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Fast-growing till over ancient ice in Beacon Valley, Antarctica

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ABSTRACT

We analyze published cosmogenic ^3He depth profiles through the till that covers relict glacier ice in Beacon Valley, Antarctica, in order to derive rigorous constraints on the till-thickness history, and on the amount and rate of ice loss by sublimation. The till is a residue of debris-laden ice that sublimed. The ^3He profiles show that the lower 80% of the till formed in the past 310–43 kyr under sublimation rates averaging $>7 \text{ m}\cdot\text{Myr}^{-1}$. Such rapid, recent growth of the till contradicts previous interpretations that it is older than 8.1 Ma at an adjacent site, where it encloses volcanic ash of this age. We question whether the ash provides a valid age constraint for the ice. Cosmogenic nuclide analysis of the till where the ash was collected for dating should resolve this question.

Keywords: Antarctica, Dry Valleys, glacial deposits, cosmogenic elements, sublimation.

INTRODUCTION

The recent history of East Antarctica is key to understanding the response of large ice sheets to climate forcing. Field evidence has spurred a debate on two conflicting scenarios advocated for this history: stable glacial conditions since the middle Miocene (Sugden et al., 1993) and ice-sheet disintegration under warming during the Pliocene (Webb et al., 1984). The ice in Beacon Valley is important in this context. It is debris-laden, thought to be the remains of an expansion of Taylor Glacier into the valley, and lies under a till layer produced by its own sublimation. Sugden et al. (1995) argued for prolonged glacial conditions because they discovered 8.1 Ma volcanic ash in the till. Under their interpretation, the ash is a direct air-fall deposit into a former frost crack in the till, and the ice, till, and crack all predate 8.1 Ma. This interpretation implies not only the oldest glacier ice on Earth, but also a low sublimation rate for its survival—and hence, a persistent cold climate—since the Miocene, with correspondingly little extra accretion of the till. In contrast, ice sublimation rates from a physical model are high, $\sim 10^3 \text{ m}\cdot\text{Myr}^{-1}$ (Hindmarsh et al., 1998). Given a reasonable initial thickness for the ice of

43 no more than a few hundred meters (Potter et al., 2003), its age should be less than 1 Ma
 44 (Van der Wateren and Hindmarsh, 1995).

45 One way to resolve this age controversy is to decipher the history of the till from
 46 cosmogenic nuclide measurements. The till is a diamict formed mainly from debris
 47 originally in the ice, although its upper part contains eolian sand and weathered rocks
 48 also. Material deep in the ice is shielded from cosmic rays, but is uncovered, becomes
 49 less shielded as the ice sublimates, and finally accretes to the base of the till, feeding its
 50 growth (Fig. 1A). In such material, the production rate of nuclides, such as ^3He , increases
 51 as the overlying ice thins; then, after the material joins the till, its depth and the
 52 production rate remain constant. We develop a model of nuclide accumulation to
 53 reexamine published data from Beacon Valley.

54 Schäfer et al. (2000), Phillips et al. (2000), and Marchant et al. (2002) analyzed
 55 cosmogenic ^3He in clasts from three vertical profiles in the till overlying the ice (Table
 56 1). The profiles are within ~ 1 km of each other. ^3He concentration N decreases rapidly
 57 with depth z . This result is expected because the production rate attenuates with depth
 58 and because, in a sublimation till, deep clasts are exposed for a shorter time compared to
 59 shallow clasts after they accrete to the till (Fig. 1). The profiles' monotonic decrease
 60 suggests that the till did not undergo cryoturbation (Phillips et al., 2000; Marchant et al.,
 61 2002), even though the ground in Beacon Valley is patterned conspicuously by
 62 contraction-crack polygons (Berg and Black, 1966; Black, 1973; Sletten et al., 2003).

63 Two arguments to support antiquity of the ice have been made using cosmogenic
 64 depth profiles: (1) Some clasts at the surface have exposure ages of 2–3 Ma, so the ice
 65 beneath is at least as old (Schäfer et al., 2000; Oberholzer et al., 2000; Marchant et al.,
 66 2002). (2) Schäfer et al. (2000) devised a method of calculating the thickness of ice that
 67 sublimed using ^3He concentrations in surficial–basal clast pairs from the till. When
 68 coupled with the till surface exposure age—a minimum age in view of weathering of the
 69 surficial clasts—their method indicates maximum (average) sublimation rates of ≤ 90
 70 $\text{m}\cdot\text{Myr}^{-1}$, which are considered to be low enough for ice survival.

