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Modeling Challenging EMC Problems Using the Method of Moments

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Abstract—A variety of numerical techniques are used in electromagnetic compatibility (EMC) for the numerical analysis of practical problems. The most important are the finite-difference time-domain technique, the finite-element method, the transmission-line-matrix method and the method of moments. All approaches have their strengths and weaknesses and cannot be applied to all kinds of problems with the same degree of efficiency, measured in memory consumption and computation time necessary. In this paper it is shown that the method of moments (MoM) can be applied to a variety of scenarios, each of which would form a challenging problem to all of the mentioned numerical techniques. It is shown that a customization of the MoM to the specific problems at hand can be accomplished. All examples shown have been computed using the academic code CONCEPT-II developed at Hamburg University of Technology (TUHH).

Index Terms—Method of moments, lightning, radar cross section (RCS), numerical modeling, H-matrices, reverberation chamber, stochastic electromagnetic fields, magnetic resonance imaging (MRI), broadband antenna optimization.

I. INTRODUCTION

NOWADAYS, numerical simulations are common-place in the daily work of an EMC engineer. Among the various techniques available the method of moments (MoM) [1, 2] has outstanding capabilities for the numerical analysis of a broad range of problems, encountered in the area of EMC. Traditionally, MoM is applied in the frequency domain. Especially for electrically large surface and/or wire structures it has advantages compared to other numerical approaches. Essentially, a radiating or scattering structure is replaced by equivalent currents, which are developed into line currents for wires and into surface currents in case of surfaces. Each of these currents is defined over electrically small neighboring segments (basis functions) and has an impact on all other surface and/or wire currents. Applying the underlying boundary condition leads to a set of linear equations expressed by a fully populated quadratic matrix with complex coefficients. Solving the corresponding equation system provides an approximation of the physical current distribution which in turn provides all quantities that might be of interest in electromagnetic engineering. MoM is one of the most widely applied numerical techniques in computational electromagnetics and has a long history since the original paper of Harrington was published. It is well known that the technique is quite sensitive to an appropriate discretization (surface patch grid), regarding both electrical size, shape, and geometrical complexity of the object under consideration. Hence, the user should have a good knowledge of the

underlying principles to quickly obtain accurate results.

To demonstrate the capabilities of the numerical technique under discussion and to get more familiar with it, this contribution gives details about a series of selected numerical problems, which have been presented in a workshop at the EMC Europe conference in Dresden, August 2015. Participants of CONCEPT-II [3] were asked to report on their experiences and surprises they made while treating specific problems.

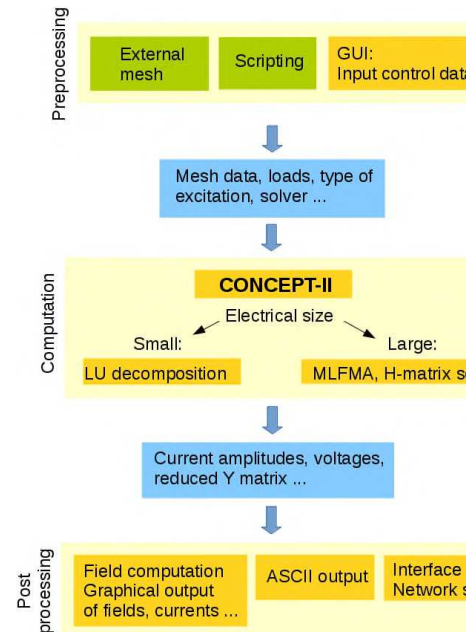


Fig. 1 Typical workflow using CONCEPT-II. Both direct and iterative solvers are available.

this code. All examples presented showed challenges not so well known in the EMC community. As a result, a broad range of interesting applications, including topics like dimensioning of magnetic resonance imaging, testing of objects in reverberation chambers, radar cross section investigations or the analysis of direct lightning strikes could be demonstrated. The intention of this paper is to summarize what can be achieved with the MoM and to provide appropriate extensions thereof.

For several decades, CONCEPT-II has been under continuous development at the Institute of Electromagnetic

Theory at the Hamburg University of Technology (TUHH), Germany. The package is based on various types of integral equations in the frequency-domain. It is designed for the computation of currents, electromagnetic fields and secondary quantities in wire and surface structures consisting of both metallic sheets and lossy dielectric materials. Various types of excitations are available. Lumped loads and generators can be placed at arbitrary locations in the grid. The package is limited to linear materials which have to be homogeneous and isotropic in all regions. The application of the code is free of charge for academic institutions that want to use it in education or public research. During recent years, its core and graphical user interface have been re-implemented with a modern code structure which is necessary for incorporating future research work. Fig. 1 illustrates the possibilities that are currently implemented.

The examples demonstrated in the following sections are difficult to solve with a basic MoM implementation, i.e., special features and enhancements are needed to obtain the results. The specific challenges being addressed by the examples include:

- stable solutions for highly resonant cavities,
- extraction and postprocessing of multipoint network parameters,
- automated and fast computation of multiple excitations,
- optimization of radiating structures,
- efficient modeling of reverberation chamber environments, and
- accurate representation of field excitations given by lightning strokes.

