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1 **THE IMPORTANCE OF THE HEEL EFFECT in X-RAY CT IMAGING of**
2 **SOILS**

3 **K. Liu¹, R. Boardman², M. Mavrogordato³, F. A. Loveridge⁴ *, and W. Powrie⁵**

4 *** corresponding author (f.a.loveridge@leeds.ac.uk; 07773346203)**

5 **ABSTRACT**

6 Non-destructive and non-invasive X-ray computed tomography (CT) is increasingly used in
7 environmental geotechnics research. As a result of recent advances in technology and image
8 processing techniques, CT with rapid scanning now has the potential to track changes in soil
9 structure or soil water conditions as they happen, rather than as previously on a specimen in
10 (temporary) stasis. Gathering meaningful data in a short scan time requires compromises to be made
11 on parameters such as exposure time, and / or the use of higher X-ray intensities and energies. Data
12 processing and imaging processing - including the removal of any artefacts, which can cause errors
13 in interpretation of soil structure or phase proportions - then become especially important. One such
14 artefact is the heel effect. It has been recognised in medical imaging, owing to its association with
15 high scan energies. However, it has not previously been identified in soil imaging, despite the trend
16 towards using higher energies. This paper presents an investigation into the potential for the heel
17 effect to affect the soil property determination. It is shown for the first time that a noticeable heel
18 effect will be present in CT images of soils and derived phase proportion data, when certain types of
19 X-ray reflection targets are used. A correction method for the heel effect is presented, use of which
20 will prevent significant errors in derived soil parameters such as water content.

21
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25
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28
29 **Notation:**

30 CT computed tomography

¹ Faculty of Engineering and Physical Sciences, University of Southampton, Southampton, SO17 1BJ, UK; 0000-0002-7533-1272

² Faculty of Engineering and Physical Sciences, University of Southampton, Southampton, SO17 1BJ, UK; 0000-0002-4008-0098

³ Faculty of Engineering and Physical Sciences, University of Southampton, Southampton, SO17 1BJ, UK; 0000-0002-3956-2866

⁴ Faculty of Engineering and Physical Sciences, University of Leeds, Leeds, LS2 9JT, UK; 0000-0002-6688-6305

⁵ Faculty of Engineering and Physical Sciences, University of Southampton, Southampton, SO17 1BJ, UK; 0000-0002-2271-0826

31 GV grey value
32 ROI region of interest
33

34 1 INTRODUCTION

35 X-ray computed tomography (CT) is increasingly used for research in soil mechanics, geotechnical
36 engineering and soil science. The technique involves passing X-rays through a soil specimen
37 mounted on the manipulator and measuring the arrival of photons at the X-ray detector. Repeating
38 this procedure for many angles of incidence as the specimen rotates generates a series of 2D
39 radiographs, one for each angle. This enables a three dimensional model to be built up that provides
40 information on density contrasts according to the differential absorption of the photons at the
41 specimen. The density information can then be used to identify the soil phases, soil structure (e.g.
42 Sleutel et al, 2008) and other geotechnical parameters including incremental strain from applied
43 loading (e.g. Ando et al, 2012a, b).

44

45 The approach has become popular as a method for non-destructive and non-invasive examination of
46 soil specimens in at least a temporary state of stasis. However, advances in technology and image
47 processing techniques mean that X-ray CT is increasingly being seen as a way to track and
48 meaningfully quantify changes in soil structure or changes in soil water conditions as they happen
49 with time (e.g. Cnudde & Boone, 2013), enhancing its utility in environmental geotechnics.

50 Applications in environmental geotechnics are varied and diverse and include for example reservoir
51 characterisation (Van Geet et al, 2000), soil water retention characterisation (Khaddour et al, 2018),
52 measurement of liquid and/or gas flow and associated deformation (Mees et al, 2003, Alvarez-
53 Borges et al, 2018, Wang et al, 2019), quantification of soil-biological interactions (Helliwell et al,
54 2013). To track changes in soil phases or structure with time, the associated rapid scanning processes
55 require a careful balance between speed and quality. Minimising the scan duration ensures that
56 temporal changes are captured and images are not blurred, for example due to a significant changes

57 in water content within the specimen or specimen movement during a single scan. At the same time,
58 it is important to maintain the image quality traditionally associated with a longer scan time. This
59 means that compromise on some scan settings may be required, for example using shorter exposure
60 times or a smaller number of projections in combination with higher X-ray intensities and energies.

61

62 If scan setting compromises lead to a lower resolution CT image, different soil phases may occur
63 within the same image voxel. This is known as the partial volume effect (Ketcham 2005). To avoid
64 errors arising from averaging of phase densities within voxels, data processing to remove artefacts
65 become especially important. Artefacts or anomalies in the data obtained may cause deterioration in
66 image quality and reduce the accuracy of image analysis, including parameters such as soil phase
67 proportions (e.g. porosity, water content), which can be derived by segmentation.

68

69 Common artefacts that can hamper reliable image segmentation include beam hardening and rings
70 (Boas & Fleischmann, 2012, Ketcham & Hannah, 2014). Beam hardening is a limitation arising from
71 polychromatic X-ray sources (Ketcham and Carlson, 2001), in which there is selective attenuation of
72 lower energy photons. This can cause dark streaks in images where there is greatest attenuation, or a
73 “cupping” artefact, where an artificially bright zone is present around the edges of a specimen. Beam
74 hardening can be easily addressed by filtration of the X-ray beam before it reaches the specimen. A
75 ring artefact is a dark ring, centred about the specimen rotation point; it results from a deficiency in
76 the X-ray detector and is a systematic error (Davis and Elliott, 2006). Rings are usually removed by
77 calibration of the detector, or by moving the specimen or detector by a small, subsequently-
78 corrected, random amount between frames.

79

80 The artefact known as the heel effect is associated with CT scanning systems in which the incident
81 X-rays are generated by means of a reflection target. It has been recognised in medical imaging due
82 to its association with high scan energies. However, it has not been previously identified in soil
83 imaging. The trend in soil mechanics towards more sensitive contrast differentiation from X-ray CT
84 images (e.g. Liu et al, 2017, Alvarez-Borges et al, 2018, Liu, 2020) places the scan energies utilised
85 into a similar category to biomedical imaging. This brings the potential for an undetected and / or
86 uncorrected heel effect leading to errors in image interpretation and derived parameters, for example
87 changes in water content associated with thermal or chemical processes, or changes in soil density
88 related to mechanical loading.

89

90 This paper presents the first investigation into the potential for the heel effect relevant to soil
91 imaging. First, we review the heel effect artefact and recent trends of increasing scan energy in
92 applications of X-ray CT to problems in soil mechanics. Then we identify the presence of the heel
93 effect in soil and control materials, before investigating which factors influence the magnitude of the
94 effect. Finally, we assess methods to avoid errors arising from the heel effect, and propose a new
95 universal correction method to remove the artefact.

96

97 **2 BACKGROUND**

98 **2.1 The Heel Effect**

99 CT scanners use targets between the electron beam source and the object being scanned to generate
100 X-rays in a pattern suitable for imaging. These targets are generally of either the reflection or the
101 transmission type (Figure 1). Reflection targets are more commonly used in laboratory microfocus

102 X-ray CT sources. The heel effect is an image artefact that results from the inherent features of a
103 reflection type target.

104

105 Reflection targets are usually formed of inclined metallic plates (Figure 1a). They are frequently
106 adopted in high energy applications as they are easily cooled and provide good flux. Electrons
107 generated by a hot tungsten filament (cathode) are accelerated to the target (anode), where
108 interactions occur at an atomic level to produce X-rays that are emitted onto the detector. Owing to
109 the angle at which the target is set, the emerging X-rays will pass through more of the target material
110 on the downslope side than on the upslope (Figure 1a). Some of the X-rays will be attenuated more
111 significantly by the extra thickness of the target material. This is because as the thickness of a given
112 material increases, the chance of a photon of a given energy being absorbed increases according to
113 the Lambert-Beer law (Swinehart, 1962). X-ray paths with different reflection angles will therefore
114 experience correspondingly different degrees of filtration. Hence in the example shown in Figure 1a,
115 the resulting projection will tend to reduce in intensity with decreasing vertical elevation. This will
116 cause a gradient in the distribution of the grey value (GV), a measure of specimen density, evaluated
117 at the detector. This is known as the heel effect.

118

119 The heel effect is often slight and may not significantly affect some types of image analysis - for
120 example, where there is no particular interest in the spatial distribution of a soil property. Owing to
121 the origin of the heel effect in the attenuation of photons, it will be more noticeable in scans where
122 high energies are used. This has typically included medical applications, but is also now starting to
123 include geotechnics as discussed in Sections 2.2 and 2.3 below.

