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# A Kinetic Study of Ni and NiO Reactions Pertinent to the Earth's Upper Atmosphere

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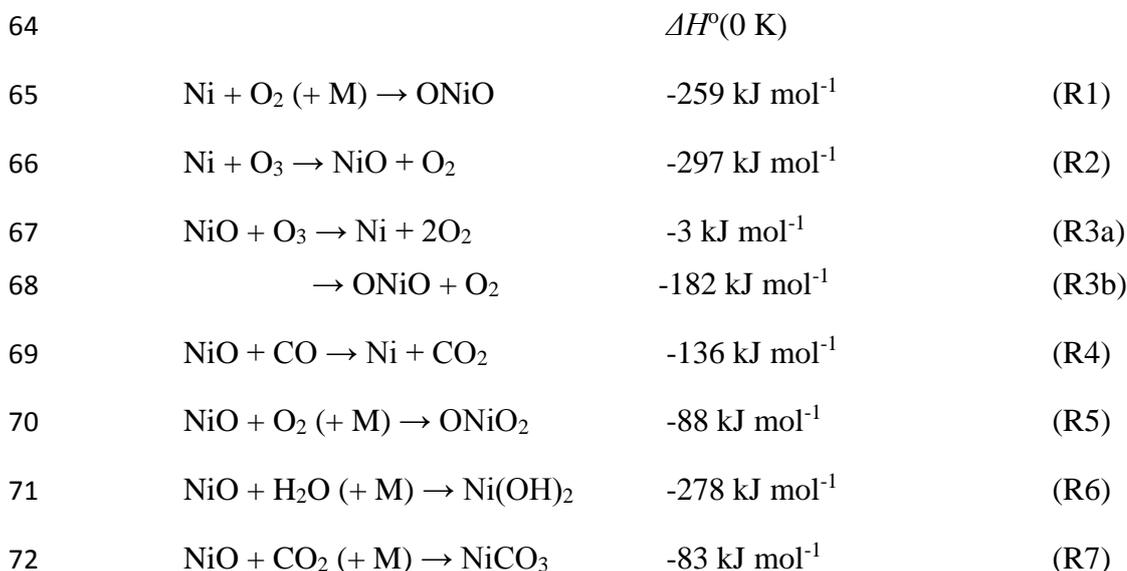
**Abstract.** Nickel atoms are injected into the Earth's mesosphere by meteoric ablation, producing a Ni layer between 70 and 105 km in altitude. The subsequent reactions of Ni and NiO with atmospherically relevant species were studied using the time-resolved pulsed laser photolysis-laser induced fluorescence technique, combined with electronic structure calculations and RRKM theory where appropriate. Results for bimolecular reactions (in  $\text{cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ):  $k(\text{Ni} + \text{O}_3, 293\text{K}) = (6.5 \pm 0.7) \times 10^{-10}$ ;  $k(\text{NiO} + \text{O}_3 \rightarrow \text{Ni} + 2\text{O}_2, 293 \text{ K}) = (1.4 \pm 0.5) \times 10^{-10}$ ;  $k(\text{NiO} + \text{O}_3 \rightarrow \text{NiO}_2 + \text{O}_2, 293 \text{ K}) = (2.5 \pm 0.7) \times 10^{-10}$ ;  $k(\text{NiO} + \text{CO}, 190 - 377 \text{ K}) = (3.2 \pm 0.6) \times 10^{-11} (\text{T}/200)^{-0.19 \pm 0.05}$ . For termolecular reactions (in  $\text{cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}$ , uncertainty  $\pm \sigma$  over the stated temperature range):  $\log_{10}(k_{\text{rec},0}(\text{Ni} + \text{O}_2 + \text{N}_2, 190 - 455 \text{ K})) = -37.592 + 7.168 \log_{10}(\text{T}) - 1.5650(\log_{10}(\text{T}))^2$ ,  $\sigma = 11\%$ ;  $\log_{10}(k_{\text{rec},0}(\text{NiO} + \text{O}_2 + \text{N}_2, 293 - 380 \text{ K})) = -41.0913 + 10.1064 \log_{10}(\text{T}) - 2.2610(\log_{10}(\text{T}))^2$ ,  $\sigma = 22\%$ ; and  $\log_{10}(k_{\text{rec},0}(\text{NiO} + \text{CO}_2 + \text{N}_2, 191 - 375 \text{ K})) = -41.4265 + 10.9640 \log_{10}(\text{T}) - 2.5287(\log_{10}(\text{T}))^2$ ,  $\sigma = 15\%$ . The faster recombination reaction  $\text{NiO} + \text{H}_2\text{O} + \text{N}_2$ , which is clearly in the fall-off region over the experimental pressure range (3 – 10 Torr), is best described by:  $\log_{10}(k_{\text{rec},0}/\text{cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}) = -29.7651 + 5.2064 \log_{10}(\text{T}) - 1.7118(\log_{10}(\text{T}))^2$ ,  $k_{\text{rec},\infty} = 6.0 \times 10^{-10} \exp(-171/\text{T}) \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ , broadening factor  $F_c = 0.84$ ,  $\sigma = 16\%$ . The implications of these results in the atmosphere are then discussed.

## 31 1. Introduction

32 The Earth's atmosphere is subjected to a continuous influx of cosmic dust particles. The global  
 33 input rate of this dust was recently estimated to be  $43 \pm 14 \text{ t d}^{-1}$  (tonnes per day), of which  $8 \text{ t}$   
 34  $\text{d}^{-1}$  ablates due to frictional heating of the particles as they collide with air molecules at  
 35 hypothermal speeds.<sup>1</sup> Meteoric ablation is the major source of the layers of metal atoms that  
 36 occur between 75 and 105 km, in a region of the atmosphere known as the mesosphere/lower  
 37 thermosphere (MLT).<sup>2</sup> The layers of Na, K, Ca,  $\text{Ca}^+$  and Fe have been studied in great detail  
 38 over several decades by the ground-based lidar technique; and the Na, K, Mg and  $\text{Mg}^+$  layers  
 39 have also been observed by space-borne optical spectroscopy.<sup>3</sup> Recently, a Ni layer was  
 40 observed for the first time by a lidar at the Poker Flat Research Range in Alaska ( $65^\circ \text{N}$ ), using  
 41 the optical transition at 337.1 nm [ $\text{Ni}(\text{z}^3\text{F}_4^0 - \text{a}^3\text{F}_4)$ ].<sup>4</sup> Surprisingly, the Ni layer exhibited  
 42 marked differences from other metals such as Fe, extending down to nearly 70 km and having  
 43 a column density roughly 20 times larger than expected based on the elemental abundance of  
 44 Ni in Carbonaceous Ivuna (CI) chondrites (1.1%).<sup>5</sup> Ni can be present in meteorites as a Fe-Ni  
 45 alloy, but is present in the majority of chondrites as Fe-Ni sulphides.<sup>6</sup>

46 Understanding the chemistry that controls the metal layers in the MLT is important as they  
 47 provide unique tracers of the chemistry and dynamics of a part of the atmosphere that can only  
 48 be accessed directly by payloads on sub-orbital rockets. The metal atoms become oxidized in  
 49 the mesosphere to form a range of oxides, hydroxides and carbonates which then aggregate  
 50 into Meteoric Smoke Particles (MSPs).<sup>3, 7</sup> MSPs are thought to act as ice nuclei for polar  
 51 mesospheric clouds<sup>8</sup> and condensation nuclei for sulphate aerosol that go on to form polar  
 52 stratospheric clouds.<sup>9</sup> The composition and concentration of MSPs may influence these  
 53 processes. To understand the characteristics of the newly-observed Ni layer and the likely Ni-  
 54 containing molecules which will contribute to MSP formation, the gas-phase chemistry of Ni  
 55 in the MLT needs to be understood.