All challenges could be met with the CONCEPT-II code.

II. EMC ANALYSIS OF HIGHLY RESONANT PC CHASSIS

In the frequency domain, a boundary element method such as the MoM and the finite-element method (FEM) are good choices in general. It is common knowledge that MoM is particularly well suited for the analysis of problems with radiation into free space. On the other hand, for 3-D closed structures, FEM is considered to be the method of choice in many cases. In the following example it is demonstrated, that MoM is capable of solving a highly resonant cavity with internal excitation yielding excellent results.

During the last decades, the compute power of personal computers (PCs) has constantly increased. With this, the clock frequencies of the PC components and especially the data buses have reached the GHz range. As a consequence, the components constitute good radiators. These effects need to be accurately considered in system level simulations when solutions to issues involving electromagnetic interference need to be solved.

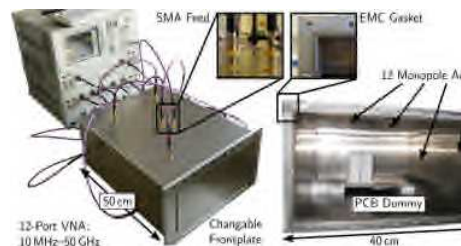


Fig. 2 Analyzed PC chassis: a rectangular chassis with the given dimensions is excited by twelve monopole probes on the walls. The interaction is measured by a network analyzer and calculated by MoM. The chassis contains a dummy fixture resembling a typical PC mainboard with a dummy extension cards. The figure has been adapted from [9], further details and exact locations are given there.

Time-domain methods have convergence problems for highly resonant, low loss cavities. Nevertheless at low frequencies, the dimensions of typical desktop workstations are in the range of several wavelengths, which leads to a large number of unknowns. This requires the utilization of iterative solvers to reduce the solution time and the required memory. Unfortunately, iterative solvers such as the multilevel multipole method (MLFMM) [5] converge slowly due to the high quality factor of the cavity. As a remedy, a fast solver based on H-Matrices has been implemented in the CONCEPT-II package. The capability to additionally include lossy dielectric materials in the simulation makes this applicable to a wide range of problems as encountered in PCs [10].

The main principle behind the H-matrix algorithm is the decomposition of far-interaction blocks in the system matrix into low-rank approximations [7]. The easiest way to do this would be to perform a singular value decomposition (SVD) of the far-interaction block. However, the SVD is costly and requires prior calculation of the complete far-interaction block. This limitation can be overcome by the adaptive cross approximation (ACA) [4], which directly calculates a low-rank approximation of the far-interaction block. Unfortunately, it is not known a priori by how far a block can be reduced. In the worst case of incompressible blocks, one would obtain two matrices of the same dimension as the original block, with twice the memory requirement and an increased computational effort. Therefore, it is essential to accurately predict the compressibility of matrix blocks and this is done in CONCEPT-II. When a minimum block size is reached, corresponding blocks are kept as full matrices. The resulting H-matrix consists of low-rank and full matrix blocks and can either be inverted for a direct solution of the equation system or used for an iterative solution [7].

This technique is applied to the setup of a PC chassis depicted in Fig. 2, originally introduced in [9]. The chassis resembles a desktop workstation and is excited by twelve monopole probes on the cavity walls. Inside the cavity a mother-board/daughter-card setup is placed on PTFE blocks and an absorbing sheet is placed on the cavity

discretized as a flat homogeneous body. The front plate contains two horizontal and two vertical slots located at 1 cm away from the walls and extending almost along the complete height and width. In [9], the correlation between simulation and measurements has been reported to be excellent, see Fig. 3.

The chassis according to Fig. 2 is modeled with a triangular mesh using the open source tool *gms* [6] the output of which can conveniently be imported into CONCEPT-II. The internal short monopole antennas are modeled as thin strips with a width of twice the monopole diameter [8]. At a maximum frequency of 3 GHz, this results in a setup of approximately 30,000 unknowns for the empty chassis and 120,000

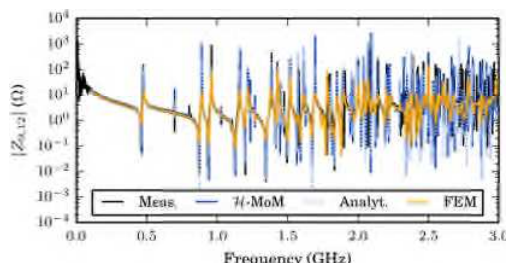


Fig. 3 Simulated and measured coupling impedance between the two closest probes on the top wall of the cavity (framed in Fig. 2). The results are in very good agreement with the analytical solution. Plot taken from [9].

unknowns for the populated one. Here, the solution times have been reduced by a factor of ten to twenty over the considered frequency range compared with the traditional LU-decomposition technique. The simulation process required a few minutes per frequency point on a desktop workstation with 32 GB of RAM and an Intel CPU (i7-4930K, one core). Due to the reduced asymptotic complexity of the H-matrix algorithm in both simulation time and memory footprint, this benefit is more pronounced for larger simulation setups. The complete PC system with internal components and absorbing sheets present could be reduced to a relative matrix size of 3% compared to the full matrix. This enabled an analysis on a desktop machine. In this case, the simulation could be performed in less than an hour per frequency point.

