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1 1. INTRODUCTION

2 Although subaqueous mass movements can develop on very gentle slopes ($\ll 2^\circ$), they
3 can run out for extremely long distances across the deep ocean floor (Talling et al.,
4 2007). In most cases the slope failures leave an empty scar in the source region with
5 almost all residual strength material evacuated (Krastel et al., 2018; Mountjoy and
6 Micallef, 2018). As source area slopes are predominantly lower than the friction angles
7 of the landslide materials (e.g. Urlaub et al., 2015), the generation of pore pressures
8 exceeding hydrostatic pressure are a likely loading mechanism. But, in some instances,
9 the debris from subaqueous slope failures only moves a short distance before coming
10 to rest despite a lack of buttressing or obvious decrease in slope gradient (Locat and
11 Lee, 2000; Micallef et al., 2013). Such small-displacement mass failures have been
12 found in a wide range of seafloor environments including passive margins (Baeten et
13 al., 2014) and active margins (Micallef et al., 2016). These indicate that there are likely
14 to be common mechanisms for the movement and arrest over a short distance of certain
15 seafloor failures.

16
17 The high pore pressures that often initiate movement can occur from widely recognised
18 processes including undrained cyclic loading during earthquakes (e.g. Sassa et al.,
19 2012), rapid sediment burial (e.g. Stigall and Dugan, 2010) and as result of focused
20 fluid flow (e.g. Dugan and Flemings, 2000; Elger et al., 2018). In addition, high pore
21 pressures may be generated by more complex processes involving for example gas
22 liberation from hydrate dissociation (Riboulot et al., 2013).

23
24 Despite the advances in understanding potential causes of submarine landslides, the
25 fundamental processes controlling their movement remain poorly constrained

26 compared with their terrestrial counterparts. Over the last century, terrestrial landslides
27 have been shown to display a wide array of movement behaviour, ranging from slow
28 creep ($\leq 1 \text{ mm a}^{-1}$) (e.g. Mansour et al., 2011), through episodic sliding and stick-slip
29 ($\approx 1 \text{ cm a}^{-1}$) (e.g. Allison and Brunsden, 1990) to acceleration to catastrophic failure
30 ($\gg 1 \text{ m s}^{-1}$) (e.g. Kilburn and Petley, 2003). Such complex movements are commonly
31 associated with pore fluid pressure-induced changes, as has been hypothesised for
32 subaqueous landslides.

33

34 The availability of high resolution landslide monitoring records onshore has allowed
35 detailed acceleration phases in multiple landslides to be examined in a variety of
36 materials (e.g. Schulz et al., 2009; Massey et al., 2013; Carey et al., 2015). These studies
37 distinguish two distinct styles of movement (Petley et al., 2002). The first movement
38 style is brittle, characterised by a distinct hyperbolic acceleration in displacement rate
39 to failure. These movement patterns can be examined by plotting in $1/v - t$ space (where
40 v is velocity) and yield a negative linear trend to failure which generally results in rapid
41 accelerations and catastrophic landsliding (e.g. Voight, 1988; Fukuzono, 1990; Petley
42 and Petley, 2006). Conversely, the second style of movement is ductile, characterised
43 by an exponential acceleration to a constant strain rate which when analysed in $1/v - t$
44 space produces a distinct asymptotic trend (Petley et al., 2005b).

45

46 Specialist laboratory testing approaches have been used to simulate landslide failure
47 conditions by increasing pore water pressures at constant normal and shear stress (e.g.
48 Brand, 1981; Anderson and Sitar, 1995; Zhu and Anderson, 1998; Dai et al., 1999;
49 Orense et al., 2004; Petley et al., 2005a, Ng and Petley, 2009; Carey and Petley, 2014).
50 These approaches confirm that movement styles are controlled by mechanisms of

51 deformation occurring within shear zones (Petley et al., 2005a; Ng and Petley, 2009;
52 Carey and Petley, 2014). In cases where brittle shear surface development occurs, the
53 hyperbolic acceleration to final failure is observed in the experiments, whilst
54 exponential acceleration occurs in landslides undergoing ductile deformation (Petley et
55 al., 2005a). Numerous conceptual models have been proposed to explain these different
56 strain responses to changes in stress state, ranging from catastrophic failure driven by
57 micro-cracking and rapid shear-surface propagation (e.g. Petley et al., 2005a; Viesca
58 and Rice, 2012) to slower, steady landslide motion in response to shear-zone dilation
59 and subsequent pore-pressure feedback (e.g. Iverson, 2005).

60

61 Given increases in pore-water pressure may drive either of movement-arrest or rapid
62 runout behaviour in terrestrial landslides, depending on the material response, similar
63 behaviour should be expected in subaqueous landslides. The mechanisms controlling
64 the transition from steady slow movement to rapid failure observed in specialist
65 laboratory tests may be key to determining the behaviour of a given subaqueous
66 landslide. However, very few small displacement or reactivated subaqueous landslides
67 have been observed and investigated.

68

69 Two subaqueous landslides on the northern Hikurangi Margin located off the east coast
70 of New Zealand have geomorphological characteristics that indicate movement-arrest
71 behaviour may be occurring (Mountjoy et al., 2009; Micallef et al., 2016) making them
72 suitable for investigating this behaviour in subaqueous slopes. In this study, we seek to
73 constrain better the potential movement mechanisms in short displacement and
74 reactivated subaqueous landslides by conducting novel laboratory experiments on
75 sediment samples collected from the shallow sedimentary sequence on the northern

76 Hikurangi Margin area. We use a Dynamic Back-Pressured Shearbox (DBPSB) to
77 accurately replicate in-situ stresses in the submarine slopes to explore the potential
78 strain response of the landslide when subject to elevated pore water and gas pressure.

79

80 2. SAMPLE AREA

81 An area of the upper continental slope on the Hikurangi Subduction Margin, off the
82 coast of Gisborne, New Zealand, hosts a number of landslides where limited
83 displacement has occurred following their initial failure (Mountjoy et al., 2009;
84 Micallef et al., 2016) or where repeated reactivation has been hypothesised (Mountjoy
85 et al., 2014) (Figure 1 A and B). The landslides occur within an active subduction zone
86 experiencing regular tectonic activity (Wallace and Bevan, 2010; Wallace et al., 2012).
87 The upper continental slope is comprised of Miocene to Recent slope basin sequences
88 (Mountjoy et al., 2009). A gravity core profile down the length of the slides (the
89 Tuaheni Landslide Complex) shows that the upper few metres of sediment are
90 dominated by mud to sand sized particles from hemipelagic drape, reworked landslide
91 debris and airfall tephra (Kuhlmann et al., 2018).