124

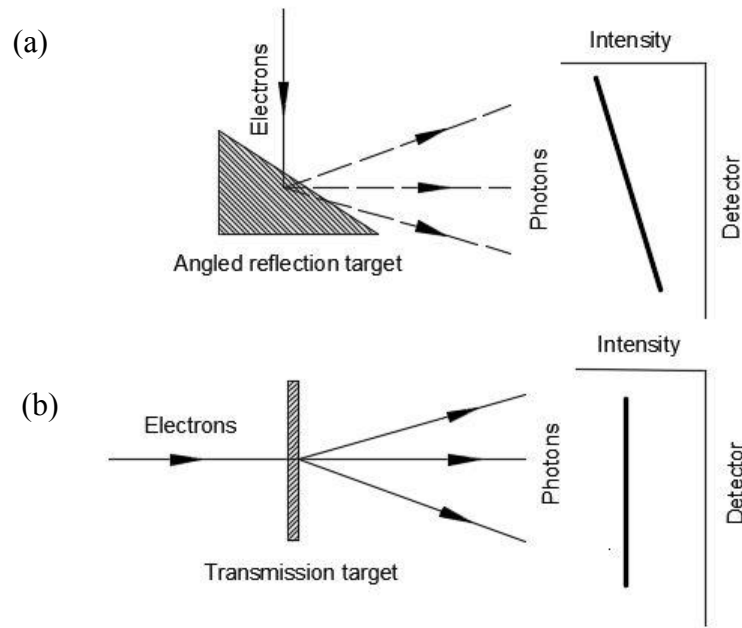


Figure 1: Schematics of X-ray (a) reflection target and (b) transmission target

125

126

127

128 As an alternative to a reflection target, a transmission target may be used. This is placed
 129 perpendicular to the electron beam (Figure 1b). X-rays are generated within a small region of
 130 substantially uniform thickness, and penetrate directly through the thin target material. Thus the
 131 intensity of the resulting X-rays is theoretically uniform and transmission targets do not show the
 132 heel effect.

133

134 2.2 The Heel Effect in Medical Applications

135 The heel effect was initially identified in the medical field, as medical CT scanners normally use
 136 reflection targets owing to the high flux required for the rapid scanning of large objects (patients).
 137 Several correction or compensation methods have been proposed. A method using a compensation
 138 filter (where the aluminium layer of either side of the target is increased in thickness on one side)
 139 was developed by Mori et al, 2005, to eliminate potential radiation damage to patients. Heel effect

140 corrections based on simple first-order beam hardening (Braun et al, 2010), and a slice-by-slice
141 background subtraction approach (Johnston et al, 2015) have also been proposed. However, these
142 approaches have limitations in that their application requires detailed understanding of the target
143 composition and geometry, hence additional specialist input.

144

145 **2.3 High Energy X-Ray CT Scanning in Soil Mechanics**

146 The heel effect will be present in all scans carried out using a reflection target. There is no threshold
147 energy above which the effect should be noticeable, as this will be a function of the image contrast.
148 Nonetheless, it is expected that higher energies would lead to a more pronounced heel effect. In this
149 section, we review recent soil scanning applications that have used higher energies.

150

151 CiSlerova et al (2002) used a medical CT scanner (Siemens SOMATOM PLUS IV) to scan soil
152 specimens with a peak scan energy of 140 kV; Farber et al (2003) used a Skyscan 1072 high-
153 resolution X-ray micro-tomography unit (Skyscan, Belgium) to study the porosity of granules; and
154 Taud et al (2005) used a peak scan energy of 130 kV on a PICKER IQ PREMIER (IQXTRA)
155 scanner to explore porosity in a rock. The energy values in all these studies are within or above the
156 range shown in this paper to cause a noticeable heel effect. However, none of these studies
157 considered the spatial distribution of specimen density, porosity etc., so any errors potentially present
158 due to the heel effect would remain unnoticed.

159

160 Some studies have derived spatial distributions of density or density related parameters without
161 explicitly identifying the influence of the heel effect. For example, Anderson et al (1990) used a
162 Siemens SOMATOM DR Version H scanner with a peak energy of 125 keV to analyse macropores

163 in undisturbed soil cores. They evaluated bulk density distributions in soil cores scanned
164 horizontally. Otani et al (2010) correlated overall material density directly with CT grey value
165 following 150 kV peak energy scanning of fine to coarse Yamazuna sand, and used the results to
166 identify shear zones. Density changes in shear zones were also assessed by Desrues (2004) for fine
167 sand scanned using a ND8000 medical scanner. Derived void ratio values were reported, but not
168 their spatial distribution. Alvarez-Borges et al (2018) examined changes in chalk density following
169 model pile penetration using a modified 225 keV Nikon/Xtek HMX device. Scan energies were 200
170 keV, and the resulting images interpreted using a GV density calibration.

171

172 Alshibli & Hasan (2008) analysed the porosity of a medium sand specimen before and after shearing,
173 using a MSFC CT facility with a peak scan energy of 335 kV. A fairly uniform distribution of void
174 ratio was obtained within their specimen before compression, based on the void ratio map of a centre
175 section. Fonesca et al (2013) reported differences between void ratios in medium sand determined by
176 gravimetric measurement and X-ray CT scans (using a micro-CT scanners, developed by phoenix|X-
177 ray (GE)) of between 9 % and 25 %. They attributed the differences to the heterogeneity of the
178 specimen reducing the accuracy of comparisons between global gravimetric data with the local
179 values determined by CT scanning, and did not explicitly consider the possibility of the heel effect.

180

181 These studies show that it is possible to scan soil specimens at high energies and not experience the
182 heel effect. This could be due to one of a number of reasons, including the nature of the soil
183 properties determined from the analysis; the highly localised nature of the results obtained; other
184 sources of error obscuring the effect; or it simply having gone unnoticed and uncorrected. We will
185 show in this paper that manifestation of the heel effect depends on energy levels and soil grain size;

186 hence it may be that the materials in these studies were coarse relative to the energy levels used. In
187 such cases, the individual grains may be distinguishable directly so that the phase proportions can be
188 calculated without error. Given these uncertainties, it is important to demonstrate that the heel effect
189 can be non-negligible in soil imaging. Awareness of the artefact can then guard against future error
190 as soil scan energies increase further.

191

192 **3 EXPERIMENTAL APPROACH**

193 A comprehensive programme of experiments was carried out to demonstrate the presence of the heel
194 effect in soil imaging and understand the importance of different factors influencing the magnitude
195 of the artefact. Three different types of specimen were used; first, empty acrylic specimen containers
196 comprising two homogeneous phases (acrylic and air) that were not susceptible to the preparation-
197 induced variability that may occur with real soils. This allowed initial carefully controlled
198 experiments to be carried out to confirm the presence of the heel effect in the scanners utilised, in
199 ideal conditions without internal specimen variability.

200

201 Secondly, granular materials were used to assess the presence of the heel effect in images of soils,
202 and to show the effect of different grain sizes on the magnitude of the artefact. Leighton Buzzard
203 sand was used for this purpose, as it is well known and understood. Different particle sizes were
204 tested, from Fraction B to Fraction E. The third material tested was a clay - London Clay was used,
205 again because its properties are well-known.

206

207 **3.1 Materials**

208 Purpose-built circular cylindrical containers, of internal diameter 5mm, 8mm and 20mm and
 209 corresponding height 10mm, 20mm and 50mm respectively, were made by 3D printing. The
 210 containers were scanned empty or containing soil specimens.

211

212 Leighton Buzzard sand specimens of different grain sizes (Table 1) were prepared by air pluviation
 213 directly into the containers. Most of the specimens were prepared dry to allow the simplest initial
 214 assessments. However, wet pluviation was also used with some specimens to facilitate comparison
 215 with the London Clay, which could not be used dry. London Clay specimens were prepared by slurry
 216 deposition.

217

218

Table 1: Grain Sizes of Soil Materials

Descriptor	<i>Leighton Buzzard sand</i>				<i>London Clay</i>
	<i>LB-B</i>	<i>LB-C</i>	<i>LB-D</i>	<i>LB-E</i>	
	Coarse sand	Medium sand	Fine to medium sand	Fine sand	Silty clay or Clay
Range of mean grain size(μm)	1180 – 600	600 – 300	300 – 150	150 – 90	1 * <5 **
% coarser than range	10	10	10	15	NA
% finer than range	10	10	15	15	NA

219 * Zhang et al, 2016; ** King (1991)

220

221 **3.2 Scanning**

222 Four different X-ray CT scanners in the University of Southampton μ -VIS X-ray Imaging Centre
223 (University of Southampton, 2017) were used for the experiments. Initially the Nikon HMX ST 225
224 (designated “HMX”), Nikon/Metris custom Hutch (225kV peak modality) (“Hutch”) and Zeiss
225 Xradia Versa 510 (“Versa”) scanners were used to confirm the presence of the heel effect. Each of
226 these machines has a different X-ray target (Table 2) and hence were expected to show different
227 results with respect to the heel effect. After these initial confirmation scans, a range of sensitivity
228 experiments related to the heel effect were carried out using the Nikon/Metrix CT Benchtop 160 Xi
229 machine (“Benchtop”), which has a vertical reflection X-ray target.

230

231 *Table 2: Features of various scanners (see University of Southampton (2017) for more details)*

<i>CT Scanner</i>	<i>Short name in text</i>	<i>Target source</i>	<i>Expected heel effect</i>
Nikon/Metrix CT Benchtop 160 Xi	Benchtop	Reflection target mounted vertically	Vertical
Nikon HMX ST 225	HMX	Reflection target mounted vertically	Vertical
Nikon/Metris custom Hutch (225kVp modality)	Hutch	Reflection target mounted horizontally	Horizontal
Zeiss Xradia Versa 510	Versa	Transmission target	None

232

233 The materials and scan settings used for all the experiments are given in Table 3. The scan settings
234 are defined by the energy, which affects image contrast; the power and exposure time, which
235 controls the number of X-ray photons; the projection count, which ensures the intactness of the
236 image features; and the number of frames per projection, which is chosen to minimise movement
237 artefacts and reduce noise. Overall, the settings were chosen to maximise image quality within the
238 constraint of scan time during which the state of the specimen would not change significantly – a key

239 factor when the specimen is not in stasis. Of the scan parameters, only the scan energy would be
240 expected to influence the magnitude of the heel effect observed. Taking the first row of Table 3 as
241 an example, two frames per projection at 500 ms exposure means that every projection is made up of
242 the mean of two acquired images, each exposed for 500 ms. Consequently, the total exposure time
243 per projection is 1000 ms.