III. MULTIPORT SOLUTION METHODS FOR MAGNETIC RESONANCE IMAGING COILS

The following structure is characterized by a resonant array antenna structure creating a homogeneous circularly polarized field inside a lossy dielectric object. In particular, the simulation utilizes many individual ports for precise and realistic adjustment of lumped components, which is achieved by combining MoM with network theory. Subsequently, an optimized excitation pattern is derived by implementing an Eigenmode decomposition.

In Magnetic Resonance Imaging (MRI), parallel reception with array coils is established since about 15 years. In the

recent past also parallel transmission became a clear trend. Clinical multi-transmit systems are commercially available since 2009. These techniques require highly sophisticated multi-channel magnetic near field antennas, also referred to as MRI transmit arrays. To keep the power deposition on the patient as small as possible, electric fields are typically eliminated to a large extent by means of multiple capacitors (typically up to a few hundred), storing the main part of the electric energy. As these capacitances need to be adjusted accurately to form resonant structures at the desired Larmor frequency (e.g. 128 MHz for a static magnetic field of 5 T), the tuning of these structures is challenging in both, sim-

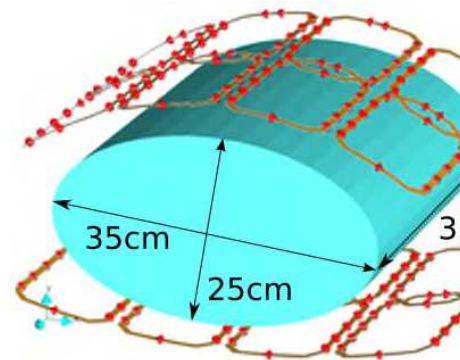


Fig. 4 MoM model (CONCEPT-II) of a 16-element transmit/receive array coil using 208 individual ports. The elements all are 10cm x 10cm in size. An overlap in z-direction compensates for mutual coupling; a decoupling is implemented by transformers (post-processing, not shown). The phantom inside mimics a well reproducible and realistic load, similar to a patient (conductivity 0.52S/m, relative permittivity 53.4).

and construction. Compared with alternative methods, MoM typically leads to small, but ill-conditioned equation systems for these near field problems, which are solved with a high degree of precision by Gaussian elimination (LU-decomposition). The simulation therefore, hardly depends on the number of degrees of freedom; hence, a multi-port model may be combined with an intelligent circuit simulation including optimization. In a second step, fields can be evaluated to describe the behav-

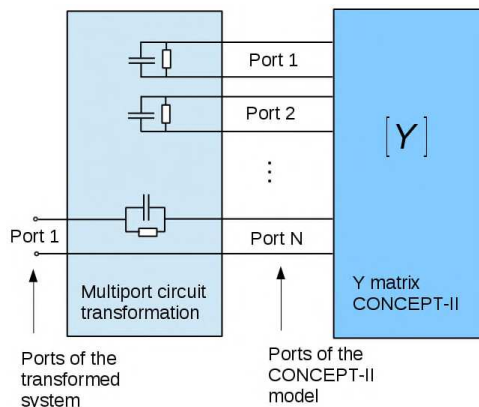


Fig. 5 Transformation of a CONCEPT-II simulated structure (linear network and field sensitivities) by means of a linear general multiport circuit, here defined by lossy resonance capacitors. The ports on the right are the ports as defined in the CONCEPT-II simulation. This yields the corresponding admittance matrix as well as the 3D electromagnetic fields as a linear combination of the port voltages. After transformation these fields as well as the residual admittance matrix are described by the port excitations on the left.

with respect to a limited number of feeding (or receiving) ports.

Fig. 4 shows a 16-element MRI coil, which was simulated for 208 individual port excitations. Most of these ports were equipped with resonance capacitors in post processing. Besides for feeding via one port per loop, some of the ports are connected to transformers, which are used to compensate for mutual coupling of neighboring elements. Both the antennas port responses and the electromagnetic fields can be written by matrix equations, typically based on the admittance or scattering matrix combined with corresponding so-called sensitivity matrices, which describe the linear relations of fields and port excitations. This behavior can be transformed by a linear network [11], which in this case is defined by matching and tuning capacitors as sketched in Fig. 5. In the given example, the 208 ports were reduced to 16 driving ports by a 224-port transformation network containing tuning capacitors, matching networks as well as transformer decoupling networks. The driving amplitudes and phases of the individual channels need to be optimized by means of a complex goal function. On the one hand, it is the aim to generate a homogeneous circularly polarized B_1 -field in the imaging region. On the other hand, the power demands need to be lowered by reducing both, RF amplifier demands and local heating by energy absorption inside the phantom (or patient). For finding a good compromise, methods based on Eigenmode analysis were proposed [12, 13] which first define modes with respect to their power efficiency and then perform an optimization by using just the most efficient modes. Fig. 6 exemplarily shows the resulting transversal field profiles for using the best 1, 2, 4, or 7 Eigenmodes. As can be seen, at

least the first seven modes are needed to generate an acceptable transmit field homogeneity. All field-maps are scaled to 10 μ T average, which demands 716 W net power for the seven modes. This is significantly less than a commercial system-integrated body coil would need to generate a similar field. Phantom experiments showed good agreement between field profiles and power demands [12], thus the described simulation technique is able to predict the performance of MRI arrays.