92

93 The Tuaheni Landslide Complex (TLC) comprises an area of approximately 145 km²
94 which is sub-divided into Tuaheni North and Tuaheni South, separated by a 2 km wide
95 spur (Figure 1 A). While Tuaheni North is characterised by multiple evacuated
96 landslide scarps, Tuaheni South is characterised by a large debris apron which has a
97 distinct scarp and bench topography and features indicative of lateral, extensional and
98 compressional deformation (Mountjoy et al., 2009; 2014). The base of gas hydrate
99 stability has been imaged beneath the TLC's extensional domain and it has been
100 suggested that the gas hydrate system may play a role in deformation of the landslide

101 mass (Mountjoy et al., 2014; Crutchley et al., 2016). Although no movement
102 measurements are available, the morphology of Tuaheni South is similar to slow-
103 moving landslide complexes observed in terrestrial environments such as earthflows
104 and mudslides (e.g. Hungr et al., 2014) which occur in similar fine-grained sediments
105 and are often subject to episodic remobilisation (e.g. Alison and Brunsden, 1990).
106 Seismic-reflection surveys across the landslide immediately north of TLC (Figure 1 C)
107 indicate that there is free gas in the proto basal failure surface, and adjacent to the
108 landslide, but gas is not observed beneath the landslide body (Micallef et al., 2016).
109 This suggests that free gas may have been present in the slope prior to the landslides
110 and that the gas migrated out of the sediment sequence during and/or after failure.
111 Micallef et al., (2016) concluded that overpressure in the slope sequence may have
112 contributed to bringing the slope to the point of failure, and that once the landslide
113 moved and dilated, the gas pressure reduced, and further failure was arrested. To further
114 test this interpretation required measurement of relevant geotechnical data, and this
115 provided an opportunity to assess different mechanisms that result in small-
116 displacement slope failures.

117

118 3. MATERIALS AND METHODS

119 We performed a suite of conventional laboratory experiments to determine the physical
120 and geomechanical characteristics of the shallow materials collected from the landslide
121 complex (Tables 1 and 2). Sediment samples were collected from the crest of the
122 landslide (Tan1404-02) and from within an extensional domain farther downslope
123 (Tan1404-10) (Figure 1 A) using a 100 mm diameter gravity corer during RV Tangaroa
124 voyage TAN1404 in April/May 2014.

125

126 These standard soil classification test results indicate that both materials have similar
127 physical properties (Table 1). Natural water contents were higher in the shallow
128 EN1285 samples (Tan1404-02) and had correspondingly higher void ratios and lower
129 dry densities when compared with deeper EN1287 samples (Tan1404-10). Particle-size
130 analyses (Figure 2) confirm that both materials are fine grained comprising of over 88
131 % silt, approximately 5 % clay and with the remaining fraction consisting of mostly
132 fine and medium grained sand. The Atterberg Limit tests performed for both samples
133 confirmed similar plastic and liquid limits and universal soil classifications at the
134 boundary of high plasticity silts and clays (Table 1). The results indicated that the
135 physical properties of the sediments were similar both within the downslope extensional
136 domain and above the current landslide crest, indicating that any sediment disturbance
137 during sampling would have negligible impact on our study.

138

139 Conventional drained direct shear tests were undertaken on 60 mm diameter circular
140 samples of both materials using a Wykeham Farrance direct shearbox WF2500.
141 Shearing was conducted at three low normal stresses (Table 2) to simulate the shallow
142 depths of burial of the samples. Shearing was initiated on completion of the
143 consolidation phase (i.e., no further significant vertical displacement at the desired
144 normal load for shearing). A slow shear rate ($1.83 \times 10^{-4} \text{ mm s}^{-1}$), was used to avoid
145 developing excess pore pressures within the specimens and a minimum of five shear
146 reversals was completed for each test to ensure a representative 'residual' shear strength
147 was reached (i.e., no further reduction in shearing stress on cyclic loading).

148

149 A suite of specialist pore-pressure reinflation experiments was performed in the
150 Dynamic Back-Pressured Shearbox (Figure 3), constructed by GDS Instruments Ltd

151 and described in detail by Brain et al. (2015) and Carey et al. (2016). Previous studies
152 have demonstrated that the DBPSB is able to induce a variety of styles of deformation
153 ranging from dynamic liquefaction (Carey et al. 2017) to creep (Carey et al. 2016). Each
154 sample was saturated prior to testing to replicate the shallow seabed conditions (Table
155 3) using methods previously described (see Carey et al., 2016). A normal effective
156 stress of 32 kPa was applied to each sample during consolidation by applying a back
157 pressure of 100 kPa and a total normal stress at 132 kPa.

158

159 After consolidation, each sample was sheared to failure at a constant strain rate ($1.83 \times$
160 $10^{-4} \text{ mm s}^{-1}$, chosen to avoid excess pore fluid generation) to form a shear zone
161 representative of the base of a shallow subaqueous landslide, and to measure initial
162 shear strength. A shear-stress of c.70 % of the undrained strength was then applied to
163 each sample, whereupon both total normal and total shear stress were held constant
164 whilst the normal effective stress was reduced by increasing pore pressure until the
165 samples failed (Figure 4). The experiments were conducted using two different pore
166 fluid conditions: water, and water plus nitrogen gas.

167

168 Pore-water pressure controlled tests (referred to as PWP) were conducted by linearly
169 increasing the back pressure applied to the sample whilst holding both the total normal
170 stress and shear stress constant (Figure 3B). De-aired water was used to ensure that the
171 pore-water pressure increases and changes in fluid movement pathways anticipated in
172 the shallow seabed were accurately simulated.

173

174 Pore-gas pressure controlled experiments (referred to as PGP) were conducted using
175 nitrogen gas because it has similar physical properties to methane for the conditions of

176 our tests (Kossel et al. 2013), but is safer to use than methane. PGP experiments were
177 performed by filling the volume controller (VC1) with nitrogen gas whilst a second
178 volume controller (VC2) was used to maintain the back pressure to the sample (Figure
179 3C). VC1 was raised to the same pressure as the vessel before being connected to the
180 sample. Once the pressures were equal, the gas volume controller was opened to the
181 base of the sample and monitored to ensure that no significant pressure change occurred
182 in either volume controller, and that no sample strain occurred prior to the test run. The
183 pressure in the gas volume controller was then increased linearly at rates of 12 kPa/hr
184 and 30 kPa/hr to replicate increasing gas pressure within a shear zone. The porous disc,
185 installed at the sample base was replaced with a specially designed gas test plug prior
186 to testing to ensure gas pressures would be applied to the sample shear zone (Figure 3
187 C). To ensure that the gas was replacing the water within the soil pores, VC1 was
188 allowed to increase in volume (to extract water) as gas pressure was increased (Figure
189 3 C). These PGP tests using the DBPSB were the first of their kind, representing a new
190 methodology for testing the impact of gas pressure on sediment failure. Both the PWP
191 and PGP testing approaches resulted in similar stress paths to failure in which
192 increasing pore-fluid pressures resulted in a reduction in mean effective stress toward
193 failure while shear stress and total normal stresses remained constant (Figure 4).

194

195 During each phase of pore-fluid pressure increase, shear strain was monitored by
196 measuring the horizontal (shear) displacement of the shear box. These experiments
197 simulated the generation of excess pore-fluid pressure, and associated shear-strain
198 response, in near surface subaqueous landslides under a representative stress state.

199

200 4. LABORATORY RESULTS

201 The conventional shear box tests indicated no notable peak strength or shear strength
202 reduction following repeated shear reversals (Table 2) consistent with ductile
203 behaviour. Shear tests at three confining pressures indicated that the sediments had a
204 linear drained failure envelope, with similar friction angles which ranged between 34°
205 and 36° with an effective cohesion of 4.0 kPa (Figure 5). The strength characteristics
206 measured in the landslide materials were comparable, with similar materials tested from
207 other fine-grained subaqueous landslide systems (e.g. Sassa et al., 2012).