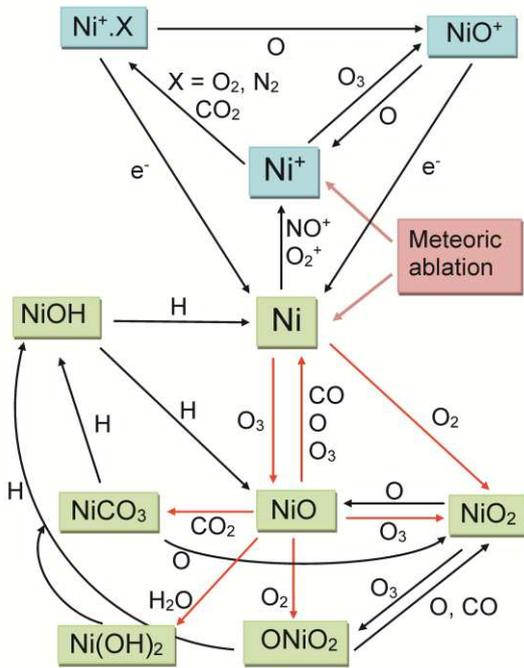
56 Figure 1 is a proposed reaction scheme for the likely neutral and ion reactions of Ni occurring  
 57 in the MLT. This is based on our previous work on Fe and Mg,<sup>3</sup> and using electronic structure  
 58 theory calculations of Ni species reacting with ambient MLT species to determine energetically  
 59 viable reaction pathways (see Section 4 for further details). Note that the ion-molecule  
 60 chemistry is included here for completeness, and is not the subject of the present study. The  
 61 only relevant reaction that has, to our knowledge, been studied previously is the recombination  
 62 of Ni with  $\text{O}_2$  (at 295 K in an Ar bath gas).<sup>10</sup> The following reactions of Ni and NiO that were  
 63 investigated in the present study (red arrows in Figure 1) are:



73

74 where M is a third body, and the standard reaction enthalpies (at 0 K) are calculated at the CBS-  
75 QB3 level of theory,<sup>11, 12</sup> apart from that for R1 which is taken from a recent multi-reference  
76 configuration interaction study.<sup>13</sup>

77



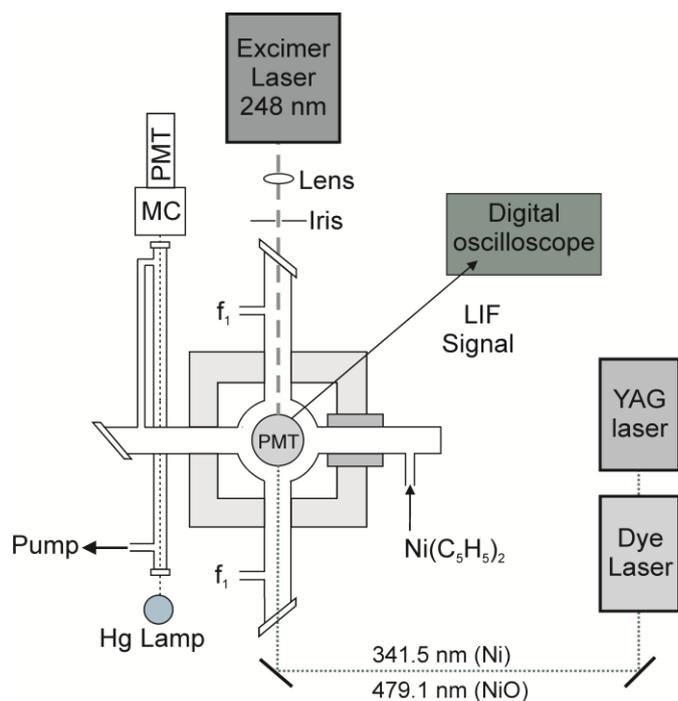
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79 Figure 1. Proposed reaction scheme for the neutral and ion-molecule chemistry of Ni in the  
80 MLT, where meteoric ablation provides the initial source of Ni and Ni<sup>+</sup>. Neutral and ionic  
81 species are shown in green and blue boxes, respectively. The kinetics of reactions indicated  
82 with red lines were measured in the present study.

83

## 84 2. Experimental

85 Reactions R1 – R7 were studied using the Pulsed Laser Photolysis-Laser Induced Fluorescence  
86 (PLP-LIF) technique. Figure 2 is a schematic diagram of the stainless steel reactor which has  
87 been used previously to study the kinetics of other mesospheric metal atoms and their oxides,  
88 including Si/SiO, Mg/MgO, Ca/CaO and Fe/FeO.<sup>3</sup>



89

90 Figure 2. Schematic diagram of the PLP-LIF system for studying the kinetics of reactions of  
 91 Ni and NiO. MC = monochromator;  $f_1$  = reactant gas flow; PMT = photomultiplier tube.

92

93 The stainless steel reactor comprises a central chamber (volume = 400 cm<sup>3</sup>) with four  
 94 horizontal side arms that are mutually orthogonal, and one vertical side arm. The reaction  
 95 chamber is surrounded by ceramic heaters housed within a stainless steel box, allowing for  
 96 stable resistive heating up to 1100 K. Alternatively, the box can be filled with dry ice to reach  
 97 a base temperature of 190 K. Reaction chamber temperatures were measured with a shielded  
 98 thermocouple inserted directly into the centre of the reactor volume.

99 Ni atoms were generated in the reactor by the multiphoton dissociation of nickelocene  
 100 (dicyclopentadienyl nickel, [Ni(C<sub>5</sub>H<sub>5</sub>)<sub>2</sub>]) vapor at 248 nm using a KrF excimer laser (Lambda  
 101 Physik COMPEX 102, typical pulse energy = 40 mJ, <5 mJ in the reactor (pulse rate = 10 Hz).  
 102 The excimer beam was loosely focused to a point 15 cm beyond the centre of the reaction  
 103 chamber using a lens (focal length = 50 cm). These experiments were limited to a maximum  
 104 temperature of 450 K because there was clear evidence in experiments at higher temperatures  
 105 that decomposition of the nickelocene precursor artificially increased reaction rates.

106 Ni atoms were probed either at 341.476 nm [Ni( $z^3F_4^0 - a^3D_3$ )] or at 339.105 nm [Ni( $z^3F_4^0 -$   
 107  $a^3F_4$ )], using a frequency-doubled Nd:YAG pumped dye laser (Sirah CBR-G-30) with Pyridine  
 108 1 laser dye. The dye and excimer laser beams counter-propagated through opposing arms of  
 109 the reactor (Figure 1). By varying the time delay between the dye laser (probe beam) and  
 110 excimer laser (photolysis beam,  $t = 0$ ), scans of the concentration of Ni/NiO with time were  
 111 obtained. In a typical experiment each LIF signal time step was the average of 5 laser shots.  
 112 The Ni LIF signal was collected through an interference filter ( $\lambda_{\text{max}} = 340$  nm, FWHM = 10  
 113 nm) using a Photomultiplier Tube (PMT, Electron Tubes, model 9816QB) attached to the  
 114 vertical side arm of the reactor, orthogonal to the laser beams. Loose focusing of the excimer  
 115 laser to achieve multi-photon dissociation was necessary to produce Ni atoms. This also  
 116 produced a population of Ni in excited states, the effect of which is discussed in Section 3.  
 117 NiO( $X^3\Sigma^-$ ) was probed at 479.108 nm<sup>5</sup> using Coumarin 102 dye, and monitored off resonance

118 using a cut-on filter at 495 nm. All the Ni and NiO transitions were calibrated using a  
119 wavelength meter (Bristol instruments 871) accurate to 0.8 pm.

120 Reactant gases (O<sub>2</sub>, O<sub>3</sub>, CO<sub>2</sub>, CO, H<sub>2</sub>O) were added via the PMT side arm and two laser beam  
121 side arms, mixing in the central chamber of the reactor. The flows in each experiment  
122 comprised the nickelocene vapor entrained in N<sub>2</sub> flow, the reactant flow and a balancing flow  
123 of the bath gas N<sub>2</sub> (total flow of 0.2 L/min). The reactant concentration was calculated using  
124 the calibrated mass flow rates and the pressure in the reactor, measured using a Baratron  
125 pressure gauge (MKS PR 4000). The total pressure in the system ranged between 3 and 10  
126 Torr. The final arm of the reactor provides the exhaust outlet leading to an O<sub>3</sub> absorption cell  
127 and vacuum pump (Edwards E28).

128 O<sub>3</sub> was produced by corona discharge of O<sub>2</sub> in an ozoniser (Fischer OZ500) with a typical  
129 conversion efficiency of ~2%. The concentration of O<sub>3</sub> was monitored by absorption  
130 spectroscopy at 253.7 nm from a Hg Pen-ray lamp, using a 1m path length absorption cell  
131 downstream of the reactor and a monochromator (Minichrom, 300 μm slits) with a PMT  
132 (Hamamatsu Type H9306-13) (Figure 1). The O<sub>3</sub> concentration was then determined using the  
133 absorption cross section at 293 K of  $1.16 \times 10^{-17}$  cm<sup>2</sup> at 253.7 nm.<sup>14</sup>

134

### 135 **Materials**

136 N<sub>2</sub> (99.9999 % pure, Air Products), O<sub>2</sub> (99.999 %, pure Air Products), CO<sub>2</sub> (99.995 % pure,  
137 Air Products), and CO (99.5 % pure, Argo International) were used without any further  
138 purification. H<sub>2</sub>O vapour was produced by freeze-pump thawing deionised water and  
139 subsequent dilution in a glass bulb on a glass vacuum gas-handling line. Nickelocene vapour  
140 (vapor pressure = 5 mTorr at 293 K<sup>15</sup>) was entrained within a flow of N<sub>2</sub> gas over a glass trap  
141 containing pure nickelocene crystals.