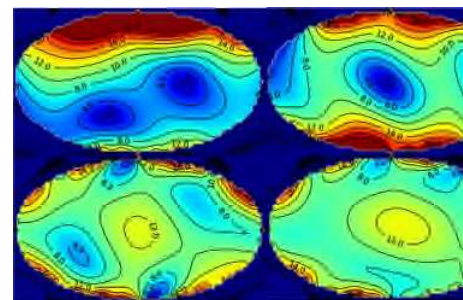


Fig. 6 Simulated B_1 -fields for using the first (top left), first two (top right), first four (bottom left) and seven (bottom right) most efficient mode(s) of a 16-element MRI coil for RF-shimming.

IV. BISTATIC RCS COMPUTATION OF AIRCRAFT

The next example includes much more unknowns related to the previous one as many right-hand sides are involved. Here, a fully populated excitation matrix is required. A solution is obtained by compressing the excitation matrix, which considerably speeds up the overall solution process.

Analyzing the scattering of electromagnetic fields at targets, such as aircraft, is of fundamental importance for the design and evaluation of radar systems. In common practice, the scattering characteristics of an aircraft is expressed in terms of a radar cross section (RCS). Besides material properties and geometry of a target, the RCS depends on frequency, incident angles associated with the illuminating field, θ , and aspect angles. Fig. 7 illustrates a typical scenario for a bistatic radar system: A small civil aircraft is approaching an airport and shall be tracked using the field of a transmitting antenna which is reflected by the aircraft and received at an observation point. As the aircraft flies along the trajectory towards the airport, the RCS is varying and the scattering analysis of such scenario requires the RCS evaluation for a large number of pairs of incident angles (θ', Φ') and observation angles (θ, Φ).

Computing such a problem with the MoM leads to a large equation with a corresponding number of right-hand sides (RHS), each representing one angle of incidence. The simulation of electrically medium-size objects can be accelerated by fast solution algorithms, such as fast iterative methods or fast direct techniques [15] based on ACA. Multiple

problems lead to a high computational effort and require the application of sophisticated techniques for iterative solvers [16], while fast decomposition techniques are well-suited for treating multiple RHS problems [15].

For the computation of the bistatic RCS scenario shown in Fig. 7 the Evektor EV-55 aircraft has been considered at a typical digital audio broadcasting (DAB) frequency of 200 MHz and there was interest in computing the RCS at 1000 equidistant points on the trajectory. The work flow for this example is as follows: A surface model has been extracted from a mechanical CAD model, it has been meshed by the external meshing tool *GiD* [17] and imported in CONCEPT-II. Incident and aspect angles were computed by a MATLAB

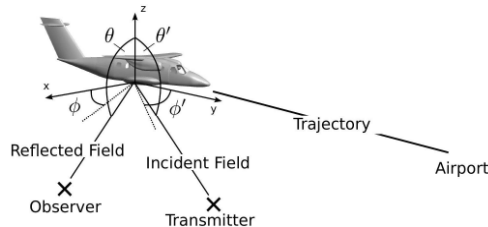


Fig. 7 Considered bistatic RCS example of the Evektor EV55 (courtesy of Evektor) associated with one observer and one transmitter for different positions on the indicated trajectory.

script, then stored in an ASCII file and finally used for the RCS excitation. An ACA treatment of the excitation was chosen with an accuracy of 10^{-5} in order to reduce the computation time of the iterative solver. For this example, two simulations were performed, one using a fast direct solver and one simulation using a fast iterative solver. The predefined solver settings were applied in both cases.

Fig. 8 shows the bistatic RCS of the aircraft for different positions on the trajectory. It is indicated that the RCS values exhibit slow variations if the aircraft is in great distance from transmitter and observer because the incident and aspect angles are slowly varying. In close vicinity to transmitter and observer, these angles are rapidly changing, leading to pronounced RCS fluctuations. Computing this problem with the fast decomposition algorithm took 3.45h. Applying a fast iterative technique with a compression of the excitation required 3.76h, whereas a computation without RHS-compression would lead to a computation time of around 48h!