208

209 The horizontal (shear) strain and vertical (axial) strain behaviour were found to be
210 broadly consistent in all three PWP tests regardless of rate of pore water pressure
211 increase (Figure 6 A and B). In each experiment horizontal deformation progressed
212 through three movement phases as the shear zone dilated; these three phases can be
213 observed by distinct changes in horizontal strain rate (Figure 7).

214

215 An initial low strain rate ($< 0.05 \mu\text{S s}^{-1}$) was observed in each sample whilst the mean
216 effective stress (p') remained high (Figure 7). The strain rate of this early phase of
217 shear-zone deformation increased exponentially in response to the reducing mean
218 effective stress and corresponding sample dilation. This initial phase was followed by
219 a distinct sliding phase which was characterised by a rapid increase and then decrease
220 in strain rate (Figure 7). This movement was not associated with a measured change in
221 bulk sample pore-water pressure and was observed in all three samples, although more
222 pronounced in experiments PWP 5 and PWP 3. The sliding phase developed at a similar
223 normal effective stress (c.10 kPa) in each experiment, regardless of the rate of pore-
224 pressure increase. The continued decrease in mean effective stress following this phase
225 resulted in a progressive increase in horizontal strain which comprised of distinct

226 periods of increasing strain rate punctuated by periods of reducing strain rate. This
227 cyclic strain phase produced an exponential increase in strain rate with reducing mean
228 effective stress (Figure 7).

229

230 The development of a short, near-instantaneous, sliding phase followed by repeating
231 cyclic strain suggested that the shear zone mobilised and sheared at a critical mean
232 effective stress without acceleration to runaway failure despite the continued reduction
233 in mean effective stress. Instead, the strain rate continued to increase exponentially with
234 pore-water pressure and therefore exhibited an asymptotic trend in $1/v-p'$ space (Figure
235 8). Despite some variability in behaviour between each experiment, the strong
236 asymptotic trend observed in $1/v-p'$ space (Figures 8 A, B and C) demonstrated that
237 each sample underwent ductile deformation after a critical mean effective stress was
238 reached, regardless of the rate of pore-water pressure increase. Similar styles of
239 behaviour have been observed in shallow terrestrial landslides (e.g. Allison and
240 Brunsden, 1990) and have been shown to be controlled by localised pore fluid changes
241 in laboratory experiments (e.g. Ng and Petley, 2009).

242

243 The PGP experiments showed broadly similar progressive strain development and
244 shear-zone dilation in response to decreasing mean effective stress as observed in the
245 PWP testing (Figure 6 A and B). However, whilst movement initiated at low horizontal
246 strain rates ($< 0.05 \mu\text{S s}^{-1}$) and high mean effective stress in both experiments (Figure
247 9 A), only PGP 12 developed the distinct sliding phase observed during the PWP
248 experiments. This sliding phase occurred at a similar critical effective normal stress
249 (c.10 kPa) to that in each of the PWP experiments indicating that pore fluid phase had

250 negligible influence on the effective stress conditions required to mobilise the shear
251 zone.

252

253 By comparison, PGP 30 experienced an exponential increase in strain rate to
254 approximately $0.5 \mu\text{S s}^{-1}$ at a higher effective stress (15 kPa, Figure 9A). Following
255 this, the strain rate remained constant whilst the mean effective stress continued to
256 reduce. The peak strain rate was then reached at a lower effective stress (c.8 kPa) and
257 resulted from a further exponential increase in strain rate. Whilst more variability was
258 observed in PGP 30, similar maximum strain rates were observed in both PGP 30 and
259 PWP 30 experiments, suggesting subtle differences in rheology may have impacted the
260 testing. Similar episodic patterns of strain-rate development were observed in both gas
261 experiments once the shear surface had mobilised, and an asymptotic trend in $1/v - p'$
262 space was observed in both PGP experiments (Figure 9 B).

263

264 The results indicated that shear-zone deformation occurred through ductile deformation
265 at a critical mean effective stress regardless of the applied rate of pore pressure increase
266 or the phase of the pore fluid. In addition, the distinct reduction in movement rates after
267 initial shear surface mobilisation and the exponential increase in displacement rate
268 observed across all experiments indicated that catastrophic failure did not develop.
269 During most of the experiments, no relationship between reduction in strain rate and
270 decreasing sample pore pressure was observed (e.g. Figure 10 A), however, a reduction
271 in sample pore gas pressure was measured during experiment PGP30, which coincided
272 with the development of peak horizontal and vertical strain rates (Figure 10 B). It was
273 inferred from this that the rapid dilation of the shear zone acted to increase its
274 permeability, which resulted in localised dissipation of pore water pressure, temporarily

275 altering the stress state within the shear zone before pore pressure increased further.
276 Slow episodic shear regulated by dilation provides a potential mechanism to explain
277 the cyclic phases of increased landslide displacement rate observed in response to
278 elevated pore-fluid pressures without the development of rapid shear failure.

279

280 5. DISCUSSION

281 Although a range of mechanisms have been suggested to explain various subaqueous
282 mass-movements, including shear surface nucleation (e.g. Viesca and Rice, 2012);
283 shear zone liquefaction and ductile extrusion (e.g. Bull et al., 2009; Sassa et al., 2012);
284 and local lateral fluid flow (Dugan and Flemings, 2000; Fleming et al., 2002), very few
285 mechanisms have been proposed to explain shallow subaqueous slope failures that
286 arrest without long runout. Based on the experimental results illustrated herein, a
287 conceptual model can be hypothesised to explain how shallow subaqueous slopes can
288 progressively deform through episodic movement when pore-fluid pressures are
289 elevated within a shear zone (Figure 11).

290

291 In the model, an increase in pore pressure generates localised dilation and strain within
292 the landslide shear zone. As a consequence, very slow pre-failure deformation initiates
293 whilst effective normal stress remains comparatively high in the slope (stage 2). Slow
294 dilation increases the permeability of the shear zone allowing pore pressures to increase
295 more rapidly which increases the landslide strain rate and drives further shear zone
296 dilation (stage 3). This progressive increase in strain rate, and the associated inter-
297 particle deformation, drives further increases in local pore-fluid pressure, within the
298 narrow shear zone. This in turn induces a further increase in strain rate (stage 4). The
299 feedback mechanism continues as pore water pressure increases to reach a critical mean

300 effective stress when the shear surface rapidly slides and dilates (stage 5). This process
301 rapidly changes the shear zone properties as the rate of permeability increase through
302 dilation exceeds the rate of pore pressure increase within the narrow shear band. This
303 leads to pore expansion and dissipation of the excess fluid pressure (stage 5), locally
304 increasing the effective stress and reducing the shear strain rate. These processes
305 prevent catastrophic acceleration (stage 6).

306

307 The process of coupling local pore-pressure increase and the development of high strain
308 rate, checked by dilation and the broadening of the shear zone, can continue while
309 external processes drive increasing pore-fluid pressure (stage 7a). Consequently, the
310 landslide will continue to display ductile deformation behaviour, characterised by
311 exponential increases in strain rate punctuated by episodic phases of decreased
312 movement rates, never leading to catastrophic failure. Should the externally derived
313 pore-fluid pressures reduce or pore fluid dissipate fully from the shear zone (Stage 7b)
314 the slope movement will arrest and expulsion of pore fluid from the pores will result in
315 re-compaction of the shear zone (stage 8).

