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Ng, F., Hallet, B., Sletten, R. et al. (1 more author) (2005) Fast-growing till over ancient ice in Beacon Valley, Antarctica. Geology, 33 (2). 121 - 124. ISSN 0091-7613

https://doi.org/10.1130/G21064.1

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1 2 Fast-growing till over ancient ice in Beacon Valley, 3 Antarctica 4 5 Felix Ng Department of Earth, Atmospheric and Planetary Sciences, Massachusetts Institute of 6 7 Technology, Cambridge, Massachusetts 02139, USA 8 Bernard Hallet 9 Ronald S. Sletten 10 John O. Stone 11 Quaternary Research Center and Department of Earth and Space Sciences, University of 12 Washington, Seattle, Washington 98195, USA 13 14 15 ABSTRACT We analyze published cosmogenic ³He depth profiles through the till that covers 16 17 relict glacier ice in Beacon Valley, Antarctica, in order to derive rigorous constraints on the till-thickness history, and on the amount and rate of ice loss by sublimation. The till is 18 19 a residue of debris-laden ice that sublimed. The ³He profiles show that the lower 80% of 20 the till formed in the past 310–43 kyr under sublimation rates averaging $>7 \text{ m}\cdot\text{Myr}^{-1}$. 21 Such rapid, recent growth of the till contradicts previous interpretations that it is older 22 than 8.1 Ma at an adjacent site, where it encloses volcanic ash of this age. We question 23 whether the ash provides a valid age constraint for the ice. Cosmogenic nuclide analysis 24 of the till where the ash was collected for dating should resolve this question. 25 26 Keywords: Antarctica, Dry Valleys, glacial deposits, cosmogenic elements, sublimation. 27 28 **INTRODUCTION** 29 The recent history of East Antarctica is key to understanding the response of large 30 ice sheets to climate forcing. Field evidence has spurred a debate on two conflicting 31 scenarios advocated for this history: stable glacial conditions since the middle Miocene 32 (Sugden et al., 1993) and ice-sheet disintegration under warming during the Pliocene 33 (Webb et al., 1984). The ice in Beacon Valley is important in this context. It is debris-34 laden, thought to be the remains of an expansion of Taylor Glacier into the valley, and 35 lies under a till layer produced by its own sublimation. Sugden et al. (1995) argued for prolonged glacial conditions because they discovered 8.1 Ma volcanic ash in the till. 36 37 Under their interpretation, the ash is a direct air-fall deposit into a former frost crack in 38 the till, and the ice, till, and crack all predate 8.1 Ma. This interpretation implies not only 39 the oldest glacier ice on Earth, but also a low sublimation rate for its survival-and 40 hence, a persistent cold climate-since the Miocene, with correspondingly little extra 41 accretion of the till. In contrast, ice sublimation rates from a physical model are high, 42 $\sim 10^3 \text{ m}\cdot\text{Myr}^{-1}$ (Hindmarsh et al., 1998). Given a reasonable initial thickness for the ice of

no more than a few hundred meters (Potter et al., 2003), its age should be less than 1 Ma
(Van der Wateren and Hindmarsh, 1995).

One way to resolve this age controversy is to decipher the history of the till from 45 46 cosmogenic nuclide measurements. The till is a diamict formed mainly from debris 47 originally in the ice, although its upper part contains eolian sand and weathered rocks 48 also. Material deep in the ice is shielded from cosmic rays, but is uncovered, becomes 49 less shielded as the ice sublimes, and finally accretes to the base of the till, feeding its 50 growth (Fig. 1A). In such material, the production rate of nuclides, such as ³He, increases 51 as the overlying ice thins; then, after the material joins the till, its depth and the 52 production rate remain constant. We develop a model of nuclide accumulation to 53 reexamine published data from Beacon Valley.

54 Schäfer et al. (2000), Phillips et al. (2000), and Marchant et al. (2002) analyzed 55 cosmogenic ³He in clasts from three vertical profiles in the till overlying the ice (Table 1). The profiles are within ~ 1 km of each other. ³He concentration N decreases rapidly 56 with depth z. This result is expected because the production rate attenuates with depth 57 58 and because, in a sublimation till, deep clasts are exposed for a shorter time compared to 59 shallow clasts after they accrete to the till (Fig. 1). The profiles' monotonic decrease 60 suggests that the till did not undergo cryoturbation (Phillips et al., 2000; Marchant et al., 2002), even though the ground in Beacon Valley is patterned conspicuously by 61 62 contraction-crack polygons (Berg and Black, 1966; Black, 1973; Sletten et al., 2003).