244

245 The confirmation scans were carried out on 5mm diameter specimen containers, either empty or
246 filled with Fraction E Leighton Buzzard Sand. These experiments were designed to show the
247 difference between the horizontal and vertical reflection targets in the Hutch and HMX machines,
248 and the presence or absence of the heel effect in the soil specimens in the HMX and Versa machines,
249 the latter of which uses a transmission target.

250

251 Subsequently, empty pots were used to assess the influence of scan energy on the magnitude of the
252 heel effect, and any impact of specimen size for a given set of scan settings. In this case, the
253 resolution changes according to the specimen size (Table 3). Finally, soil specimens of different
254 grain sizes and saturations were tested, again at the same scan settings.

255

256

Table 3: Experiment Details and Scan Settings

<i>Purpose</i>	<i>Container Size Dia/Ht (mm)</i>	<i>Specimen</i>	<i>Scanner</i>	<i>Peak Energy kV</i>	<i>Power W</i>	<i>Exposure ms</i>	<i>Frames per projection</i>	<i>Projection count</i>	<i>Resolution μm</i>
Confirm presence of heel effect	5/10	Empty	HMX	75	7	500	2	3142	5.8
	5/10	Empty	Hutch	75	7	500	1	1601	6.3
Confirm presence of heel effect in soils	5/10	LB-E dry	HMX	85	7	500	1	3142	3.9
	5/10	LB-E dry	Versa	85	7	4	1	3201	5.7
Demonstrate effect of scan energy	5/10	Empty	Benchtop	60	6	1067	4	1905	10.4
	5/10	Empty	Benchtop	80	6	1067	4	1905	10.4
	5/10	Empty	Benchtop	100	6	1067	4	1905	10.4
	5/10	Empty	Benchtop	120	6	1067	4	1905	10.4
	5/10	Empty	Benchtop	140	6	1067	4	1905	10.4
Investigate effect of absolute specimen size	5/10	Empty	Benchtop	80	7.5	534	4	1000	9.1
	8/20	Empty	Benchtop	80	7.5	534	4	1000	17.3
	20/50	Empty	Benchtop	80	7.5	534	4	1000	43.2
Investigate effect of soil grain size and saturation	5/10	LB-B dry	Benchtop	100	6	1067	4	1905	9.5
	5/10	LB-C dry	Benchtop	100	6	1067	4	1905	9.5
	5/10	LB-D dry	Benchtop	100	6	1067	4	1905	9.5
	5/10	LB-E dry	Benchtop	100	6	1067	4	1905	9.5
	5/10	LB-E saturated	Benchtop	80	6	1067	4	1905	9.5
	5/10	London Clay saturated	Benchtop	80	6	1067	4	1905	9.1

259 **3.3 Data Processing**

260 *3.3.1 Image Reconstruction*

261 During scanning, the specimens were rotated and radiographic projections (“X-ray images”)
262 taken with an equiangular spacing of 360 degrees divided by the projection count. The raw data
263 were then reconstructed using vendor-specific implementations of the Feldkamp Davis Kress
264 algorithm (Feldkamp et al, 1984) in Nikon X-TEK CT Pro 3D (Version XT 2.2 Service Pack
265 11), to provide a stack of 2D horizontal plane images together giving a representation of the 3D
266 image. The stack of images was used to evaluate the vertical heel effect using the grey value
267 (GV) and derived porosity distributions.

268 To evaluate the heel effect in the horizontal plane (when using the Hutch machine), the GV was
269 measured directly using the original 2D radiographic projections, before any reconstruction was
270 carried out. This is because reconstruction reduces the magnitude of the horizontal heel effect
271 when the reflection target is arranged horizontally. In this specific case, the artefact is largely
272 averaged out between reciprocal projections in radiographs oriented at 180 degrees.

273

274 *3.3.2 Image Processing and Thresholding*

275 After reconstruction, image processing for each soil specimen was limited to a region of interest
276 (ROI) centred on the axis of the specimen, to eliminate boundary or edge effects. Analyses of
277 empty specimen containers focused on small regions of the acrylic wall.

278

279 For the larger image data from the HMX, Hutch and Versa machines, the GV range of each
280 image stack was scaled to enhance the contrast between the different phases. Before further
281 image processing, the reconstructed raw image data were converted into 8-bit unsigned integer
282 format to reduce the computing time. Analysis of the relatively smaller image data from the
283 Benchtop scans was based on the raw 32-bit floating point data. For greyscale featured image
284 data, a lower GV corresponds to less dense and a higher GV to denser material. For example in
285 an 8-bit integer image, the range of greyscale is between 0 (completely “black”) and 255
286 (completely “white”), representing the least and the most dense materials, respectively. For
287 presentation of some images in this paper, contrast enhancement was applied to help illustrate
288 the heel effect.

289

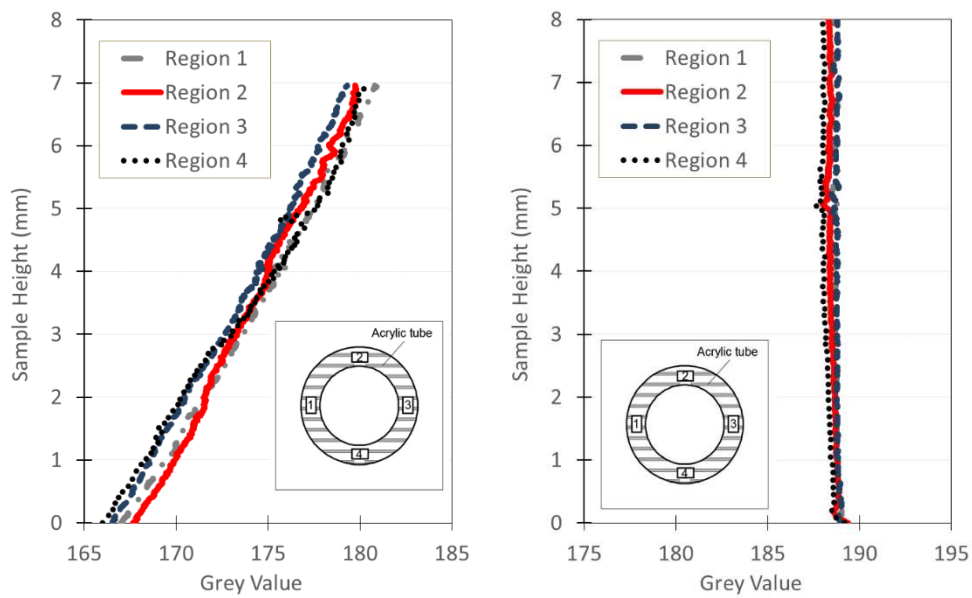
290 For quantitative assessment of the soil phase relationships, the scan data were thresholded to
291 differentiate the sand grains from the voids based on their density, as reflected in their GV. In
292 demonstrating the presence of the heel effect and presenting the results of the sensitivity analysis
293 (in Sections 4 and 5 respectively), the primary method of thresholding used is the Otsu method
294 (Otsu, 1979). This popular method was chosen because it is simple and straightforward to apply,
295 has a very fast analysis time (in the order of seconds), and is robust as it does not require the
296 selection of additional fitting parameters. The approach searches for the threshold that minimises
297 the intra-class variance based on the shape of the overall GV intensity histogram, and has been
298 shown to work well for two phase soil systems (e.g. Watanabe et al., 2012; Zhao et al, 2015)
299 where there is sufficient contrast between those phases. In Section 6, the Otsu method is
300 compared with more sophisticated adapted thresholding methods (e.g. Bernsen, 1986; Niblack,
301 1986) offering a possible way of removing the heel effect artefact.

302

303 4 DEMONSTRATION OF THE HEEL EFFECT

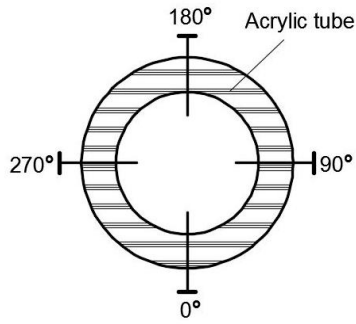
304 4.1 Vertical Heel Effect

305 Four regions of similar area at different orientations on the empty specimen container were
306 investigated (Figure 2a). As the density of the acrylic container was homogeneous, the GV
307 distribution should have been uniform. However, a marked gradient is apparent in the vertical
308 direction (Figure 2a) for the data gathered using the HMX. In contrast, the vertical distribution of
309 GV obtained using the Hutch is uniform when shown at the same scale (Figure 2b). This
310 confirms the presence of a vertical heel effect in the HMX scanner and the absence of the
311 vertical heel effect in the Hutch machine, which is equipped with a horizontal reflection target.