316

317 Our study has focused specifically on the movement mechanisms in subaqueous slopes
318 using a linear increase in pore-fluid pressure within a pre-defined shear zone. Other
319 potential mechanisms hypothesised to influence landslide motion such as state-and-
320 rate-variable friction affects (e.g. Helmsetter et al., 2004), complex perturbations in
321 effective stress (e.g. Hangwerger et al., 2016) and shear-surface geometry (e.g. Aryal
322 et al., 2015) have not been analysed. The coupling of shear deformation and shear-zone
323 dilation observed in our experiments has been shown to promote steady landslide
324 motion, particularly within clay-rich landslide shear zones (Iverson, 2005). Similar

325 movement patterns are also observed in terrestrial landslides on shallow slopes, (e.g.
326 Ng and Petley, 2009) and slowly deforming mudslides (e.g. Allison and Brunsten,
327 1990). In addition, similar behaviour has been used to describe ice flow dynamics (e.g.
328 Damsgaard et al., 2016) and the pore fluid driven cyclic fault-valve model proposed for
329 seismic slip (Sibson, 1992), indicating that this behaviour can be expected across a
330 broad range of geological processes.

331

332 Despite the different physical properties of nitrogen gas and liquid water, similar
333 patterns of behaviour were observed across all the experiments, demonstrating that
334 either fluid can generate similar movement characteristics with increasing pore
335 pressure. The results explain how unconstrained subaqueous landslides can episodically
336 move downslope when pore-fluid pressures at the landslide shear zone are elevated by
337 external factors such as the injection of water or gas from below.

338

339 The experiments presented herein provide credible support for the hypothesis that over-
340 pressuring by free gas can result in episodic/slow movements required to produce
341 submarine spreading failures observed in different parts of the world (e.g. Micallef et
342 al., 2007; Mountjoy et al. 2009; Micallef et al. 2016). Cyclic changes in strain rate
343 driven by a negative feedback mechanism associated with shear-zone dilation provides
344 a credible mechanism through which a subaqueous landslide can accumulate large
345 strain without catastrophic failure. This cyclic process, therefore, determines the long-
346 term behaviour of subaqueous mass movement in shallow sedimentary sequences, and
347 as similar materials commonly form submarine slopes, it is likely to be a widespread
348 seafloor process.

349

350 CONCLUSION

351 Geomorphological evidence suggests that some shallow subaqueous slopes can
352 accumulate substantial amounts of downslope deformation without transitioning to
353 catastrophic failure, even though their downslope terminations are unconstrained
354 (Mountjoy et al., 2009; Micallef et al., 2016). Novel laboratory experiments have been
355 used to explore this behaviour by examining the response of such slope materials to
356 states of low effective normal stress associated with high pore-fluid pressure induced
357 by either gas or water injected from below. The strain behaviour observed in the
358 experiments were found to be similar to the movement patterns measured in terrestrial
359 landslides which deform along ductile shear zones. In such circumstances movement
360 develops when pore water pressures are sufficiently elevated in the slope and movement
361 rates increase exponentially with increasing pore pressure. Given that this response is
362 dominated by the effective stress conditions operating within the shear zone, such
363 behaviour can be expected in submarine slopes regardless of the pore fluid phase (gas
364 or liquid). The behaviour we observed provides a mechanism through which
365 subaqueous landslides may accumulate strain without undergoing catastrophic failure.
366 This has important implications when assessing marine geohazards since these types of
367 slope failure will not be tsunamigenic.

368

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601 FIGURE CAPTIONS

602

603 Figure 1. Study area, east of New Zealand's North Island. (A) 'Hillshade' plot of
604 Tuaheni Basin bathymetry, showing the Tuaheni Landslide Complex (TLC) and a
605 shallow subaqueous slope failure (Inset B) which both show evidence of limited
606 displacement despite having unconstrained toes. Yellow dots = location of Sample
607 Tan1404-10; yellow contours = metres below sea level adapted from Mountjoy et al.
608 (2014). (B) Bathymetric map of shallow subaqueous slope failure draped over slope
609 gradient map and showing key morphological features including arrested debris (ad)
610 and evacuated debris (ed) after Micallef et al. (C) Seismic section across the subaqueous
611 slope failure headscarp after Micallef et al. (2016).

612

613 Figure 2. The particle size distribution of samples EN1285 (Tan1404-01) and EN1287
614 (Tan1404-10) collected from the Tuaheni South section of the TLC.

615

616 Figure 3. Laboratory testing apparatus and procedures. (A) Schematic diagram of the
617 Dynamic Back Pressure Shearbox apparatus. (B) Experimental procedure used for pore
618 water pressure testing. (C) Experimental procedure used for pore gas pressure testing.

619

620 Figure 4. Testing parameters used during the specialist pore pressure reinflation testing.
621 (A) Applied pore pressure (Back pressure) against mean effective stress. (B) Stress
622 paths followed during the pore water pressure (PWP) and pore gas pressure (PGP)
623 experiments in relation to the conventional failure envelope.

624

625 Figure 5. Residual strength envelopes constructed from conventional drained shear tests
626 carried out on samples EN1285 and EN1287.

627

628 Figure 6. Change in horizontal and vertical strain measured during PWP and PGP
629 experiments. (A) Horizontal strain against normal effective stress. (B) Vertical strain
630 against normal effective stress.

631

632 Figure 7. Horizontal strain rate against normal effective stress illustrating three distinct
633 movement phases measured during PWP experiments.

634

635 Figure 8. Analysis of movement styles observed during each PWP experiment using 1/
636 Horizontal strain rate (V) against normal effective stress (p'). (A) Asymptotic trend
637 calculated during PWP 30 experiment. (B) Asymptotic trend calculated during PWP 12
638 experiment. (C) Asymptotic trend calculated during PWP 5 experiment.

639

640 Figure 9. Analysis of movement styles observed during each PGP experiment. (A)
641 Horizontal strain rate against normal effective stress. (B) 1/ Horizontal strain rate (v)
642 against normal effective stress (p'). (C) Asymptotic trend calculated during PGP 30
643 experiment. (D) Asymptotic trend calculated during PGP 12 experiment.

644

645 Figure 10. Comparison of behaviour observed during the PWP and PGP experiments.
646 (A) Change in pore water pressure and strain rate against time measured during
647 experiment PWP 30. (B) Change in pore gas pressure and horizontal strain rate against
648 time measured during experiment PGP 30.

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651 Figure 11. Conceptual model of the development of slow movement in the shallow
652 subaqueous slopes driven by elevated pore-fluid pressure.