71 Here we reach different conclusions. We argue that the ^3He profiles constrain
 72 minimum, not maximum, sublimation rates; that the surficial clasts are unreliable
 73 indicators of age. Moreover, new constraints on the history of till thickness suggest that
 74 the ash was not emplaced in the way Sugden et al. (1995) envisaged. These results
 75 emerge when we analyze how the profiles record the sublimation and accretion
 76 processes.

77

78 **MODEL OF NUCLIDE CONCENTRATION**

79 Consider first a model for simulating the ^3He profiles from clast-exposure history
 80 (Fig. 1). We assume a nondeforming till of porosity ϕ . We measure the depth z relative to
 81 the lowering surface and let $\ell(T)$ be the till thickness, where T denotes age. If the
 82 sublimation rate is $S(T)$ and the debris concentration of the subliming ice (by volume) is c
 83 ($\ll 1$), then the till thickens at a rate

84

$$85 \quad -\frac{d\ell}{dT} = \frac{cS}{1-\phi}. \quad (1)$$

86

87 The debris concentration c varies with T if debris in the ice is not uniformly distributed;
 88 we return to the consequences of this situation later.

89 Cosmogenic dating models that are used widely to constrain exposure age and
 90 erosion rate of rock surfaces (Lal, 1991) do not adequately describe our system. Although
 91 the ice may be likened as being eroded as it sublimates, the till is a lag that has no analogue
 92 in such models. Here we follow the depth history of each clast, $z = h(T)$, to calculate its
 93 exposure history. Given its depth today, z_0 , we reconstruct h by backtracking (Fig. 1B)—
 94 observing that h is constant after the clast accretes to the till; that the age of accretion, T_A ,
 95 satisfies $\ell(T_A) = z_0$; and that, although h differs from z_0 prior to accretion, the clast,
 96 contained then by the ice, approaches the surface at velocity $S + d\ell/dT$. These
 97 considerations yield
 98

$$\begin{aligned}
 & h(T) = z_0 && \text{for } 0 \leq T \leq T_A, \\
 99 \quad & h(T) = \ell(T) + \int_{T_A}^T S(\xi) d\xi && \text{for } T > T_A, \quad (2)
 \end{aligned}$$

100
 101 in which the integral represents the overlying ice thickness (ξ is the variable of
 102 integration). We distinguish three stages in the clast's exposure history: *inheritance* ($T \geq$
 103 T_{AS}), *preaccretion* ($T_{AS} > T > T_A$), and *postaccretion* ($T_A > T \geq 0$), where T_{AS} is the age of
 104 the till surface ($= T_A$ for $z_0 = 0$; Fig. 1). Inheritance thus comprises nuclide contributions
 105 before the till layer develops. We separate inheritance from preaccretion, because it
 106 includes exposure contributions before the clast was incorporated into the ice, which are
 107 unknown. This uncertainty makes it difficult to determine how the stages partition the
 108 nuclide concentration N measured for a given clast.

109 For a stable cosmogenic nuclide such as ^3He , we model its accumulation rate in
 110 the clast (using Lal's (1991) formulation) as
 111

$$112 \quad -\frac{dN}{dT} = P_0 e^{-\frac{\rho_I}{\Lambda}[h(T) - \ell(T)]_{0+}} e^{-\frac{\rho_S}{\Lambda}(1 - \phi)z_0}, \quad (3)$$

113
 114 where P_0 is the surface production rate, ρ_I is ice density, ρ_S is sediment density, Λ is
 115 absorption mean free path, and $[x]_{0+} = \max(x, 0)$. In equation 3, the first exponential factor
 116 describes shielding of the clast by ice; the second exponential factor describes shielding
 117 of the clast by overlying debris, which remains above the clast after enclosing ice
 118 sublimates away. Equation 3 ignores ^3He production by muon-induced reactions, whose
 119 rate at the surface has not been calibrated but is estimated at $\sim 3\%$ of the corresponding
 120 rate by spallation (Lal, 1987). We expect muon-induced production to dominate at depths
 121 $> 4\text{--}5$ m. Including its effect in our (spallation-only) model leads to a slight increase in the
 122 ^3He accumulated in clasts prior to accretion that lowers the bound $T_{A,\text{max}}$, raises the
 123 bounds S_{min} and $\Delta_{I,\text{min}}$ derived below, and strengthens the conclusions of this paper.

