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The STELLA Apparatus for Particle-Gamma Coincidence Fusion Measurements with Nanosecond Timing

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Abstract

The STELLA (STELlar LAboratory) experimental station for the measurement of deep sub-barrier light heavy-ion fusion cross sections has been installed at the Andromède accelerator at the Institut de Physique Nucléaire, Orsay (*France*). The setup is designed for the direct experimental determination of heavy-ion fusion cross sections as low as tens of picobarn. The detection concept is based on the coincident measurement of emitted gamma rays with the UK FATIMA (FAst TIMing Array) and evaporated charged particles using a silicon detector array. Key developments relevant to reaching the extreme sub-barrier fusion region are a rotating target mechanism to sustain beam intensities above 10μA, an ultra-high vacuum of 10^{-8} mbar to prevent carbon built-up and gamma charged-particle timing in the order of nanoseconds sufficient to separate proton and alpha particles.

Keywords: rotating target, $LaBr_3$ self-calibration, coincidence technique, proton-alpha separation, fusion

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1. Introduction

Heavy-ion fusion reactions involving ^{12}C and ^{16}O nuclei such as the ^{12}C+^{12}C reaction play a key role in the evolution of massive stars and in explosive astrophysical scenarios such as type Ia supernovae and super-bursts in binary systems. Since the 1950s the ^{12}C+^{12}C system was well known to exhibit strongly resonant behaviour [1] which also manifests in the fusion cross-section with prominent resonances, at energies ranging from a few MeV per nucleon, down to the Coulomb barrier and below [2]. Such resonances have been attributed to the formation of long-lived ^{12}C+^{12}C molecular configurations.

The presence of these resonances will inevitably have a strong impact on the carbon burning reaction rates under the different astrophysical scenarios. Direct cross-section measurements are therefore needed down into the Gamow window corresponding to carbon burning in massive stars. These experiments are hugely challenging as the relevant cross-sections are well below the nanobarn level. Reaction rates presently rely on extrapolations of cross-section data from higher energies. These data are largely based either on the detection of evaporated charged particles (α, p) or the characteristic gamma decay of the α and p evaporation residues ^{23}Na and ^{20}Ne [3]. The former technique suffers from the presence of low-level deuterium contamination in the carbon target as the reaction ^{12}C(d,p) has a large cross section and the resulting protons are at similar energies to the far weaker evaporated charged particles from ^{12}C+^{12}C fusion. Gamma-ray detection is challenging at the level of the cross-sections of interest from the point of view of discriminating signal from background. A clear way to achieve a system with strong background suppression is to measure coincidences between evaporated charged particles and their associated gamma rays since this is a unique signature. This technique has been realised by Jiang *et al.* [4, 5] using the Gammasphere germanium detector array and an annular silicon detector array at Argonne National Laboratory. Their initial results are very promising but the full potential of the technique is limited by the available beam currents (of the order of 1 μA) and the potential running period (of the order of one week). To extend this approach to the energies of astrophysics interest, such experiments will need beam currents in the microampere range and extended running periods of many weeks duration. This is the challenge addressed by the STELar LAboratory described in this paper. The key elements of STELLA are:

- rotating targets which can sustain high beam intensities,
- high-efficiency particle and gamma-ray detection systems, and
- employment of a coincident technique which allows the extraction of the relevant fusion signal from the dominant background.

2. Apparatus

The scattering chamber of the STELLA system is presented in Figure 1. The chamber contains several annular DSSSDs (Double-Sided Silicon-Strip-