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656 TABLE CAPTIONS

657 Table 1. Physical properties of the Tuaheni sediments

658 Table 2. Summary of conventional drained shear tests

659 Table 3. Summary of pore pressure reinflation tests

660

TABLE 1. PHYSICAL PROPERTIES OF THE TUAHENI SEDIMENTS

Sample Reference	EN1285 (Tan1404-02)		EN1287 (Tan1404-10)	
	Range	Average	Range	Average
Sampling depth (*mbsb)	0.32-0.49		2.23-2.40	
Moisture content (%)	77.0-80.0	78.5	70.0-72.0	70.5
Void ratio	3.11-3.17	3.14	2.00-2.18	2.06
Dry density	0.65-0.66	0.65	0.85-0.90	0.88
Atterberg limits:				
Plastic limit (%)		75		76
Liquid limit (%)		35		32
Plasticity index (%)		40		44
Particle size distribution:				
Clay (%)		5.50		4.18
Silt (%)		88.18		88.48
Sand (%)		6.32		7.33

*mbsb = meters below sea bed

TABLE 2. SUMMARY OF CONVENTIONAL DRAINED SHEAR TESTS

Sample number	Sample depth (mbsl)	Initial water content w_n (%)	Void Ratio	Dry density ρ_d (t/m ³)	Normal stress (kPa)	Shear stress (kPa)	
						Peak	Residual
EN1285a	0.40-0.44	78	3.14	0.65	32	27	27
EN1285c	0.36-0.40	77	3.11	0.66	76	59	59
EN1285d	0.32-0.36	79	3.17	0.65	10	11	11
EN1287a	2.31-2.35	70	2.02	0.89	32	31	31
EN1287c	2.27-2.31	72	2.18	0.85	10	12	12
EN1287e	2.23-2.27	70	2.00	0.90	76	66	66

TABLE 3. SUMMARY OF PORE PRESURE REINLFATION TESTS

Test Reference	Sample Number	Normal effective stress (kPa)	Initial shear strength (kPa)	Post failure shear stress/ Percentage of initial shear strength (kPa) / (%)	Pore fluid increase rate (kPa/ hr)	Pore fluid
PWP 5	EN1287i*	32*	22	15 (68%)	5	water
PWP 12			20	15 (75%)	12	water
PWP 30	EN1287p	32	20	13 (65%)	30	water
PGP 12	EN1287q	32	20	13 (65%)	12	nitrogen
PGP30	EN1287k	32	19	13 (68%)	30	nitrogen

*Test PWP 5 and PWP 12 conducted on the same sample

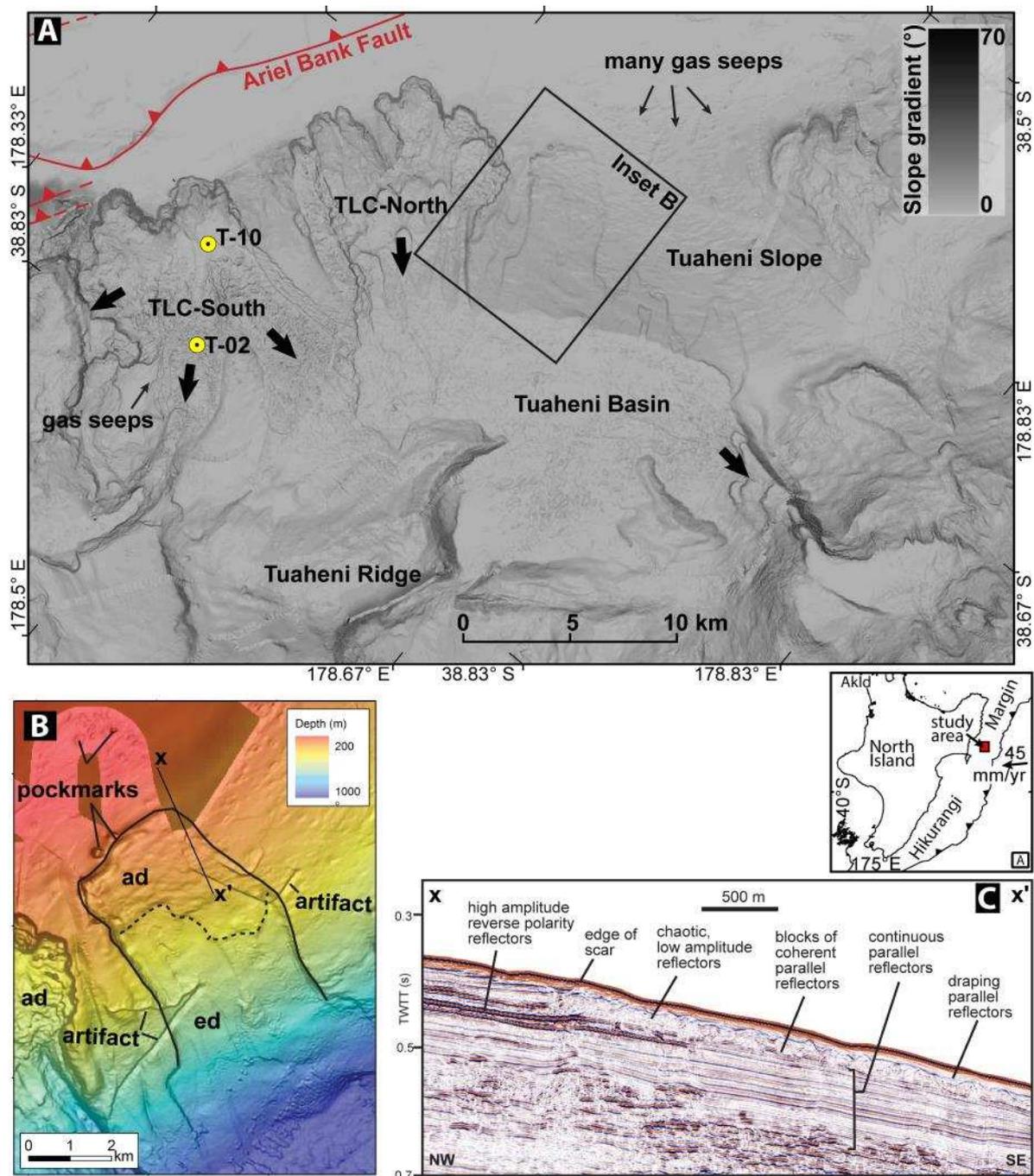


Figure 1. Study area, east of New Zealand's North Island. (A) 'Hillshade' plot of Tuaheni Basin bathymetry, showing the Tuaheni Landslide Complex (TLC) and a shallow subaqueous slope failure (Inset B) which both show evidence of limited displacement despite having unconstrained toes. Yellow dots = location of Sample Tan1404-10; yellow contours = metres below sea level adapted from Mountjoy et al. (2014). (B) Bathymetric map of shallow subaqueous slope failure draped over slope gradient map and showing key morphological features including arrested debris (ad) and evacuated debris (ed) after Micallef et al. (C) Seismic section across the subaqueous slope failure headscarp after Micallef et al. (2016).