312

313 *Figure 2: Grey value (GV) distribution of the specimen container acrylic wall from (a) the HMX with*
314 *vertical reflection target and (b) the Hutch with horizontal reflection target.*



315

316

Figure 3: Sketch of the 2D slice showing the projections at four orthogonal angles

317

318 **4.2 Horizontal heel effect**

319

Four mutually orthogonal radiographs (at $\theta = 0^\circ, 90^\circ, 180^\circ$ and 270° ; Figure 3), over a 3 mm

320

height positioned above the horizontal centreline, were used to assess a potential horizontal heel

321

effect. The plots in Figure 4 show the GV profile for the ROI in each projection as a function of

322

the distance along the projection (in pixels). In each case, the GV for the air phase is high in the

323

centre of the plot, then falls where the projection intercepts the acrylic wall of the container. The

324

GV reaches a minimum at the internal boundary of the wall, and increases when the air phase

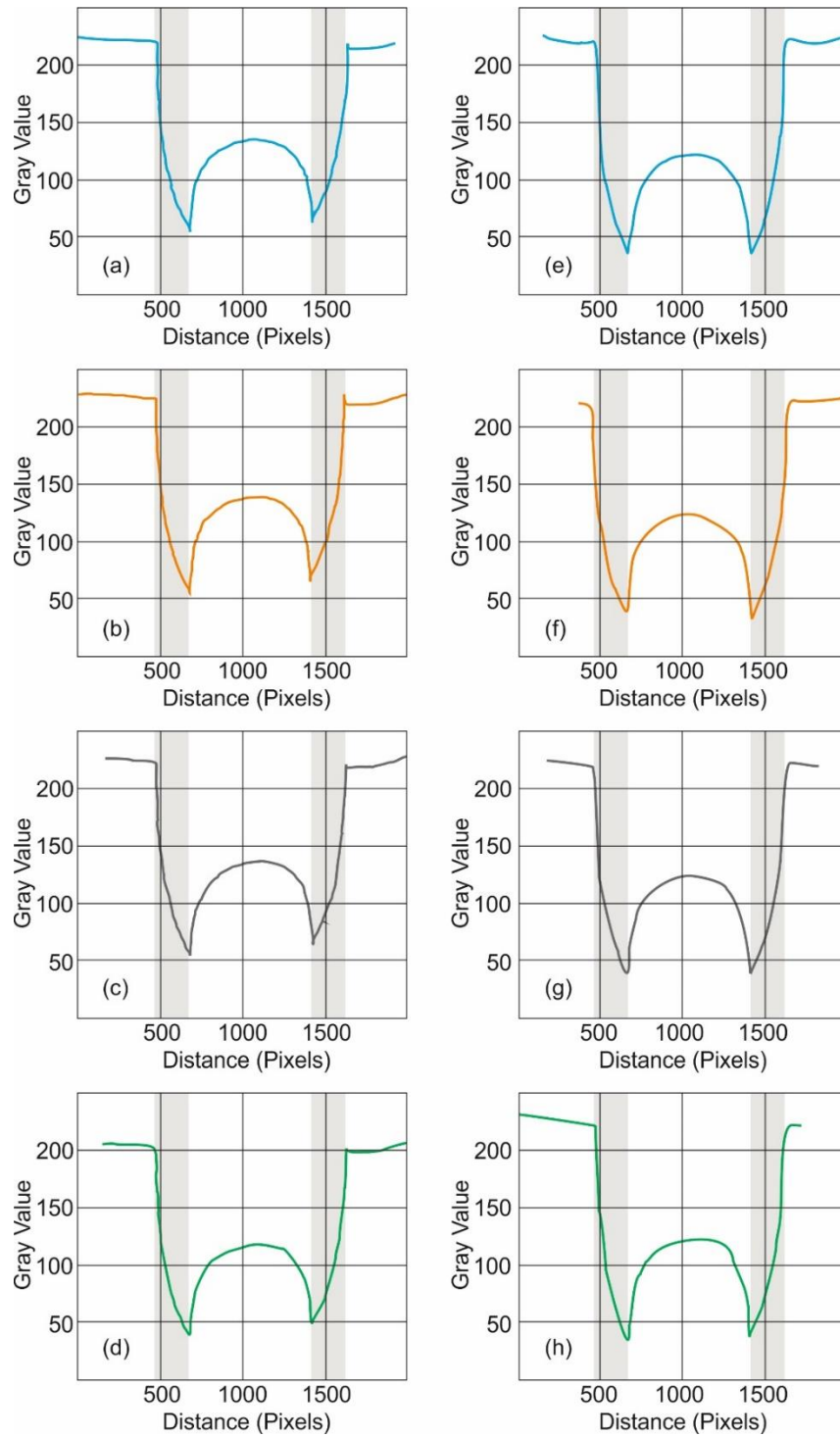
325

outside the container is reached. In the absence of the heel effect, each radiograph would be

326

expected to show reasonable symmetry.

327



328

329

330

331

332

Figure 4: Grey value profiles for the ROIs from four orientations: (a) 0° (b) 90° (c) 180° (d) 270° using data from the Hutch machine with horizontal reflection target; (e) 0° (f) 90° (g) 180° (h) 270° using data from the HMX machine with vertical reflection target. The shaded regions represent the acrylic wall of the specimen containers.

333 In the Hutch results (Figure 4a to Figure 4d), the minima on the left are always lower than those
334 on the right, regardless of the orientation of the scan. This is a manifestation of the horizontal
335 heel effect. In comparison, the HMX data (Figure 4e to Figure 4h) show a much more consistent
336 result. In Figure 4e and Figure 4g, the left hand minima are the same as those on the right. This is
337 because (e) and (g), and (f) and (h) (and also (a) and (c), and (b) and (d)) are pairs of mirror
338 projection images at the same orientation. The consistency of these pairs demonstrates that the
339 horizontal heel effect is not present in the HMX.

340

341 **4.3 The Heel Effect in Sand**

342 *4.3.1 Effect on Grey Value*

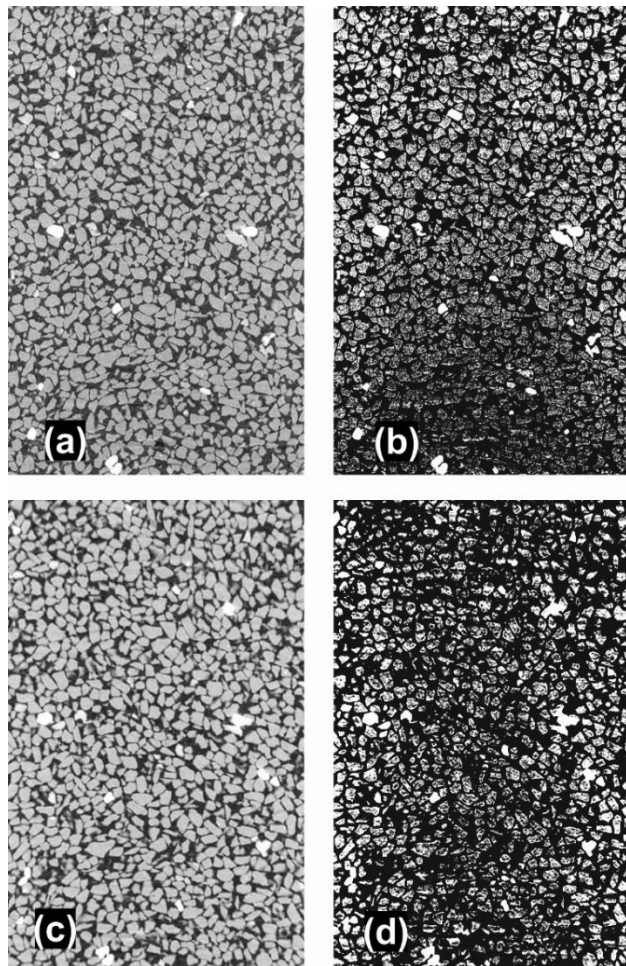
343 A specimen of dry Fraction E Leighton Buzzard sand in a 5mm diameter container was used to
344 assess the implications of the heel effect for imaging in soils. The HMX (vertical reflection
345 target) and Versa (transmission target) machines were used. To avoid the influence of cone-beam
346 artefacts (Zbijewski & Beekman, 2006; Hsieh et al, 2007) and reduce computing time, a section
347 about 4 mm in height around the middle of a 10 mm sand specimen was selected for analysis.
348 Figure 5 shows the resulting orthogonal images from the two scans, as originally processed
349 (Figure 5a, c) and with contrast enhancement (Figure b, d). The contrast enhancement reveals
350 the vertical heel effect as a gradient in greyscale. In Figure 5c there is a clear top to bottom
351 darkening of the image from the HMX. This effect is not noticeable with Versa.

352

353 Taking a slice from the HMX dataset at mid height and near the base (Figure 6) shows clearly
354 how the heel effect has effected the GV. Figure 6c also includes the GV histogram for the ROI

355 in each of the two slices. The two phase peaks are both shifted to the left in the lower darker
356 slice; and in addition, the overall range of GV is also reduced. This has implications for
357 thresholding as a reduced range may make it harder to separate phases. This is discussed below
358 and in Section 5, where the impact of specimen grain size is investigated.