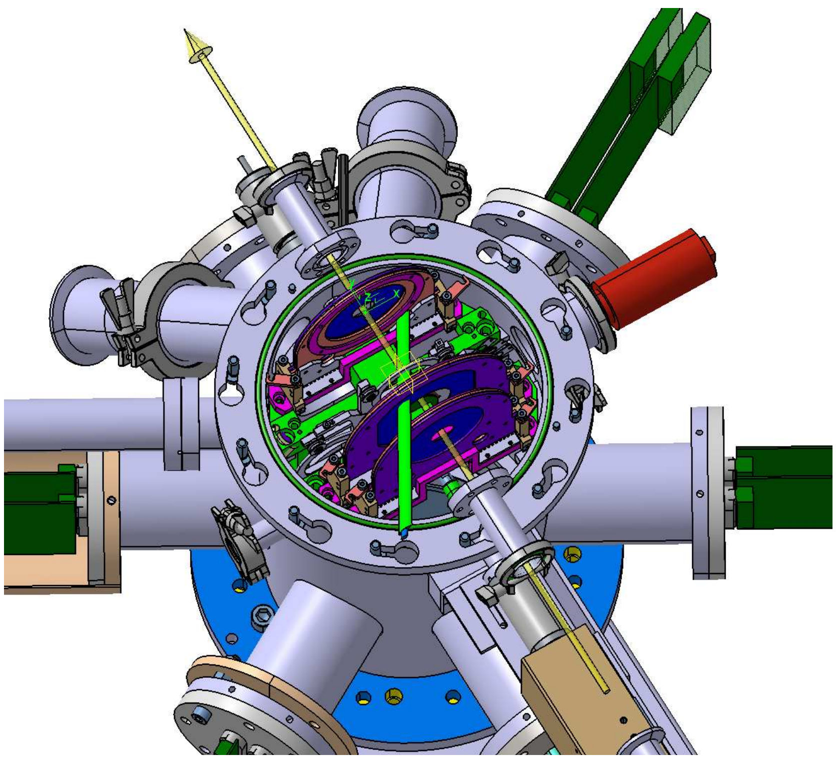

Figure 1: [color online] View into the target chamber that is mounted on top of the cryogenic pump and that is closed by a this Al dome. The annular particle detectors shown in dark blue are aligned along the beam axis around the target at the center of the reaction chamber. The extensions serve as feed-throughs for detector signals and host two surface barrier silicon monitor detectors at 45° with respect to the beam line.

Detector), described in detail in section 2.2, for high-efficiency particle measurements. The detectors are aligned along the beam axis around the target (see section 2.1) at the center of the reaction chamber. All support structures and signal cables are directed towards the bottom of the chamber where a cryogenic ultra-high vacuum pump is located providing a vacuum of 10^{-8} mbar. The gamma-ray detectors comprising an array of lanthanum bromide ($LaBr_3(Ce)$) scintillators are supported from above and surround the 2.5 mm thick aluminum dome-shaped target chamber with a diameter of 20 cm. The gamma-detection array is introduced in section 2.3. All data are time-stamped with sampling times of 1 ns and 8 ns, respectively, for gamma ray and charged-particle detection. The synchronization of coincident gamma-particle events in $^{12}C+^{12}C$ fusion reactions as well as a first background reduction estimate is discussed in section 2.4. Two surface barrier silicon detectors for the measurement of scattered beam particles as well as a Faraday integrator are used for the precise determination of the beam intensity during the measurements. The beam particle monitors are located 23 cm from the target in extensions of the reaction chamber that form an angle of 45° with the beam line.

2.1. Rotating Target

With ^{12}C beam intensities in the order of pμA and a beam spot diameter of 2 mm, a heat input of Watts may be estimated based on energy loss for targets of a few tens $\mu g/cm^2$ thickness. In order to avoid breaking of the targets, it was necessary to develop the rotating target mechanism displayed in Figure 2. It is designed to increase the effective beam spot size to distribute the thermal load. The wheel hosts three rotating target frames and seven slots for fixed target experiments. A magnetic feed-through connects the rotation-driving motor outside the vacuum with the central wheel to spin *via* friction with the target frame bearing in contact. This bearing transmits the rotation to the rotating target frames, each with a diameter of 6.3 cm. In total, three bearings guide each target frame. The axis of the target revolver mechanism used to change the target is slightly off the target rotation axis. In this way, only the uppermost target can spin, because the other bearings are not in contact with the drive shaft.