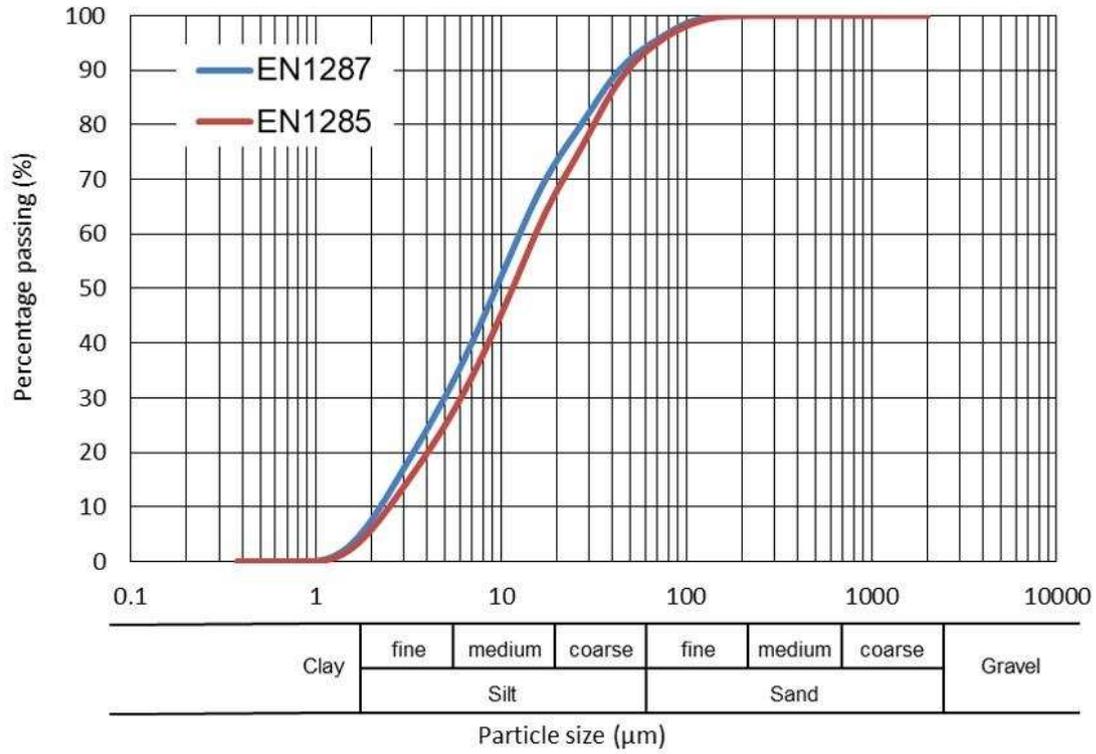


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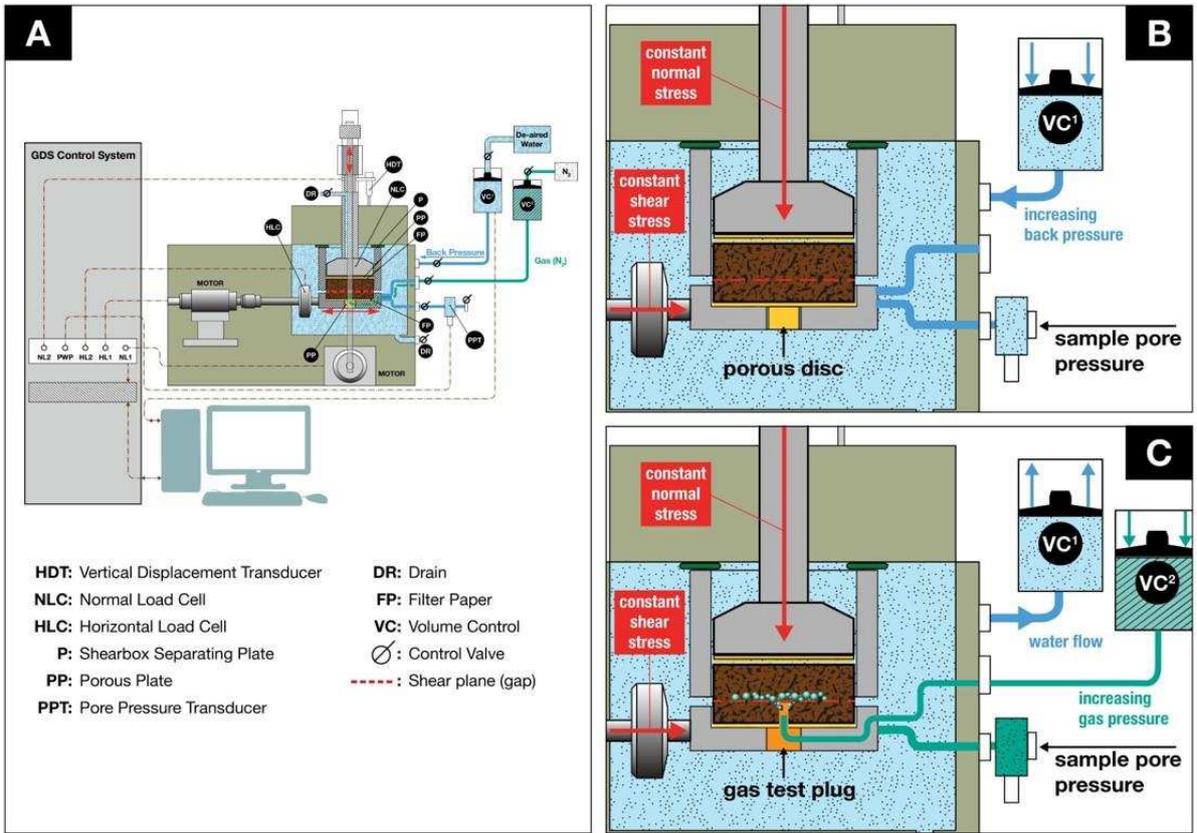


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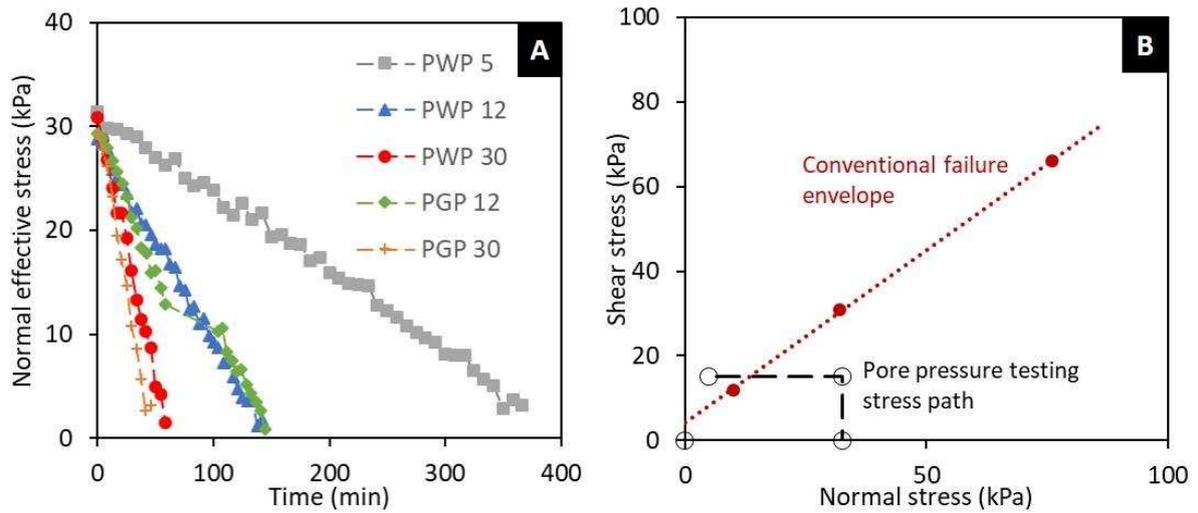


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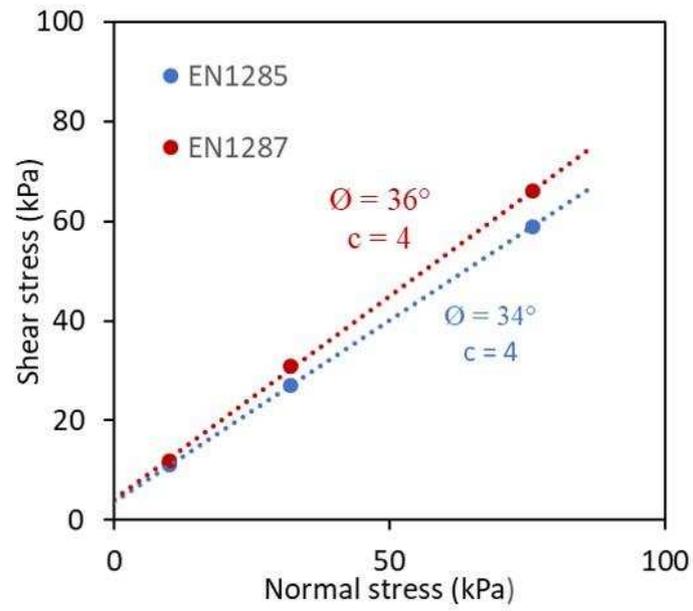


Figure 5. Residual strength envelopes constructed from conventional drained shear tests carried out on samples EN1285 and EN1287.

