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3	Drained cavity expansion analysis with a unified state
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#### 21 ABSTRACT

22 This paper presents an analytical solution for drained expansion in both spherical and 23 cylindrical cavities with a unified state parameter model for clay and sand (CASM) (Yu, 1998). The solution developed here provides the stress and strain fields during the expansion of a 24 25 cavity from an initial to an arbitrary final radius. Small strains are assumed to the elastic region 26 and large strains are applied for soil in the plastic region by using logarithmic strain definitions. 27 Since its development, the unified CASM model has been demonstrated by many researchers 28 to be able to capture the overall soil behaviour for both clay and sand under both drained and 29 undrained loading conditions. In this study, the CASM model is used to model soil behaviour 30 whilst we develop a drained cavity expansion solution with the aid of an auxiliary variable. This is an extension of the undrained solution presented by the authors (Mo and Yu, 2017). The 31 32 parametric study investigates the effects of various model constants including the stress-state 33 coefficient and the spacing ratio on soil stress paths and cavity expansion curves. Both London 34 clay and Ticino sand are modelled under various initial stress conditions and initial state 35 parameters. The newly-developed analytical solution highlights the potential applications in 36 geotechnical practice (e.g. for the interpretation of cone penetration test (CPT) data) and also 37 serves as useful benchmarks for numerical simulations of cavity expansion problems in critical 38 state soils.

39

#### 40 KEYWORDS

41 Cavity expansion analysis, analytical solution, drained analysis, unified state parameter model,

42 cone penetration test

43

44 List of notations provided on Page 3

# 46 NOTATION

а	radius of cavity
С	radius of the elastic/plastic boundary
е	void ratio of granular material
m	parameter to combine cylindrical $(m = 1)$ and spherical $(m = 2)$ analysis
n	stress-state coefficient for CASM
p',q	mean stress and deviatoric stress
$p_0'$	initial mean effective stress
$p_{y0}^{\prime}$	preconsolidation pressure
r	radial position of soil element around the cavity
$r^*$	spacing ratio for the concept of state parameter
G	elastic shear modulus
K	elastic bulk modulus
$R_0$	isotropic overconsolidation ratio, defined as $p'_{y0}/p'_0$
X	auxiliary independent variable, defined as $u/r$
δ,γ	volumetric and shear strains
$arepsilon_p$ , $arepsilon_q$	volumetric and shear strains
$arepsilon_r$ , $arepsilon_ heta$	radial and tangential strains
η	stress ratio, defined as $q/p'$
μ	Poisson's ratio of soil
ν	specific volume, defined as $1 + e$
$\sigma_r'$ , $\sigma_ heta'$	radial and tangential stresses
ξ	state parameter
$\xi_R$	reference state parameter
Μ, κ, λ, Γ, Λ	critical state soil parameters

#### 48 INTRODUCTION

49 The cavity expansion method and its applications to geotechnical problems have been extensively developed in the last five decades (e.g., Yu 2000). While early research works-was 50 51 mainly focused on the expansion in elastic materials, analytical solutions have been developed 52 using increasingly more sophisticated constitutive soil models (e.g., Palmer and Mitchell 1971; Vesic 1972; Carter et al. 1986; Yu and Houlsby 1991; Collins and Yu 1996; Chen and 53 Abousleiman 2012, 2013, 2016, 2017; Mo et al. 2014; Vrakas and Anagnostou 2014; Mo and 54 Yu 2017). As a result, the solutions have been particularly of interest to geotechnical 55 56 engineering problems, such as in-situ soil testing, pile foundations, and tunnelling, largely due 57 to their successful applications in providing simple but useful geotechnical solutions.

58 Perfect plasticity was initially adopted for cavity expansion in soils under either undrained or 59 drained conditions. Total stress analysis of cohesive soil is typically used for the Tresca and 60 von Mises materials, whereas the drained behaviour of soil is modelled by the effective stress analysis for the Mohr Coulomb material. Among the solutions in elastic-perfectly plastic soils, 61 62 one of the milestones in cavity expansion solutions was provided by Yu and Houlsby (1991), 63 who derived a unified analytical solution of cavity expansion in dilatant elastic-plastic soils, 64 using the Mohr-Coulomb yield criterion with a non-associated flow rule. The large strain analysis in the plastic region, with the aid of a series expansion, was used to derive a rigorous 65 66 closed-form solution for both cylindrical and spherical cavities. However, to account for the 67 variation of soil strength during cavity expansion, a solution using a strain-hardening/softening plasticity model was clearly necessary. 68