142

## 143 **3. Results**

144

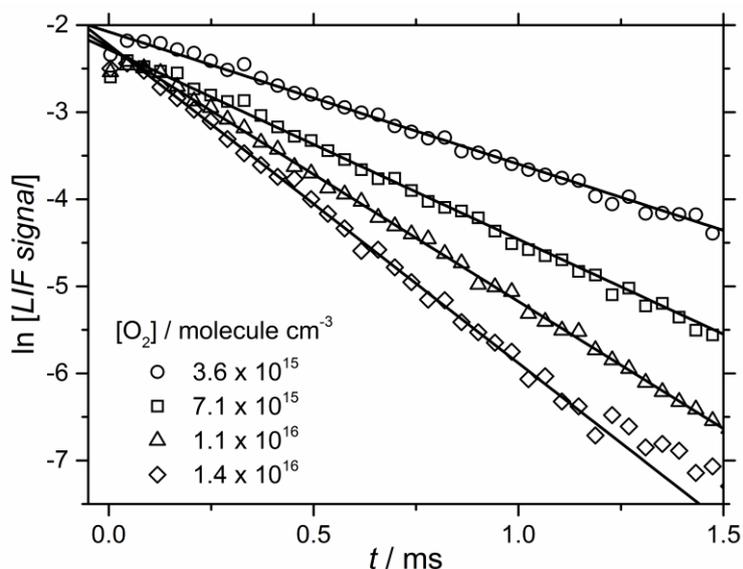
### 145 **3.1 Ni + O<sub>2</sub>**

146 The time-resolved LIF signals are described by a single exponential form  $A \times \exp(-k' t)$ , where  
147 the first-order decay coefficient  $k'$  is given by

$$148 \quad k' = k_{\text{diff}} + k_1[\text{O}_2][\text{M}] \quad (1)$$

149  $k_{\text{diff}}$  is the rate of diffusion of Ni atoms out of the volume of the excimer laser that is within the  
150 line of sight of the PMT, and has a value  $< 800$  s<sup>-1</sup>. Figure 3 illustrates plots of the natural  
151 logarithm of the LIF signal against time. Note the increase in Ni(<sup>3</sup>F) concentration within the  
152 first 100 μs of each kinetic trace, which is due to the quenching of electronically excited Ni  
153 atoms produced during photolysis of the nickelocene precursor. These excited atoms were  
154 quenched to the ground-state by the N<sub>2</sub> bath gas in ~ 100 μs,<sup>16</sup> and kinetic traces were therefore  
155 always fitted at longer times when there was no further evidence of highly excited metastable  
156 Ni atoms (i.e. above the low-lying Ni(<sup>3</sup>D) state) playing a role.

157



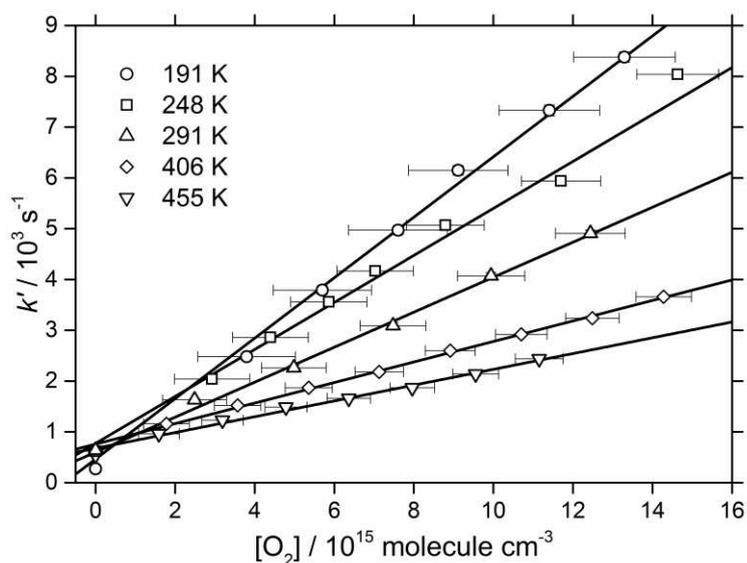
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159 Figure 3. Kinetic decays of Ni(<sup>3</sup>F) reacting with [O<sub>2</sub>] (total pressure = 3 Torr). The lines are  
 160 linear regression fits through each decay.

161

162 Figure 4 shows that plots of *k'* against [O<sub>2</sub>] over the temperature range 191 – 455 K are linear  
 163 with a slope providing the second-order rate coefficient at a particular pressure of N<sub>2</sub>. The  
 164 resulting third-order rate coefficients are plotted as function of temperature in Figure 5.

165



166

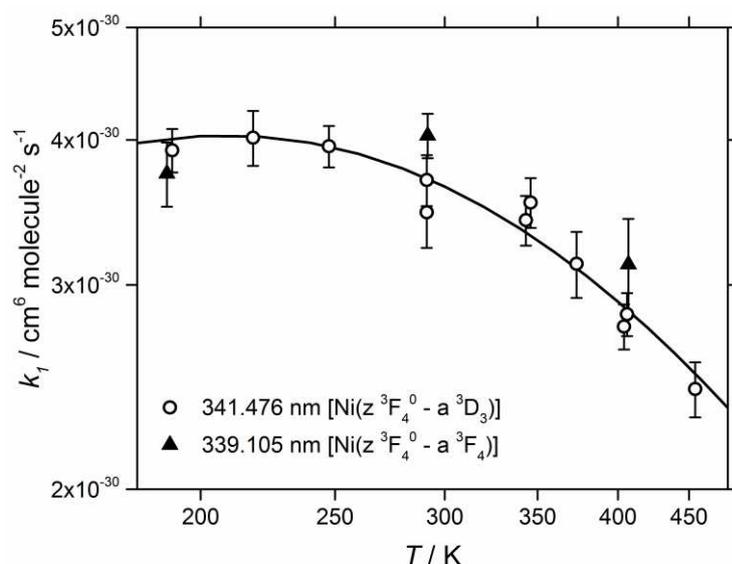
167 Figure 4. Plots of *k'* against [O<sub>2</sub>] for R1 (3 Torr total pressure). The slopes of the linear  
 168 regression fits yield the second-order rate coefficients, illustrating the negative temperature  
 169 dependence of the reaction from 191 to 455 K.

170

171 Nickel has two low lying states, the ground state Ni(<sup>a</sup>3F<sub>4</sub>, configuration [Ar].3d<sup>8</sup>.4s<sup>2</sup>) and the  
 172 excited Ni(<sup>a</sup>3D<sub>3</sub>, configuration [Ar].3d<sup>9</sup>.4s<sup>1</sup>) state which is only 204 cm<sup>-1</sup> above the ground state  
 173 so that the spin-orbit manifolds of these configurations overlap.<sup>17</sup> Note that the prefix in front  
 174 of the atomic term symbol is conventionally used to distinguish terms of the same type and  
 175 multiplicity from each other in complex atomic spectra.<sup>18</sup> In the present study most reactions

176 were investigated by probing the Ni( $a^3D_3$ ) excited state, as it produced significantly more LIF  
 177 signal since the Einstein coefficient for the Ni( $z^3F_4^0 - a^3D_3$ ) transition is  $5.5 \times 10^7 \text{ s}^{-1}$ , compared  
 178 with  $1.8 \times 10^7 \text{ s}^{-1}$  for the Ni( $z^3F_4^0 - a^3F_4$ ) transition.<sup>17</sup> As shown in Figure 5, R1 was studied at  
 179 several temperatures between 190 and 410 K by using both Ni transitions, and the resulting  
 180 rate coefficients agreed within error. This is likely to be because these states are rapidly  
 181 thermally equilibrated on the time-scale of reaction. This is to be expected given their close  
 182 separation in energy (the Boltzmann fraction of Ni in the  $a^3D_3$  state varies from 14.3 to 26.5%  
 183 between 191 and 400 K) and the relatively high pressure in the reactor ( $> 1 \text{ Torr N}_2$ ). Thus, the  
 184 reaction kinetics of both states are tightly coupled, and either transition can be used as a  
 185 spectroscopic marker for the overall kinetics which will be governed by the more reactive of  
 186 the two states.

187



188

189 Figure 5. Third-order rate coefficient for R1 as a function of temperature where  $\text{N}_2$  is the third  
 190 body. Measurements were made using LIF transitions from the Ni ( $a^3F_4$ ) ground state and the  
 191 low-lying Ni( $a^3D_3$ ) excited state. The solid line is an RRKM fit to the experimental data.