This example shows that arbitrary incident and aspects can be treated, making an investigation of aircraft trajec

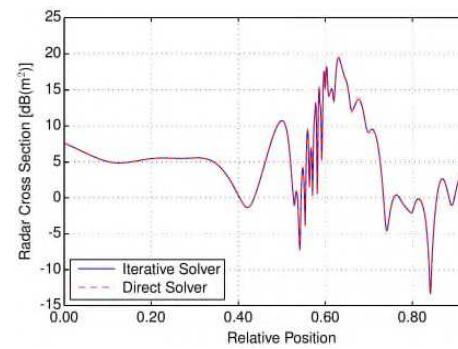


Fig. 8 Bistatic RCS of the EV55 at 200MHz for the example shown in Fig. 7. The RCS is depending on the position of aircraft on the indicated trajectory.

feasible. Computing the bistatic RCS of such scenarios is carried out in an efficient way thanks to the application of the fast direct solver and an efficient treatment of the excitation matrix, respectively.

V. ANALYSIS AND OPTIMIZATION OF A HYBRID BROADBAND REVERBERATION CHAMBER ANTENNA

Optimizing a structure with regard to specific electromagnetic properties is an important goal in EMC and RF engineering. The following example shows that the MoM is a suitable method for achieving excellent antenna designs. The motivation was to produce an antenna for use in a reverberation chamber over a wide frequency range. Specifically, this antenna should be usable in a chamber of $4.7 \text{ m} \times 3.0 \text{ m} \times 2.4 \text{ m}$ size over a frequency range 200 MHz to 20 GHz.

Specifically, the goal was to maximize efficiency while minimizing size. Therefore the aim was to keep the reflection coefficient below 0.316 (-10 dB) over the entire frequency range. As the direction in which the antenna radiates into the RC is not important, the radiation pattern was not specified, but was an interesting output of the simulation. The solution was a dual mode antenna, whose structure is a hybrid of a monopole and an exponential taper like a Vivaldi antenna [18]. This taper is formed from a curved metal strip of height 305 mm, aligned vertically over a ground plane with width 300 mm and length 375 mm. Its mode of operation at low frequency is like a monopole (Fig. 9) with a resonance frequency of approximately 250 MHz. This is determined by the effective minimum length of the monopole, as defined by the curved edge above the feed point. However it behaves more like a simple exponential taper from 400 MHz upwards. The key is that the taper takes over before the half-wave resonance of the monopole at about 500 MHz. At 200 MHz the antenna pattern is close to omnidirectional at 200 MHz. The antenna shows higher gain above 500 MHz: Fig. 10 shows the radiation pattern at 1GHz.

For optimization of the design, the antenna was modelled with CONCEPT-II using the external tool *gms* for the meshing

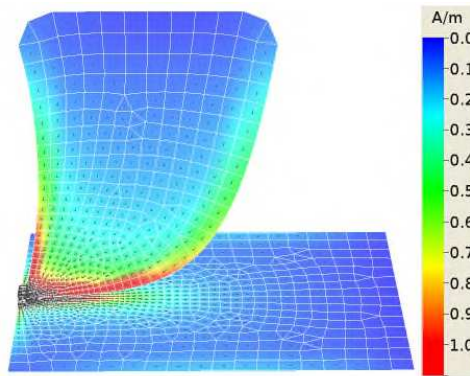


Fig. 9 Surface currents on a blade antenna acting like a monopole at 200MHz as simulated by CONCEPT-II. Dimensions: height 0.305m, width 0.3m, length 0.375m. Feed voltage: 1V at zero phase.

[6]. The target was $S_{11} < -10$ dB from 200-1000 MHz; prior experience showed that if this were achieved, the antenna would also perform well at gigahertz frequencies.

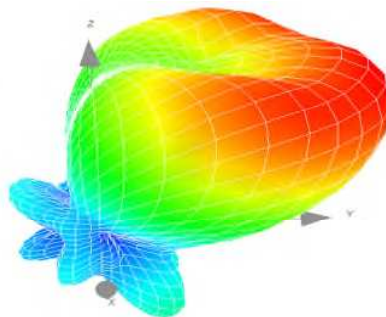


Fig. 10 Blade antenna pattern at 1 GHz as simulated by CONCEPT-II.

Design parameters were varied including the length and shape of the exponential taper, the size of feed block, and also the effect of adding top-loading formed by a horizontal plate at top of the taper. The mesh size was around 10 mm at edges away from the feed, and finer (around 0.5 mm) near the feed [19]. The automated optimization process used a genetic algorithm (GA). Programs were implemented in a portable subset of GNU Octave and MATLAB, using the MATLAB GA toolbox or an in-house compatible Octave GA. The MATLAB/Octave functions also automatically wrote the CONCEPT-II input files. In the first step of the optimization process the genotype had to be decoded to get the design parameters. A *gms* (.geo) file was created with the required

parameters. Next CONCEPT-II input files were created templates and a CONCEPT-II run was started. The processing step provided the reflection coefficient $|S_{11}|$, the cost function has to be evaluated as area between $|S_{11}|$ the upper mask (see (Fig. 11)).

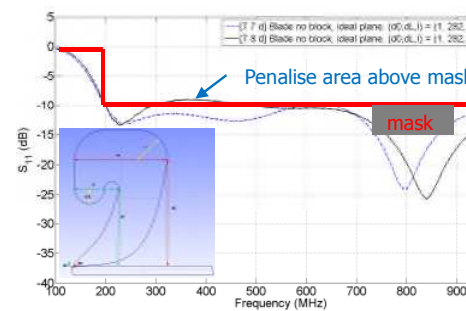


Fig. 11 Cost function for the GA. Inset: parametric computer aided design (CAD) for *gms*.