69 As the most widely used strain-hardening or softening models in soil mechanics, critical state 70 soil models (Schofield and Wroth 1968) have been used to derive cavity expansion solutions 71 under both drained and undrained conditions in the last two decades (e.g., Collins and Yu 1996; 72 Cao et al. 2001; Chen and Abousleiman 2012, 2013, 2016; Mo and Yu 2017). It should be noted 73 that drained cavity expansion solutions in critical state soils are very limited due to the unknown 74 stress paths and variations of the specific volume during the cavity expansion process. Palmer 75 and Mitchell (1971) were the first to derive an approximate small-strain analytical solution for cylindrical cavity expansion in normally consolidated clay. Similarity solutions for drained 76 77 cavities from zero initial radius in critical state soils were presented by Collins et al. (1992) and 78 Collins and Stimpson (1994), who provided the limit cavity pressures for both spherical and 79 cylindrical cavities. However, the asymptotic solutions are only valid for large cavity expansion 80 due to the approach of geometric self-similarity. Other similarity solutions were also developed 81 by Russell and Khalili (2002) using the conventional Mohr-Coulomb failure criterion and a 82 state parameter sand behaviour model with a non-linear critical state line. More recently, semianalytical solutions for crushable granular materials were proposed by Jiang and Sun (2012)
using a new critical state line, with a state-dependent dilantancy and a bounding surface
plasticity model. Again, similarity transformation was introduced for the cavity expansion
solutions, and plastic deformation was assumed as zero for constant stress ratio.

87 By abandoning the assumption of similarity, drained solutions for the expansion of cylindrical 88 cavities in the Modified Cam-clay and bounding surface plasticity soils were reported by Chen 89 and Abousleiman (2013, 2016), with the aid of an auxiliary variable in the plastic region, which 90 aims to convert the Eulerian formulation into Lagrangian form. The approach of auxiliary 91 variable is also applied to the proposed drained solutions for the general shear strain 92 hardening/softening Drucker-Prager models (Chen and Abousleiman, 2017) and for the unified 93 hardening parameter-based critical state model (Li et al. 2017). However, as pointed out by Yu 94 (1998) among others, it is also true that the conventional critical state models are less suitable 95 for modelling sand behaviour and heavily overconsolidated clays. Hence existing solutions for 96 cavity expansion for a unified critical state soil model for clay and sand are still limited.

97 In the present paper, an analytical solution for the expansion of both spherical and cylindrical 98 cavities with a unified state parameter model for clay and sand (CASM) (Yu, 1998) is 99 developed. This is an extension of the undrained cavity expansion solutions of Mo and Yu (2017) to drained loading conditions. After introducing the unified state parameter model 100 101 CASM, the small strain theory is applied in the elastic region, and the large strain assumption is used for soil in the plastic region. The approach of auxiliary variable used by Chen and 102 103 Abousleiman (2013) is employed for our drained analysis, which is valid for the expansion of 104 either a spherical or a cylindrical cavity in clay or sand material. In this paper, the results of 105 cavity expansion in both London clay and Ticino sand are presented for stress paths and cavity 106 expansion curves. A parametric study is also provided to investigate the effects of the stress-107 state coefficient and the spacing ratio, as well as the effects of initial stress condition and initial 108 state parameter of the soil. The interpretation of CPT data using the proposed solution is also 109 compared with data from relevant calibration chamber tests.

110

#### 111 PROBLEM DESCRIPTION

112 A spherical or cylindrical cavity with initial radius  $a_0$  in an infinite soil (Fig. 1a) is assumed to 113 be expanded under fully drained conditions. As reported in Mo and Yu (2017), Fig. 1b 114 schematically illustrates the geometry and kinematics of cavity expansion. The initial stress 115 state is assumed as isotropic, with  $\sigma'_{r,0} = \sigma'_{\theta,0} = p'_0$ . For the cylindrical case,  $\sigma'_{z,0}$  is equal to  $p'_0$ , 116 and the effect of  $\sigma'_z$  is not included in this study. For soil with an overconsolidated stress history, the preconsolidation pressure is referred to as  $p'_{y0}$ , and  $R_0 = p'_{y0}/p'_0$  represents the isotropic overconsolidation ratio in terms of the mean effective stress. The initial specific volume is referred to as  $v_0$ , and the specific volume varies during the process of expansion for the drained analysis. Note that a compression positive notation is used throughout this paper, consistent with the undrained solution of Mo and Yu (2017).

For cavity expansion problems, the stresses of soil must satisfy the following quasi-staticequilibrium equation:

124 
$$\sigma_{\theta}' - \sigma_{r}' = \frac{r}{m} \frac{d \, \sigma_{r}'}{d \, r} \tag{1}$$

where the parameter 'm' is used to integrate both spherical (m = 2) and cylindrical (m = 1)scenarios (following Yu and Houlsby 1991, Collins and Yu 1996, and Mo and Yu 2017);  $\sigma'_r$ and  $\sigma'_{\theta}$  are the effective radial and tangential stresses, and r is the radius of the material element  $(r_0$  is the initial position before cavity expansion). The symbol 'd' denotes the Eulerian derivative for every material particle at a specific moment.