192

### 193 3.2 Ni + O<sub>3</sub> → NiO + O<sub>2</sub>

194 The loss of Ni in the presence of O<sub>3</sub> and O<sub>2</sub> should be described by:

$$195 \quad k' = k_{\text{diff}} + k_2[\text{O}_3] + k_1[\text{O}_2] \quad (2)$$

196 The points with square symbols in Figure 6 are a plot of  $k'$  as a function of  $[\text{O}_3]$ . This shows  
 197 curvature at increasing  $[\text{O}_3]$  (the dashed line is provided to guide the eye), which is caused by  
 198 NiO being recycled back to Ni by reaction with O<sub>3</sub> (reaction R3a). The result is that the decay  
 199 of Ni due to R2 is retarded, an effect that we have previously observed for reactions such as  
 200 Na and Mg<sup>+</sup> with O<sub>3</sub>.<sup>19, 20</sup> In order to prevent recycling of NiO back to Ni, HCl was added in  
 201 large excess (by a factor of 25 over the maximum  $[\text{O}_3]$  i.e.  $5 \times 10^{14} \text{ molecules cm}^{-3}$ ). NiO reacts  
 202 with HCl via channel R8a, which is effectively thermoneutral:



205 As shown in Figure 6, the decay rate of Ni is now significantly increased at the same  $[\text{O}_3]$ . The  
 206 contribution of R1 (up to 12%) has been subtracted from  $k'$ . The linear regression fit in the

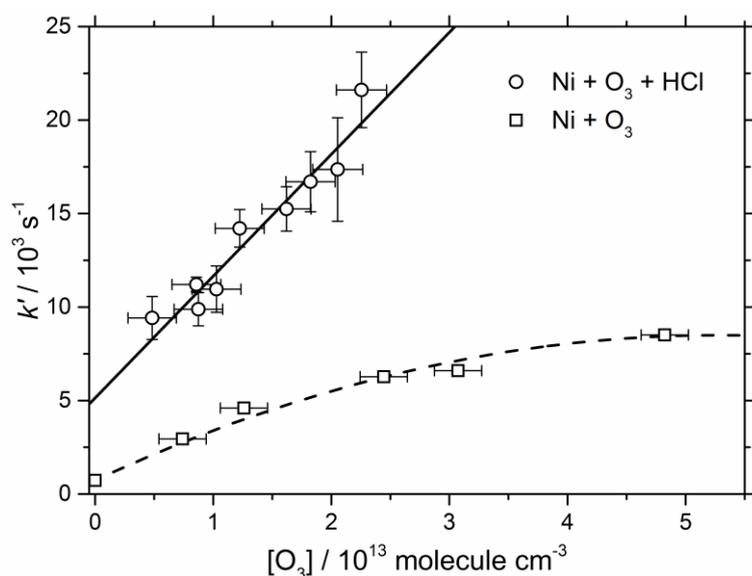
207 presence of HCl yields  $k_2(293\text{ K}) = (6.5 \pm 0.7) \times 10^{-10}\text{ cm}^3\text{ molecule}^{-1}\text{ s}^{-1}$ . Kinetic decays  
 208 recorded over a range of excimer pulse energies (nominal energy 30-70 mJ, maximum fluence  
 209 in the reactor  $6 \times 10^{15}\text{ photons cm}^{-2}$ ) did not display a discernable change in the Ni first-order  
 210 decay rate, suggesting that  $\text{O}_3$  photolysis was not significant at the relatively low excimer  
 211 fluence in the reactor. Indeed, as we show in Section 4,  $k_2$  is very close to the long-range capture  
 212 rate coefficient.

213 The large intercept of the plot in Figure 6 of  $5100\text{ s}^{-1}$  is due to the reaction of Ni with HCl:



215 which therefore has a rate coefficient  $k_9(293\text{ K}) \sim 9 \times 10^{-12}\text{ cm}^3\text{ molecule}^{-1}\text{ s}^{-1}$ . There is a  
 216 significant barrier on the PES for R9, calculated at the CBS-QB3 level of theory<sup>12</sup> to have a  
 217 height of  $23\text{ kJ mol}^{-1}$  (probably an overestimate by around  $10\text{ kJ mol}^{-1}$  to be consistent with the  
 218 magnitude of  $k_9$ ).

219



220

221 Figure 6. Plots of  $k'$  against  $[\text{O}_3]$  in the presence and absence of HCl at 293 K. The added HCl  
 222 ( $5 \times 10^{14}\text{ molecule cm}^{-3}$ ) reacts with NiO to prevent recycling back to Ni via reaction R3a.

223

### 224 3.3 Reactions of NiO

#### 225 NiO + O<sub>3</sub>

226 When studying reactions R3-R7, we were not aware of a direct photolytic source of NiO.  
 227 Furthermore, when we explored using  $\text{N}_2\text{O}$  or  $\text{NO}_2$  to produce NiO from Ni, both these oxidants  
 228 were found to produce NiO too slowly to be useful for studying NiO kinetics (in agreement  
 229 with a previous study of  $\text{Ni} + \text{N}_2\text{O}$ <sup>16</sup>). This meant that NiO had to be produced by reaction of  
 230 Ni with  $\text{O}_3$  via reaction R2. The resulting time-resolved behavior of NiO in the presence of  $\text{O}_3$ ,  
 231  $\text{O}_2$  and other reactants then had to be analysed through kinetic modeling.

232 The model comprised a set of differential equations accounting for the removal of Ni and NiO  
 233 by  $\text{O}_2$ ,  $\text{O}_3$  and other reagents (reactions R4 - R7), as well as diffusional loss out of the volume  
 234 created by the excimer laser and within the field-of-view of the PMT. This volume is essentially  
 235 a cylinder of radius  $r = 0.6\text{ cm}$ . The diffusion coefficients of Ni, NiO and ONiO were  
 236 calculated<sup>21</sup> from the dipole-induced dipole and long-range dispersion forces between the Ni

237 species and N<sub>2</sub>, using the parameters in Table 1. The calculated diffusion rate of Ni out of the  
 238 excimer laser volume, which is given by  $5.81D/r^2$ ,<sup>22</sup> is in accord with the measured  $k_{\text{diff}}$  rates  
 239 of 600 – 700 s<sup>-1</sup> at 3 Torr. Ni has the largest diffusion coefficient because it does not have a  
 240 dipole moment, and its polarizability is smaller than NiO and ONiO. In the model a T<sup>1.83</sup>  
 241 dependence for the diffusion coefficients was assumed.<sup>19, 21</sup>

242

243 Table 1. Parameters and estimated diffusion coefficients for Ni species in N<sub>2</sub> at 293 K.

Species	Dipole moment Debye	Polarizability 10 <sup>-24</sup> cm <sup>3</sup>	Ionization energy / eV	Diffusion coefficient <sup>a</sup> Torr cm <sup>2</sup> s <sup>-1</sup>
Ni	-	6.8 <sup>b</sup>	7.64 <sup>b</sup>	128
NiO	4.43 ± 0.04 <sup>c</sup>	8.4 <sup>d</sup>	9.83 <sup>d</sup>	105
ONiO	-	11 <sup>d</sup>	8.82 <sup>d</sup>	100

244 <sup>a</sup> Calculated using the standard formalism in Maitland et al.<sup>21</sup>

245 <sup>b</sup> Handbook of Physics and Chemistry<sup>23</sup>

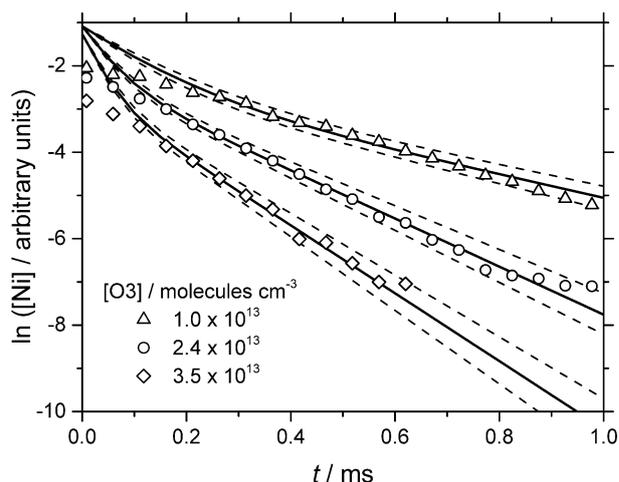
246 <sup>c</sup> Experimental value.<sup>24</sup>

247 <sup>d</sup> Calculated at the B3LYP/6-311+g(2d,p) level of theory.<sup>11</sup>

248

249 Reaction 3 was studied by measuring the time-resolved variation of the measured Ni and NiO  
 250 LIF signals when Ni was produced in the presence of O<sub>3</sub> (and O<sub>2</sub>, but without added HCl). The  
 251 rate coefficients for both channels of R3,  $k_{3a}$  and  $k_{3b}$ , were included in the model as floating  
 252 parameters. A non-linear least-squares fitting procedure was then used to scale and fit the  
 253 model simulation to the experimental data after reaction times longer than 150 μs, when  
 254 quenching of excited Ni states was complete (see above). The best-fit value for R3a (NiO + O<sub>3</sub>  
 255 → Ni + 2O<sub>2</sub>) is  $k_{3a}(293 \text{ K}) = (1.4 \pm 0.5) \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ , and R3b (NiO + O<sub>3</sub> → Ni +  
 256 2O<sub>2</sub>) is  $k_{3b}(293 \text{ K}) = (2.5 \pm 0.7) \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ . The uncertainties were estimated  
 257 using Monte Carlo sampling of the uncertainties of the kinetic parameters in the model,  
 258 assuming a “top-hat” probability distribution, and repeating the fit to each experimental kinetic  
 259 trace 10<sup>3</sup> times. The simulated fits to three experimental decays of Ni are shown in Figure 7,  
 260 with the uncertainty envelope depicted by the dashed lines. The uncertainty mostly arises from  
 261 uncertainties in [O<sub>3</sub>] and  $k_2$ . Note that the model satisfactorily captures the departure from  
 262 single exponential decay behavior.