63 Two arguments to support antiquity of the ice have been made using cosmogenic 64 depth profiles: (1) Some clasts at the surface have exposure ages of 2–3 Ma, so the ice 65 beneath is at least as old (Schäfer et al., 2000; Oberholzer et al., 2000; Marchant et al., 2002). (2) Schäfer et al. (2000) devised a method of calculating the thickness of ice that 66 sublimed using ³He concentrations in surficial–basal clast pairs from the till. When 67 68 coupled with the till surface exposure age-a minimum age in view of weathering of the 69 surficial clasts—their method indicates maximum (average) sublimation rates of ≤ 90 70 $m \cdot Myr^{-1}$, which are considered to be low enough for ice survival.

Here we reach different conclusions. We argue that the ³He profiles constrain minimum, not maximum, sublimation rates; that the surficial clasts are unreliable indicators of age. Moreover, new constraints on the history of till thickness suggest that the ash was not emplaced in the way Sugden et al. (1995) envisaged. These results emerge when we analyze how the profiles record the sublimation and accretion processes.

77 78

MODEL OF NUCLIDE CONCENTRATION

79 Consider first a model for simulating the ³He profiles from clast-exposure history 80 (Fig. 1). We assume a nondeforming till of porosity ϕ . We measure the depth *z* relative to 81 the lowering surface and let $\ell(T)$ be the till thickness, where *T* denotes age. If the 82 sublimation rate is *S*(*T*) and the debris concentration of the subliming ice (by volume) is *c* 83 (<<1), then the till thickness at a rate

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85
$$-\frac{d\ell}{dT} = \frac{cS}{1-\phi}.$$
 (1)

86

The debris concentration c varies with T if debris in the ice is not uniformly distributed; we return to the consequences of this situation later.

Cosmogenic dating models that are used widely to constrain exposure age and erosion rate of rock surfaces (Lal, 1991) do not adequately describe our system. Although the ice may be likened as being eroded as it sublimes, the till is a lag that has no analogue in such models. Here we follow the depth history of each clast, z = h(T), to calculate its exposure history. Given its depth today, z_0 , we reconstruct *h* by backtracking (Fig. 1B) observing that *h* is constant after the clast accretes to the till; that the age of accretion, T_{A_2}

satisfies $\ell(T_A) = z_0$; and that, although h differs from z_0 prior to accretion, the clast,

96 contained then by the ice, approaches the surface at velocity $S + d\ell/dT$. These

- 97 considerations yield
- 98

$$h(T) = z_0 \qquad \text{for } 0 \le T \le T_A,$$

$$h(T) = \ell(T) + \int_{T_A}^T S(\xi) d\xi \qquad \text{for } T > T_A,$$
(2)

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101 in which the integral represents the overlying ice thickness (ξ is the variable of 102 integration). We distinguish three stages in the clast's exposure history: *inheritance* ($T \ge$ T_{AS}), preaccretion ($T_{AS} > T > T_A$), and postaccretion ($T_A > T \ge 0$), where T_{AS} is the age of 103 the till surface (= T_A for $z_0 = 0$; Fig. 1). Inheritance thus comprises nuclide contributions 104 105 before the till layer develops. We separate inheritance from preaccretion, because it 106 includes exposure contributions before the clast was incorporated into the ice, which are 107 unknown. This uncertainty makes it difficult to determine how the stages partition the 108 nuclide concentration N measured for a given clast.

For a stable cosmogenic nuclide such as ³He, we model its accumulation rate in the clast (using Lal's (1991) formulation) as

111

 $-\frac{dN}{dT} = P_0 e^{-\frac{\rho_{\rm I}}{\Lambda} [h(T) - \ell(T)]_{0+}} e^{-\frac{\rho_{\rm S}}{\Lambda} (1 - \phi) z_0}, \qquad (3)$