130 According to Collins and Yu (1996), the mean and deviatoric effective stresses (p'; q) for 131 cavity expansion problems can be defined as follows:

132 
$$p' = \frac{\sigma'_r + m \cdot \sigma'_{\theta}}{1 + m} \qquad (2)$$
$$q = \sigma'_r - \sigma'_{\theta}$$

133 Accordingly, the volumetric and shear strains ( $\delta$ ;  $\gamma$ ) can be written as:

134 
$$\delta = \varepsilon_r + m \cdot \varepsilon_\theta \\ \gamma = \varepsilon_r - \varepsilon_\theta$$
(3)

As stated in Mo and Yu (2017), the definitions of 'p'', 'q' provided in eq. (2) and ' $\delta$ ', ' $\gamma$ ' in eq. (3) are used consistent with the solution of Collins and Yu (1996), which can contribute to the simplification of the analytical solutions. For the problem with an isotropic in-situ stress state, the possible error introduced by this simplification has been shown to be negligible by a rigorous numerical (finite element) simulation (Sheng et al. 2000), which has also been reported by Chen and Abousleiman (2012).

141 Considering plastic soil behaviour, the strains are decomposed additively into elastic and plastic 142 components. The superscripts 'e' and 'p' are used to distinguish the elastic and plastic 143 components of the total strains. According to Collins and Stimpson (1994), the deformation in 144 the elastic region is in fact isochoric with no volumetric change, although the material is 145 compressible. Thus, the small strain analysis is used for soil in the elastic region, as expressed:

146 
$$\varepsilon_r = -\frac{d u}{d r}$$

$$\varepsilon_\theta = -\frac{u}{r}$$
(4)

147 where u is the radial displacement. Conversely, to accommodate the effect of large deformation 148 in the cavity expansion process, the large strain analysis is adopted for the plastic regions by 149 assuming logarithmic strains (which are also termed true strains or Hencky strains):

$$\varepsilon_{r} = -\ln\left(\frac{d\,r}{d\,r_{0}}\right)$$

$$\varepsilon_{\theta} = -\ln\left(\frac{r}{r_{0}}\right)$$
(5)

151

150

### 152 UNIFIED STATE PARAMETER MODEL

The unified state parameter model (CASM, developed by Yu 1998) is briefly described in this
section, which was also provided in Mo and Yu (2017). The critical state line is fully defined
as:

156 
$$q = M p'$$

$$\nu = \Gamma - \lambda \ln p'$$
(6)

where *q* and *p'* are the deviatoric and mean effective stresses; *M* is the slope of the critical state line in p' - q space;  $\nu = 1 + e$  is the specific volume, and *e* is the void ratio;  $\lambda$ ,  $\kappa$  and  $\Gamma$  are the critical state constants.

160 The state parameter  $\xi$  is defined by Wroth and Bassett (1965) and Been and Jefferies (1985) as 161 the vertical distance between the current state and the critical state line in  $\ln p' - \nu$  space (see 162 Fig. 2a):

163 
$$\xi = \nu + \lambda \ln p' - \Gamma \tag{7}$$

With benefits of the concept of state parameter, Yu (1998) proposed a unified state parameter
model for clay and sand, which is referred to as CASM. The state boundary surface of the
CASM is described as:

167 
$$\left(\frac{\eta}{M}\right)^n = 1 - \frac{\xi}{\xi_R} \tag{8}$$

168 where  $\eta = q/p'$  is known as the stress ratio; *n* is the stress-state coefficient, which is a new 169 material constant and typically ranges between  $1.0 \sim 5.0$ ;  $\xi_R = (\lambda - \kappa) \ln r^*$ , is the reference 170 state parameter; and  $r^*$  is the spacing ratio, defined as  $p'_{\nu}/p'_{x}$  (Fig. 2a). Equation (8) also represents the stress-state relation and the yield function. In terms of the preconsolidation pressure  $p'_{\nu}$ , the yield surface can be rewritten as follows:

173 
$$\left(\frac{\eta}{M}\right)^n = -\frac{\ln(p'/p_y')}{\ln r^*} \tag{9}$$

The variation of state boundary surfaces (eq. (9)) with the stress-state coefficient are shown in
Fig. 2b, with normalisation of the preconsolidation pressure. Rowe's stress-dilatancy relation
(Rowe 1962), as expressed by:

177 
$$\frac{D \,\delta^p}{D \,\gamma^p} = \frac{9 \,(M - \eta)}{9 + 3 \,M - 2 \,M \,\eta} \times \frac{m}{m + 1} \tag{10}$$

is adopted to define the plastic potential, which has been widely accepted with greatest success in describing the deformation of sands and other granular media. The symbol 'D' denotes the Lagrangian derivative for a given material particle. The hardening law is then adopted based on a typical isotropic volumetric plastic strain hardening, as shown to be:

182 
$$D p'_{y} = \frac{v p'_{y}}{\lambda - \kappa} D \delta^{p}$$
(11)

It should be noted that the adopted soil model CASM after Yu (1998) could be taken as a basis 183 184 for further extensions; e.g. to include shear hardening, to include viscoplasticity, for unsaturated soils, for bounded geomaterials, etc. (see Yu, 2006). In terms of a general three-dimensional 185 186 stress state, M value varying with Lode's angle (proposed by Sheng et al., 2000) could also be 187 included in the yield function, capturing more realistic soil behaviour under various loading 188 paths. This paper, however, focuses on the derivation of drained cavity expansion with the 189 original proposed soil model CASM, largely owing to the simple stress paths of spherical and 190 cylindrical cavity expansion.

