-3} , for a smaller group of functionals. Figure 4 illustrates the energy per particle as a function of Z . We show a re-scaled energy per particle, e_s , defined as

$$e_s(n_b) = e(n_b) - e_{PNM}(n_b) - e_{n,pair}(n_b), \quad (5)$$

where $e_{n,pair}$ is the neutron pairing energy per particle. This is purely for visual reasons, so that all values lie within a similar energy range. Since e_{PNM} is independent of Z , and $e_{n,pair}$ is roughly constant with respect to

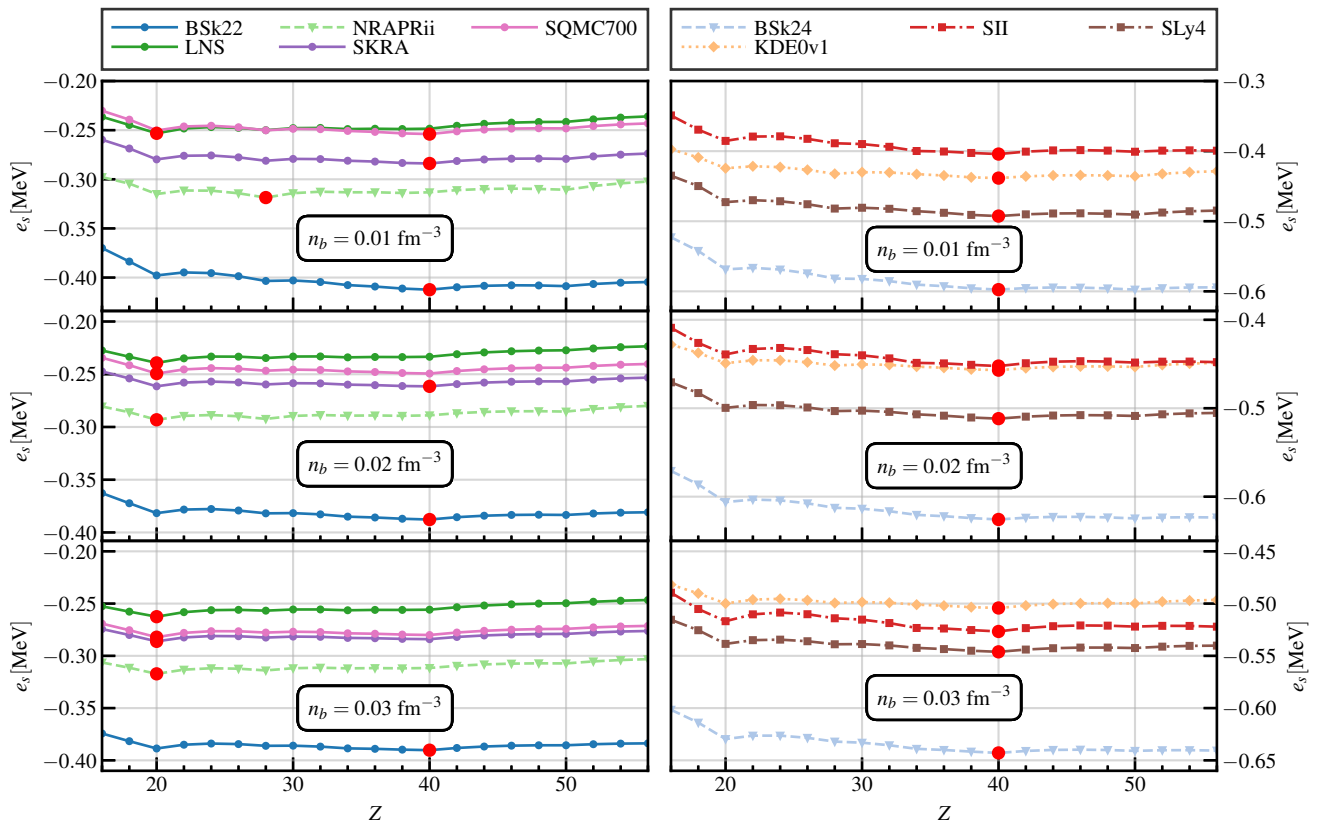


FIG. 4. (Colors online) The re-scaled energy per particle e_s , defined in Eq.5, for selected functionals from the Soft (left column) and Stiff (right column) sets, at three values of baryonic density n_b . The large dot on each curve represents the position of the energy minimum. See text for details.

Z [42], this results in a simple shift in e for a given functional. In the left column are 5 from the Soft set: **BSk22**, **LNS**, **NRAPRii**, **SKRA**, and **SQMC700**; in the right column are 4 functionals from the Stiff set: **BSk24**, **KDE0v1**, **SII**, and **SLy4**, as explained in Sec. III.