The layout is optimized for heat dissipation using the MATHCAD15® package where the temperature distribution at the beam spot position is calculated solving the heat-flow equation with a radiative heat loss term. This follows the Stefan-Boltzmann law at high excess temperatures and the net-radiative heat loss over time is obtained with $P_{rad} = \epsilon\sigma S\cdot(T^4 - T_s^4)$, where

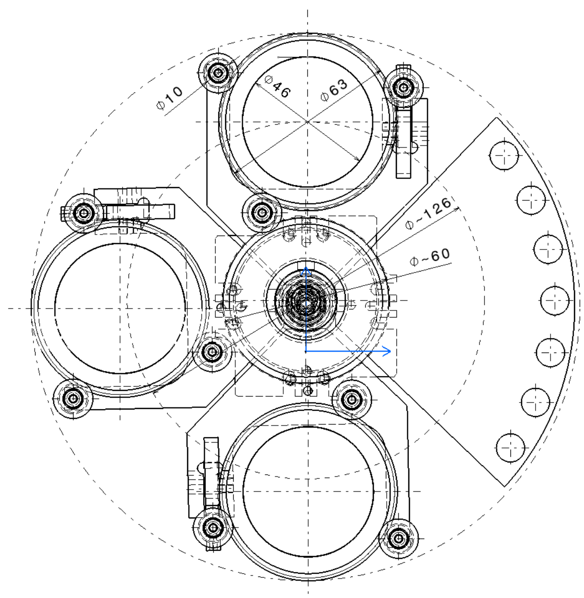

Figure 2: Front view of the target wheel where the uppermost quadrant is exposed to the beam. The wheel hosts a quadrant with seven fixed target slots (on the right) and three rotating targets with a diameter of 6.3 cm. The central wheel transmits the rotation from the external motor to the target frames.

the emissivity $\epsilon = 0.8$, σ is the Stefan-Boltzmann constant, S is the surface area during one turn of the target, T the temperature of the environment is 20°C and T_s is the target foil temperature.

The voxels of the target material are heated when exposed to the beam and they cool *via* radiation when off the beam axis during the rotation. Taking into account these effects, the time-dependent profile of the target temperature may be calculated per turn of the target frame. The resulting temperature converges towards a maximum T_{max} within seconds for the chosen parameters with a saw-tooth like cooling modulation ΔT between two heating pulses. The dynamics are mainly dependent on the beam spot size, the radius of the beam track on the target, and the rotation velocity at a given beam intensity. An example of the multi-parameter study is given in Table 1 for a rotation speed of 1000 rpm. The maximum temperature T_{max} is

d [mm]	T_{max} [C°]	d [mm]	T_{max} [C°]	d [mm]	T_{max} [C°]
$P = 1$ W		$P = 2$ W		$P = 3$ W	
2	550	2	920	2	1240
3	410	3	670	3	910
4	340	4	550	4	730
5	290	5	470	5	620

Table 1: Temperature dependence T_{max} from the beam spot diameter d for various heat input power P at a rotation velocity of 1000 rpm.

listed depending on the beam spot diameter d for various heat input power P. A higher rotation speed leads to more efficient cooling as the heat input per voxel decreases. The radiative cooling is most efficient right after the heat is deposited and the shortened rotation cycle has a negligible impact. The dependency of maximum temperature on the beam spot diameter is presented in Figure 3 for different rotation velocities at a heat input power of 2 W. It can be seen that for small beam diameters a rotation with 1000 rpm is necessary to keep the target material below a temperature of 1000°C. This value serves as an empirical benchmark to ensure that the material properties of the carbon target remain unchanged during irradiation. With a target frame diameter of 6.3 cm, the trajectory of a beam spot close to the frame becomes around 14 cm long which is sufficient to keep the maximum temperature of the target material below the benchmark value. The Figure 4 demonstrates