191

### 192 ANALYTICAL SOLUTION

193 The drained analytical solution is provided in this section, for a cavity expanded from  $a_0$  to a. 194 After a certain expansion, the soil medium around the cavity becomes plastic, and the plastic 195 region develops from the cavity wall. The symbol 'c' is the radius of the elastic-plastic 196 boundary; thus, for r > c, soil is in the elastic region, and the plastic region is for soil at a <197 r < c (see Fig. 1).

198 Solution for soil in the elastic region

To describe the stress-strain relationship in the elastic region, the elastic strain rates areexpressed as follows:

201  
$$D \delta^{e} = \frac{1}{\kappa} D p'$$
$$D \gamma^{e} = \frac{1}{2G} D q$$
(12)

where *K* is the elastic bulk modulus, which is equal to  $\nu p'/\kappa$ ; *G* is the elastic shear modulus for an isotropic linear elastic material as defined by Collins and Stimpson (1994), which is determined as:

205 
$$G = \frac{(1+m)(1-2\mu)\nu p'}{2[1+(m-1)\mu]\kappa}$$
(13)

Based on the assumption of small strains, the distributions of effective stresses in the elasticregion can be expressed as follows, according to the solution of Yu and Houlsby (1991):

$$\sigma_{r}' = p_{0}' + B_{1} \times \frac{1}{r^{1+m}}$$

$$\sigma_{\theta}' = p_{0}' - B_{1} \times \frac{1}{m r^{1+m}}$$
(14)

where  $B_1$  is a constant of integration. And the distributions of strains in the elastic region can be solved as:

211 
$$\delta = 0$$
  

$$\gamma = B_2 \times B_1 \times \frac{1+m}{\nu_0 p'_0 m r^{1+m}}$$
(15)

where  $B_2 = [1 + (m - 1) \mu] \kappa / [(1 + m) (1 - 2 \mu)]$ . For the elastic stage (i.e. there is no plastic region),  $B_1$  can be derived based on the boundary condition:  $\varepsilon_{\theta}|_{r=a} = -(a - a_0)/a$ , which results in  $B_1 = v_0 p'_0 m a^m (a - a_0)/B_2$ . However, for the plastic stage, the elasticplastic boundary is located at r = c, and the initial yielding deviatoric stress can be found from the initial yield surface:  $q_c = (\ln R_0 / \ln r^*)^{1/n} M p'_0$ . The boundary condition at r = c gives that  $B_1 = q_c m c^{1+m}/(1 + m)$  for the plastic stage, and the size of the plastic region *c* needs to be determined based on the solution for the plastic region.

219

208

#### 220 Solution for soil in the plastic region

Note that for soil in the plastic region (a < r < c), the elastic moduli (*K* and *G*) are not constants but functions of the mean effective stress p'. The volumetric strain is related to the specific volume:  $\delta = -\ln(\nu/\nu_0)$ . In order to convert the Eulerian formulation (e.g. eq. (1)) to the Lagrangian description, a suitable auxiliary independent variable,  $\chi = u/r = (r - r_0)/r$ , is introduced according to Chen and Abousleiman (2013). For the exact solution in the plastic region, numerical integration is required from the elastic-plastic boundary (r = c), where the initial yielding conditions are known with  $p' = p'_0$ ,  $q = q_c$ ,  $\nu = \nu_0$ , and  $\chi = (c - c_0)/c =$  228  $B_2 q_c / [(1 + m) v_0 p'_0]$ . For a given derivative  $D \chi$ , three formulations need to be established 229 to relate  $D \chi$  with D p', D q, and D v, which will be derived from the equilibrium equation, the

volumetric strain rate, and the deviatoric strain rate, respectively.

Together with the assumption of large strains (eq. (5)), the expression of strains can be converted into the forms of  $\chi$ , as follows:

233

$$\varepsilon_{\theta} = -\ln\left(\frac{r}{r_{0}}\right) = \ln(1-\chi)$$

$$\varepsilon_{r} = \delta - m \varepsilon_{\theta} = -\ln\left(\frac{\nu}{\nu_{0}}\right) - m \ln(1-\chi) = -\ln\left[\frac{\nu}{\nu_{0}}(1-\chi)^{m}\right]$$

$$\gamma = -\ln\left[\frac{\nu}{\nu_{0}}(1-\chi)^{m+1}\right]$$
(16)

#### • Equilibrium equation

By using the auxiliary independent variable, the equilibrium equation (eq. (1)) can thus berewritten as:

237 
$$-q = \frac{r}{m} \frac{D\left(p' + \frac{m}{m+1}q\right)}{D\chi} \frac{d\chi}{dr}$$
(17)