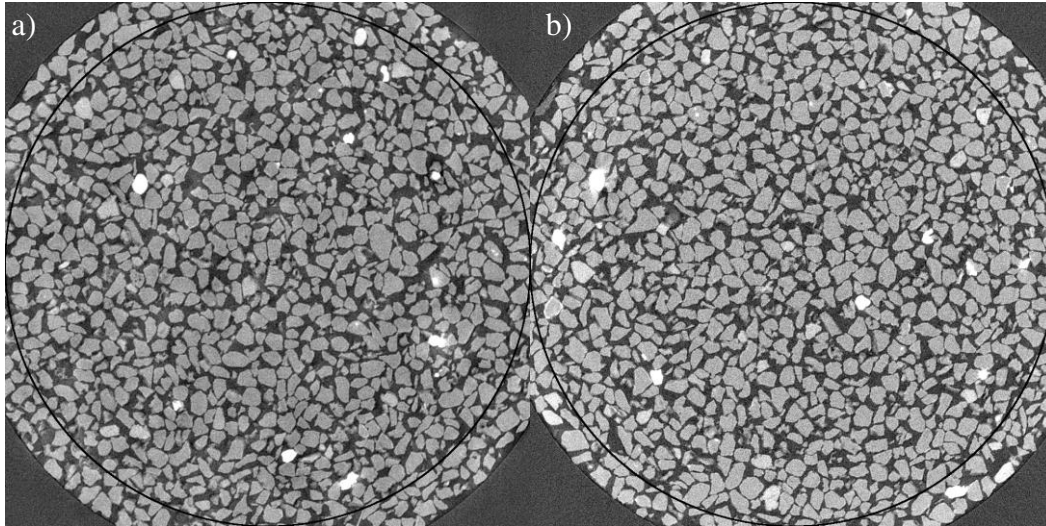
359



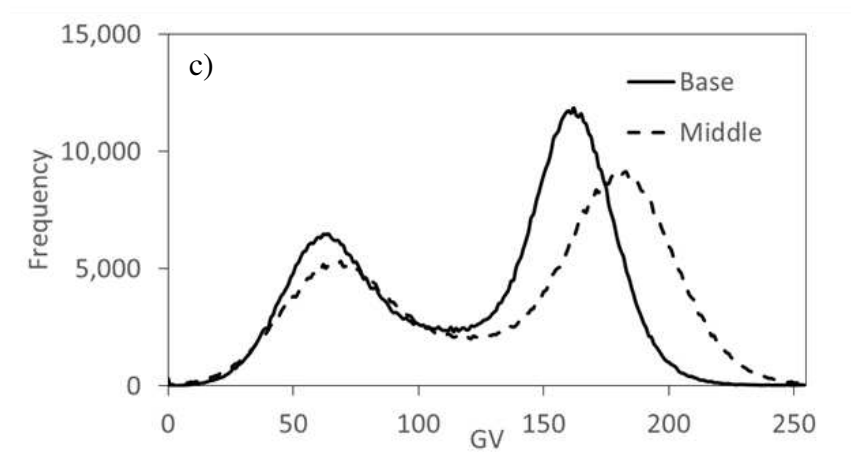
360

361 *Figure 5: Orthogonal images for dry Fraction E Leighton Buzzard sand a) scanned using the HMX; b)*
362 *scanned using the HMX with contrast enhancement, c) scanned with the Vera; d) scanned with the Vera*
363 *with contrast enhancement*

364



365



366

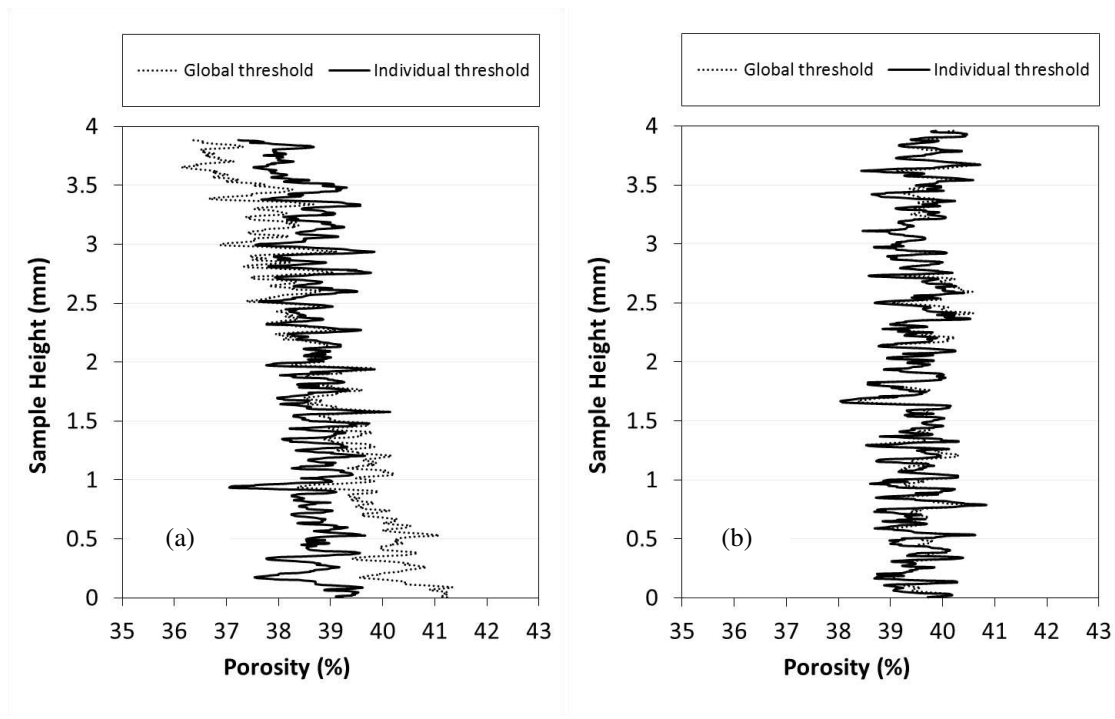
367 *Figure 6: Processed CT images for Fraction E Leighton Buzzard sand taken from the HMX for: a) a slice*
 368 *near the base; b) a slice near the centre of the specimen; c) GV histograms for the ROI in the two slices.*

369 **4.3.2 Effect on Thresholding and Phase Determination**

370 The magnitude of the heel effect may be quantified. Instead of working in terms of grey value (as
 371 in Sections 4.1 and 4.2 when considering only the specimen container), in this case the scan data
 372 were thresholded to determine the phase proportions. Data were processed using both global
 373 thresholding and local individual slice thresholding. Global thresholding is applied to each image
 374 slice based the GV thresholds determined from the overall GV intensity histogram for the entire
 375 specimen. In individual slice thresholding, each image slice is assessed with reference to the GV

376 thresholds determined from its own GV intensity histogram. Comparison of the results of the two
377 thresholding approaches can be used to evaluate the influence of the heel effect. Figure 7a shows
378 that with global thresholding, the data obtained using the HMX exhibit an apparent reduction in
379 porosity of about 5 % over the imaged specimen height (4 mm ROI). In contrast, the Versa data
380 (Figure 7b) show a uniform distribution of porosity with a variation of less than 1 % over the
381 same length scale.

382



383

384 *Figure 7: Porosity distributions from (a) HMX and (b) Versa scans of a 5 mm diameter dry Fraction E*
385 *Leighton Buzzard sand specimen*

386

387 The apparent decrease in porosity with height shown by the HMX is an artefact resulting from
388 the increasing gradient of the GV distribution caused by the vertical heel effect. It is shown later

389 that the degree of error introduced by the heel effect (about 5 % in this case) can be affected by
390 the applied energy and the extent of the partial volume effect.

391

392 In contrast, the porosities determined using individual thresholding are uniform (Figure 7a).

393 While this may at first sight seem to offer a suitable way of correcting for the heel effect, making

394 reference to only the GV data from each slice without knowledge of the remainder of the data

395 sets leaves the process open to errors and the potential to fail to identify real height dependent

396 trends. For example, when applying individual thresholding, a small number of high density

397 grains could cause the threshold value to be chosen at a GV range adequate for these particles

398 but inappropriate for the remaining grains. A similar effect could occur with partial saturation.

399 Hence individual thresholding is not recommended in practice. Individual thresholding is also

400 more computationally expensive than global thresholding, and is not always possible in cases

401 with a strong partial volume effect, where the GV must be calibrated for density or other

402 techniques applied (e.g. see Liu et al, 2017). Since global thresholding is also a more common

403 approach in practice, this means that the heel effect may be important in geotechnical analysis

404 for certain scan settings and soil specimen conditions. The magnitude of this effect and how it

405 may be countered are discussed in Sections 5 and 6 respectively.

406

407 The overall global porosity of the specimen was also determined from the HMX and Versa data

408 (Table 4). These results are consistent, with a variation of less than 1 %. Thus the heel effect

409 seems mainly to influence the apparent distribution of the specimen density, rather than the

410 overall value. This is another reason why the effect may not have been noticed in previous

411 research (see Section 2.3). In the present case, the overall porosity determined from gravimetric
 412 data was 39.3 %, which is in close agreement with the CT data.

413

414 It can therefore be concluded that the heel effect exists independently of specimen materials. If a
 415 reflection target is used in soil scanning, the artefact will influence the resulting image analyses
 416 of geotechnical phase relationships where knowledge of the spatial distribution is required. The
 417 heel effect can only be prevented directly by avoiding the use of reflection targets, especially
 418 those oriented vertically.