124 Now, the integral of equation 3 from $T = T_{AS}$ to $T = 0$ represents the ^3He
 125 accumulated in the clast since the till layer began forming. We substitute for h from

126 equation 2 and, by replacing z_0 with z , generalize this integral for all clasts. If we include
 127 the inheritance stage, the outcome is an expression for today's depth profile:
 128

$$\begin{aligned}
 N(z) = & N_{\text{inh}}(z) && \text{(inheritance, by } T_{\text{AS}} \text{ years ago)} \\
 & + P_0 e^{-\frac{\rho_s(1-\phi)z}{\Lambda}} \int_{T_A(z)}^{T_{\text{AS}}} \exp\left[-\frac{\rho_l}{\Lambda} \int_{T_A(z)}^T S(\xi) d\xi\right] dT && \text{(preaccretion)} \\
 & + P_0 e^{-\frac{\rho_s(1-\phi)z}{\Lambda}} T_A(z) && \text{(postaccretion),}
 \end{aligned} \tag{4}$$

130
 131 in which we identify each exposure stage and N_{inh} denotes the inherited concentration in
 132 material at depth z today. In a forward simulation $S(T)$ and $c(T)$ are specified, and
 133 equation 4 is evaluated with the accretion age distribution $T_A(z)$ (or $\ell(T)$, its inverse)
 134 found from equation 1.

135 136 THE INVERSE MODEL

137 The challenge is the opposite: to find the sublimation and till-thickness histories
 138 $S(T)$ and $\ell(T)$, given $N(z)$. Equations 4 and 1 cannot be solved for these histories uniquely
 139 because of the extra unknowns N_{inh} and c . In particular, the debris concentration $c(T)$ of
 140 the sublimed ice may differ from c for the relict ice today. The measured profiles also are
 141 discrete. Here we seek constraints instead of solution.

142 We first raise a caveat on the method by Schäfer et al. (2000) that explains also
 143 our apparent reversal of their maximum bound on sublimation rate in this paper. They
 144 assumed a constant rate of sublimation S_c and inheritance-free clasts ($N_{\text{inh}} = 0$). In this
 145 case, the ratio of N for two clasts from the surface and base of the till can be used to find
 146 the initial ice thickness between the clasts, because the overall shielding effect of the ice
 147 as it sublimed is predictable. For the clasts, equation 4 reduces to

$$148 \quad N(0) = P_0 T_{\text{AS}}, \quad N(\ell_0) = P_0 e^{-\frac{\rho_s(1-\phi)\ell_0}{\Lambda}} \int_0^{T_{\text{AS}}} \exp(-\rho_l S_c T / \Lambda) dT, \tag{5}$$

150
 151 where ℓ_0 is the till thickness today, and the ratio of N can be written in the form

$$152 \quad \frac{N(\ell_0)}{N(0) e^{-\rho_s(1-\phi)\ell_0/\Lambda}} = \frac{1}{Z} \int_0^Z e^{-\rho_l \xi / \Lambda} d\xi = \frac{1 - e^{-\rho_l Z / \Lambda}}{\rho_l Z / \Lambda}, \tag{6}$$

154
 155 where $Z = S_c T_{\text{AS}}$ is the sublimed ice thickness in the model. Schäfer et al. (2000) used
 156 equation 6 to determine Z from the end data of a profile, and the sublimation rate from S_c
 157 $= Z / T_{\text{AS}} = P_0 Z / N(0)$. They claimed that in the last step, surface erosion would render T_{AS}
 158 (denominator) a minimum age, making S_c a maximum sublimation rate. The caveat is that
 159 Z (numerator) is not an upper-bound estimate: the actual sublimed ice thickness could
 160 exceed Z if unsteady sublimation (e.g., due to climate change) had violated the

161 assumption that S was constant. Therefore, the value S_c does not constrain sublimation
 162 rates and cannot be used to dismiss the model results of Hindmarsh et al. (1998). (But, as
 163 expected, S_c satisfies our constraint below where we allow for all possible sublimation
 164 histories. For profiles I, II, and III, Marchant et al. [2002] and Schäfer et al. [2000]
 165 obtained $S_c \approx 20, 90, \text{ and } 6 \text{ m} \cdot \text{Myr}^{-1}$, respectively.)