238 and

239 
$$\frac{r d \chi}{d r} = -\frac{u}{r} + \frac{d u}{d r} = -\chi + \frac{d u}{d r}$$
(18)

where d u/d r can be obtained from the expression of  $\varepsilon_r = ln(1 - d u/d r)$  together with eq. (16), i.e.  $d u/d r = 1 - v_0/[v (1 - \chi)^m]$ . Therefore, the formulation based on the equilibrium equation is derived as:

243 
$$-q = \frac{D p' + \frac{m}{m+1} D q}{m D \chi} \left[ 1 - \chi - \frac{\nu_0}{\nu (1-\chi)^m} \right]$$
(19)

#### • Volumetric strain rate

245 The volumetric strain rate in the plastic region indicates the rate of specific volume (i.e.  $D \delta = -D v / v$ ), which is also a combination of elastic and plastic components:

247 
$$D \delta = -D \nu / \nu = D \delta^{e} + D \delta^{p} = \kappa \times \frac{D p'}{\nu p'} + \frac{\lambda - \kappa}{\nu} \frac{D p'_{y}}{p'_{y}}$$
(20)

The integration together with the yield criterion (eq. (9)) is equivalent to the expression of thestate parameter (eq. (7)), which gives:

250 
$$\nu = \nu_0 - \lambda \ln \frac{p'}{p'_0} + (\lambda - \kappa) \left[ \ln R_0 - \left(\frac{\eta}{M}\right)^n \ln r^* \right] = C_1 + C_2 \ln p' + C_3 \eta^n$$
(21)

251 where

$$C_{1} = v_{0} + \lambda \ln p'_{0} + (\lambda - \kappa) \ln R_{0}$$

$$C_{2} = -\lambda$$

$$C_{3} = -(\lambda - \kappa) \ln r^{*} / M^{n}$$
(22)

253 The derivative form can then be rewritten as:

254 
$$D v = C_2 \frac{1}{p'} D p' + C_3 n \eta^{n-1} \left( \frac{1}{p'} D q - \frac{q}{{p'}^2} D p' \right)$$
(23)

**255** • Deviatoric strain rate

256 Similarly, the deviatoric strain rate is thus further expressed as:

257 
$$D \gamma = -\frac{D \nu}{\nu} + \frac{m+1}{1-\chi} D \chi = D \gamma^{e} + D \gamma^{p} = B_{2} \frac{D q}{\nu p'} + \frac{\lambda - \kappa}{\nu} \frac{D p'_{y}}{p'_{y}} \frac{9+3 M-2 M \eta}{9 (M-\eta)} \frac{m+1}{m}$$
(24)

Therefore, the three formulations (eqs. (19), (23), and (24)) provide the increments of D p', D q, and D v for a given  $D \chi$  from  $\chi|_{r=c}$  to  $\chi|_{r=a} = (a - a_0)/a$ . Thus, the distributions of v,  $\chi$ , stresses and strains in the plastic region are obtained from the numerical integration. The equivalent location of a material particle around the cavity r corresponding to the auxiliary variable  $\chi$  is revived by integration from a to r:

263 
$$\int_{a}^{r} \frac{dr}{r} = \ln \frac{r}{a} = \int_{\chi|_{r=a}}^{\chi} \frac{d\chi}{1 - \chi - \nu_0 / [\nu (1 - \chi)^m]}$$
(25)

The elastic/plastic boundary *c* is also obtained from eq. (25) by integration from *a* to *c*, which is used to determine  $B_1$  and the distributions in the elastic region (eqs. (14) and (15)).

266

#### 267 RESULTS AND DISCUSSION

#### 268 Validation of the analytical solution

269 After examining the state boundary surface and the stress-state relation, the Modified Cam-clay model could be accurately recovered by choosing  $r^* = 2.0$  and a suitable value of  $n \approx 1.5$  – 270 271 2.0, as noted by Yu (1998). The validation of the proposed solution is performed by the comparisons of the cylindrical cavity expansion between the recovered Modified Cam-clay 272 273 analysis and the results of exact analytical solution for the Modified Cam-clay model, which 274 were reported by Chen and Abousleiman (2013) in conjunction with their drained analysis. The test with an isotropic in-situ stress condition was adopted for  $R_0 = 3$ . The parameters were 275 276 selected to be equivalent to those in Chen and Abousleiman (2013), as summarised in Table 1. 277 The stress paths, the distributions of stresses and specific volume are presented in Fig. 3, with 278 comparisons of data from Chen and Abousleiman (2013), which was also verified by the finite 279 element simulation. Note that all stress paths presented in this paper are provided for the soil element at the cavity wall. As the solution is quasi-static and time-independent, all soil elements 280 281 follow the same stress path, but at any stage of the cavity expansion those elements closer to 282 the cavity boundary are further along that path. The present analytical solution is thus validated 283 by the close agreement between the calculated behaviour of the cavity expansion and the 284 verified analytical results, although the Modified Cam-clay model is assumed by matching the 285 state boundary surface and the stress-state relation using the CASM and the differences on the 286 flow rules.