419 *Table 4: Overall porosity from HMX and Versa scans of a 5 mm diameter dry Fraction E Leighton*
 420 *Buzzard sand specimen*

<i>Porosity (%)</i>	<i>HMX data</i>	<i>Versa data</i>	<i>Corrected HMX data</i>
Otsu global threshold	38.8	39.6	38.9
Otsu individual threshold	38.6	39.5	38.8
Adaptive threshold (Bernsen)	38.4		
Adaptive threshold (Niblack)	38.6		
Gravimetric measurement		39.3%	

421

422

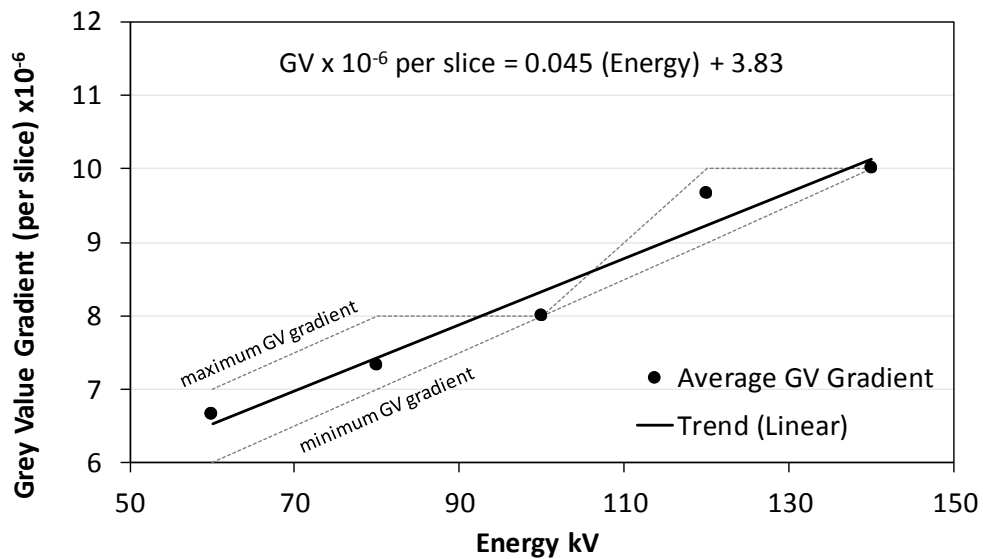
423 **5 RESULTS OF SENSITIVITY ANALYSIS**

424 **5.1 Influence of Scan energy**

425 As the heel effect is related to X-rays, energy is a key parameter and the heel effect is expected
 426 to be greater at higher energies. This was demonstrated by scans of an empty 5 mm internal
 427 diameter acrylic specimen container. Five different peak scan energies (60 kV, 80 kV, 100 kV,
 428 120 kV and 140 kV) were applied using the Benchtop machine. Three scans were carried out for
 429 each energy level, giving fifteen scans in total. Four ROIs on the acrylic wall of the container

430 from each set of data were used for analysis. The mean gradient of the four GV curves was
 431 determined in each case, with the average of the results from the three scans at the same energy
 432 being taken as representative. The analysis was based on the raw 32-bit data to avoid potential
 433 inconsistencies arising from data conversion. 800 slices without any cone-beam artefacts were
 434 selected for the analysis of the gradient of GV distribution.

435



436

437 *Figure 8: Average and range of grey value gradients for a 5mm internal diameter specimen container*
 438 *using different scan energies in the Benchtop machine.*

439

440 The grey value gradient induced by the heel effect is plotted against the scan energy in Figure 8,
 441 in which the data points represent the average gradient for each energy, the solid line is the trend,
 442 and the dashed lines give the minimum and maximum GV gradients determined. As expected a
 443 positive linear relationship is obtained. This is because as the scanning energy increases, the X-
 444 rays will penetrate deeper into the target material with more generated photons being attenuated,

445 causing a wider GV magnitude range to be reflected on the X-ray detector. In this case doubling
446 the scan energy results in 145% of the initial GV gradient.

447

448 **5.2 Influence of Specimen size**

449 To assess the effect of specimen size on the heel effect, empty acrylic containers having internal
450 diameters of 5 mm (10 mm in height), 8 mm (20 mm in height) and 20 mm (50 mm in height)
451 (with a wall thickness of 1 mm in each case) were scanned under the same imaging conditions
452 and with the same scan settings (Table 3). When the containers were scanned, they filled entirely
453 the field of view of the radiograph. This means that larger containers were scanned to a
454 correspondingly lower resolution, with the container wall represented by fewer voxels.

455

456 Analysis of the empty containers was based on the same four ROIs on the acrylic wall as before,
457 to obtain the gradient of GV distribution from the raw 32-bit data. The gradients of the GV
458 curves are the same in each case (2×10^{-6} in terms of GV/slice), showing that the magnitude of
459 the heel effect is not dependent on the overall specimen size. It should be noted that the GV
460 gradients are slightly less than those shown in Figure 8 for the same scan energy. This is because
461 the other scan conditions (power, exposure, resolution) are not identical (Table 3).

462

463 **5.3 Influence of Soil Grain Size**

464 The potential influence of soil grain size on the heel effect was assessed in dry and wet
465 conditions using the Benchtop machine. Leighton Buzzard sands (Fractions B, C, D and E) and
466 London Clay were used. Dry sand specimens of different grain sizes and wet soil specimens of

467 Fraction E sand and London Clay were tested separately at two different peak scan energies, 100
468 kV and 80 kV, respectively (Table 3). A consistent ROI was chosen within each specimen,
469 covering 8 mm of the total 10 mm specimen depth. Global and individual thresholding
470 approaches were used to determine the porosity within the ROI. While it is accepted that
471 thresholding will not be strictly appropriate for the clay specimen owing to the partial volume
472 effect (mean grain size $<5\mu\text{m}$ and scan resolution $9.5\mu\text{m}$), the same approach was applied
473 throughout for consistency.

474

475 The measured vertical porosity gradients are given in Table 5. The errors induced by the heel
476 effect are quantified by the ratio (Column 6) of the porosity gradients obtained using global
477 (Column 5) and individual (Column 4) thresholding. As individual thresholding is unaffected by
478 the heel effect, this ratio provides a consistent measure of the magnitude of the artefact. The
479 results in Columns 5 and 6 show that the influence of the heel effect increases with decreasing
480 grain size. The effect is apparent in both the dry and the wet specimens, although to a lesser
481 extent in the latter, for a given (smaller) grain size.

482

483 The changes in GV gradient (Column 7) for specimens composed of smaller grains are more
484 varied, but there is a clear trend of greater gradients at smaller grain sizes for the same or similar
485 scan energies. The GV gradient for the finer specimens is at least twice that for the coarsest. The
486 data in Table 5 demonstrate that the heel effect is magnified by the thresholding process,
487 particularly when the grain size is small such that it becomes comparable with the spatial
488 resolution.

489

490

Table 5: Porosity gradients for different soil grain sizes

Specimen Material	Comments	Featured peak scan energy	Porosity Gradient (mm^{-1})		Ratio Global: Individual	GV Gradient ($\times 10^{-6}$ GV per slice)
			Individual threshold	Global threshold		
Dry LB-B		100 kV	0.416	0.416	1	17.4
Dry LB-C		100 kV	-0.377	-0.842	2.2	45
Dry LB-D		100 kV	-0.120	-0.787	6.6	33.3
Dry LB-E		100 kV	-0.064	-1.797	28.1	51.6
Saturated LB-E		80 kV	-0.483	-1.729	3.58	44.5
Saturated London Clay		80 kV	-0.076	-1.950	25.66	30.6
Dry LB-E	Original Data (Figure 7)	85 kV	-0.1152	-0.9615	8.4	18807
Dry LB-E	Corrected Data (Figure 13)	85 kV	-0.0944	-0.0635	0.7	56

491

492

493 The observed trend is due to two factors. As the grain sizes reduce with respect to the scan
494 resolution, additional errors will be introduced in thresholding owing to the partial volume effect.
495 These errors will be compounded by the heel effect, which increases the underlying apparent GV
496 gradient and hence further impedes the ability of thresholding to resolve the phase proportions
497 accurately. Referring to Figure 6, it can be seen that the heel effect has compressed the GV
498 histogram of a slice near the base of specimen where the image is darkened. This makes
499 thresholding more difficult and will introduce additional errors. For the different grain sizes
500 scanned, the porosity difference over the height of the sample (10mm) due to the heel effect
501 could be over 20%, i.e. an error in excess of $\pm 10\%$. The systematic nature of this error makes it
502 especially important to identify.

503

504 Also included in Table 5 are the porosity gradients from the confirmation scans (the data shown
505 in Figure 7). This case shows a lower global to individual thresholding ratio than the
506 corresponding grain size in the sensitivity analysis, which is to be expected given the better
507 resolution, lower scan energy and other differences in scan settings, but nonetheless the errors
508 are in the same order of magnitude. The differences in scan settings mean that the GV gradient
509 should also not be compared directly with the sensitivity scans, but with the corrected data. The
510 corrected data (discussed in Section 6) show a substantially reduced heel effect, with the artefact
511 eliminated almost entirely. Owing to the use of a different machine (HMX) for these scans, the
512 GV gradients are not directly comparable with the sensitivity scans carried out using the
513 Benchtop machine. Nonetheless the same pattern is shown, with the GV gradient almost
514 eliminated by the correction method.

515

516 **6 CORRECTION FOR THE HEEL EFFECT**

517 Section 4 has demonstrated the presence of the heel effect in sand and other materials. Section 5
518 has showed how the effect is influenced by scan energy and the grain size of the soil scanned,
519 and that the heel effect can be of significance in geotechnical problems. The simplest way to
520 avoid the heel effect is to use transmission target CT scanners. This will prevent the introduction
521 of the heel effect into the resulting radiographs. However, depending on available equipment,
522 this option will not always be possible. Therefore two approaches to correcting for the heel
523 effect, and the results of their application, are presented in this Section. First, the role of
524 advanced adaptive thresholding is considered. This could potentially give derived soil phase
525 proportions unaffected by the heel effect, although the original artefact will remain in the CT

526 images and GV data. Secondly, a universal correction technique, termed the “self-wedge”, is
527 proposed. This correction is applied to the radiographic projections directly to provide a
528 completely clean dataset for subsequent analysis.