113

where P_0 is the surface production rate, ρ_I is ice density, ρ_S is sediment density, Λ is 114 115 absorption mean free path, and $[x]_{0+} = \max(x, 0)$. In equation 3, the first exponential factor 116 describes shielding of the clast by ice; the second exponential factor describes shielding 117 of the clast by overlying debris, which remains above the clast after enclosing ice sublimes away. Equation 3 ignores ³He production by muon-induced reactions, whose 118 119 rate at the surface has not been calibrated but is estimated at $\sim 3\%$ of the corresponding 120 rate by spallation (Lal, 1987). We expect muon-induced production to dominate at depths 121 >4–5 m. Including its effect in our (spallation-only) model leads to a slight increase in the 122 ³He accumulated in clasts prior to accretion that lowers the bound $T_{A,max}$, raises the bounds S_{\min} and $\Delta_{I,\min}$ derived below, and strengthens the conclusions of this paper. 123 124 Now, the integral of equation 3 from $T = T_{AS}$ to T = 0 represents the ³He

accumulated in the clast since the till layer began forming. We substitute for *h* from

equation 2 and, by replacing z_0 with z, generalize this integral for all clasts. If we include the inheritance stage, the outcome is an expression for today's depth profile:

128

$$N(z) = N_{\text{Inh}}(z) \qquad (\text{inheritance, by } T_{\text{AS}} \text{ years ago}) + P_0 e^{-\frac{\rho_{\text{S}}}{\Lambda}(1-\phi)z} \int_{T_{\text{A}}(z)}^{T_{\text{AS}}} \exp\left[-\frac{\rho_{\text{I}}}{\Lambda} \int_{T_{\text{A}}(z)}^{T} S(\xi) d\xi\right] dT \qquad (\text{preaccretion}) + P_0 e^{-\frac{\rho_{\text{S}}}{\Lambda}(1-\phi)z} T_{\text{A}}(z) \qquad (\text{postaccretion}),$$
(4)

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in which we identify each exposure stage and N_{Inh} denotes the inherited concentration in material at depth *z* today. In a forward simulation S(T) and c(T) are specified, and equation 4 is evaluated with the accretion age distribution $T_A(z)$ (or $\ell(T)$, its inverse)

134 found from equation 1.

136 THE INVERSE MODEL

137 The challenge is the opposite: to find the sublimation and till-thickness histories 138 S(T) and $\ell(T)$, given N(z). Equations 4 and 1 cannot be solved for these histories uniquely 139 because of the extra unknowns N_{Inh} and c. In particular, the debris concentration c(T) of 140 the sublimed ice may differ from c for the relict ice today. The measured profiles also are 141 discrete. Here we seek constraints instead of solution.

We first raise a caveat on the method by Schäfer et al. (2000) that explains also our apparent reversal of their maximum bound on sublimation rate in this paper. They assumed a constant rate of sublimation S_c and inheritance-free clasts ($N_{Inh} = 0$). In this case, the ratio of N for two clasts from the surface and base of the till can be used to find the initial ice thickness between the clasts, because the overall shielding effect of the ice as it sublimed is predictable. For the clasts, equation 4 reduces to

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149
$$N(0) = P_0 T_{\rm AS}, \quad N(\ell_0) = P_0 e^{-\frac{\rho_{\rm S}}{\Lambda}(1-\phi)\ell_0} \int_0^{T_{\rm AS}} \exp(-\rho_{\rm I} S_{\rm c} T / \Lambda) \, dT, \quad (5)$$

150

151 where ℓ_0 is the till thickness today, and the ratio of N can be written in the form

152

$$\frac{N(\ell_0)}{N(0)e^{-\rho_{\rm S}(1-\phi)\ell_0/\Lambda}} = \frac{1}{Z} \int_0^Z e^{-\rho_{\rm I}\xi/\Lambda} d\xi = \frac{1-e^{-\rho_{\rm I}Z/\Lambda}}{\rho_{\rm I}Z/\Lambda}, \quad (6)$$

154

153

155 where $Z = S_c T_{AS}$ is the sublimed ice thickness in the model. Schäfer et al. (2000) used 156 equation 6 to determine Z from the end data of a profile, and the sublimation rate from S_c 157 $= Z/T_{AS} = P_0 Z/N(0)$. They claimed that in the last step, surface erosion would render T_{AS} 158 (denominator) a minimum age, making S_c a maximum sublimation rate. The caveat is that 159 Z (numerator) is not an upper-bound estimate: the actual sublimed ice thickness could 160 exceed Z if unsteady sublimation (e.g., due to climate change) had violated the 161 assumption that *S* was constant. Therefore, the value S_c does not constrain sublimation 162 rates and cannot be used to dismiss the model results of Hindmarsh et al. (1998). (But, as 163 expected, S_c satisfies our constraint below where we allow for all possible sublimation 164 histories. For profiles I, II, and III, Marchant et al. [2002] and Schäfer et al. [2000]

165 obtained $S_c \approx 20, 90$, and $6 \text{ m} \cdot \text{Myr}^{-1}$, respectively.)

