



Deposited via The University of Sheffield.

White Rose Research Online URL for this paper:

<https://eprints.whiterose.ac.uk/id/eprint/77887/>

---

**Monograph:**

Banks, S.P. and Chanane, B. (1987) On Frequency Response for Nonlinear Systems. Research Report. Acse Report 323 . Dept of Automatic Control and System Engineering. University of Sheffield

---

**Reuse**

Items deposited in White Rose Research Online are protected by copyright, with all rights reserved unless indicated otherwise. They may be downloaded and/or printed for private study, or other acts as permitted by national copyright laws. The publisher or other rights holders may allow further reproduction and re-use of the full text version. This is indicated by the licence information on the White Rose Research Online record for the item.

**Takedown**

If you consider content in White Rose Research Online to be in breach of UK law, please notify us by emailing [eprints@whiterose.ac.uk](mailto:eprints@whiterose.ac.uk) including the URL of the record and the reason for the withdrawal request.

Q 629.8 (S)



ON FREQUENCY RESPONSE FOR NONLINEAR SYSTEMS

by

S. P. Banks and B. Chanane

Department of Control Engineering

University Of Sheffield

Mappin Street

SHEFFIELD, S1 3JD

Research Report No. 323

December 1987

### Abstract

In this paper we define a 'generalized frequency response' for a nonlinear input-output map  $S_{x_0}$  as the mapping  $\mathcal{F} \circ \mathcal{F} \circ S_{x_0} \circ \mathcal{F}^{-1} \circ \mathcal{F}^{-1}$  from  $\ell^2$  to  $\ell^2$  where  $\mathcal{F}$  is the Fourier transform and  $\mathcal{F}$  the usual isomorphism from  $L^2[-\infty, \infty]$  to  $\ell^2$ . Realization results relative to linear and bilinear systems are presented. Also sufficient conditions for  $\mathcal{L}^2$ -stability of bilinear systems are derived.

200081498



equivalent.

The paper is organised as follows: In section 2 we shall consider sufficient conditions for a bilinear system to be  $\varrho^2$ -stable. In section 3 we define a 'GFR' for a nonlinear input-output map. We shall illustrate the 'GFR' for linear and bilinear systems. Finally a realization theory will be presented in section 4. We shall make use in the sequel of the following notation: The Fourier transform  $F$  of a square integrable function  $f$  is defined by

$$F(i\omega) = \frac{1}{\sqrt{2\pi}} \int_0^{\infty} f(t) e^{-i\omega t} dt$$

If  $x$  is an  $n$ -vector and  $A$  is an  $n \times n$ -matrix, we take as compatible norms:

$$\|x\| = \max_{i=1, n} |x_i| \quad \text{and} \quad \|A\| = \max_{i=1, n} \sum_{k=1}^n |a_{i,k}|$$

$\|\cdot\|_1$  and  $\|\cdot\|_{2w}$  are the norms associated with the standard  $L^1$  and  $L^2_w$  spaces; the space of absolutely integrable functions and the space of square integrable functions (with weight  $w$ ) respectively. Whereas,  $\ell^2$  designates the space of square summable sequences.

### 2. $\varrho^2$ -stability of bilinear systems:

In this section we shall present sufficient conditions for the  $\varrho^2$ -stability of bilinear systems. By  $\varrho^2$ -stability we understand that:  $\forall u \in \Omega \subset L^2_w[0, \infty] \longrightarrow y \in L^2_w[0, \infty]$  where  $u$  and  $y$  are respectively the input and the output of the system and  $\Omega$  is defined by the sufficient conditions. The weight function  $w$  is such that  $w(t)e^{-2\alpha t} \in L^1[0, \infty]$  for some  $\alpha \geq 0$ ,  $w(t) \geq 0$  for  $t \geq 0$ .

Consider for simplicity the single input-single output system

described by

$$\begin{cases} \dot{x} = Ax + uNx + Bu & , \quad x(0) = 0 \\ y = Cx \end{cases} \quad (2.1)$$

where A, N, B, C are constant matrices of suitable dimensions.