529

530 **6.1 Adaptive Thresholding**

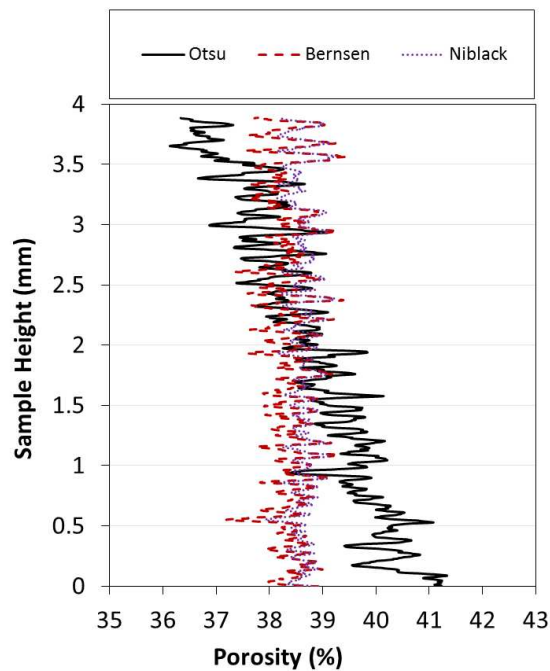
531 Adaptive thresholding computes the threshold for each voxel in an image in accordance with GV
532 information from neighbouring voxels. As such, there are similarities to local thresholding and
533 the application of the Otsu method on a slice by slice basis (e.g. the results presented in Section
534 4.3.2). Adaptive thresholding uses a local area (either circular or rectangular) with the target
535 voxel at its centre. As well as defining the size of this area, several of the more sophisticated
536 approaches (e.g Bernsen, 1986; Niblack, 1986; Phansalskar et al, 2011) also require specification
537 of one or two additional fitting parameters. Owing to the use of local GV information to
538 compute thresholds, adaptive thresholding may only be suitable if there are not expected to be
539 other in situ characteristics or changes within a soil specimen that would be masked by the
540 adaptive process. It also takes longer to carry out than global thresholding.

541

542 A number of adaptive thresholding methods were tested using the original HMX scans on
543 Fraction E Leighton Buzzard sand, to see if they could offer a way to overcome the heel effect.
544 Methods based on using adjusted local average GV for the threshold were found to remove the
545 artificial GV gradient induced by the heel effect, but tended to substantially over-estimate the
546 porosity. However, two techniques, one based on local contrast thresholds (Bernsen, 1986) and
547 one that uses the standard deviation as well as the mean in determining the local threshold
548 (Niblack, 1986), were found to work well.

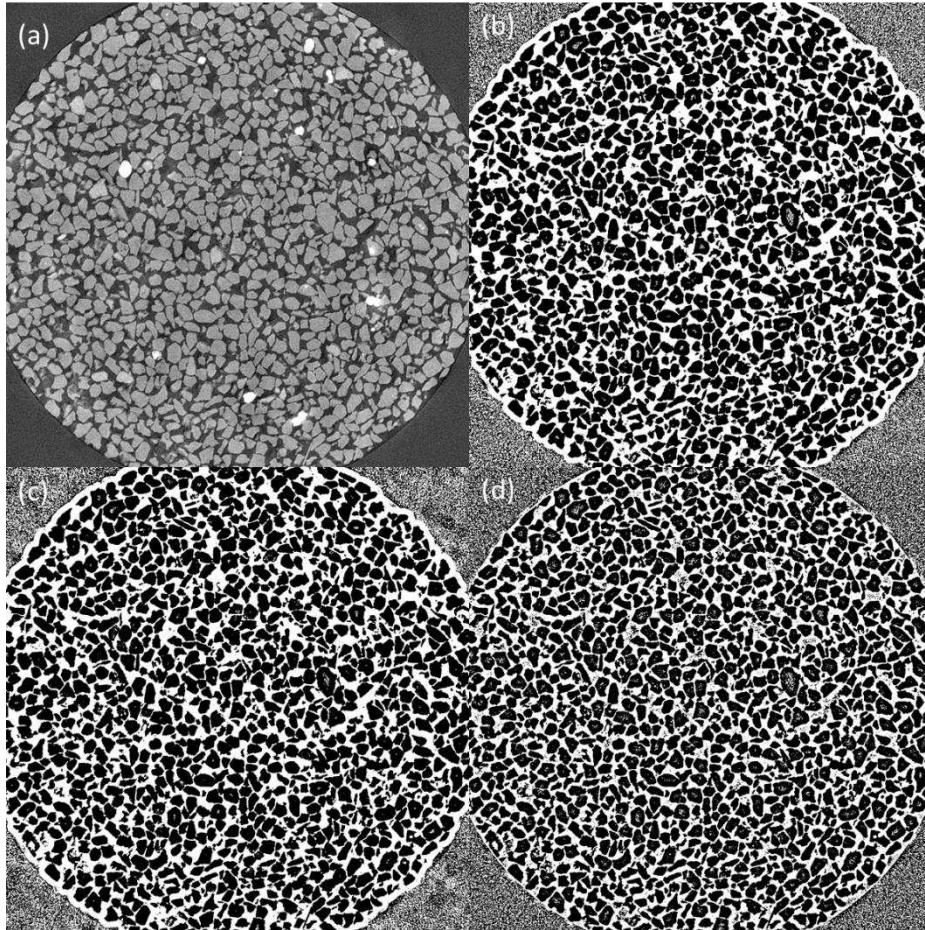
549

550 Figure 9 compares the porosity distributions determined using the Otsu method (global
551 thresholding) with those using the adaptive methods of Bernsen (1986) and Niblack (1986). The
552 segmented images using the different approaches are shown in Figure 10. It can be seen that the
553 two adaptive thresholds produce sensible porosity outcomes that are very similar to the
554 individual thresholding using Otsu shown in Figure 7. Furthermore, the average porosities for
555 the two adaptive methods are 38.4 % and 38.6 % respectively, which compares favourably with
556 the Otsu and gravimetric measurement results given in Table 4. This both confirms Otsu as a
557 sensible approach in these types of soils and shows more generally that adaptive thresholding is
558 capable of negating the heel effect.



559

560 *Figure 9: Porosity distributions from the Fraction E Leighton Buzzard sand scanned using the HMX,*
561 *illustrating the effect of different thresholding methods.*



562

563 *Figure 10: Comparison of images of Fraction E Leighton Buzzard sand from the HMX using different*
564 *thresholding approaches: (a) original image (b) binarised image using Otsu (individual slice)*
565 *thresholding (c) Bernsen adaptive thresholding (d) Niblack adaptive thresholding.*

566

567 However, two other adaptive thresholding methods (Phansalskar et al, 2011; Sauvola &
568 Pietaksinen, 2000) based on the approach of Niblack (1986) were found to underestimate the
569 specimen porosity, despite also removing the artificial GV gradient. This shows that care is
570 always required in selecting thresholding methods and their associated fitting parameters. Given
571 this, it is also desirable to have a universal method for correcting the heel effect, which acts on

572 the original GV data. This will enable the application of global thresholding and implementation
573 of the fullest range of data and image analysis approaches.

574

575 **6.2 Self-Wedge Correction**

576 *6.2.1 Approach*

577 As there are limitations to the use of local and adaptive thresholding methods, a correction to
578 eliminate the errors introduced by the heel effect in the GV data is needed. This will permit the
579 application of a wider range of thresholding and other image processing techniques. Existing
580 correction methods (Section 2.2) may require specific knowledge about the nature of the target,
581 or the placement of compensation filters within the scanner itself. However, a more universal
582 approach would be beneficial in that it could be applied with no special prior knowledge and
583 carried out any time after the scan data had been obtained. The “self-wedge” correction
584 proposed in this paper fulfils this need. It is based on proposals by Ketcham & Carlson (2001)
585 for use with beam hardening. However, application of the technique to the scan data for the 5
586 mm diameter specimen of Fraction E dry sand in the HMX machine will be the first use of the
587 self-wedge correction for the purpose of countering the heel effect.

588

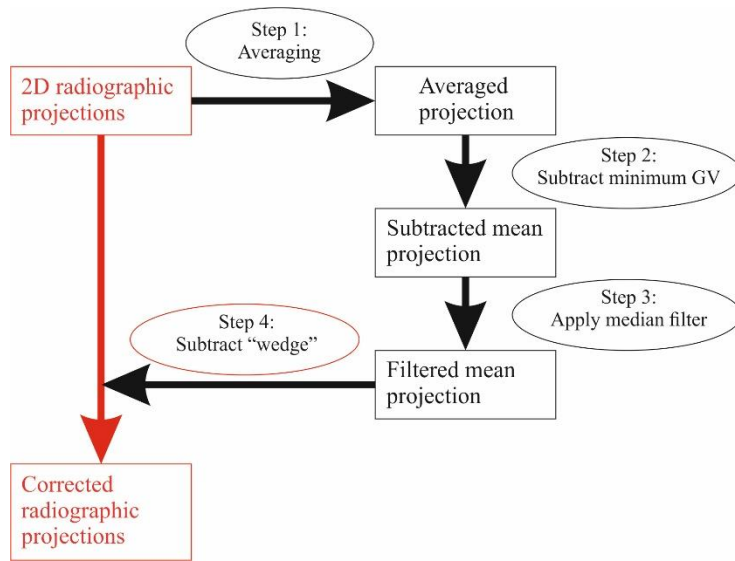
589 The self-wedge correction is based on an X-ray signal calibration method known as a “wedge”
590 (Ketcham & Carlson, 2001), wherein some CT protocols scan a wedge of material of known
591 dimensions or uses the specimen itself (i.e. a “self-wedge”) to provide the correction by taking
592 the mean across all projections to calibrate the signal. In this study, the correction method
593 initially averages the 2D radiographic projections for all angles (Figure 11, Step 1). Then the

594 minimum GV in the average projection is subtracted from the averaged projection itself (Figure
595 11, Step 2), meaning that the average projection, which will be used to correct the original
596 radiographs, has a minimum value of 0. This minimises the deviation from the original grey
597 values.