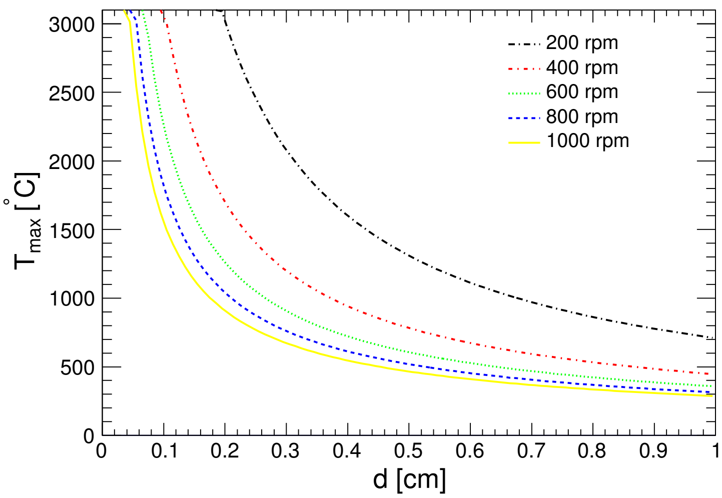

Figure 3: [color online] Dependency of the maximum temperature T_{max} from beam spot diameter d at a heat input power of 2 W for different rotation velocities.

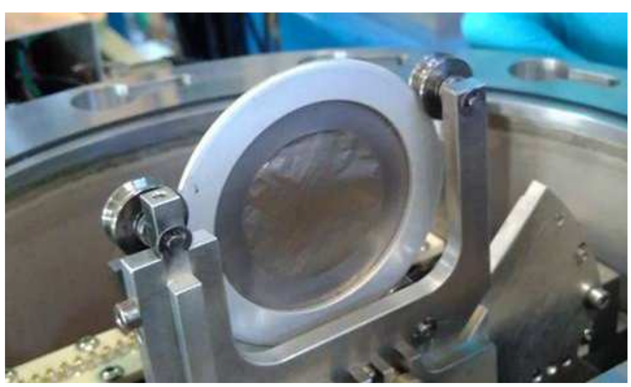

Figure 4: [color online] Photography of a 35 μg/cm^2 carbon target after 50 h of exposure to 2.5 pμA beam. The beam focus is at the upmost position so the beam spot forms a track along the outer area of the target foil close to the frame.

the effect of 50 h of 2.5 pμA beam exposure on the target material. The beam spot forms a track along the outer area of the target foil during the rotation.

2.2. Charged Particle Detection

Light charged particles from the reaction are detected by a set of annular S1- and S3-type DSSSD (Double-Sided Silicon Strip Detector) based on chips manufactured by Micron Semiconductor Ltd., where the S1 (S3) chips are segmented in 16 (24) rings on the junction side and 16 (32) sectors on the ohmic side. In the design (see Figure 1) developed by the IPHC Mechanics and Microtechnique Department at CNRS Strasbourg, the chips are sandwiched between low outgassing RO4003C Rogers® ceramics which serve to replace a regular PCB in its role of providing detector polarization and signal readout. The same design permits to fit in chips of 500 μm (S1/S3) or 1000 μm (S3) thickness where the incomplete rings of S1 are connected to a closed circle with an adapted PCB cabling.

The signal connection to the front-end electronics is via a series of contacts at the base of the PCB connected via spring-like pins on the detector support inside the reaction chamber. This connection system is integrated into the vertical slots of the sliding system for the PCBs that are kept in position with a precision better than 1 mm using clamps. Low-outgassing Kapton® insulated cables feed the electronics signals into sets of MPR-16D differential Mesytec® preamplifier cards outside the reaction chamber before processing towards the digitizers.

Integrated aluminum absorber foils in front of the silicon detectors protect them from delta electrons and radiation damage from scattered beam particles under experimental conditions. The thickness of the foils is adapted to minimize the degradation of the proton and alpha-particle energies. The system is grounded to mitigate the effects of charging. The junction side of the DSSSD is biased with a negative potential, while the ohmic side facing the target is grounded, thus guarding against damage due to possible sparking from charge depositions of beam induced particles between the protecting aluminum foil and the detector surface.