287

### 288 Drained cavity expansion in clay

289 This section describes the results of drained cavity expansion in clay using the CASM, for both 290 spherical and cylindrical scenarios. Unless stated otherwise, all results are presented by 291 choosing the material constants similar to those of London clay, as suggested by Yu (1998). 292 The soil model parameters and the initial conditions for London clay are listed in Table 2. Note 293 that the frictional constant M is determined by the critical state friction angle, using M =294  $2(m+1) \sin \phi_{cs}/[(m+1)-(m-1) \sin \phi_{cs}]; \phi_{cs}$  is also assumed based on the triaxial critical state friction:  $\phi_{cs} = \phi_{tx}$  for spherical scenario and  $\phi_{cs} = 1.125 \phi_{tx}$  for cylindrical 295 scenario, as suggested by Wroth (1984). 296

Fig. 4 shows the stress paths in normalised p' - q space for  $a/a_0 = 1$  to 10 with the variation 297 298 of overconsolidation ratio  $R_0$ , keeping the initial specific volume constant as 2.0. The critical state lines and initial yield surfaces for the tests with different values of  $R_0$  overlap in 299 normalised p' - q space, and all stress paths start from q = 0 and gradually approach the 300 critical state line. The critical state is reached only when the conditions are satisfied: q/p' =301 Dq/Dp' = M. It can be seen that the normalised stresses (i.e.  $p'/p'_{y0}$ ,  $q/(M \cdot p'_{y0})$ ) increase 302 with the overconsolidation ratio, and slightly higher normalised stresses are found for the 303 304 spherical tests comparing to the cylindrical tests.

The cavity expansion curves for  $a/a_0 = 1$  to 10 are presented in Fig. 5 for both spherical and cylindrical scenarios, respectively; while the variations of the elastic-plastic radius *c* with the overconsolidation ratio  $R_0$  are shown in Fig. 6. It is clear that the normalised cavity pressure  $(\sigma'_r/p'_0)$  increases with the overconsolidation ratio, whereas the elastic-plastic radius appears to be smaller for the test with a higher overconsolidation ratio. The limiting cavity pressure and the constant ratio of c/a are obtained after expansion of approximately 4 times of the initial cavity size, while the cylindrical tests seem to require larger expansion before reaching the
limiting values. In addition, comparing to the spherical scenario, the cylindrical tests have lower
normalised cavity pressure but larger elastic-plastic radius.

314 With benefits of the CASM which can be recovered to the Original Cam-clay (n = 1 and  $r^* =$ 315 2.7183), the effects of model constants n and  $r^*$  are investigated by comparing the modelled London clay and the Original Cam-clay. The results of stress paths and cavity expansion curves 316 317 for both  $R_0 = 1$  and 16 are shown in Figs. 7-8, respectively. The difference on the yield surfaces results in the loci of stresses and cavity expansion curves for both London clay and the Original 318 319 Cam-clay. Higher normalised stresses and limiting cavity pressure are found for London clay with  $R_0 = 1$ , whereas the tests of the Original Cam-clay show higher values of normalised 320 stresses and limiting cavity pressure for heavily overconsolidated clay. It is clear that the 321 322 analytical solution with the CASM can be used for materials with different softening/hardening 323 responses, by modifying the values of stress-state coefficient n and spacing ratio  $r^*$ .

324

#### 325 Drained cavity expansion in sand

326 Similarly, the results of drained cavity expansion in sand using the CASM are described in this

section, which are presented by choosing the material constants similar to those of Ticino sand, as suggested by Yu (1998). The soil model parameters for Ticino sand and the initial conditions under  $p'_0 = 200 \ kPa$  are listed in Table 3.

330 To investigate the effect of initial state parameter,  $\xi_0$  from -0.075 to 0.075 is examined under a 331 constant initial mean stress of 200 kPa. Note that  $\xi_0 = 0.075$  indicates the initial condition at the normal compression line, since the reference state parameter  $\xi_R = 0.075$ . The results of the 332 cavity expansion curves and stress paths in  $\ln p' - \nu$  space are presented in Figs. 9-10, 333 respectively. It is shown that the increase of initial state parameter reduces the limiting cavity 334 335 pressure and increases the limiting specific volume on the critical state line. Comparing to the 336 spherical tests, the value of limiting cavity pressure for the cylindrical scenario is about half of 337 that of the spherical scenario, which also results in a higher specific volume in Fig. 10.