As is well known [3], the input-output map is given by:

$$\begin{aligned} y(t) = & Ce^{At}x_0 \\ & + \sum_{j \geq 1} \int_0^t \int_0^{\sigma_j} \dots \int_0^{\sigma_2} v_j(t, \sigma_1, \dots, \sigma_j) u(\sigma_1) \dots u(\sigma_j) d\sigma_1 \dots d\sigma_j \\ & + \sum_{j \geq 1} \int_0^t \int_0^{\sigma_j} \dots \int_0^{\sigma_2} w_j(t, \sigma_1, \dots, \sigma_j) u(\sigma_1) \dots u(\sigma_j) d\sigma_1 \dots d\sigma_j \end{aligned} \quad (2.2)$$

where

$$v_1(t, \sigma_1) = Ce^{A(t-\sigma_1)}B$$

$$v_j(t, \sigma_1, \dots, \sigma_j) = Ce^{A(t-\sigma_j)}Ne^{A(\sigma_j-\sigma_{j-1})}N \dots Ne^{A(\sigma_2-\sigma_1)}B \quad , \quad j > 1$$

and

$$w_1(t, \sigma_1) = Ce^{A(t-\sigma_1)}x_0$$

$$w_j(t, \sigma_1, \dots, \sigma_j) = Ce^{A(t-\sigma_j)}Ne^{A(\sigma_j-\sigma_{j-1})}N \dots Ne^{A(\sigma_2-\sigma_1)}Ne^{A\sigma_1}x_0 \quad , \quad j > 1.$$

### Theorem 1 :

A sufficient condition for the system (2.1) to be  $\mathcal{L}^2$ -stable is that the following holds:

- (i) the eigenvalues of A have real parts less than  $-\alpha$  ,
- (ii)  $u$  ,  $e^{\alpha t}u \in L^1[0, \infty]$  .

### Proof :

It follows from (i) as a standard result in the theory of differential equations that there exists a positive constant K such that

$$\|e^{At}\| \leq Ke^{-\alpha t} \quad , \quad t \geq 0 .$$

We have,

$$\begin{aligned}
\int_0^{\infty} w(t)y^2(t)dt \leq & \int_0^{\infty} w(t) \cdot \left[ \|C\| \cdot \|x_0\| \cdot K \cdot e^{-\alpha t} \right. \\
& + \|C\| \cdot \|B\| \cdot K \cdot \int_0^t e^{-(t-\sigma_1)\alpha} |u(\sigma_1)| d\sigma_1 \\
& + \sum_{j \geq 2} \int_0^t \int_0^{\sigma_j} \dots \int_0^{\sigma_2} \|C\| \cdot \|B\| \cdot \|N\|^{j-1} e^{-(t-\sigma_1)\alpha} \\
& \cdot K^j |u(\sigma_1)| \dots |u(\sigma_j)| d\sigma_1 \dots d\sigma_j \\
& + \|C\| \cdot \|x_0\| \cdot K \cdot \int_0^t e^{-(t-\sigma_1)\alpha} |u(\sigma_1)| d\sigma_1 \\
& \left. + \sum_{j \geq 2} \int_0^t \int_0^{\sigma_j} \dots \int_0^{\sigma_2} \|C\| \cdot \|x_0\| \cdot \|N\|^j \cdot K^j \cdot e^{-\alpha t} \right. \\
& \left. \cdot |u(\sigma_1)| \dots |u(\sigma_j)| d\sigma_1 \dots d\sigma_j \right]^2 dt \quad (2.3)
\end{aligned}$$

therefore,

$$\begin{aligned}
\int_0^{\infty} w(t)y^2(t)dt \leq & \int_0^{\infty} w(t)e^{-2\alpha t} dt \cdot \left[ \|C\| \cdot \|x_0\| \cdot K \right. \\
& + \|C\| \cdot \|B\| \cdot K \cdot \int_0^{\infty} e^{\sigma_1 \alpha} |u(\sigma_1)| d\sigma_1 \\
& + \sum_{j \geq 2} \|C\| \cdot \|B\| \cdot \|N\|^{j-1} \cdot K^j \cdot \int_0^{\infty} e^{\sigma_1 \alpha} |u(\sigma_1)| d\sigma_1 \cdot \frac{\|u\|_1^{j-1}}{(j-1)!} \\
& + \|C\| \cdot \|x_0\| \cdot K \cdot \int_0^{\infty} e^{\sigma_1 \alpha} |u(\sigma_1)| d\sigma_1 \\
& \left. + \sum_{j \geq 2} \|C\| \cdot \|x_0\| \cdot \|N\|^j \cdot K^j \cdot \frac{\|u\|_1^{j-2}}{j!} \right]^2 \quad (2.4)
\end{aligned}$$