In this figure, we clearly see the different behaviour of the Soft and Stiff sets. The Stiff set has a persistent minimum at $Z = 40$, while the Soft set has minima that favour lower Z , or that shift towards lower Z as n_b increases, with all favouring $Z = 20$ at $n_b = 0.03 \text{ fm}^{-3}$. The exception in the Soft set is **BSk22**, which transitions to $Z = 20$ just above $n_b = 0.03 \text{ fm}^{-3}$, and which has one of the stiffest PNM EoS in this set. **NRAPRii** is unique among all the functionals investigated, in that it favours $Z = 28$ at $n_b = 0.01 \text{ fm}^{-3}$. This is likely related to an issue involving the spin-orbit parameter, W_0 . See the discussion in Ref. [73] for more details. Below $n_b = 0.01 \text{ fm}^{-3}$, nearer the transition region between the outer and inner crust, other finite-size effects not considered in our ETFSI+pairing method may become significant, changing the results. The HFB method is preferable in this very low density range.

To better understand the origin of the different minima, we plot in Fig. 5 the *smooth* part of the total energy

per particle, e_{smooth} , defined as

$$e_{smooth} = e - e_{n,pair} - (e_p^{SI} + e_{p,pair}) \quad (6)$$

where the last term is the sum of the SI correction with BCS pairing for protons, and the proton pairing energy, as explained in Ref. [40]. This results in a smooth parabola-like shape, with a single minimum. We select 2 representative functionals **SQMC700** and **BSk24**, from the Soft and Stiff sets respectively. The parabola shifts significantly in Z going between the two functionals. Furthermore, as n_b increases, we observe that for **SQMC700** the minimum moves from $Z = 32$ to $Z = 28$, while for **BSk24** it stays at $Z = 46$.

When the SI correction is included, as in Fig. 4, each functional at each density displays local minima at (semi-)magic Z values between 20 and 50. However, Fig. 5 shows how the global minimum in each case is governed more by the stiffness of the PNM EoS. In this work, we strictly consider no temperature effects. Since the various energy minima are quite close to each other, the inclusion of such effects may change this picture. See discussion in Ref. [47] for more details.

A more systematic study of the evolution of the proton

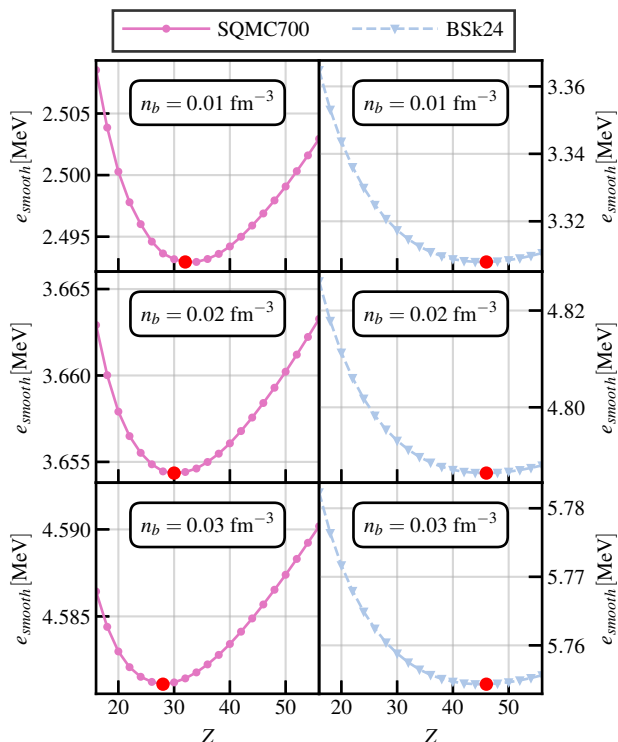


FIG. 5. (Colors online) Total energy per particle, e_{smooth} , obtained without using shell or pairing corrections, for two functionals at three representative baryonic densities n_b . See text for details.

content of the clusters in the inner crust is illustrated in Fig. 6, for all functionals in the Soft and Stiff sets. For clarity, small vertical offsets have been added to each line, and the Soft and Stiff sets in are shown in separate panels. At very low n_b several values of Z are observed, but above $n_b = 0.02 \text{ fm}^{-3}$ only two minima are observed: $Z = 20$ and $Z = 40$. The minimum at $Z = 40$ seems to be favoured by most functionals at around $n_b = 0.01 \text{ fm}^{-3}$, but at $n_b = 0.03 \text{ fm}^{-3}$ the majority of the Soft set have transitioned to the $Z = 20$ configuration. The only exception to this rough classification is SQMC650, whose e_{PNM} is the highest out of the Soft set, as seen in Fig. 2, and it maintains a $Z = 40$ minimum up to quite a high baryonic density of $n_b \approx 0.04 \text{ fm}^{-3}$.