The annular charged-particle detectors, placed along the beam axis, are displayed in blue in Figure 1 in the top view of the scattering chamber. Upstream a pair of S3 and S1 detectors located 5.6 cm and 3.1 cm, respectively, from the target. The relative positioning is chosen to avoid shadowing of the target vertex. At the same time, the angular coverage is maximized. Downstream, an S3 is at around 6 cm from the target. In this configuration, the angular acceptance is 30% of the solid angle. The angular coverage per strip ranges from 10.0 mrad (outer ring S3) to 27.8 mrad (inner ring S1) due to the compact geometry of the system. The relative energy resolution obtained with the α-emitter ^{239}Pu is 0.5% FWHM at 5154 keV. During the commissioning campaign, the detector in forward direction is shielded against beam induced background, most importantly secondary protons from the ^{12}C(d,p)^{13}C reaction, with aluminum foils of up to 1 μm thickness. For this extreme case, this results in an energy resolution of 2.7% FWHM in the energy regime of the ^{239}Pu source. An energy spectrum with light charged particles from the ^{12}C+^{12}C reaction at a beam energy of 11 MeV is shown in Figure 5, for the S3 detec-

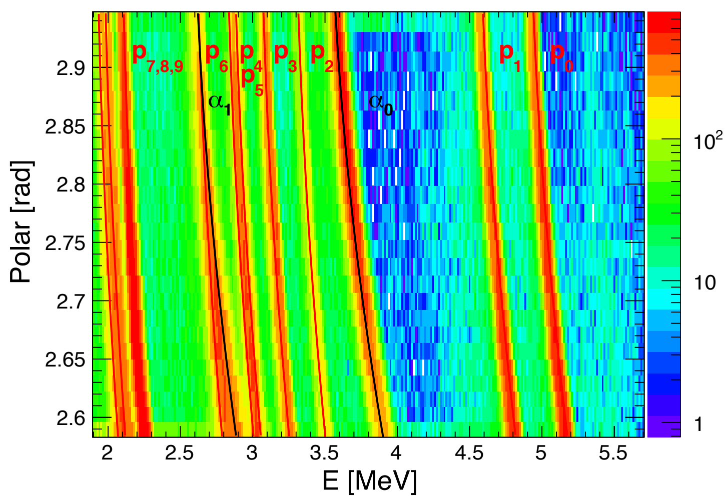

Figure 5: [color online] Angular distribution of protons and alphas associated with various excitation levels i of the corresponding fusion evaporation nucleus ^{23}Na (p_i) and ^{20}Ne (α_i), respectively, for the S3 in backward direction at a beam energy of 11 MeV.

tor in backward direction. In the angular range, various bands with protons and alphas, where the index stands for the excitation level of the associated daughter nucleus ^{23}Na and ^{20}Ne, respectively, are labeled. The solid lines are kinematics calculations for the emitted particles. The angular resolution is sufficient to cleanly distinguish the exit channels. For clean separation of protons and alphas, additional selection criteria based on timing are used (see section 3).

E [MeV]	0.01	0.44	1.0	1.63	2.0	3.0	4.0	5.0	6.0	7.0
ϵ_{sing} [%]	23.1	8.0	3.5	2.2	1.8	1.1	0.7	0.5	0.4	0.4
ϵ_{sum} [%]	23.1	8.6	4.1	2.6	2.1	1.4	1.0	0.7	0.6	0.5

Table 2: Full-energy peak gamma detection efficiency ϵ of 36 $LaBr_3$ detectors in percent. The efficiency obtained from the analysis of single detector spectra ϵ_{sing} is compared to the value reflecting the total energy deposit in the array ϵ_{sum}.

higher gamma energies.

$LaBr_3$(Ce) as a scintillator is well known to contain appreciable levels of self activity from the decay of ^{138}La and the chemically similar ^{227}Ac isotope [7]. The former is the main source of the background with around 100 Hz per detector. The ^{138}La nucleus decay comprises two gamma lines at 789 keV and 1436 keV from the de-excitation of the daughter nuclei ^{138}Ce and ^{138}Ba, respectively. The former decay is accompanied by a beta particle with an end-point energy of 258 keV while the latter gamma line is broadened due to X-rays from the electron capture escaping the crystal. These features can be well reproduced in simulation [10]. For the STELLA setup, the ^{138}La decay pattern is implemented for all $LaBr_3$ assembled in the detection array and compared to packages of experimental data to obtain energy calibration correction parameters (see [9] for details). The quadratic energy-response term is suppressed by 10^{-10} with respect to the linear term for gamma energies lower than 1.5 MeV, characterized with multiple emission lines of a ^{152}Eu source. In the fit of experimental data to the nominal energies in the simulated self-decay spectrum, an exponential background is also taken into account. The corrections with data samples of 45 min are illustrated in Figure 9 for the peak position of the 1.436 MeV line of the ^{138}La decay, which