166 In contrast, an approach is now developed to give robust *minimum* mean
 167 sublimation rates (S_{\min}). The crux is to derive, for any pair of clasts in a profile, a *lower*
 168 bound on the original thickness of ice that separated them ($\Delta_{I,\min}$) and an *upper* bound on
 169 the time over which this ice sublimed (t_{\max}). The result $S_{\min} = \Delta_{I,\min}/t_{\max}$ is rigorous.

170

171 **Constraint on Ice Thickness**

172 Suppose the clasts are numbered 1 (lower) and 2 (upper) and have concentrations
 173 N_1 and N_2 , depths z_1 and z_2 , respectively (Fig. 1A). We can constrain their original
 174 separation in the ice (Δ_I) because the concentrations reflect different depth histories. The
 175 clasts' separation today is $\Delta_S = z_1 - z_2$, so the intervening sediment thickness is $(1 - \phi)\Delta_S$.
 176 Given the shielding by this sediment, we can predict what the ratio N_2/N_1 should be, but
 177 the data show that the ratio is always larger, which could only have resulted because of
 178 intervening ice that has disappeared. If we neglect ^3He inheritance before the clasts were
 179 incorporated into the ice, then the minimum intervening ice thickness, $\Delta_{I,\min}$, can be
 180 computed from

181

$$182 \quad e^{[\rho_I \Delta_{I,\min} + \rho_S (1 - \phi) \Delta_S] / \Lambda} = \frac{N_2}{N_1}. \quad (7)$$

183

184 The value $\Delta_{I,\min}$ is the minimum initial ice thickness, because the ice could only have
 185 thinned: for a smaller initial thickness, past ^3He production rates in the clasts would have
 186 been too similar for us to explain the data. We calculate $\Delta_{I,\min}$ from N_1 , N_2 , and Δ_S (Table
 187 1). Equation 7 holds regardless of sublimation rate changes and does not depend on P_0 .
 188 ^3He production by muon-induced reactions, which have large attenuation lengths,
 189 effectively increases Λ used in our model, making $\Delta_{I,\min}$ an underestimate.

190

191 **Constraint on Sublimation Time**

192 Next, we deduce a maximum sublimation time t_{\max} for the ice between clasts 1
 193 and 2. This ice began subliming after clast 2 (the upper clast) accreted to the till and none
 194 of it remains today (Fig. 1A), so the maximum accretion age of clast 2 suffices as our
 195 choice for t_{\max} . For any clast, its maximum accretion age ($T_{A,\max}$) is simply the maximum
 196 duration of its postaccretion stage, which we can calculate by attributing all of its
 197 measured N -value to exposure at its current depth z in the till; thus,

198

$$199 \quad T_A(z) \leq T_{A,\max}(z) = \frac{N(z)}{P_0 e^{-\rho_S (1 - \phi) z / \Lambda}}. \quad (8)$$

200

201 Accordingly we put $t_{\max} = T_{A,\max}(z_2)$. In Table 1, dividing $\Delta_{I,\min}$ by $T_{A,\max}(z_2)$ gives S_{\min} ,
 202 our minimum sublimation rate.

203 The bound t_{\max} cannot be tightened, for we cannot deduce from the profiles the
 204 most recent time at which the lower clast (clast 1) could have joined the till (i.e., a
 205 minimum T_A) without making assumptions. Consequently, for a given depth profile, we
 206 cannot resolve the different sublimation periods for ice that existed between successive
 207 clast pairs. For any two clasts, the time over which S_{\min} is defined (and constrains the
 208 sublimation rate) is fixed by the upper clast—it begins no earlier than the age $T_{A,\max}(z_2)$
 209 and ends at the present, regardless of where in the profile the lower clast is taken. Hence
 210 we pick the lower clast always from the base of the till, to ensure the largest admissible
 211 $\Delta_{I,\min}$ for calculating S_{\min} .