It is interesting to note that $Z = 20$ is favoured by functionals whose PNM EoS is compatible with *ab-initio* calculations. This result is also in good agreement with the findings of Ref. [58]. Although further analysis is necessary, our results suggest that a better understanding of the EoS in low-density PNM may help to clarify the chemical composition of the inner crust.

Having calculated the chemical composition of the inner crust, we now study its general properties with the different functionals. In Fig. 7, we show the total energy per particle e (Eq. 1) for the WS cells in the inner crust, obtained through the energy minimisation, as a function

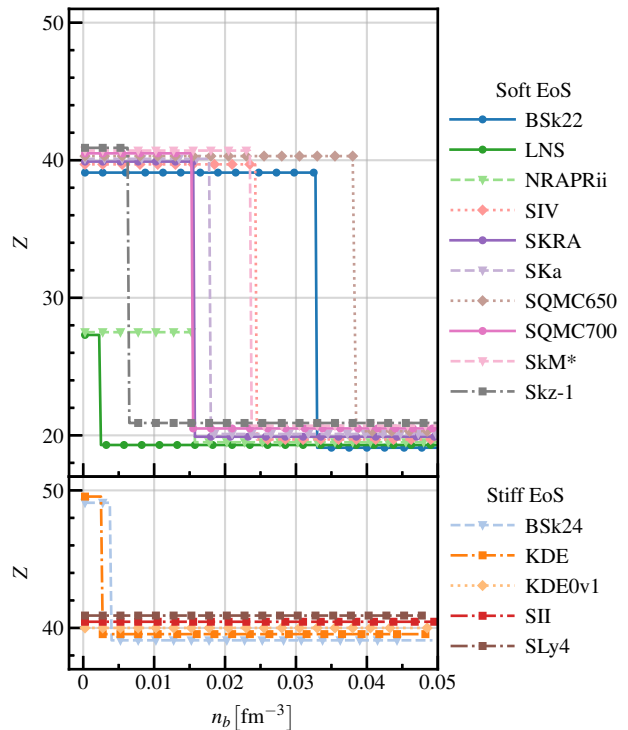


FIG. 6. (Colors online) Proton content of WS cells in the inner crust as a function of the baryonic density n_b , for all functionals in the Soft and Stiff sets. A small vertical offset has been added to each line to make them more visible; all lines only take values from one of $Z = 20, 28, 40, 50$. See text for details.

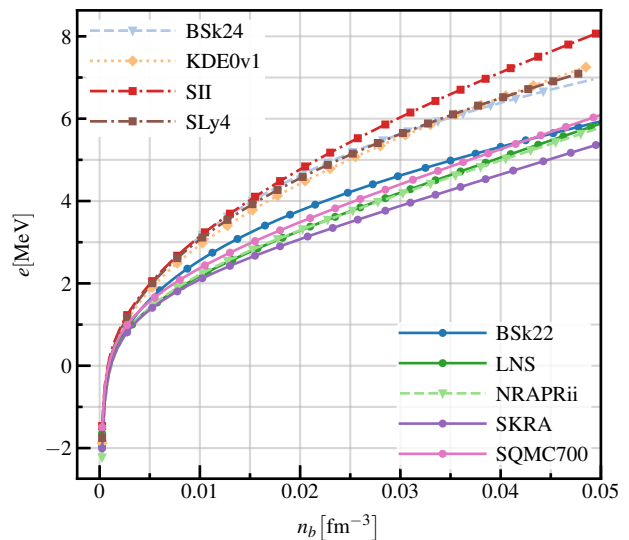


FIG. 7. (Colors online) The energy per particle e in the inner crust, as a function of the baryonic density n_b . 4 functionals from the Stiff set are shown, labelled in the upper left, and 5 from the Soft set, labelled in the lower right. See text for details.