166 In contrast, an approach is now developed to give robust *minimum* mean 167 sublimation rates (S_{min}). The crux is to derive, for any pair of clasts in a profile, a *lower* 168 bound on the original thickness of ice that separated them ($\Delta_{I,min}$) and an *upper* bound on 169 the time over which this ice sublimed (t_{max}). The result $S_{min} = \Delta_{I,min}/t_{max}$ is rigorous.

170

171 Constraint on Ice Thickness

172 Suppose the clasts are numbered 1 (lower) and 2 (upper) and have concentrations 173 N_1 and N_2 , depths z_1 and z_2 , respectively (Fig. 1A). We can constrain their original 174 separation in the ice (Δ_I) because the concentrations reflect different depth histories. The clasts' separation today is $\Delta_{\rm S} = z_1 - z_2$, so the intervening sediment thickness is $(1 - \phi)\Delta_{\rm S}$. 175 176 Given the shielding by this sediment, we can predict what the ratio N_2/N_1 should be, but 177 the data show that the ratio is always larger, which could only have resulted because of 178 intervening ice that has disappeared. If we neglect ³He inheritance before the clasts were 179 incorporated into the ice, then the minimum intervening ice thickness, $\Delta_{I \min}$, can be 180 computed from

181

$$e^{\left[\rho_{\rm I}\Delta_{\rm I,min} + \rho_{\rm S}(1-\phi)\Delta_{\rm S}\right]/\Lambda} = \frac{N_2}{N_1}.$$
 (7)

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184 The value $\Delta_{I,min}$ is the minimum initial ice thickness, because the ice could only have 185 thinned: for a smaller initial thickness, past ³He production rates in the clasts would have 186 been too similar for us to explain the data. We calculate $\Delta_{I,min}$ from N_1 , N_2 , and Δ_S (Table 187 1). Equation 7 holds regardless of sublimation rate changes and does not depend on P_0 . 188 ³He production by muon-induced reactions, which have large attenuation lengths, 189 effectively increases Λ used in our model, making $\Delta_{I,min}$ an underestimate.

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191 **Constraint on Sublimation Time**

192 Next, we deduce a maximum sublimation time t_{max} for the ice between clasts 1 193 and 2. This ice began subliming after clast 2 (the upper clast) accreted to the till and none 194 of it remains today (Fig. 1A), so the maximum accretion age of clast 2 suffices as our 195 choice for t_{max} . For any clast, its maximum accretion age ($T_{A,max}$) is simply the maximum 196 duration of its postaccretion stage, which we can calculate by attributing all of its 197 measured *N*-value to exposure at its current depth *z* in the till; thus,

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199
$$T_{A}(z) \leq T_{A,\max}(z) = \frac{N(z)}{P_{0}e^{-\rho_{S}(1-\phi)z/\Lambda}}.$$
 (8)

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201 Accordingly we put $t_{\text{max}} = T_{A,\text{max}}(z_2)$. In Table 1, dividing $\Delta_{I,\text{min}}$ by $T_{A,\text{max}}(z_2)$ gives S_{min} ,

202 our minimum sublimation rate.

The bound t_{max} cannot be tightened, for we cannot deduce from the profiles the 203 204 most recent time at which the lower clast (clast 1) could have joined the till (i.e., a 205 minimum T_A) without making assumptions. Consequently, for a given depth profile, we 206 cannot resolve the different sublimation periods for ice that existed between successive 207 clast pairs. For any two clasts, the time over which S_{\min} is defined (and constrains the sublimation rate) is fixed by the upper clast—it begins no earlier than the age $T_{A,max}(z_2)$ 208 209 and ends at the present, regardless of where in the profile the lower clast is taken. Hence 210 we pick the lower clast always from the base of the till, to ensure the largest admissible 211 Δ_{Lmin} for calculating S_{min} .

212

213 **DISCUSSION**

Our results (Table 1) shed new light on the evolution of the ice and overlying till in Beacon Valley. Mean sublimation rates have not necessarily been low. Profiles I, II, and III indicate minimum mean rates S_{\min} of ~4, 23, and 2 m·Myr⁻¹, respectively, within the past 1.1 Myr, 170 kyr, and 1.6 Myr, causing at least several meters of ice loss at all three sites. Erosion of the surficial clasts can invalidate these results, but not the higher S_{\min} values for the more recent past indicated by buried clast pairs.