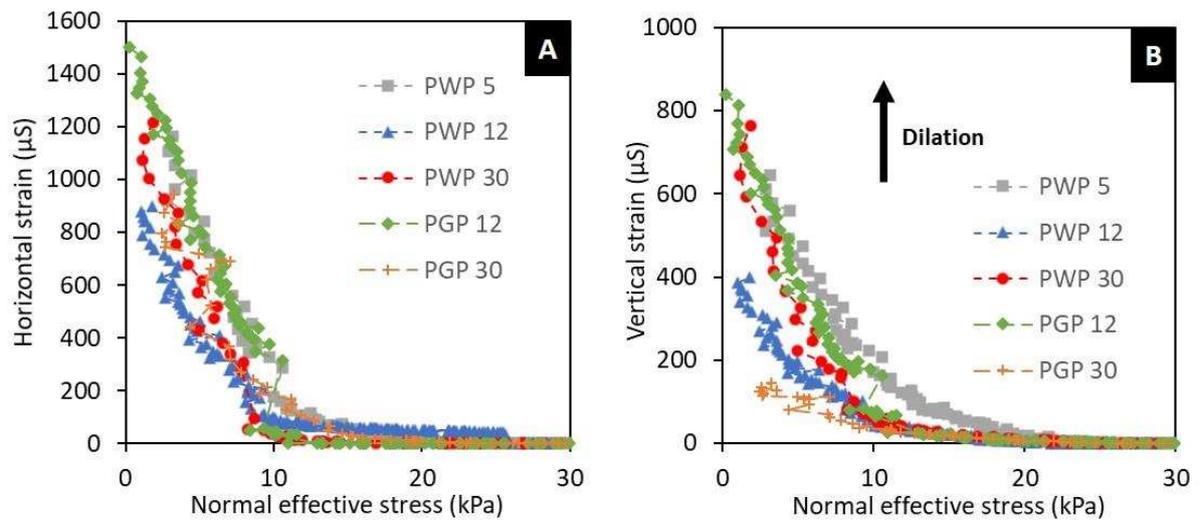


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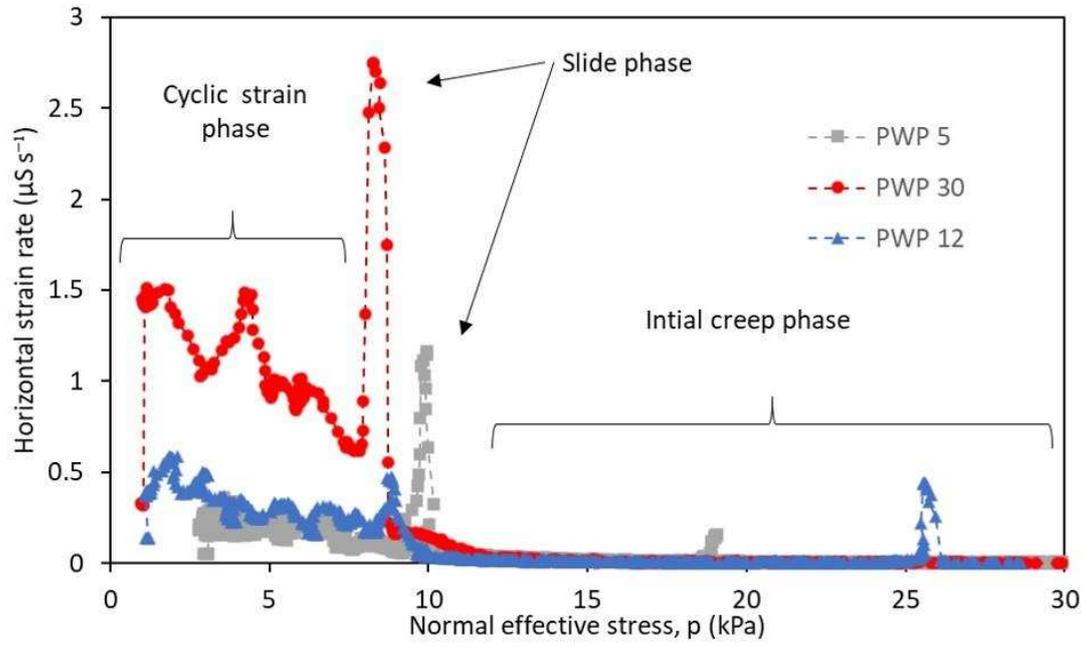


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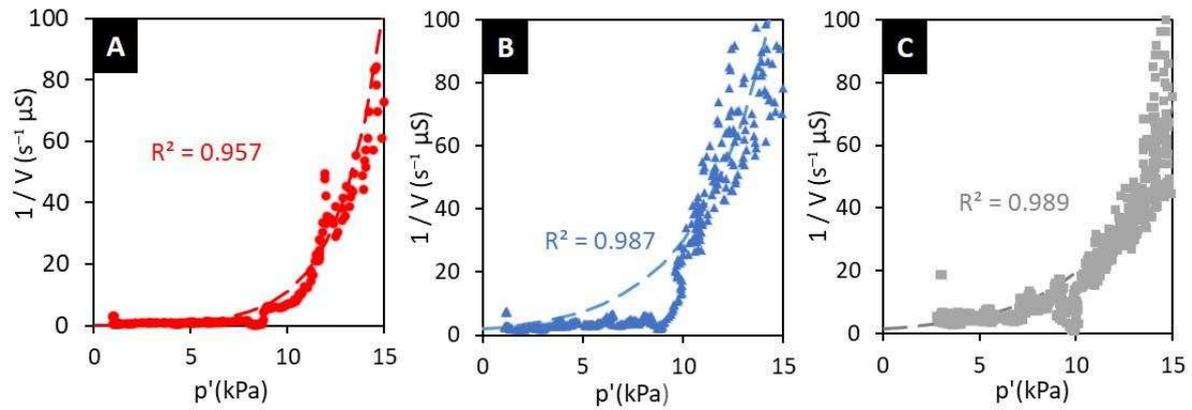


Figure 8. Analysis of movement styles observed during each PWP experiment using $1/V$ Horizontal strain rate (V) against normal effective stress (p'). (A) Asymptotic trend calculated during PWP 30 experiment. (B) Asymptotic trend calculated during PWP 12 experiment. (C) Asymptotic trend calculated during PWP 5 experiment.