263



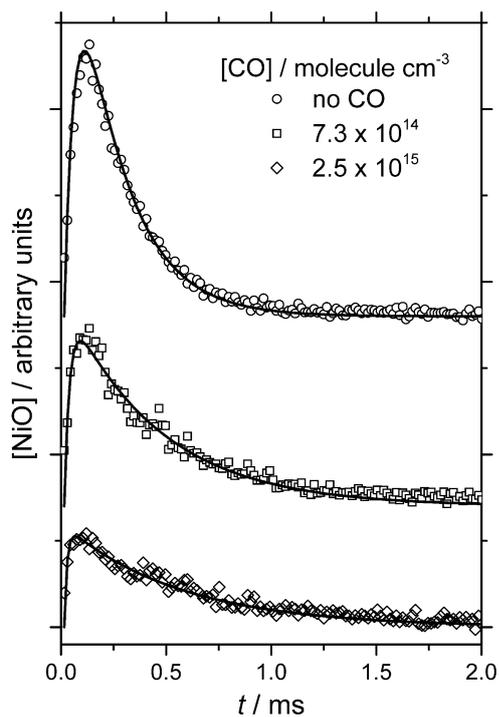
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265 Figure 7. Time-resolved decays of the Ni concentration (in arbitrary units) in the presence of  
 266 O<sub>3</sub> and O<sub>2</sub> at 295 K. The symbols are experimental values and the solid lines fitted simulations  
 267 using the kinetic model. The dashed lines indicate the uncertainty envelope.

268

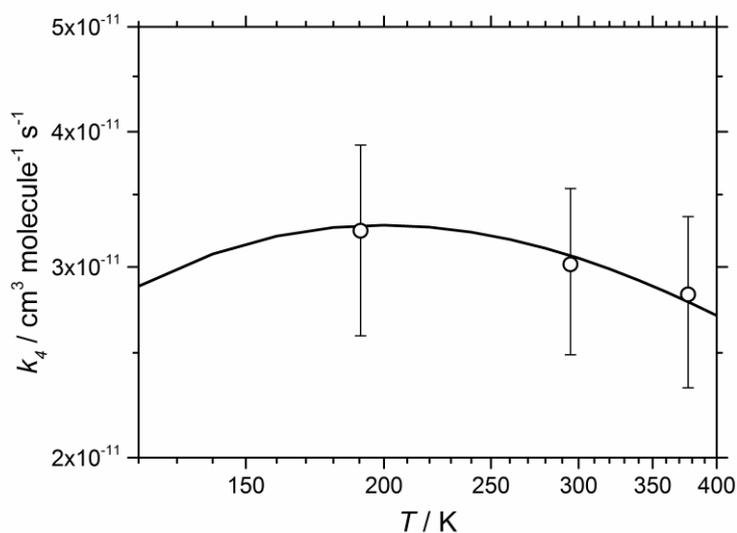
#### 269 NiO + CO, O<sub>2</sub>, H<sub>2</sub>O and CO<sub>2</sub>

270 The kinetic model, including the new rate coefficients for R3a and R3b, was then used to study  
 271 R4 - R7 by fitting to time-resolved kinetic traces of NiO when Ni was produced in the presence  
 272 of fixed [O<sub>3</sub>] and varying concentrations of CO, O<sub>2</sub>, H<sub>2</sub>O or CO<sub>2</sub>. Examples of fitted  
 273 simulations to experimental data for R4 are shown in Figure 8. Background fits in the absence  
 274 of the additional reactant (in this case CO) were first used to constrain the model [Ni] (taken  
 275 before and after runs with CO added). The excellent fit in Figure 8 (top trace, [CO] = 0)  
 276 demonstrates that the kinetic model with  $k_{3a}$  and  $k_{3b}$  optimised to fits of Ni decay profiles can  
 277 equally well simulate the time-resolved NiO profile. Kinetic traces for experiments including  
 278 CO, where  $k_4$  is a floating parameter, were then fitted using the same least-squares method. As  
 279 shown in the middle and lower traces in Figure 8, increasing [CO] decreases the NiO peak,  
 280 while increasing the NiO tail at longer times, due to the recycling reaction  $\text{NiO} + \text{CO} \rightarrow \text{Ni} +$   
 281  $\text{CO}_2$ . Figure 9 illustrates the resulting variation of  $k_4$  over the temperature range 190 – 376 K.  
 282 A linear regression fit yields,  $k_4(190 - 376 \text{ K}) = (3.2 \pm 0.6) \times 10^{-11} (T/200)^{-0.19 \pm 0.05} \text{ cm}^3 \text{ molecule}^{-1}$   
 283  $\text{s}^{-1}$  (uncertainty at the 95% confidence level).



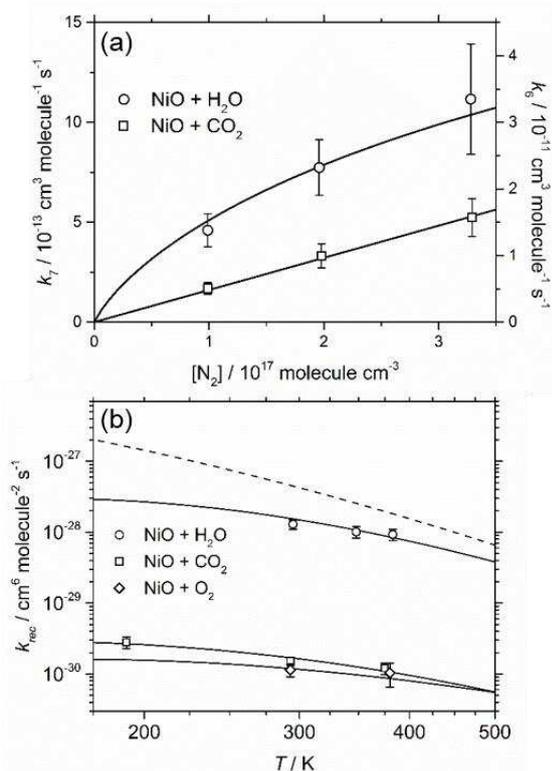
284

285 Figure 8. Time-resolved variation of the NiO concentration (in arbitrary units) in the presence  
 286 of  $O_3$  and varying CO (see figure legend). The solid lines are fitted model simulations to the  
 287 experimental data (open symbols). Conditions:  $T = 295$  K,  $P = 3$  Torr,  $[O_3] = 2 \times 10^{13}$  molecules  
 288  $cm^{-3}$ ,  $[O_2] = 1.2 \times 10^{15}$  molecule  $cm^{-3}$ . The data for each experiment is offset vertically for  
 289 clarity.



290

291 Figure 9. Rate coefficient for the reaction  $NiO + CO \rightarrow Ni + CO_2$  as a function of  $T$ . The solid  
 292 line is a fit using RRKM theory.



293

294 Figure 10. (a) Pressure dependence of R6 ( $\text{NiO} + \text{H}_2\text{O}$ , right-hand ordinate) and R7 ( $\text{NiO} +$   
 295  $\text{CO}_2$ , left-hand ordinate) at 293 K. The symbols are experimental points, and the solid lines are  
 296 RRKM fits. (b) Temperature dependence of R5, R6 and R7. The symbols are experimental  
 297 values, and the solid lines are RRKM fits. Note that the R6 experimental points and RRKM fit  
 298 are at a pressure of 3 Torr, and the dashed line shows the true low-pressure limiting rate  
 299 coefficient for the reaction.

300 The recombination reactions R5-7 with  $\text{N}_2$  as the third body were investigated using the same  
 301 kinetic model to fit time-resolved traces of  $\text{NiO}$ . Reaction R5 ( $\text{NiO} + \text{O}_2$ ) proved challenging  
 302 to fit because  $\text{O}_2$  also recombines with  $\text{Ni}$  (R1), so that varying  $[\text{O}_2]$  caused a comparatively  
 303 small change in the time-resolved variation of  $\text{Ni}$ , which was only apparent above the signal-  
 304 to-noise of the  $\text{NiO}$  LIF signal at temperatures between 293 and 380 K. Figure 10a shows the  
 305 pressure dependence of the second-order rate coefficients for R6 and R7 at 293 K. Note that  
 306 R6 is in the fall-off region over the pressure range 3 – 10 Torr. Figure 10b illustrates the  
 307 negative temperature dependences of  $k_5$ ,  $k_6$  and  $k_7$ , with R6 proceeding  $\sim 100$  times faster than  
 308 R5 and R7. R5 and R7 are close to the low pressure limit in the experimental pressure range,  
 309 yielding  $k_5(293 - 380 \text{ K}) = (1.1 \pm 0.3) \times 10^{-30} (T/293)^{-0.4 \pm 0.1}$  and  $k_7(191 - 375 \text{ K})$   
 310  $= (1.55 \pm 0.22) \times 10^{-30} (T/293)^{-1.28 \pm 0.34} \text{ cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}$  (quoted uncertainties at the 95%  
 311 confidence level). A theoretical extrapolation of  $k_6$  to the low-pressure limit is presented below.