220 Rapid sublimation (Hindmarsh et al., 1998) could be considered likely, if one is 221 prepared to make assumptions about the ice that sublimed. Its maximum average debris 222 concentration can be calculated from our results as the ratio of sediment thickness to 223 minimum ice thickness: $c_{\text{max}} = (1 - \phi)\Delta_{\text{S}}/\Delta_{\text{I,min}}$ (Table 1). c_{max} is several times c_0 (~3%) 224 for the relict ice. In contrast, one might expect the ice that sublimed to contain less debris 225 than the relict ice, if the latter is basal ice from Taylor Glacier, as assumed by Sugden et 226 al. (1995). Thus our bounds may be overconservative. By assuming ice no dirtier than today's, i.e., $c(T) \le c_0$, alternative minimum bounds can be found from $\Delta_{\text{Lmin}}^{R} = (1 - 1)^{R}$ 227 ϕ) $\Delta_{\rm S}/c_0$ (for sublimed ice thickness) and $S_{\rm min}^{\rm R} = \Delta_{\rm I,min}^{\rm R}/T_{\rm A,max}$ (for sublimation rate). 228 These bounds indicate mean sublimation rates exceeding $\sim 10-100 \text{ m} \cdot \text{Myr}^{-1}$ (Table 1), 229 consistent with an independent estimate of 50 m·Myr⁻¹ from ¹⁰Be analysis of the ice and 230 231 of debris within the ice (Stone et al., 2000) in the part of Beacon Valley where profiles I to III were measured. 232

233 Equation 8 constitutes a powerful constraint on the till accretion history. On the 234 depth vs. age plot of Figure 2A, the accretion history $T = T_A(z)$ is confined to the region 235 right of the line representing the maximum accretion age $T = T_{A,max}(z)$. Consequently the 236 line also limits the till thickness: the apparent exposure age of a clast, $T_{A,max}$ (calculated 237 on the basis of current shielding), implies that the till was, at that age, no thicker than the 238 till above the clast today. Prior to $T_{A,max}$ the clast must have still been in the ice and below 239 the till. For discrete depth profiles, this constraint takes the form of a staircase (Figs. 2B, 240 2C) provided that the till had not thinned over time.

241 We stress that, according to Figures 2B and 2C, all but the topmost 20% of till at 242 the sites measured by Phillips et al. (2000) and Marchant et al. (2002) formed within the 243 past 310 kyr (profile I) and 43 kyr (profile II). Prior to these times the till was 244 exceptionally thin: ≤ 14 cm (profile I) and ≤ 9 cm (profile II), and by these times there 245 were relatively old clasts aged 800 ka (I) and 130 ka (II) at the surface. These surficial 246 clasts have uncertain provenance; unlike subsurface clasts released by ice, they might 247 have originated via rockfall onto Taylor Glacier. Prior exposure may account for most of 248 their ³He concentration, so that they may not be used to infer a minimum age for the ice,

which could be as little as several hundred thousand years. Although the old exposure age of the surficial clasts can be explained in other ways (e.g., the ice that originally separated them from the next lower clast in the profile was very thick, or sublimed very slowly), we caution against using them to support the case for ancient ice.

253 An outstanding conundrum is the past relationship between ash and till. The 254 interpretation advanced by Sugden et al. (1995) is that the ice in Beacon Valley was 255 already mantled by \sim 50 cm of till at 8.1 Ma, when ash filled a frost crack, and that the till 256 has thickened little since. In contrast, our analysis shows that no more than a thin veneer 257 of till existed prior to 310 ka, and that the bulk of the till has accreted since. The 3 He 258 profiles examined here are not located at the "ash site", and their differences reflect some 259 spatial variability in till evolution. Nevertheless, the profiles are close enough spatially 260 and in stratigraphic context for our interpretation of them to challenge the antiquity of the 261 till enclosing the ash. Our results show that the ash may not be a reliable stratigraphic 262 indicator. The case for Miocene ice is likely to remain unsettled until a profile similar to 263 the ones already discussed is measured at a site containing old ash.

264

265 ACKNOWLEDGMENTS

We thank S. Byrne and H. Conway for helpful discussions, and A. Fountain and R. C. A. Hindmarsh for critical review. This study was supported by a Massachusetts Institute of Technology Leavitt Research Fellowship (to Ng) and is based on work supported by the U.S. National Science Foundation under Grant Nos. 9726139 and 0124824.