212 DISCUSSION

213 Our results (Table 1) shed new light on the evolution of the ice and overlying till
 214 in Beacon Valley. Mean sublimation rates have not necessarily been low. Profiles I, II,
 215 and III indicate minimum mean rates S_{\min} of ~ 4 , 23, and 2 $\text{m}\cdot\text{Myr}^{-1}$, respectively, within
 216 the past 1.1 Myr, 170 kyr, and 1.6 Myr, causing at least several meters of ice loss at all
 217 three sites. Erosion of the surficial clasts can invalidate these results, but not the higher
 218 S_{\min} values for the more recent past indicated by buried clast pairs.

219 Rapid sublimation (Hindmarsh et al., 1998) could be considered likely, *if* one is
 220 prepared to make assumptions about the ice that sublimed. Its maximum average debris
 221 concentration can be calculated from our results as the ratio of sediment thickness to
 222 minimum ice thickness: $c_{\max} = (1 - \phi)\Delta_S/\Delta_{I,\min}$ (Table 1). c_{\max} is several times c_0 ($\sim 3\%$)
 223 for the relict ice. In contrast, one might expect the ice that sublimed to contain less debris
 224 than the relict ice, if the latter is basal ice from Taylor Glacier, as assumed by Sugden et
 225 al. (1995). Thus our bounds may be overconservative. By assuming ice no dirtier than
 226 today's, i.e., $c(T) \leq c_0$, alternative minimum bounds can be found from $\Delta_{I,\min}^R = (1 -$
 227 $\phi)\Delta_S/c_0$ (for sublimed ice thickness) and $S_{\min}^R = \Delta_{I,\min}^R/T_{A,\max}$ (for sublimation rate).
 228 These bounds indicate mean sublimation rates exceeding ~ 10 – $100 \text{ m}\cdot\text{Myr}^{-1}$ (Table 1),
 229 consistent with an independent estimate of $50 \text{ m}\cdot\text{Myr}^{-1}$ from ^{10}Be analysis of the ice and
 230 of debris within the ice (Stone et al., 2000) in the part of Beacon Valley where profiles I
 231 to III were measured.

232 Equation 8 constitutes a powerful constraint on the till accretion history. On the
 233 depth vs. age plot of Figure 2A, the accretion history $T = T_A(z)$ is confined to the region
 234 right of the line representing the maximum accretion age $T = T_{A,\max}(z)$. Consequently the
 235 line also limits the till thickness: the apparent exposure age of a clast, $T_{A,\max}$ (calculated
 236 on the basis of current shielding), implies that the till was, at that age, no thicker than the
 237 till above the clast today. Prior to $T_{A,\max}$ the clast must have still been in the ice and below
 238 the till. For discrete depth profiles, this constraint takes the form of a staircase (Figs. 2B,
 239 2C) provided that the till had not thinned over time.

240 We stress that, according to Figures 2B and 2C, all but the topmost 20% of till at
 241 the sites measured by Phillips et al. (2000) and Marchant et al. (2002) formed within the
 242 past 310 kyr (profile I) and 43 kyr (profile II). Prior to these times the till was
 243 exceptionally thin: $\leq 14 \text{ cm}$ (profile I) and $\leq 9 \text{ cm}$ (profile II), and by these times there
 244 were relatively old clasts aged 800 ka (I) and 130 ka (II) at the surface. These surficial
 245 clasts have uncertain provenance; unlike subsurface clasts released by ice, they might
 246 have originated via rockfall onto Taylor Glacier. Prior exposure may account for most of
 247 their ^3He concentration, so that they may not be used to infer a minimum age for the ice,
 248

249 which could be as little as several hundred thousand years. Although the old exposure age
 250 of the surficial clasts can be explained in other ways (e.g., the ice that originally separated
 251 them from the next lower clast in the profile was very thick, or sublimed very slowly), we
 252 caution against using them to support the case for ancient ice.

253 An outstanding conundrum is the past relationship between ash and till. The
 254 interpretation advanced by Sugden et al. (1995) is that the ice in Beacon Valley was
 255 already mantled by ~50 cm of till at 8.1 Ma, when ash filled a frost crack, and that the till
 256 has thickened little since. In contrast, our analysis shows that no more than a thin veneer
 257 of till existed prior to 310 ka, and that the bulk of the till has accreted since. The ³He
 258 profiles examined here are not located at the “ash site”, and their differences reflect some
 259 spatial variability in till evolution. Nevertheless, the profiles are close enough spatially
 260 and in stratigraphic context for our interpretation of them to challenge the antiquity of the
 261 till enclosing the ash. Our results show that the ash may not be a reliable stratigraphic
 262 indicator. The case for Miocene ice is likely to remain unsettled until a profile similar to
 263 the ones already discussed is measured at a site containing old ash.