hence,

$$\begin{aligned}
\|y\|_{2w}^2 \leq & \int_0^{\infty} w(t)e^{-2\alpha t} dt \cdot \|C\|^2 \cdot \left[ K \cdot \|x_0\| + K \cdot \|B\| \cdot \int_0^{\infty} e^{\sigma_1 \alpha} |u(\sigma_1)| d\sigma_1 \right. \\
& + K \cdot \|B\| \cdot \int_0^{\infty} e^{\sigma_1 \alpha} |u(\sigma_1)| d\sigma_1 \cdot (e^{K \cdot \|N\|} \cdot \|u\|_{1-1}) \\
& + K \cdot \|x_0\| \cdot \int_0^{\infty} e^{\sigma_1 \alpha} |u(\sigma_1)| d\sigma_1 \\
& \left. + \|x_0\| \cdot (e^{K \cdot \|N\|} \cdot \|u\|_{1-1} - K \cdot \|N\| \cdot \|u\|_1) \right]^2 < \infty \quad (2.5)
\end{aligned}$$

### 3. Generalized frequency response of Nonlinear input-output

maps:

Consider a system S given in terms of an input-output map

$$S : \mathbb{R}^n \times L_W^2[0, \infty] \longrightarrow L_W^2[0, \infty] \quad (3.1)$$

$$\text{defined by} \quad y(t) = S(x_0, u(\cdot))(t) \quad (3.2)$$

where  $u$  and  $y$  are respectively the input and the output of the system and  $x_0$  is the initial state in some given state-space realization. For each fixed initial state  $x_0$ , we have a map

$$S_{x_0} \triangleq S(x_0, \cdot) : L_W^2[0, \infty] \longrightarrow L_W^2[0, \infty] \quad (3.3)$$

For simplicity, we have assumed scalar input and scalar output.

In a recent paper [1], we introduced the notion of a 'Generalized Frequency Response' by using the natural isomorphism  $\phi$  between  $L_W^2[0, \infty]$  and  $\ell^2$  in the time-domain, and then defining the 'G.F.R.' as the induced map from  $\ell^2$  to  $\ell^2$  such that the diagram

$$\begin{array}{ccc} L_W^2[0, \infty] & \xrightarrow{S_{x_0}} & L_W^2[0, \infty] \\ \phi \downarrow & & \phi \downarrow \\ \ell^2 & \xrightarrow{s_{x_0}} & \ell^2 \end{array} \quad (3.4)$$

commutes.

Alternatively, we can use the natural isomorphism between  $L^2[-\infty, \infty]$  and  $\ell^2$  in the frequency domain (i.e., we operate on the Fourier transforms of the input and output) and then defining the 'G.F.R.' as the induced map from  $\ell^2$  to  $\ell^2$  such that the diagram

$$\begin{array}{ccc}
L^2[0, \infty] & \xrightarrow{S_{x_0}} & L^2[0, \infty] \\
\mathcal{F} \downarrow & & \mathcal{F} \downarrow \\
L^2[-\infty, \infty] & \xrightarrow{\hat{S}_{x_0}} & L^2[-\infty, \infty] \\
\mathcal{F} \downarrow & & \mathcal{F} \downarrow \\
\ell^2 & \xrightarrow{s_{x_0}} & \ell^2
\end{array} \tag{3.5}$$

commutes.

Let  $\{e_j\}_{j \geq 0}$  be a basis of  $L^2[0, \infty]$  and  $E_j = \mathcal{F}\{e_j\}$ ,  $j \geq 0$ . Recalling the fact that the scalar product is invariant under the Fourier transform, we deduce that  $\{E_j\}_{j \geq 0}$  is a basis of  $L^2[-\infty, \infty]$ . Let  $\mathcal{F}$  denotes the usual isomorphism

$$\mathcal{F}: L^2[-\infty, \infty] \longrightarrow \ell^2$$

given by

$$\mathcal{F}(F) = \{F_j\}_{j \geq 0} \tag{3.6}$$

where

$$F \in L^2[-\infty, \infty], \quad F = \sum_{j \geq 0} F_j E_j$$

explicitly  $s_{x_0}$  is given by

$$s_{x_0} = \mathcal{F} \circ \mathcal{F} \circ S_{x_0} \circ \mathcal{F}^{-1} \circ \mathcal{F}^{-1} \tag{3.7}$$

that is

$$s_{x_0}(\{U_k\}_{k \geq 0}) = \{Y_j\}_{j \geq 0} \tag{3.8}$$

where  $U_k = \langle \mathcal{F}(u), E_k \rangle$ ,  $Y_j = \langle \mathcal{F}(S_{x_0}(u)), E_j \rangle$

and  $\langle \cdot, \cdot \rangle$  denotes the scalar product in  $L^2[-\infty, \infty]$ . We have

$$U_k = \langle u, e_k \rangle \text{ and } Y_j = \langle S_{x_0}(u), e_j \rangle$$