312

## 313 4. Discussion

314

315 The individual second-order rate coefficients measured for R1 - R7 are listed as a function of  
 316 temperature and pressure in Table S1 (in the Supporting Information, SI). In order to  
 317 understand the unusual behaviour of several of these reactions, and also to extrapolate to the  
 318 low pressure conditions of the MLT, electronic structure calculations were combined (where  
 319 appropriate) with Rice-Ramsperger-Kassel-Markus (RRKM) theory. The hybrid density

320 functional B3LYP method was employed from within the Gaussian 16 suite of programs,<sup>11</sup>  
321 combined with the 6-311+G(2d,p) triple- $\zeta$  basis set. Molecular geometries were first optimized  
322 and checked for wave function stability, and their respective vibrational frequencies were  
323 calculated. More accurate energies were then determined using the Complete Basis Set (CBS-  
324 QB3) method of Petersson and co-workers.<sup>12</sup> The resulting geometries, rotational constants,  
325 vibrational frequencies and heats of formation of NiO, ONiO<sub>2</sub>, NiCO<sub>3</sub> and Ni(OH)<sub>2</sub> are listed  
326 in Table S2 in the SI, and the geometries of these molecules are illustrated in Figure S1. Table  
327 S2 also contains relevant data for the stationary points on the NiO + CO and Ni + O<sub>2</sub> reaction  
328 potential energy surfaces (PES).

329 RRKM theory is used to calculate the pressure-dependent recombination rate coefficient  
330 between the low- and high-pressure limiting rate coefficients i.e.  $k_{\text{rec},0}$  and  $k_{\text{rec},\infty}$ , respectively.  
331 Here we use a solution of the Master Equation (ME) based on the inverse Laplace transform  
332 method.<sup>25</sup> We have applied this formalism previously to reactions of metal-containing species  
333 where stable intermediates are present on the PES,<sup>26,27</sup> and so only a brief description is given  
334 here. These reactions proceed via the formation of an excited adduct, which can either  
335 dissociate back to reactants or forward to products, or be stabilized into one or more energy  
336 wells by collision with a third body. The internal energy of the adduct is divided into a  
337 contiguous set of grains (width typically 30 – 70 cm<sup>-1</sup>), each containing a bundle of  
338 rovibrational states. Each grain is then assigned a set of microcanonical rate coefficients for  
339 dissociation back to the reactants, and to products (if appropriate). These rate coefficients are  
340 determined using inverse Laplace transformation to link them directly to  $k_{\text{rec},\infty}$ . For these neutral  
341 reactions,  $k_{\text{rec},\infty}$  should be close to the typical capture rate coefficient with a small positive  
342 temperature dependence characteristic of a long-range potential governed by the dispersion and  
343 dipole-induced dipole force (the measured dipole moment of NiO is  $4.43 \pm 0.04$  D<sup>24</sup>). The  
344 density of states of the adduct is calculated using the Beyer–Swinehart algorithm<sup>28</sup> for the  
345 vibrational modes (without making a correction for anharmonicity) and a classical densities of  
346 states treatment for the rotational modes (the vibrational frequencies and rotational constants  
347 are listed in Table S2 of the SI). The probability of collisional transfer between grains was  
348 estimated using the exponential down model, where the average energy for downward  
349 transitions  $\langle \Delta E \rangle_{\text{down}}$  was set to between 200 and 300 cm<sup>-1</sup> which is typical of M = N<sub>2</sub> at 300 K,  
350 with a  $T^\alpha$  temperature dependence where  $\alpha$  is a small number between -1 and 1.<sup>28</sup> The ME,  
351 which describes the evolution with time of the grain populations of the adduct, is expressed in  
352 matrix form and then solved to yield the rate coefficients to form stable intermediates or  
353 products, at a specified pressure and temperature.

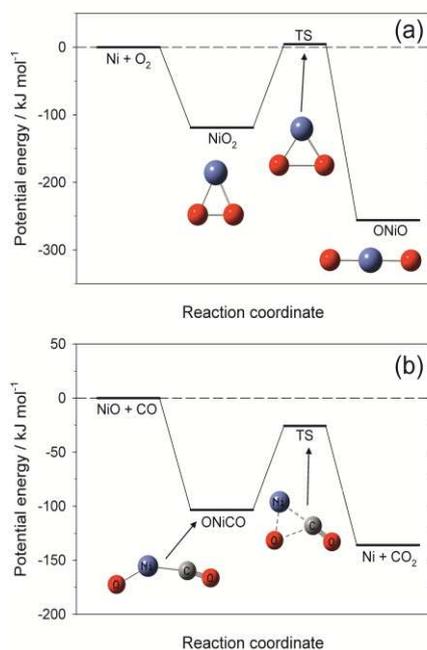
#### 354 **4.1 Ni + O<sub>3</sub> and NiO + O<sub>3</sub>**

355 R2 is fast with  $k_2(293 \text{ K}) = (6.5 \pm 0.7) \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ . The ionization energy of Ni  
356 is too large (7.64 eV<sup>23</sup>) for a simple electron jump mechanism<sup>29</sup> to operate, since the avoided  
357 crossing between the ionic and covalent surfaces would be at only 2.6 Å with a rate coefficient  
358 of  $1.0 \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ . Instead, the rate coefficient can be compared with long-range  
359 capture theory, where the reaction is controlled by the dipole-induced dipole and dispersion  
360 forces. Using the data for Ni in Table 1, and a polarizability for O<sub>3</sub> of  $3.21 \times 10^{-24} \text{ cm}^3$  and  
361 ionization potential of 12.43 eV,<sup>23</sup> the ion-induced dipole capture rate is  $1.1 \times 10^{-10} \text{ cm}^3$   
362  $\text{molecule}^{-1} \text{ s}^{-1}$ , and that due to the dispersion force is  $4.8 \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ . Following  
363 the recommendation of Georgievskii and Klippenstein,<sup>30</sup> a good estimate of the overall capture  
364 rate is given by 1.3 times the larger of these terms. This gives  $6.2 \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ,  
365 which is in good agreement with the measured  $k_2$ , and provides further evidence that O<sub>3</sub>  
366 photolysis was not significant in the experiment. A small  $T^{1/6}$  dependence for  $k_2$  is then  
367 expected.<sup>30</sup>

368 In the case of R3, the overall rate coefficient  $k_{3a} + k_{3b} = (3.9 \pm 0.9) \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ,  
369 which is 60% of the capture rate coefficient estimated in the same way as for R2. The branching  
370 ratio to form Ni is 36%.

#### 371 4.2 Ni + O<sub>2</sub> + N<sub>2</sub>

372 The previous measurement of  $k_1$  with M = Ar was  $(1.7 \pm 0.3) \times 10^{-30} \text{ cm}^6 \text{ molecules}^{-2} \text{ s}^{-1}$  at 296  
373 K.<sup>10</sup> This is slower by a factor of  $0.46 \pm 0.15$  than the result of the present study, which may  
374 be explained by the relative efficiencies of Ar and N<sub>2</sub> as third bodies.<sup>28</sup> As shown in Figure 5,  
375 the rate coefficient for R1 increases with decreasing T down to around 210 K, and then plateaus  
376 and starts to decrease slightly. The PES for this reaction is shown in Figure 11(a), based on the  
377 recent high level theory study.<sup>13</sup> The reaction produces the triangular NiO<sub>2</sub> initially (Figure S1  
378 in the SI); an asymmetric Ni-O<sub>2</sub> with a Ni-O-O bond angle of 123.5° is about 80 kJ mol<sup>-1</sup> less  
379 stable, and linear Ni-O-O is not stable. The Ni can then insert into the O<sub>2</sub> to form the much  
380 more stable linear ONiO. Reaction to ONiO is therefore favoured by the much higher density  
381 of rovibrational states, but only if the low barrier between these two isomers (3.7 kJ mol<sup>-1</sup> with  
382 respect to the reactants) can be cleared. At the lowest experimental pressure of 3 Torr, the  
383 reaction largely produces NiO<sub>2</sub> at very low temperatures, and ONiO at higher temperatures: for  
384 example, 54% of the product is NiO<sub>2</sub> at 100 K, but this decreases to only 1.8% at 600 K. The  
385 product ratio is somewhat pressure dependent: at a pressure of 0.003 Torr which is typical of  
386 85 km in the atmosphere, 35% of the product is NiO<sub>2</sub> at 100 K, and 1.6% at 600 K. Note that  
387 both triangular NiO<sub>2</sub> and ONiO (but not linear or asymmetric Ni-O<sub>2</sub>) have been observed in  
388 matrix isolation studies, with relative yields depending on the nature of the matrix.<sup>13</sup> The  
389 RRKM fit through the data is shown by the solid line in Figure 5, using the optimized  
390 parameters listed in Table S3 in the SI. The fitted temperature dependence is then:  
391  $\log_{10}(k_5(T)/\text{cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}) = -37.592 + 7.168\log_{10}(T) - 1.5650(\log_{10}(T))^2$ , with an  
392 uncertainty of  $\pm 11\%$  within the experimental temperature range (191 – 455 K).