264

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271

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323

324 **FIGURE CAPTIONS**

325 Figure 1. Model of subliming ice and accreting till with no deformation. A: Processes in a
 326 reference frame fixed to the ice. B: Depth vs. age plot shows processes in a reference
 327 frame fixed to till surface $z = 0$. Heavy dashed line denotes till-thickness history ℓ . Solid
 328 arrowed line is depth history h of clast at $z = z_0$ today; sublimation uncovers clast until it
 329 accretes to the till at age T_A , whose value depends on (and is a function of) z_0 .
 330 Trajectories of several other clasts are shown dotted.

331 Figure 2. Constraint on past till thickness using ^3He depth profiles. A: On depth vs. age
 332 plot (right panel), the till-thickness history $z = \ell(T)$ or equivalently the accretion age
 333 distribution $T = T_A(z)$ (dashed line) must lie outside hatched region, to the right of the
 334 boundary $T = T_{A,\text{max}}(z)$. This boundary (solid line), given by data (left panel) via equation
 335 8, indicates the maximum till thickness at a given time. B, C: Application of model in A
 336 to profiles I and II. In these cases the boundary $T_{A,\text{max}}(z)$ is step-like.

337
338
339
340
341TABLE 1. COSMOGENIC ^3He IN CLASTS FROM BEACON VALLEY TILL AND MODEL RESULTS

Data		Results							
z (cm)	N ($\times 10^6$ atoms $\cdot\text{g}^{-1}$)	Clast pair (cm)	Δ_S (cm)	$\Delta_{I,\min}$ (m)	$T_{A,\max}$ (Ma)	S_{\min} (m $\cdot\text{Myr}^{-1}$)	c_{\max} (%)	$\Delta_{I,\min}^R$ (m)	S_{\min}^R (m $\cdot\text{Myr}^{-1}$)
<u>Profile I</u>									
0	612	0–70	70	4.52	1.123	4.02	10.3	15.6	13.9
14	140	14–70	56	2.37	0.310	7.66	15.8	12.4	40.2
21	85	21–70	49	1.69	0.206	8.21	19.3	10.9	52.8
59	28	59–70	11	0.69	0.113	6.10	10.6	2.44	21.7
70	16	70–70	—	—	0.075	—	—	—	—
<u>Profile II</u>									
0	93	0–38	38	3.90	0.171	22.9	6.5	8.44	49.5
9	21	9–38	29	1.62	0.043	37.3	11.9	6.44	148.3
25	8.9	25–38	13	0.54	0.023	23.9	16.0	2.89	126.8
38	5.4	38–38	—	—	0.016	—	—	—	—
<u>Profile III</u>									
0	880	0–70	70	3.44	1.615	2.13	13.6	15.6	9.63
70	44	70–70	—	—	0.205	—	—	—	—

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Note: Symbols: z = depth of clast sample; N = ^3He concentration; Δ_S = clast-pair separation; $\Delta_{I,\min}$ = minimum original interclast ice thickness; $T_{A,\max}$ = maximum accretion age of upper clast of pair; S_{\min} = minimum sublimation rate of interclast ice; $c_{\max} = (1 - \phi)\Delta_S/\Delta_{I,\min}$ = maximum debris concentration of ice that sublimed; $\Delta_{I,\min}^R = (1 - \phi)\Delta_S/c_0$ (see discussion); $S_{\min}^R = \Delta_{I,\min}^R/T_{A,\max}$ (see discussion). Data sources: Phillips et al. (2000) and Marchant et al. (2002) for profiles I and II and Schäfer et al. (2000) for profile III. Deepest clast of each profile is located at the base of till. In the $\Delta_{I,\min}$ column, subtracting two values gives $\Delta_{I,\min}$ for the two clasts appearing on the same row as the values. Model does not correct for the (unknown) sampling position on each clast. Model constants: $\rho_i = 0.9 \text{ g}\cdot\text{cm}^{-3}$, $\rho_s = 3.0 \text{ g}\cdot\text{cm}^{-3}$, $\phi = 1/3$, $\Lambda = 150 \text{ g}\cdot\text{cm}^{-2}$, $c_0 = 0.03$, and (following Marchant et al., 2002) $P_0 = 545 \text{ atoms}\cdot\text{g}^{-1}$ per year.