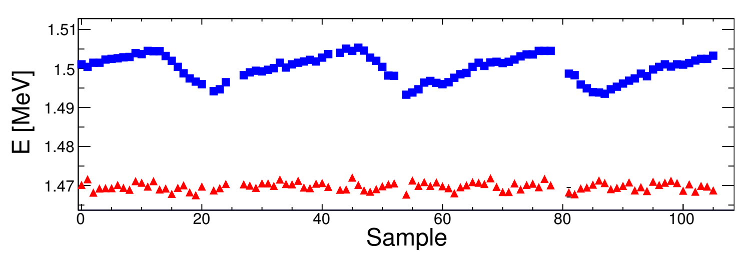

Figure 9: [color online] Correction (red triangles) of the temperature drift (blue squares) of a $LaBr_3$ detector over three days using data samples of around 45 min. The offset of uncorrected data reflects the strong drift since the detector calibration.

is accompanied by 37 keV barium x-rays, for a three-day data set. The day-night cycling of ± 7 keV (blue squares) is corrected with a precision of a few keV (red triangles). It is noteworthy to mention that the ^{138}La decay can also be used to synchronize the time offsets of the gamma detector matrix. Here, the time stamp difference of pairs of $LaBr_3$ is analyzed around the 1.436 MeV line where Compton events generate unambiguous coincident signatures used to extract the time calibration parameters.

2.4. Data Acquisition

The $LaBr_3$ signals (FATIMA) are processed by a 1 GHz VME-based Caen® V1751 card that accepts external triggers and clocks to synchronize with additional devices. The QDCs are remotely controlled by MIDAS (Multi Instance Data Acquisition System), developed in the Daresbury laboratory [11].

For the charged-particle signals, commercial ABACO® 125 MHz μTCA compatible FMC112 cards hosting 12 ADC channels are used for digitization. In the design developed by the IPHC-SMA group at CNRS Strasbourg, two FMC112s are grouped using a FC7 AMC (Advanced Mezzanine Card) [12] with 4 GB DDR3 memory around a Xilinx Kintex7 FPGA and a communication protocol/framework based on the IPbus communication scheme. Each STELLA acquisition card provides 24 single-ended DC coupled input channels with 2 V range and a programmable DC offset correction of ± 1.25 V. The digital triggering system supports TTL compatible I/O used for the synchronization with the gamma detection system. The data readout is through a μTCA crate with Gbit ethernet communication providing remote control. The PC interface is based on the TNT corpus [13] which allows for the online analysis of single signals with a trace acquisition mode in addition to the time-stamped energy acquisition features. The Java-based software is substantially expanded for the STELLA experiment for a comprehensive handling of the DAQ with DGIC (Distributed Glibex IPbus Control), for the setup of single STELLA cards with GIC (Glibex IPbus Control), and for the merging of all data streams onto tape alongside offline analysis functionality with TAN (Tnt ANalysis) servers.

The time alignment of the gamma ray and particle detection is based on reference signals from a μTCA compatible GLIB (Gbit Link Interface Board) card [14]. Figure 10 illustrates the distribution of a 10 MHz signal to all clocks on the respective cards for synchronization as well as the time reference signal to dedicated readout channels. The FC7 boards pass the