Validation showed good agreement between the model measurements, both in free space and over a ground plane (Fig. 12).

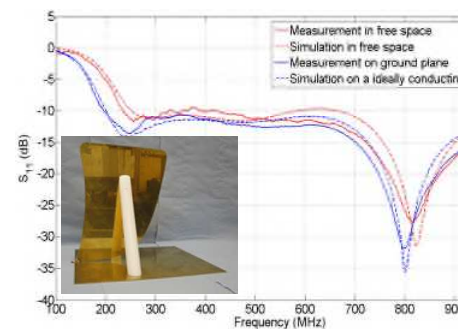


Fig. 12 Validation of model (S_{11} continues to be good up to 20 GHz final fabricated antenna).

A smooth interface between the MoM solver, MATLAB optimiser and the *gms* meshing program has been achieved. The CONCEPT-II simulations revealed a successful trade-off between antenna modes, and showed the initial design was already close to optimal. There were design trade-offs between feed-point location controls S_{11} against power handling capability and top loading controls S_{11} against low frequency performance. The final antenna has acceptable performance from 200 MHz to 25 GHz, maybe higher [18]. Its frequency range was adjusted by dimensional scaling.

VI. COUPLING OF STOCHASTIC ELECTROMAGNETIC FIELDS INTO TRANSMISSION LINE STRUCTURES

The numerical computation of a reverberation chamber including excitation, stirrer and equipment under test (EUT) is a challenging task (Fig. 13). There are several reasons for this.

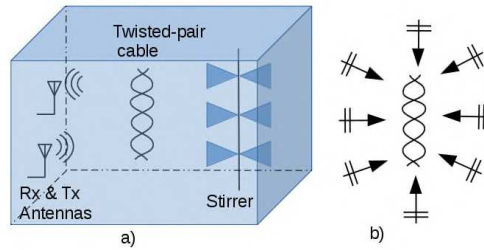


Fig. 13 a) Schematic view of a reverberation chamber with a twisted-pair cable as an EUT. b) One advantage of the presented methodology is that the chamber itself as well as the antennas and the stirrer do not have to be discretized but only the twisted pair cable.

Firstly, the chamber walls have to be included in the computation leading to huge amounts of unknowns making a numerical solution unfeasible especially in the higher frequency range. Secondly, a reverberation chamber is a completely closed cavity with a large Q factor leading to special demands both for surface-based and volume-based solvers.

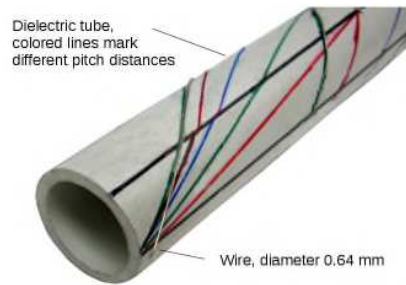


Fig. 14 Open-circuited end of the practical twisted cable that was analyzed in a reverberation chamber.

In the following, a fast and efficient way to simulate the coupling of stochastic electromagnetic fields to general transmission line structures using MoM is explained. Such fields develop in reverberation chambers, where the statistical analysis is done over a range of different stirrer positions, or in overmoded cavities like aircraft fuselages or car bodies, where an ensemble of different spatial positions or nearby frequencies can be analyzed statistically. If such fields couple into transmission lines the terminal voltages or currents also become stochastic values. Analytical calculations (integral over waves) or plain numerical simulations (sum over waves) are very fast, but cannot handle general transmission line

structures. Full wave solvers such as the MoM can handle more arbitrary geometries of the line and the EUT itself, but are very slow for stochastic simulations.

To overcome this problem, CONCEPT-II includes a technique to use stochastic fields as an excitation by simulating the source of this field itself [21]. The technique is based on a plane wave approach. A certain number of waves form one stochastic sample. Although there is a stirrer present in the simulation, but just a random number generator, the stochastic sample is called *stirrer position*, consistent with its practical meaning to the simulation in reverberation chambers. Many of these stirrer positions are generated and simulated. The main advantage is that the transmission line under test is treated conventionally, which leads to a smaller number of unknowns and a smaller computational effort.

The geometry of the reverberation chamber (including the stirrer and the antennas) or the overmoded cavity is explicitly included, so no resonance problems occur and only one system matrix which refers to the structure under test

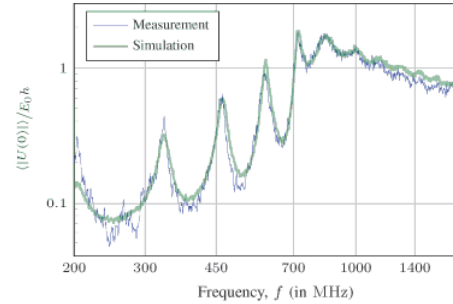


Fig. 15 Average magnitude of the coupled voltage at the beginning of a twisted double wire transmission line as a function of the frequency. The simulation was done by CONCEPT-II.

be used for all different stirrer positions. The approach can also be parallelized very easily, because different independent stirrer positions can be simulated on different CPUs.