The effect of initial mean stress is also investigated by varying  $p'_0$  from 200 kPa to 800 kPa for  $\xi_0$  of both -0.075 and 0.075. The corresponding soil parameters and the initial conditions are provided in Table 4, and the stress paths in  $\ln p' - \nu$  space are illustrated in Fig. 11 for both spherical and cylindrical scenarios, respectively. Clearly, apart from the initial state parameter, the initial stress condition has a large influence on the stress-strain relationship for soil around the cavity. 344 Furthermore, the effects of the model constants n and  $r^*$  are illustrated in Figs. 12-13, for the results of cavity expansion curves and stress paths in  $\ln p' - \nu$  space, respectively. By varying 345 the stress-state coefficient n between 2 and 4, and the spacing ratio  $r^*$  between 108.6 and 1000, 346 different softening responses of sand can be satisfactorily modelled, as suggested by Yu (1998). 347 Thus the responses of cavity expansion in Fig. 12 show that the increase of either n or  $r^*$  can 348 reduce the limiting cavity pressure for  $\xi_0 = -0.075$ , while the limiting cavity pressure 349 350 increases with n and  $r^*$  for  $\xi_0 = 0.075$ . The stress paths in Fig. 13 present different loci of 351  $\ln p' - \nu$  relation, while the difference of loci for  $\xi_0 = 0.075$  is significantly larger than that of  $\xi_0 = -0.075$ . Correspondingly, the limiting state of specific volume decreases with n and  $r^*$ 352 for  $\xi_0 = 0.075$ , and the reverse trends are found for  $\xi_0 = -0.075$ . 353

354

#### 355 Potential geotechnical applications

Note that the proposed solution provides a general approach for drained cavity expansion/contraction problems using the critical state soil models, with the concept of state parameter and two additional soil parameters. The current solution with an arbitrary cavity expansion has major potential applications, including cone penetration tests, pressuremeter tests, pile foundations, tunnelling, and wellbore instability. Moreover, the solution serves as a benchmark for validating numerical simulations of boundary value problems.

362 A simple example for application to the interpretation of CPT data has been provided here using the developed analytical solution. The cone penetration testing in the calibration chambers is 363 364 widely accepted as a versatile tool for interpretation between penetration resistance and soil properties. The cone tip resistance  $q_c$  is one of the main test measurements, which is usually 365 related to the in situ effective stress and soil density. The approach of spherical cavity expansion 366 367 idealises the cone penetration as an analogy of the expanded cavity under the same conditions by Vesic (1977) and Yu and Mitchell (1998) amongst many others. The cone resistance can 368 369 therefore be predicted based on the calculated cavity pressure (Ladanyi and Johnson, 1974):

370 
$$q_c = \sigma'_r|_{r=a} \times \left(1 + \sqrt{3}\tan\phi\right)$$
(26)

where  $\phi$  is assumed as the critical state friction angle. Thus the relationship between the normalised cone tip resistance Q, defined as  $(q_c - p'_0)/p'_0$ , and the in situ state parameter  $\xi_0$ is provided. The tests with Ticino sand (soil parameters can be found in Table 3) are conducted at an initial effective stress of  $p'_0 = 74$  kPa (after a test of Ghafhazi and Shuttle 2008). The initial state parameter  $\xi_0$  varies from -0.3 to 0.0, indicating an initial specific volume from 1.58 to 1.88. The results are shown in Fig. 14, with a good comparison with data from the calibration 377 chamber tests (Shuttle and Jefferies 1998; Ghafghazi and Shuttle 2008). The calibration 378 chamber tests cover the initial mean stress in the range 50 kPa  $< p'_0 < 500$  kPa, and the initial specific volume between 1.5 and 1.9. The results show that the normalised cone tip resistance 379 380 decreases with the value of initial state parameter, whereas the stress level was found to have 381 little effect on the  $Q - \xi_0$  curve. It should be noted that, for application of the proposed solution, 382 further study is required for the back-analysis of CPT data. To estimate the properties of soils 383 based on the limited measured data, other techniques (e.g. probabilistic identification, Wang et 384 al. 2013; statistical characterization, Niazi et al. 2011) are desired to be incorporated into the 385 solution developed in this paper.

386

#### 387 CONCLUSIONS

388 A new analytical solution for drained expansion of both spherical and cylindrical cavities with a unified state parameter model for clay and sand (CASM) (Yu, 1998) is proposed in this paper. 389 390 CASM is a critical state soil model with two additional material constants, which has the ability 391 to capture the overall behaviour of either clay or sand under both drained and undrained loading 392 conditions. The developed cavity expansion solution with large strain analysis provides the 393 entire stress-strain histories of soils in the elastic and plastic regions. The approach of auxiliary 394 variable is employed for our drained analysis, which unifies the spherical/cylindrical scenarios 395 and clay/sand models.

396 As an illustration, both London clay and Ticino sand are modelled under various initial stress 397 conditions and initial state parameters. The parametric study investigates the effects on stress 398 paths and cavity expansion curves. Higher normalised cavity pressure  $(\sigma'_r/p'_0)$  is obtained for 399 the test with a higher overconsolidation ratio, which also results in a smaller elastic-plastic 400 radius. The increase of initial state parameter reduces the limiting cavity pressure but increases 401 the limiting specific volume on the critical state line. The results also show the ability of this 402 solution for modelling materials with different softening/hardening responses by modifying the 403 values of the stress-state coefficient and the spacing ratio. In addition, this analytical solution 404 provides a general analytical approach for drained cavity expansion problems using other 405 sophisticated critical state soil models. A simple application to the interpretation of CPT data 406 using the proposed solution shows a good comparison with data from the calibration chamber 407 tests. As shown by Yu (2000), it is expected that the new cavity expansion solution developed in this paper can also be applied with success to other relevant geotechnical problems such as 408 409 pressuremeter tests, pile foundations and tunnelling in clay and sand under drained loading 410 condition.