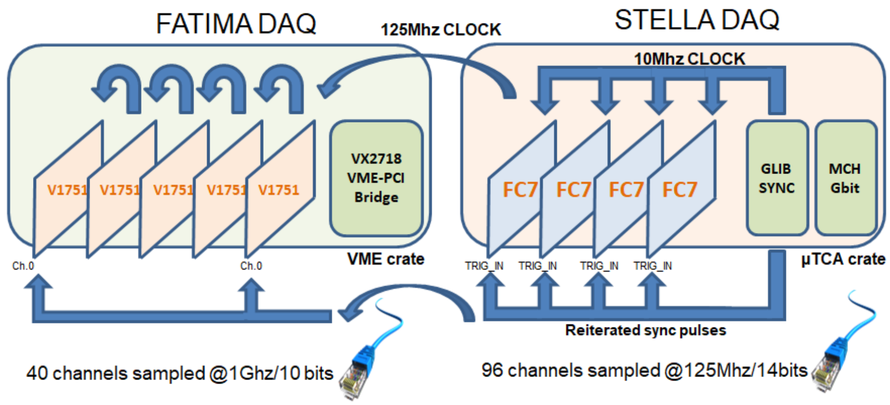

Figure 10: [color online] Time stamp synchronization of the particle (STELLA) and gamma (FATIMA) data acquisition. The clock (10 MHz) and occasional time signals are distributed by the GLIB card.

clock to their daughter cards and distribute one 125 MHz signal from a FMC112 to the Caen V1751 modules, where it is daisy chained among the internal clocks. The time reference signals are used to determine the time-stamp offsets and to detect drifts between individual clocks.

Using four FC7 boards and five V1751 modules, 96 channels for particle and 38 channels for gamma ray detection are established. The stand-alone time-stamped trigger-less data acquisition is synchronized with reference signals from a GLIB card with a precision of a few nanoseconds.

3. Particle-Gamma Coincidences

The STELLA experiment is installed at the Andromède accelerator [15] in Orsay, *France*, providing ^{12}C beam intensities of particle-micro-ampere. Fusion reactions are measured by the coincident detection of gamma rays and light charged particles broadly following the methodology of Jiang *et al.* [4, 5]. Gamma detector spectra and gamma-particle coincidence spectra can be found in [9] and [16]. The excellent timing resolution of the $LaBr_3$(Ce) detectors used in STELLA as compared to the germanium detectors employed by Jiang *et al.* combined with the time-stamped data acquisition with sampling times of 8 ns for the digital triggering [13], provides a new functionality, namely that particle-gamma timing can be used to further improve background suppression and to cleanly discriminate evaporated protons from alpha particles.

The achievable separation is shown for the ^{12}C+^{12}C reaction at a beam energy $E = 11$ MeV in correlation with particle detector energy entries in Figure 11. The gamma-particle time stamp difference Δt reveals contribu-

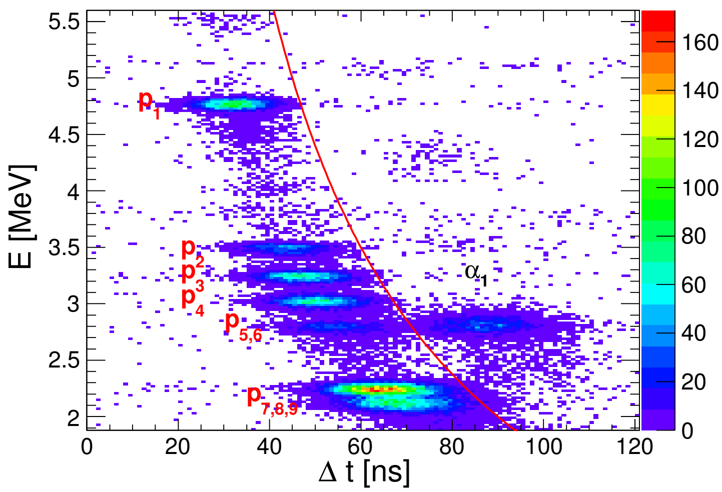

Figure 11: [color online] Time correlation of particle energies in ^{12}C+^{12}C reactions with coincident gamma rays at a beam energy $E_{beam} = 11$ MeV. The gamma-particle time stamp difference Δt allows to distinguish protons (left of red line) and alphas (right of red line).

tions from protons (left of red line) and alphas (right of red line), where the time-of-flight difference cannot be resolved. The distributions are separated due to different electronic pulse shapes based on the respective energy deposition characteristics in the silicon detector substrate. These processes depend on the particle velocity that reflects in the quadratic trend of the distributions with respect to the energy of the alphas and protons [17].