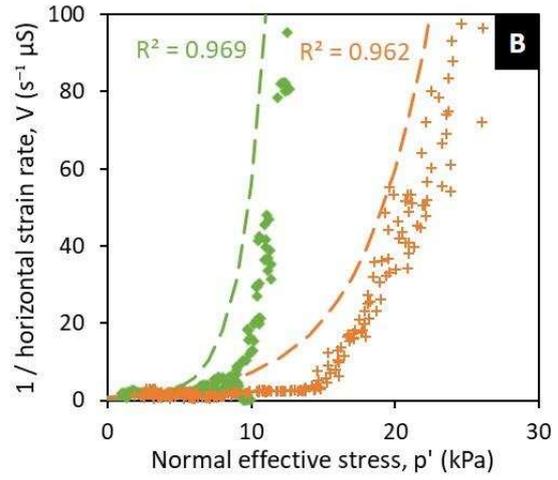
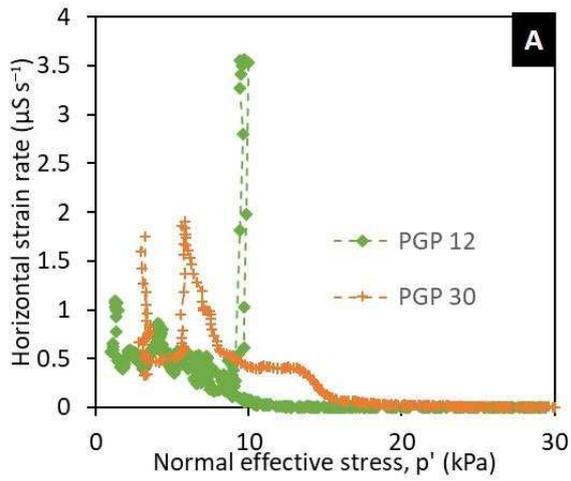


Figure 9. Analysis of movement styles observed during each PGP experiment. (A) Horizontal strain rate against normal effective stress. (B) 1/ Horizontal strain rate (v) against normal effective stress (p'). (C) Asymptotic trend calculated during PGP 30 experiment. (D) Asymptotic trend calculated during PGP 12 experiment.

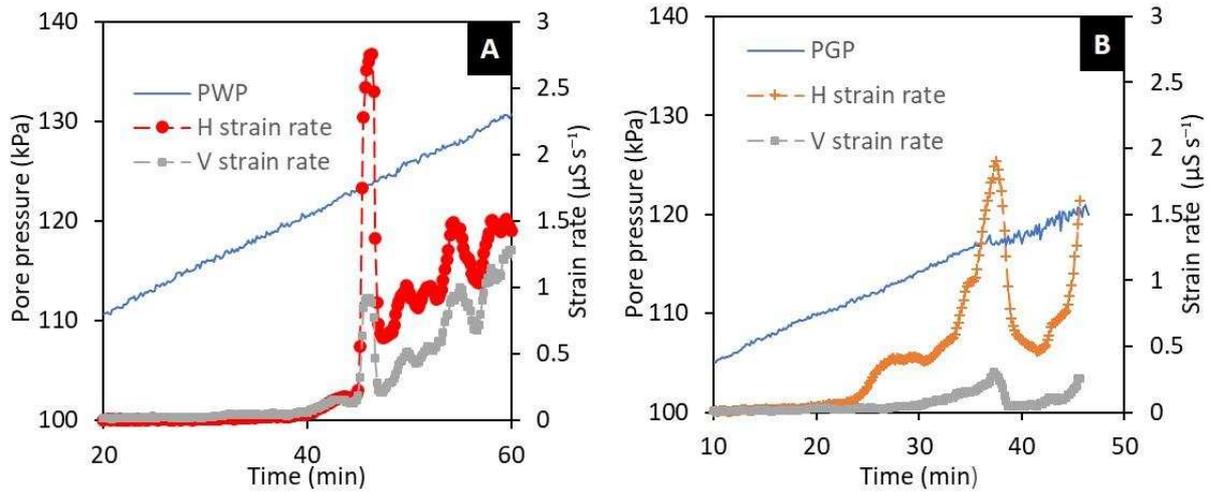


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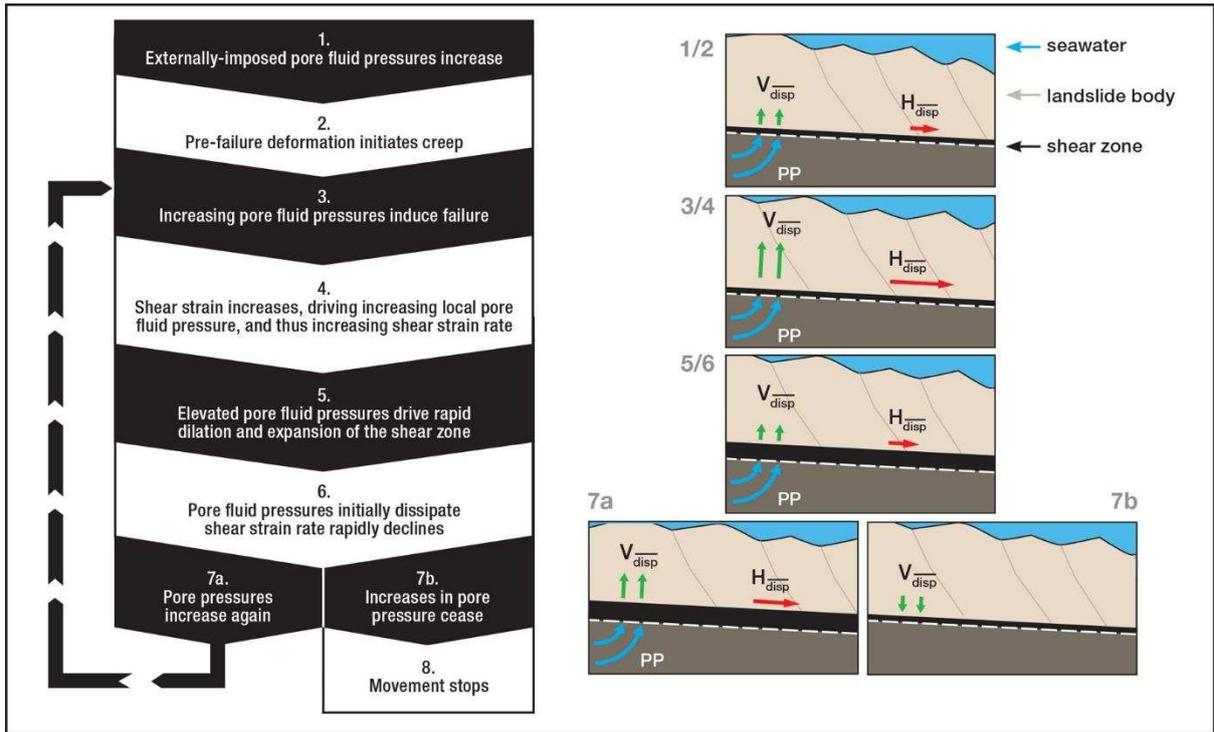


Figure 11. Conceptual model of the development of slow movement in the shallow subaqueous slopes driven by elevated pore-fluid pressure.