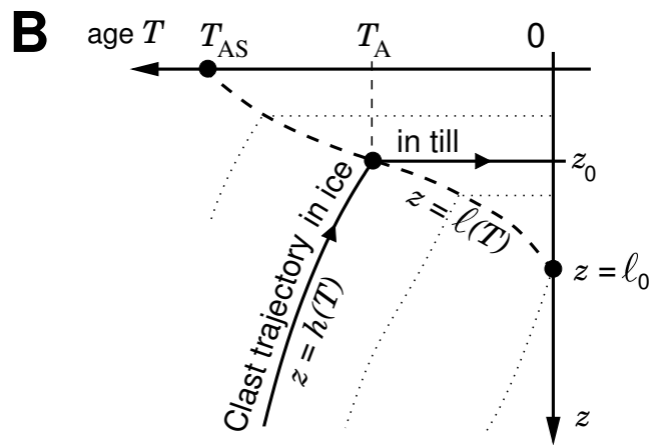
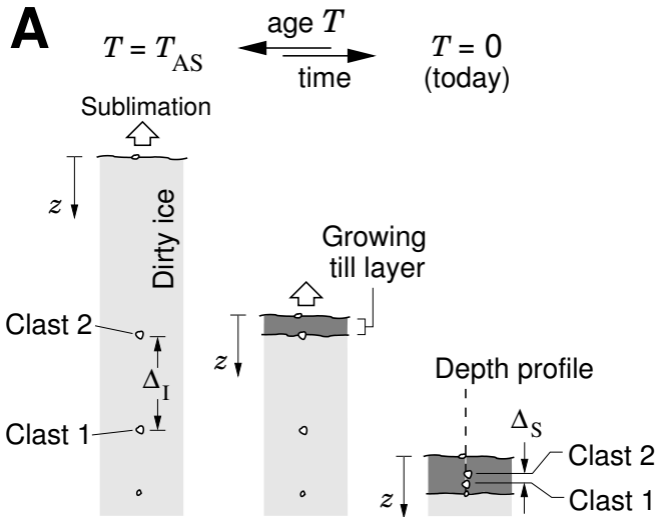


Figure 1 (65%)

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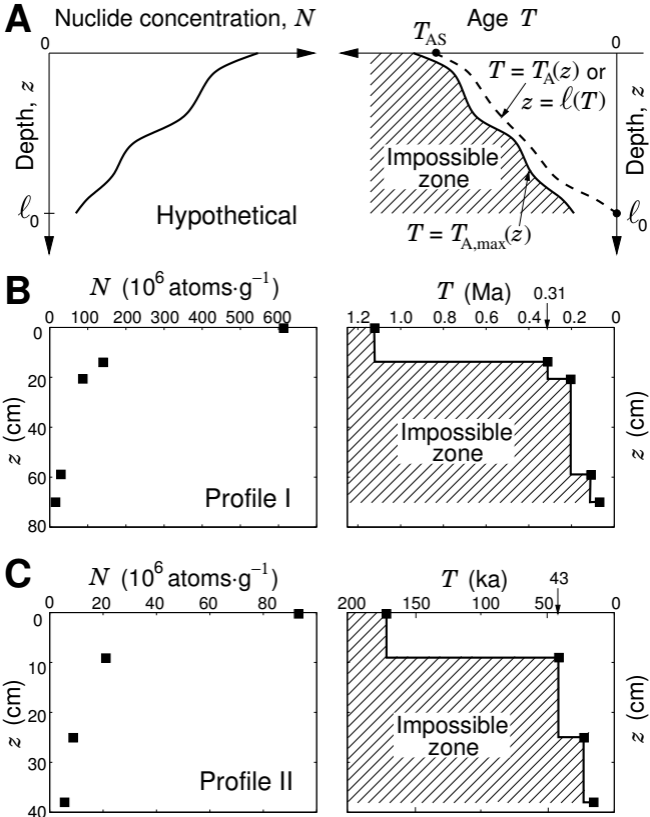


Figure 2 (44.5 %)

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