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323	
324	FIGURE CAPTIONS
325	Figure 1. Model of subliming ice and accreting till with no deformation. A: Processes in a
326	reference frame fixed to the ice. B: Depth vs. age plot shows processes in a reference
327	frame fixed to till surface $z = 0$. Heavy dashed line denotes till-thickness history ℓ . Solid
328	arrowed line is depth history h of clast at $z = z_0$ today; sublimation uncovers clast until it
329	accretes to the till at age T_A , whose value depends on (and is a function of) z_0 .
330	Trajectories of several other clasts are shown dotted.
331	Figure 2. Constraint on past till thickness using ³ He depth profiles. A: On depth vs. age
332	plot (right panel), the till-thickness history $z = \ell(T)$ or equivalently the accretion age
333	distribution $T = T_{\lambda}(z)$ (dashed line) must lie outside batched region to the right of the
334	boundary $T = T_{A(2)}$ (dustice line) must be obtained findence region, to the right of the
335	8 indicates the maximum till thickness at a given time R C: Application of model in Δ
336	to profiles L and II. In these cases the boundary $T_{1} = (7)$ is sten-like
550	o promos i uno ii. In those cuses the obtinuity I A,max(2) is step like.

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TABLE 1. COSMOGENIC ³He IN CLASTS FROM BEACON VALLEY TILL AND MODEL RESULTS

Data		Results							
z (cm)	N (× 10 ⁶	Clast pair	Δ_{s}	$\Delta_{I,min}$	T _{A,max} (Ma)	S _{min} (m·Mvr ^{−1})	c _{max} (%)	$\Delta_{I,min}^{R}$	S _{min} ^R (m⋅Mvr ⁻¹)
	atoms ·g ^{−1})	(cm)	(0)	()	()	((,,,,)	()	(),)
Profile I									
0	612	0–70	70	4.52	1.123	4.02	10.3	15.6	13.9
14	140	14–70	56	2.37	0.310	7.66	15.8	12.4	40.2
21	85	21–70	49	1.69	0.206	8.21	19.3	10.9	52.8
59	28	59–70	11	0.69	0.113	6.10	10.6	2.44	21.7
70	16	70–70	—	—	0.075	—	—	—	—
Profile II									
0	93	0–38	38	3.90	0.171	22.9	6.5	8.44	49.5
9	21	9–38	29	1.62	0.043	37.3	11.9	6.44	148.3
25	8.9	25–38	13	0.54	0.023	23.9	16.0	2.89	126.8
38	5.4	38–38	_	—	0.016	—	—	_	—
Profile III									
0	880	0–70	70	3.44	1.615	2.13	13.6	15.6	9.63
70	44	70–70	_	—	0.205	—	—	_	_

342 343 344 345 346 347 348 349 *Note*: Symbols: z = depth of clast sample; $N = {}^{3}$ He concentration; $\Delta_{s} =$ clast-pair separation; $\Delta_{Lmin} =$ minimum original interclast ice thickness; $T_{A,max}$ = maximum accretion age of upper clast of pair; S_{min} = minimum sublimation rate of interclast ice; $c_{max} = (1 - \phi)\Delta_S/\Delta_{l,min}$ = maximum debris concentration of ice that sublimed; $\Delta_{l,min}^{R} = (1 - \phi)\Delta_S/c_0$ (see discussion); $S_{min}^{R} = \Delta_{l,min}^{R}/T_{A,max}$ (see discussion). Data sources: Phillips et al. (2000) and Marchant et al. (2002) for profiles I and II and Schäfer et al. (2000) for profile III. Deepest clast of each profile is located at the base of till. In the $\Delta_{l,min}$ column, subtracting two values gives $\Delta_{l,min}$ for the two clasts appearing on the same row as the values. Model does not correct for the (unknown) sampling position on each clast. Model constants: $\rho_l = 0.9 \text{ g} \cdot \text{cm}^{-3}$, $\rho_s = 3.0 \text{ g} \cdot \text{cm}^{-3}$, $\phi = 1/3$, $\Lambda = 150 \text{ g} \cdot \text{cm}^{-2}$, $c_0 = 0.03$, and (following Marchant et al., 2002) $P_0 = 545$ atoms g^{-1} per year.

350





Figure 1 (65%) Ng et al. G21064



Figure 2 (44.5 %) Ng et al. G21064