598

599 The process of rotational averaging described above has the potential to introduce artificial ring
600 artefacts. To counter this a median filter is applied (Figure 11, Step 3). A significant kernel size
601 can be applied to the average angular projection, with the kernel size is chosen such that small,
602 high-density artefacts in the mean image can be eliminated. Typically, a kernel value of 5 would
603 suffice, but this may be larger. If a median filter is insufficient then a polar transformation
604 coupled with a Fast Fourier Transform bandpass filter can be applied for the elimination of rings.

605 The filtered average projection is then subtracted from each original projection in turn (Figure
606 11, Step 4), ensuring that any grey values falling outside the grey level boundaries are set to that
607 boundary value (for example, if the grey value range is $0 < GV < 65535$, then a value of -3
608 would be set to 0, a value of 420 would remain at 420, and a value of 65577 would become
609 65535). If the resulting corrected images appear too dark, it is also possible to scale them at this
610 stage. For subsequent analysis dependent on the absolute grey values, a calibration step can also
611 be included after Step 4 to restore the original grey values. Calibration can be carried out by
612 comparing the corrected image with a single slice from the original, or by a complete Hounsfield
613 unit calibration scan with the same style specimen container (Feeman, 2015).



614

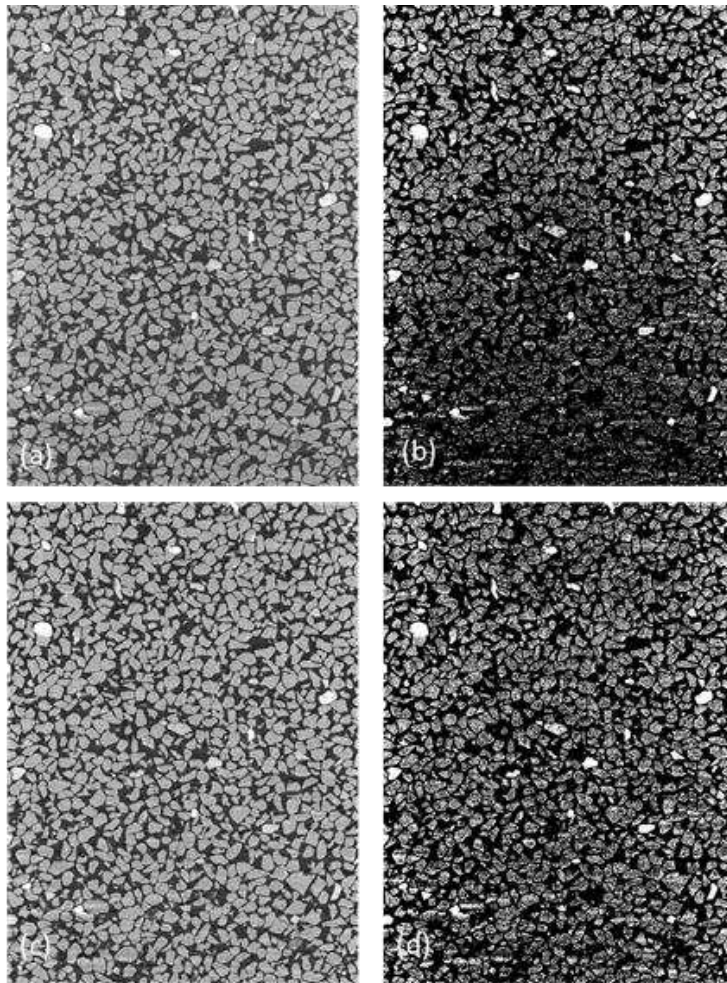
615

Figure 11: Flowchart for self-wedge correction process

616

617 6.2.2 Results

618 Applying the correction to the 2D projection scan data for the Fraction E dry sand specimen
 619 scanned in the HMX results in a small visible change in absolute GV in the reconstructed image
 620 data. This is illustrated in Figure 12a and Figure 12c for an orthogonal image. Also shown in
 621 Figure 12 are contrast enhanced versions of the images before and after the self-wedge
 622 correction. Figure 12b shows the GV gradient resulting from the heel effect, which is absent in
 623 the corrected image of Figure 12d.



624

625 *Figure 12: Comparison of orthogonal images; (a) before correction (b) before correction with contrast*
626 *enhanced (c) after correction (d) after correction with contrast enhanced.*

627

628 Thresholding was carried out on the HMX data after the self-wedge correction had been applied.

629 Individual and global thresholding of the corrected data using the Otsu method resulted in near

630 identical porosity profiles with negligible heel effect as shown in Table 5. The resulting porosity

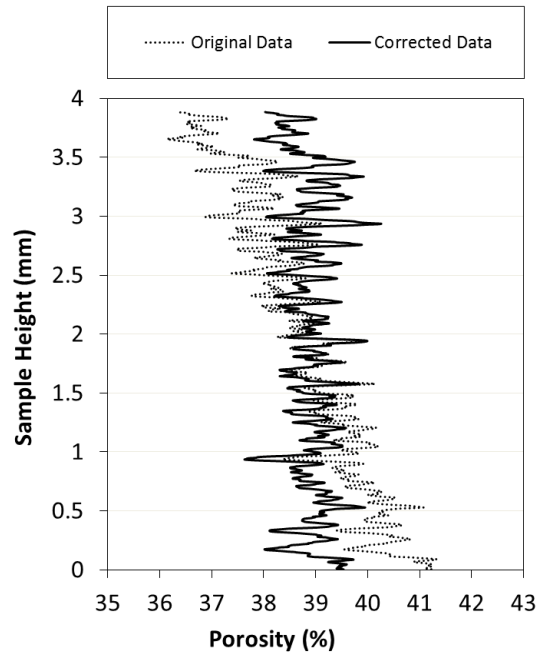
631 distribution profiles from global thresholding for the original and corrected data are shown in

632 Figure 13. There is a clear difference between the original and corrected data. The corrected data

633 compare favourably with both the results from the Versa machine (Figure 7), which uses a

634 transmission target, and the local and adaptive thresholding (Figure 9). The overall porosity

635 obtained from the corrected data is unchanged at 38.9 %, which is consistent with the results
636 obtained from the Versa (39.6 %) and gravimetric assessment (39.3 %) shown in Table 4. Thus
637 the self-wedge correction method is shown to have successfully eliminated the heel effect.



638
639 *Figure 13: Porosity distributions using global thresholding for the original and wedge correct data.*
640 *Specimen is 5 mm diameter Fraction E Leighton Buzzard dry sand scanned using the HMX machine.*

641

642 7 SUMMARY AND CONCLUSIONS

643 The heel effect, an artefact of CT imaging that occurs with scanners using a reflection target for
644 X-ray generation, has been identified in images of soil specimens for the first time. The heel
645 effect results in a gradient in the distribution of the grey value, which is used to determine
646 specimen density, at the detector. Therefore, it may also manifest as an artificial gradient in
647 derived geotechnical parameters such as porosity. Owing to the nature of image reconstruction,
648 the effect will be more significant for reflection targets orientated vertically, and seen especially

649 as a vertical parameter gradient. The heel effect is more significant at high scan energies, hence
650 is relevant to soil mechanics research as this field moves into higher energy temporal scanning;
651 for example for imaging changes in water content or density in response to thermal, chemical or
652 mechanical loading in specimens not in a state of stasis.

653

654 The heel effect has been demonstrated to occur in specimens comprising different fractions of
655 Leighton Buzzard sand and London Clay. For a dry Fraction E sand specimen scanned using the
656 HMX machine, the heel effect gave rise to an apparent gradient in porosity, reducing by about
657 5 % in absolute terms from the top to the bottom of a 4 mm high region, even though the sand
658 specimen was in reality essentially uniform. Subsequent sensitivity scans using the Benchtop
659 machine over a range of grain sizes showed that porosity gradient errors of up to at least $\pm 10\%$
660 could be easily generated in a 10 mm high specimen. Thus the heel effect is potentially
661 significant in geotechnical applications of X-ray CT techniques and may give rise to errors, for
662 example in the measurement of the distribution of phase proportions in soils.

663

664 It has been shown that specimen size does not significantly influence the heel effect. However
665 for a given scan energy, the soil grain size does. Finer, more densely packed soils are subject to
666 greater heel effect errors, especially when using a thresholding approach for phase determination.
667 The effect will increase in magnitude as the grain size reduces with respect to the scan resolution
668 and the partial volume effect becomes more significant.

669

670 To prevent the heel effect entirely reflection targets, especially those orientated vertically, should
671 be avoided. In some circumstances the effect may be negated by the use of local or adaptive
672 thresholding. Alternatively, a simple-to-apply self-wedge correction technique has been
673 developed to remove the error in GV data caused by the heel effect. Corrected HMX porosity
674 profiles compare favourably with those from the Versa machine, which has a transmission target.
675 Importantly, the proposed new correction can be applied universally post-scanning to any data
676 for which the projections are available.

677

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680 Engineering, the Doctoral Training Centre at University of Southampton and EPSRC
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682 Southampton repository (DOI to be confirmed). The authors are also grateful for the contribution
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- Figure 1: Schematics of X-ray (a) reflection target and (b) transmission target**
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833 **Table 1: Grain Sizes of Soil Materials**

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