393

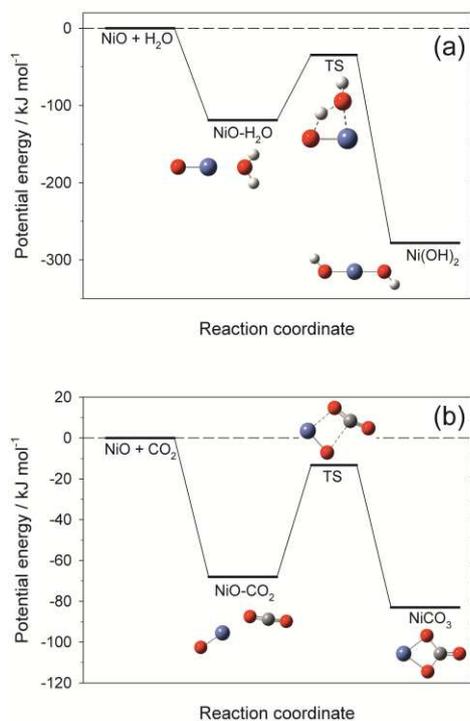
394 Figure 11. (a). Potential energy surface for the Ni + O<sub>2</sub> reaction forming NiO<sub>2</sub> and ONiO.  
395 Energies from Hübner and Himmel.<sup>13</sup> (b) Potential energy surface for the NiO + CO → Ni +  
396 CO<sub>2</sub> reaction. Calculated at the CBS-QB3 level of theory. Atom colors: oxygen (red); carbon  
397 (grey); nickel (blue)

398 **4.3 NiO + CO**

399 As shown in Figure 9,  $k_4$  is around  $3 \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ , with a possibly slightly negative  
 400 T-dependence. This result can be compared with three other metal oxide + CO reactions for  
 401 which kinetic data is available:  $k(\text{NaO} + \text{CO}, 296 \text{ K}) = (2.1 \pm 0.5) \times 10^{-10} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$   
 402 <sup>31</sup>;  $k(\text{MgO} + \text{CO}, 295 \text{ K}) = (1.1 \pm 0.3) \times 10^{-11} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ,<sup>32</sup> and  $k(\text{FeO} + \text{CO}, 295 \text{ K}) =$   
 403  $1.5 \times 10^{-13} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ , obtained by extrapolating the expression derived by Smirnov<sup>33</sup>  
 404 from detailed balance with high temperature (1180 - 2380 K) shock tube measurements of the  
 405 reverse reaction  $\text{Fe} + \text{CO}_2 \rightarrow \text{FeO} + \text{CO}$ . Thus, while being within a factor of 3 and 7 of the  
 406 MgO and NaO + CO reactions, respectively, the NiO reaction is ~200 times faster than the FeO  
 407 reaction. A likely reason for this is discussed below.

408 Figure 11(b) illustrates the PES for NiO + CO. The reaction involves forming a fairly stable  
 409 OC-NiO intermediate that is bound by  $104 \text{ kJ mol}^{-1}$ . However, this can promptly dissociate  
 410 over a barrier to the Ni + CO<sub>2</sub> products, since the barrier is submerged by  $26 \text{ kJ mol}^{-1}$  with  
 411 respect to the reactants. The RRKM fit to this reaction (parameters in Table S3) is shown in  
 412 Figure 9. The very slight temperature dependence is captured well, with  $k_4$  increasing slightly  
 413 between 400 and 200 K, and then decreasing slowly at lower temperatures. The barrier on the  
 414 PES is far enough below the surface that an insignificant fraction ( $8 \times 10^{-8}$  at 295 K and 10  
 415 Torr) of the OC-NiO complex forms compared to the Ni + CO<sub>2</sub> product. Nevertheless, because  
 416 of the relatively large uncertainty in the experimental data, we recommend a simple linear fit  
 417 through the three data points:  $k_4(T) = (3.2 \pm 0.6) \times 10^{-11} (T/200)^{-0.19 \pm 0.05} \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ,  
 418 which agrees within 20% with the RRKM fit between 120 and 400 K.

419 In contrast to R4, FeO + CO has a barrier that is calculated at the CBS-QB3 level to be  $21.6 \text{ kJ}$   
 420  $\text{mol}^{-1}$  above the reactants, which explains why the reaction is so much slower.



421  
 422 Figure 12. (a) The potential energy surface for the recombination of NiO and H<sub>2</sub>O. (b) The  
 423 potential energy surface for the recombination of NiO and CO<sub>2</sub>. Calculated at the CBS-QB3  
 424 level of theory. Atom colors: oxygen (red); hydrogen (white); carbon (grey); nickel (blue).

425

#### 426 **4.4 NiO + O<sub>2</sub>, H<sub>2</sub>O and CO<sub>2</sub>**

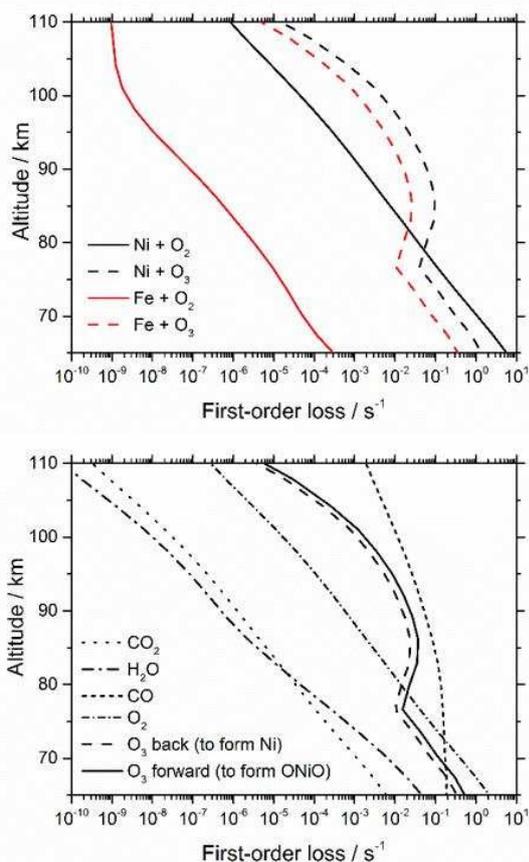
427 Figure 10(a) shows that R6 (NiO + H<sub>2</sub>O) is already in the fall-off region over the experimental  
428 pressure range, and so an RRKM extrapolation to the low pressures that characterise the MLT  
429 is essential. The reaction PES is illustrated in Figure 12(a). NiO forms a fairly strongly-bound  
430 ONi-H<sub>2</sub>O adduct, which then rearranges over a submerged barrier to form the di-hydroxide,  
431 Ni(OH)<sub>2</sub>. Since the overall reaction is strongly exothermic ( $\Delta H^{\circ}(0\text{ K}) = -278\text{ kJ mol}^{-1}$ ), the  
432 reaction is comparatively fast and hence in the fall-off region at 1 Torr of N<sub>2</sub>. The RRKM fit is  
433 shown with a solid line in Figure 10(a) (see Table S3 for the optimized RRKM parameters).  
434 Figure 10(b) shows the measured temperature dependence of  $k_6$  at 3 Torr, which compares well  
435 with the RRKM fit. The dashed line illustrates the low-pressure limiting rate coefficient,  $k_{6\text{rec},0}$ ,  
436 which is significantly larger (e.g. by a factor of 3 at 300 K), as expected. The rate coefficient can  
437 then be expressed by a Lindemann expression modified by a broadening factor  $F_c$ :<sup>34</sup>  $\log_{10}(k_{6\text{rec},0}$   
438  $/ \text{cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}) = -29.7651 + 5.2064\log_{10}(T) - 1.7118(\log_{10}(T))^2$ ;  $k_{6\text{rec},\infty} = 6.0 \times 10^{-10} \exp(-$   
439  $171/T) \text{ cm}^3 \text{ molecule}^{-1} \text{ s}^{-1}$ ;  $F_c = 0.84$ , with an uncertainty of  $\pm 16\%$  within the experimental  
440 temperature range (191 – 375 K).

441 The PES for reaction R7 is illustrated in Figure 12(b). This shows that the reaction involves  
442 first forming an ONi-OCO adduct, which rearranges over a submerged barrier to form NiCO<sub>3</sub>.  
443 Although the overall reaction is not particularly exothermic ( $\Delta H^{\circ}(0\text{ K}) = -83\text{ kJ mol}^{-1}$ ), the large  
444 number of low vibrational frequencies in the carbonate product (Table S2) ensures that the rate  
445 coefficient is reasonably large. The ONi-OCO adduct is only bound by 67 kJ mol<sup>-1</sup>, and so the  
446 two lowest frequencies (35 and 121 cm<sup>-1</sup>) were treated as a 2-dimensional hindered rotor, with  
447 a barrier  $V_0 = 14\text{ kJ mol}^{-1}$  (see Table S3 for the other fitted RRKM parameters). The pressure-  
448 and temperature-dependent RRKM fits to the experimental data-points are shown in Figure 10,  
449 yielding  $\log_{10}(k_7 / \text{cm}^6 \text{ molecule}^{-2} \text{ s}^{-1}) = -41.4265 + 10.9640\log_{10}(T) - 2.5287(\log_{10}(T))^2$ , with  
450 an uncertainty of  $\pm 15\%$  within the experimental temperature range (191 – 375 K).