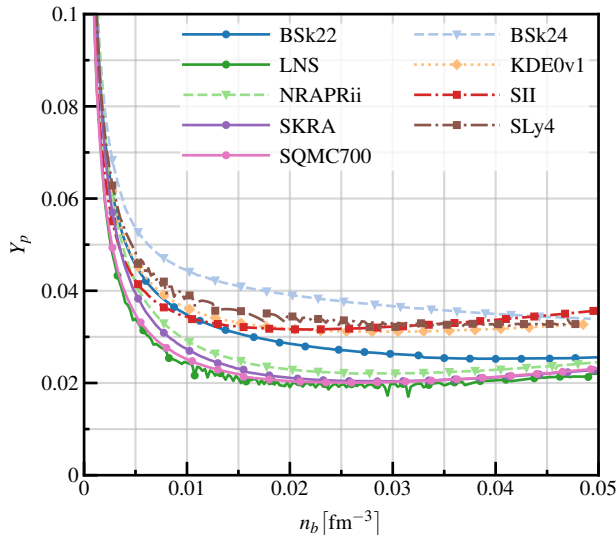


FIG. 8. (Colors online) Proton fraction Y_p in the inner crust, as a function of the baryonic density n_b . 4 functionals from the Stiff set are shown, labelled in the upper right, and 5 from the Soft set, labelled in the upper left. See text for details.

of the baryonic density n_b . The pattern formed by the various Soft and Stiff functionals is almost identical to that seen for e_{PNM} , shown in Fig. 2. This further supports the PNM EoS being the major driver of the inner crust EoS.

Another relevant astrophysical quantity is the proton fraction in the inner crust, $Y_p = Z/A$. It plays a role in the transport properties of a NS and in the determination of the neutrino mean free path in the stellar medium [85]. In Fig. 8, we report Y_p as a function of n_b for the same smaller selection of functionals shown in Fig. 4. The same feature already observed in Fig. 6 is clear: the Soft functionals predict WS cells with fewer protons, and have a lower Y_p , while the contrary is true for the Stiff functionals. At around $n_b = 0.02 \text{ fm}^{-3}$, Y_p varies by up to a factor of 2 between the two sets. When Z in a WS cell changes, the baryon content of each WS cell follows this change quite closely. As a consequence, the factor of 2 observed between the two dominant minima in Fig. 6 ($Z = 20, 40$) is conserved here.

It is worth commenting on the case of BSk24, a functional specifically created to be able to provide a unified description of the NS EoS. This was adjusted under many constraints, including the requirement that it reproduce the LS2 equation of state [86] in PNM, which was calculated using the microscopic Brueckner-Hartree-Fock method. The authors note in Ref. [62] that they focused on supernuclear densities when constraining BSk24 to the LS2 EoS. As a result, the BSk24 PNM EoS is significantly stiffer than LS2 at inner crust densities. Around $n_b = 0.04 \text{ fm}^{-3}$, LS2 is better approximated by SQMC650, SQMC700, and NRAPRii, than by BSk24.

In Ref. [43, 44], the authors claim that the constraining

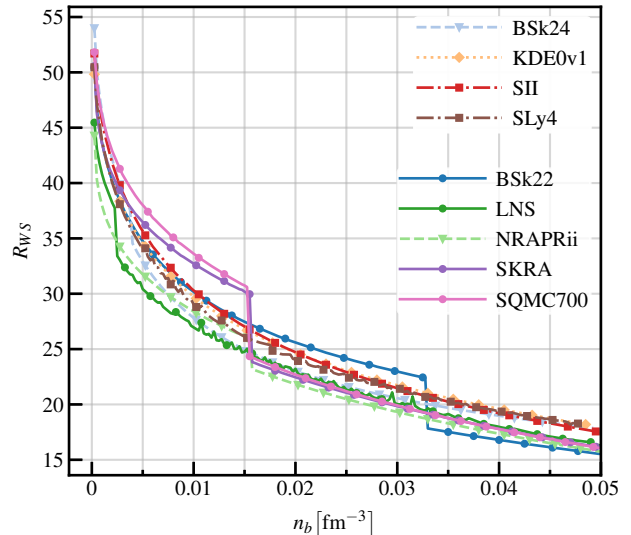


FIG. 9. (Colors online) Radius of WS cells R_{WS} in the inner crust, as a function of the baryonic density n_b . 4 functionals from the Stiff set are shown, labelled in the upper right, and 5 from the Soft set, labelled just below on the right. See text for details.

PNM EoS has little effect on Y_p in the inner crust, and show that J is an important INM quantity. However, by inspecting carefully Figs. 2 and 8, we see that BSk24 ($J = 30 \text{ MeV}$), SLy4 ($J = 32 \text{ MeV}$), and KDE0v1 ($J = 34.6 \text{ MeV}$) have a Y_p largely in the range $0.03 - 0.04$. The functionals with a softer PNM EoS — BSk22, LNS, NRAPRii, SKRA, and SQMC700 — have a Y_p mostly in the range $0.02 - 0.025$, despite having J ranging from 31.3 MeV to 33.4 MeV . We therefore conclude that J is not the driving factor leading to large variations of Y_p , but it is the equation of state of PNM at low density, or equivalently the symmetry energy S at low density as shown in Fig. 3 and discussed in Sec. III.