The field strength of the stochastic field is normalized to the average squared magnitude of the total electric field strength of $E_0^2 = 1 \text{ V}^2 \text{ m}^{-2}$. The simulation results are written to a text file that includes a line for each frequency and stirrer position. Each line, in turn, includes the results for all different stirrer positions. Such file can be parsed very easily with sophisticated mathematical software tools like MATLAB, Python with NumPy and SciPy for the stochastic analysis and post-processing.

The outlined methodology is presented using a twisted double wire transmission line as an example. Such transmission line structures are used for many applications (e.g. Ethernet, CAN bus) to enhance the immunity against external electromagnetic interference. The corresponding model of the twisted-pair as used in the simulation is shown in Fig. 14. The line is 75 cm long and features 15 twists, which corresponds to a pitch distance of 15 cm. The configuration was measured in a reverberation chamber.

A plastic tube was used to fix the geometry of the line. This tube also determines the line spacing of 33.7 mm between the conductors.

In CONCEPT-II, each conductor is simulated by 50 straight wire segments. To include the dielectric effect of the plastic tube, a dielectric coating of the wires with an equivalent relative permittivity of 4.4 and a thickness of 1.3 mm was used. The measured and simulated number of stirrer positions is 360. In the simulation one hundred incoming waves for each stirrer position were used, which was found to be enough to get reasonable results. The measurement took about 45 min, the simulation time was approximately 3 h on a standard PC.

As an exemplary result, Fig. 15 shows the average magnitude of the coupled voltage at the beginning of the line. For an easier comparison between measurement and simulation, the coupled voltage has been normalized to the chamber constant E_0 and to half of the spacing between the conductors. For small frequencies, the protection due to the twisting is effective and the coupling to the line is very low. The coupling rises with frequency, shows the expectable transmission line resonances and reaches a maximum shortly above 700 MHz, where the wavelength equals the 2.8-fold of the pitch distance, i.e. the length of a single twist. Here, the twisting loses its protective effect. For higher frequencies and smaller wavelengths, the coupling decreases again due to radiation losses. The overall agreement between the experimental and the simulated results is very good. The simulation in CONCEPT-II is able to predict the correct amplitude, the transmission line resonances and the general frequency behavior. The superimposed noise on both curves is mainly due to the statistical uncertainty of the measurement or simulation itself that could be lowered by an even higher number of (independent) stirrer positions, but would require more effort.

VII. TRANSIENT RESPONSE OF THE TOP STRUCTURE OF THE PEISSENBERG TOWER TO LIGHTNING

As has been mentioned earlier, the applied program package is based on integral equations that are formulated in the frequency domain. For a time-domain analysis, an inverse Fourier transform has to be carried out. Such an undertaking is particularly challenging in the case of a lightning stroke. Reasons are very low frequencies and the incorporation of the impact of the spatially distributed lightning current. Here, nonlinear effects are not considered.

The mountain called "Hoher Peissenberg" is an isolated ridge topping the surrounding terrain by about 250 m. The mountain is located in the south of Germany close to the mountains of the Alps, about 60 km far from Munich. On this mountain, the 160 m high Peissenberg Tower is located in an altitude of about 940 m above mean sea level. Fig. 16 (a) shows a picture of the tower.

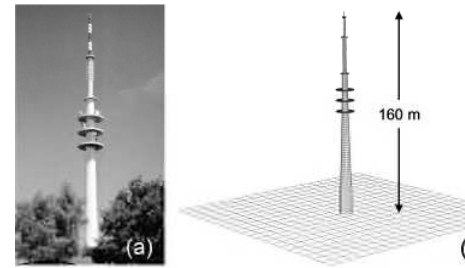


Fig. 16 (a) Picture of the Peissenberg Tower, and (b) model of the Peissenberg Tower without the attached lightning channel.

The subject that matters here is the analysis of the frequency oscillation of the lightning currents which have been measured at the top of the tower [22, 23]. Oscillations appeared only during the initial period of rising currents. Fig. 17 presents the current record of a negative lightning return stroke which occurred on March 5th, 1998. The inset shows that the current oscillation is related to the initial period of the current rise.

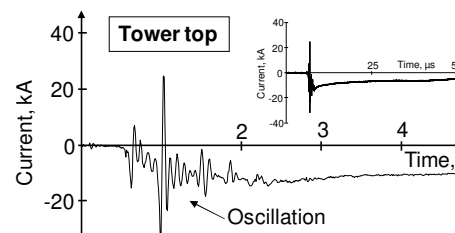


Fig. 17 Current record (tower top) of a negative return stroke which occurred on March, 5th, 1998.