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- 415

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Fig. 14. Prediction of the relationship between normalised cone tip resistance and initialstate parameter.

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Fig. 1. Geometry and kinematics of cavity expansion.







524 
$$p'/p'_{\nu} - q/(M \cdot p'_{\nu})$$
 space.



Fig. 3. Comparisons between the proposed solution and results after solution of Chen and
Abousleiman (2013) for the Modified Cam-clay model.





Fig. 4. Stress paths for  $a/a_0 = 1$  to 10 with variation of overconsolidation ratio of  $R_0$ : (a) spherical scenario; (b) cylindrical scenario.





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Fig. 7. Effect of model constants n and  $r^*$  on stress paths for clay: (a) spherical scenario; (b) cylindrical scenario.





542Fig. 8. Effect of model constants n and  $r^*$  on cavity expansion curves for clay: (a) spherical543scenario; (b) cylindrical scenario.



544

545 Fig. 9. Cavity expansion curves for  $a/a_0 = 1$  to 10 with variation of initial state parameter 546  $\xi_0$ : (a) spherical scenario; (b) cylindrical scenario.





548 Fig. 10. Stress paths in  $\ln p' - \nu$  space for  $a/a_0 = 1$  to 10 with variation of initial state 549 parameter  $\xi_0$ : (a) spherical scenario; (b) cylindrical scenario.





Fig. 11. Stress paths in  $\ln p' - \nu$  space for  $a/a_0 = 1$  to 10 with variation of initial mean stress  $p'_0$ : (a) spherical scenario; (b) cylindrical scenario.







Fig. 13. Effect of model constants n and  $r^*$  on stress paths in  $\ln p' - \nu$  space for sand: (a) spherical scenario; (b) cylindrical scenario.



560 Fig. 14. Prediction of the relationship between normalised cone tip resistance and initial

state parameter.

### 564 TABLES:

$\Gamma = 2.74; \ \lambda = 0.15; \ \kappa = 0.03; \ \mu = 0.278; \ M = 1.2; \ R_0 = 3; \ \nu_0 = 1.97$							
This study Chen and Abousleiman (2013)							
Spacing ratio $r^*$	2.0	-					
Stress-state coefficient n	1.5	-					
Initial stress $p'_0$ (kPa)	122.6	120					
$G_0$ (kPa)	3575	4113					

Table 1. Soil model parameters and initial conditions for validation of the proposed solution.

566

567 Table 2. Soil model parameters and initial conditions for London clay.

$\Gamma = 2.759; \ \lambda = 0.161; \ \kappa = 0.062; \ \mu = 0.3; n = 2.0; r^* = 3.0$							
$\phi_{tx} = 22.75^{\circ}: M = 0.8879 \text{ (spherical)}, M = 0.8640 \text{ (cylindrical)}$							
Overconsolidati	1	2	4	6			
Initial specific v	2.0	2.0	2.0	2.0			
Initial stress $p'_0$	219.15	143.11	93.45	39.84			
Initial state para	0.1088	0.0401	-0.0285	-0.1657			
Ga (kPa)	spherical	3263	2131	1391	593		
υ <sub>0</sub> (π u)	cylindrical	2828	1847	1206	514		

568

569

Table 3. Soil model parameters and initial conditions for Ticino sand under  $p'_0 = 200 \, kPa$ .

$\Gamma = 1.986; \ \lambda = 0.024; \ \kappa = 0.008; \ \mu = 0.3; n = 2.0; r^* = 108.6$	
$\phi_{tx} = 32.0^{\circ}: M = 1.2872 \ (spherical), M = 1.1756 \ (cylindrical)$	
	_

Initial state parameter $\xi_0$		-0.075	-0.025	-0.005	0.005	0.025	0.075
Initial stress $p'_0$ (kPa)		200	200	200	200	200	200
Overconsolidati	11792	518.1	148.4	79.5	22.8	1.0	
Initial specific volume $v_0$		1.7838	1.8338	1.8538	1.8638	1.8838	1.9338
Go (kPa)	spherical	20583	21160	21390	21506	21737	22314
00 (ni u)	cylindrical	17838	18338	18538	18638	18838	19338

Table 4. Soil model parameters and initial conditions for Ticino sand under  $p'_0 =$ 400,600,800 kPa.

* • • •	~							
Initial state pa	trameter $\xi_0$	-0.07	$5 (R_0 = 11)$	.792)	$0.075 \ (R_0 = 1)$			
Initial stress p	0' (kPa)	400	600	800	400	600	800	
Initial specific volume $v_0$		1.7672	1.7575	1.7506	1.9172	1.9075	1.9006	
Go (kPa)	spherical	40782	60836	80796	44243	66028	87719	
u) (iii u)	cylindrical	35344	52724	70023	38344	57224	76023	