The softest functional investigated in Ref. [43, 44], BSk22, which more closely follows LS2 at inner crust densities, gives results favouring an energy minimum at $Z = 20$ at intermediate inner crust densities, and lower values for Y_p , e , and P . This means that the results presented in this article are consistent with those in Ref. [43, 44], but given the small variations in the properties of the functionals used, the authors were not able to observe such an interesting correlation.

It is also interesting to note from Fig. 8 that almost all of the CSKP* set, selected in Ref. [46] according to several criteria based on INM properties, tend to favour a very small Y_p within the crust. The only exception is KDE0v1, from the Stiff set.

In Fig. 9, we illustrate the evolution of the radius of the Wigner-Seitz cell, R_{WS} , as a function of the baryonic density, again for the smaller selection of functionals. While in the low density region, $n_b \in [0.00025, 0.015] \text{ fm}^{-3}$, we observe significant varia-

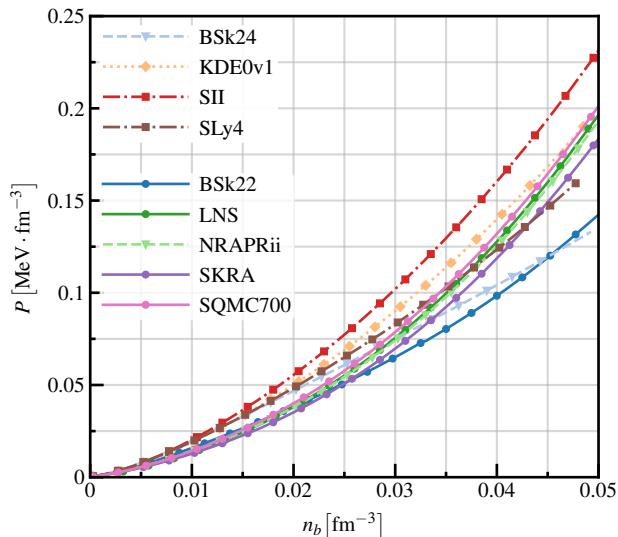


FIG. 10. Pressure P in the inner crust, as a function of the baryonic density n_b . 4 functionals from the Stiff set are shown, labelled in the upper left, and 5 from the Soft set, labelled just below on the left. See text for details.

tions in the sizes of the cells, at higher baryonic densities, we notice that the size of the cell is almost independent of the functional. In particular, we clearly identify again the Stiff set and the Soft set, giving values of R_{WS} very close to each other. The sudden jumps shown in this figure correspond to abrupt changes in the chemical composition of the crust, as seen in Fig. 6.

Finally, in Fig. 10, we illustrate the evolution of the pressure P as a function of the baryonic density n_b , for the same 8 functionals already discussed in Figs. 7 and 8. In this case, the separation into the Stiff and Soft sets is only visible up to around $n_b = 0.025 \text{ fm}^{-3}$. We recall

that the pressure depends on the derivatives of the EoS, rather than its value.

V. CONCLUSIONS

In this article we have performed a systematic analysis of the structure of the inner crust of a non-accreting neutron star (NS), using the ETFSI+pairing method. By selecting 15 different functionals with different infinite nuclear matter properties, we have illustrated a possible correlation between the a functional's pure neutron matter (PNM) equation of state (EoS) at low density, and the proton content and pressure of the Wigner-Seitz cells. In particular, we have shown that functionals with a *soft* PNM EoS at low densities tend to favour the $Z = 20$ minimum, as well as a lower proton fraction, and a lower pressure up to around $n_b = 0.025 \text{ fm}^{-3}$, while functionals with a *stiff* PNM EoS show the opposite behaviour, and favour instead the $Z = 40$ minimum. This result clearly shows the importance of the constraining the PNM EoS more tightly at subnuclear densities in the adjustment protocol of Skyrme functionals used to study NS. This is not the *standard* procedure, since authors tend to focus more on the high density trend of the EoS; indeed, this is necessary to prevent the collapse of the NS, and to reproduce its global properties such as the maximum mass and radius [87]. More stringent constraints on functionals obtained by more recent *ab-initio* calculations at low density should pave the way toward a more reliable unified EoS for neutron star matter.

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