In the computer simulation with CONCEPT-II the Peissenberg Tower according to Fig. 16 (b) is modeled by about 3600 triangular and rectangular ideal conducting plates and about 500 wires with the conductivity of aluminum ($\sigma = 3.8 \times 10^6$ S/m). The time-domain solutions are obtained by an inverse Fourier transformation. Three different frequency ranges are chosen in order to minimize the number of frequencies. Starting with a lowest frequency of 2 kHz the frequency is increased in steps of $\Delta f = 4$ kHz up to 20 kHz. Then the frequency step is increased to $\Delta f = 20$ kHz in the second frequency range up to 4 MHz. In the highest frequency range between 4 MHz and 40 MHz, the frequency step is further increased to $\Delta f = 80$ kHz.

Fig. 18 (a) shows the numerical model for the top section of the Peissenberg Tower. The lightning strike is assumed to occur at one of the small Franklin rods located at the top of the ring construction. The current measuring point was located at the central down conductor. The return stroke process is taken into account by the well-known transmission-line (TL) model [24]. In the frequency domain, this model uses a

source, where the phase velocity is given by the progress velocity of the return stroke. The inverse Fourier transformation gives a time-varying current wave starting from the point of strike and traveling along a straight wire with a constant return stroke velocity v corresponding to the chosen phase velocity (see Fig. 18b). In the simulations, the return velocity was chosen to $v = 100 \text{ m}/\mu\text{s}$. Fig. 19 shows the calculated waveform of the current at the tower top. The channel-base current is normalized to the maximum and is taken into account by a unit-step function, but with a linear rising front. The front time is chosen to 10 ns. The current at the tower top is superimposed by strong pulsation with the frequency of 12 MHz being about the same as in the current records (see Fig. 17).

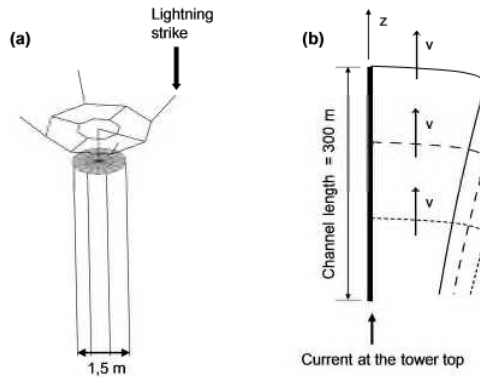


Fig. 18 (a) CONCEPT-II model for the tower, and (b) simulation of the return stroke channel with the TL model.

Oscillations in the recorded waveforms are often an indication for electromagnetic disturbance which affects the measuring system due to either electromagnetic coupling or electromagnetic radiation emitted from an external source. However, with the computer code under discussion, the oscillations could be reproduced and it is therefore very likely that the oscillations are caused by the response of the top structure of the Peissenberg Tower.

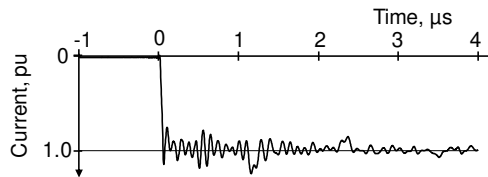


Fig. 19 With CONCEPT-II calculated current at the tower top as response of the unit-step function, but with a linear rising front. The front time is chosen to 10 ns.

VIII. CONCLUSION

In this paper it was shown that MoM can be applied to a wide range of applications. Although MoM is typically associated

with open-bounded, freely radiating problems, the examined the investigated PC chassis example clearly demonstrated the method is also a powerful tool for the analysis of closed-box or almost closed-cavity problems. To overcome the necessity to compute the current distribution inside an electrically large cavity of a reverberation chamber, a surface formulation was chosen. Here, an efficient method for simulating the stochastic field coupling based on a plane wave approach was presented. It can be used for transmission problems and even for more general structures and drastically reduces the computational effort. For the example of a twisted pair cable, a good agreement between theoretical and experimental results has been observed. A further strength of the MoM is that it can easily be combined with network theory. The inverse Fourier transform even lightning phenomena can be investigated with regard to direct and indirect coupling. Ensuring a proper grid is essential for the quality of the numerical solution. Here, it has been demonstrated that external tools can be applied with great success to fast optimization antenna design. Far field computations, e.g., through calculation, are the classical domain of the MoM. The inclusion of advanced matrix techniques provides a significant acceleration of the solution process. Of course, the disadvantages of the MoM should not be omitted. Inherently it does not provide simple means for the investigation of non-linear problems and/or non-uniform material for example. In case of electrically large problems, the available computing power sets a practical limit for the frequencies that can be investigated, in the case of very high frequencies the electric field integral equation becomes numerically unstable without a low-frequency stabilization. For the classical MoM, this low-frequency limit depends on the geometrical size of the structure and is somewhere in the range of 50 Hz to 1 kHz in many cases of practical interest. In any case, the user has to validate the numerically computed results, this takes time and is non-optional – a statement that is not only true for the MoM but to any kind of numerical technique.

In summary it has been shown that our MoM software is ideally suited for a fast investigation of a variety of electromagnetic problems. The CONCEPT-II suite can be used to investigate large scale problems, but it also serves as a sandboxing tool for getting insight into the electromagnetic behavior of complex problems and setups. Due to its open interface users may find quick answers to open questions or even new ways to solve their problems.

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