451 Reaction 5 involves simple addition of O<sub>2</sub> to NiO to form the most stable isomer, which is  
452 ONiO<sub>2</sub> (Figure S1) with triplet spin multiplicity ( $\Delta H^{\circ}(0\text{ K}) = -145\text{ kJ mol}^{-1}$ ). Singlet and quintet  
453 ONiO<sub>2</sub> are less stable than the triplet by 51 and 57 kJ mol<sup>-1</sup>, respectively. The most stable  
454 ozonide NiO<sub>3</sub> isomer is the triplet, which is 124 kJ mol<sup>-1</sup> less stable than the triplet ONiO<sub>2</sub>  
455 isomer. The RRKM fit (Figure 10b; fitted parameters in Table S3), yields  $\log_{10}(k_5(T) / \text{cm}^6$   
456  $\text{ molecule}^{-2} \text{ s}^{-1}) = -41.0913 + 10.1064\log_{10}(T) - 2.2610(\log_{10}(T))^2$  with an uncertainty of  $\pm 22\%$   
457 within the experimental temperature range (293 – 380 K).

458 In comparison with the analogous FeO reactions at room temperature,<sup>35</sup> the NiO reactions with  
459 H<sub>2</sub>O and CO<sub>2</sub> are faster by factors of 1.5 and 5, respectively, whereas the O<sub>2</sub> recombination is  
460 3 times slower.

#### 461 **4.5 Atmospheric implications**



462

463 Figure 13. First-order removal rates in the MLT between 65 and 110 km, at midday during  
 464 January at 40°N. Top panel: removal rates of Ni and Fe through reaction with O<sub>3</sub> and O<sub>2</sub>, where  
 465 the Fe rate coefficients are from a recent review.<sup>36</sup> Bottom panel: removal rates of NiO.

466

467 In order to ascertain the relative importance in the MLT of reactions R1 – R7, the rate  
 468 coefficients were combined with profiles of the relevant reactant species, temperature and  
 469 pressure. The profiles of these parameters between 65 and 110 km were taken from the Whole  
 470 Atmosphere Community Climate Model (WACCM).<sup>37</sup> Figure 13 shows the first-order loss  
 471 rates of Ni and NiO as a function of height. The top panel compares the Ni and Fe loss rates.  
 472 In contrast to Fe, the Ni + O<sub>2</sub> loss pathway becomes competitive with the Ni + O<sub>3</sub> pathway  
 473 around 80 km, which is below the 87 km peak of the Ni layer.<sup>4</sup> This difference is due to the  
 474 ~1000 times faster reaction of Ni with O<sub>2</sub> compared to Fe<sup>38</sup> under these conditions. Note that  
 475 O<sub>3</sub> increases in the MLT by around 1 order of magnitude at night,<sup>3</sup> when the formation rates of  
 476 NiO and ONiO would be essentially the same.

477 Inspection of the bottom panel of Figure 13 shows that the recycling reaction of NiO with CO  
 478 dominates, due to the relatively fast k<sub>4</sub> and the presence of significant CO in the MLT (formed  
 479 primarily by CO<sub>2</sub> photolysis). The O<sub>3</sub> reaction recycling NiO back to Ni is also significant. In  
 480 the future we plan to investigate the reaction NiO + O which, by analogy with the other  
 481 meteoric metals,<sup>3</sup> may also make a major contribution to NiO recycling.

482 In terms of the recombination reactions, R5 (NiO + O<sub>2</sub>) is the most competitive because of  
 483 the high concentration of O<sub>2</sub>. Although k<sub>6</sub> is around 2 orders of magnitude faster (Figure  
 484 10(b)), the low concentration of H<sub>2</sub>O (just a few ppm) means that R6 is the least important  
 485 recombination reaction above 80 km. However, given that Ni(OH)<sub>2</sub> is so much more stable

486 than ONiO<sub>2</sub> or NiCO<sub>3</sub>, it is likely to be the major Ni reservoir, formed from ONiO<sub>2</sub> and  
 487 NiCO<sub>3</sub> through these exothermic reactions:



490 The potential energy surfaces for R10 and R11 are shown in Figure S2 (SI). In both cases, an  
 491 initial complex with H<sub>2</sub>O forms which then dissociates over a submerged barrier to Ni(OH)<sub>2</sub>.

492

## 493 5. Conclusions

494 The reactions of Ni and NiO with atmospherically relevant species have been experimentally  
 495 investigated using the PLP-LIF technique, all but one of them (R1) for the first time. The  
 496 reaction kinetics are then explained using electronic structure theory combined with RRKM  
 497 calculations. Reactions R1 and R4 exhibit unusual temperature dependences, which are  
 498 explained by shallow submerged barriers on their potential energy surfaces. Theoretical fits to  
 499 the experimental data were then used to extrapolate to the temperatures (120 – 240 K) and  
 500 pressures (< 5 Pa) of the MLT region. A summary of the rate coefficients is given in Table 2.

501

502 Table 2. Summary of reaction rate coefficients measured in the present study.

No.	Reaction	Rate coefficient <sup>a</sup>
R1	Ni + O <sub>2</sub> (+N <sub>2</sub> ) → NiO <sub>2</sub>	$k_{rec,0} = 10^{-37.592 + 7.168 \log_{10}(T) - 1.5650(\log_{10}(T))^2}$
R2	Ni + O <sub>3</sub> → NiO + O <sub>2</sub>	$(6.5 \pm 0.7) \times 10^{-10} (T/293)^{1/2}$
R3a	NiO + O <sub>3</sub> → Ni + 2O <sub>2</sub>	$(1.4 \pm 0.5) \times 10^{-10} (T/293)^{1/2}$
R3b	NiO + O <sub>3</sub> → NiO <sub>2</sub> + O <sub>2</sub>	$(2.5 \pm 0.7) \times 10^{-10} (T/293)^{1/2}$
R4	NiO + CO → Ni + CO <sub>2</sub>	$(3.2 \pm 0.6) \times 10^{-11} (T/200)^{-0.19 \pm 0.05}$
R5	NiO + O <sub>2</sub> (+N <sub>2</sub> ) → ONiO <sub>2</sub>	$k_{rec,0} = 10^{-41.0913 + 10.1064 \log_{10}(T) - 2.2610(\log_{10}(T))^2}$
R6	NiO + H <sub>2</sub> O (+N <sub>2</sub> ) → Ni(OH) <sub>2</sub>	$k_{rec,0} = 10^{-29.7651 + 5.2064 \log_{10}(T) - 1.7118(\log_{10}(T))^2}$ $k_{rec,\infty} = 6.0 \times 10^{-10} \exp(-171/T)$ <sup>b</sup> ; F <sub>c</sub> = 0.84
R7	NiO + CO <sub>2</sub> (+N <sub>2</sub> ) → NiCO <sub>3</sub>	$k_{rec,0} = 10^{-41.4265 + 10.9640 \log_{10}(T) - 2.25287(\log_{10}(T))^2}$

503 <sup>a</sup> Units for termolecular reactions: cm<sup>6</sup> molecule<sup>-2</sup> s<sup>-1</sup> (see Section 4 for uncertainties). Units for  
 504 bimolecular reactions: cm<sup>3</sup> molecule<sup>-1</sup> s<sup>-1</sup>.

505

506 In the MLT, Ni reacts most rapidly with O<sub>3</sub> above 80 km, and at a similar rate with O<sub>2</sub> below  
 507 80 km. Interestingly, the NiO + CO reaction rapidly recycles NiO to Ni, which may explain  
 508 the unexpected observation of a significant concentration of Ni atoms below 80 km, unlike  
 509 other metals such as Fe.<sup>4</sup>

510

511 **Supporting Information.** Contains the following: Table S1 listing second-order rate  
 512 coefficients as a function of temperature and pressure; Table S2 listing the molecular properties  
 513 and heats of formation of Ni-containing molecules; Table S3 of parameters used in the RRKM

514 fits to reaction kinetics; Figure S1 illustrating the geometries of Ni-containing molecules; and  
515 Figure S2 showing the potential energy surfaces for R10 and R11.

516

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522 Natasha Aylett (University of Leeds) for supplying WACCM output. All the second-order rate  
523 coefficients, and the *ab initio* geometries, rotational constants, vibrational frequencies and  
524 heats of formation for the relevant Ni-containing molecules are listed in the Supplementary  
525 Information.

526

527

528 **References**

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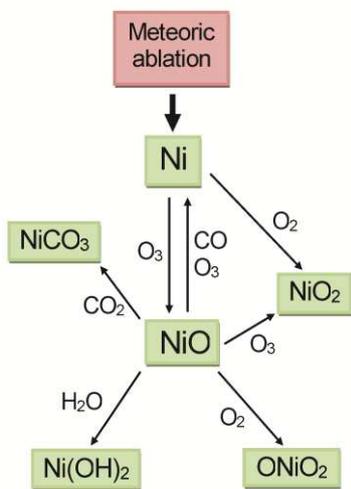
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627 TOC graphic

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