The selection criterion indicated by the red line in the picture is used to resolve the distributions from different particle types in the energy spectrum in Figure 12. The particle spectrum (blue line) at a beam energy $E_{beam} = 11$ MeV is decomposed into the proton- (red line) and alpha-channel (black line) contributions in coincidence with gammas. Several transitions i from excited states of ^{23}Na (p_i) and ^{20}Ne (α_i) are labeled. Based on the timing, contributions from different particles can be resolved as can be seen for the entries around 3 MeV in the spectrum. Note that the background contamination *e.g.* at the p_0 or α_0 energies in the coincidence spectrum, where no associated gamma ray is emitted, is based predominantly on random co-

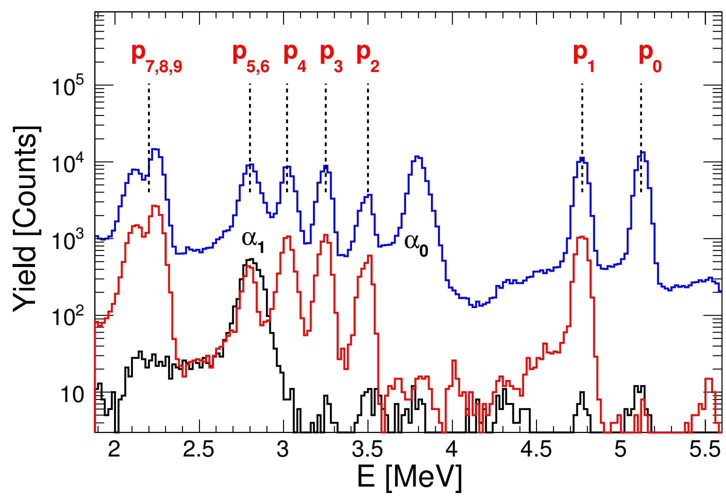

Figure 12: [color online] Particle energies at a beam energy $E_{beam} = 11$ MeV where proton and alpha contributions are selected based on the timing. The distribution of all particles (blue line) can be resolved into alphas (black line) and protons (red line) coincident with gamma rays. The excitation levels i of ^{20}Ne (α_i) as well as ^{23}Na (p_i) are labeled and the positions of the proton energies are indicated to guide the eye.

incidence with the self-activity of the $LaBr_3$ and can be well determined in the time domain (compare Figure 11). Beyond this example case, the technique is also utilized to extract the relevant signals from the large overall background, essentially due to the the ubiquitous contamination of hydrogen and deuterium in the target.

4. Summary

The STELLA experiment has been commissioned at the Andromède accelerator with long-running measurements of the ^{12}C+^{12}C reaction using fixed targets as well as the rotating target mechanism. The used beam energy ranges from around the Coulomb barrier of ^{12}C fusion towards deep sub-barrier energies. The beam intensity was gradually increased from 100 pnA to 2.5 pμA, where for the latter settings exclusively the rotating targets are utilized. These targets are accessible for the measurement of the thickness off the illuminated area and beside it to determine the effect of the beam exposure. During the commissioning, the S3-type DSSSDs and various gamma detector configurations of around 30 $LaBr_3$ detectors are installed. The accuracy of the performance of the gamma ray detection system is guaranteed by the instant-calibration routine. It is based on the comparison of the simulated ^{138}La decay with experimental data and has an accuracy of a few keV. Repeated alpha-source runs in the course of the campaign are utilized to ascertain the correctness of the particle detection performance.

The STELLA-FATIMA data acquisition systems are synchronized with frequently distributed time stamp pulses to dedicated readout channels. The reliability during long measurements is validated using coincident gamma-particle events from ^{12}C fusion reactions. An enormous background reduction is achieved with the measurement of synchronous events in the gamma and particle detection system. Beyond this, reaction channels with different species of charged particles are well separated based on the timing. This guarantees a reliable measurement of deep sub-barrier partial fusion cross sections with the STELLA station.

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