Let assume the systems at hand are  $\mathcal{L}^2$ -stable. We shall illustrate the expression (3.7) for the linear and the bilinear input-output maps.

### Example 1: Linear systems

Consider the linear system

$$\begin{cases} \dot{x} = Ax + Bu \\ y = Cx \end{cases}, \quad x(0) = x_0 \tag{3.9}$$

then the input-output map is given by

$$y(t) = Ce^{At}x_0 + \int_0^t Ce^{A(t-\tau)}Bu(\tau)d\tau \quad (3.10)$$

In this case

$$S(x_0, u(\cdot))(t) = g_0(t) + (g*u)(t) \quad (3.11)$$

where  $g_0(t) = Ce^{At}x_0$  ,  $g(t) = Ce^{At}B$

and  $*$  denotes the convolution operator.

Taking the Fourier transform of both sides of (3.10) we obtain

$$Y(i\omega) = G_0(i\omega) + G(i\omega) \cdot U(i\omega) \quad (3.12)$$

Let  $x_0 = 0$  and introduce a basis  $\{E_j\}_{j \geq 0}$  of  $L^2[-\infty, \infty]$ .

$$\text{We obtain, } Y_1 = \sum_{j \geq 0} G_{1j} \cdot U_j, \quad j \geq 0 \quad (3.13)$$

$$\text{where } Y(i\omega) = \sum_{l \geq 0} Y_l \cdot E_l(i\omega), \quad U(i\omega) = \sum_{j \geq 0} U_j \cdot E_j(i\omega)$$

$$G_{1j} = \langle G(i\omega) \cdot E_j(i\omega), E_1(i\omega) \rangle$$

We therefore see that the matrix representation of the linear operator  $s : \ell^2 \longrightarrow \ell^2$  for the linear system above, with respect to the basis  $(1, 0, 0, \dots), (0, 1, 0, \dots), \dots$  of  $\ell^2$  is

$$(G_{1j})_{1, j \geq 0}$$

### Example 2: Bilinear systems

Consider the bilinear system (2.1). Let  $N = N_1 \cdot N_2$  where the dimensions of  $N, N_1, N_2$  are respectively  $n \times n, n \times m, m \times n$ , and

$$\begin{aligned} g_0(t) &= Ce^{At}B, & g_1(t) &= Ce^{At}N_1 \\ g_2(t) &= N_2 e^{At}N_1, & g_3(t) &= N_2 e^{At}B \\ g(t) &= Ce^{At}x_0, & g_4(t) &= N_2 e^{At}x_0 \end{aligned} \quad (3.14)$$

Equation (2.2) can therefore be rewritten as

$$\begin{aligned} y(t) &= g(t) + (g_0 * u)(t) \\ &+ \sum_{j \geq 2} (g_1 * [u(g_2 * [u(\dots u(g_2 * [u(g_2 * [u(g_3 * u)]]) \dots))]])](t) \\ &+ (g * u)(t) \\ &+ \sum_{j \geq 2} (g_1 * [u(g_2 * [u(\dots u(g_2 * [u(g_2 * [u(g_2 * (g_4 \cdot u)]]) \dots))]])](t) \end{aligned} \quad (3.15)$$

where  $g_2$  appears  $(j-2)$  times in the first summation and  $(j-1)$  times in the second.

Taking the Fourier transform of both sides of (3.15) and using its properties we obtain

$$\begin{aligned}
 Y(i\omega) = & G(i\omega) + G_0(i\omega) \cdot U(i\omega) + \sum_{j \geq 2} \frac{1}{\sqrt{2\pi}^{j-1}} (G_1 \cdot [U * (G_2 \cdot [U * (\dots \\
 & \dots U * (G_2 \cdot [U * (G_2 \cdot [U * (G_3 \cdot U)])]) \dots]))](i\omega) + G(i\omega) \cdot U(i\omega) \\
 & + \sum_{j \geq 2} \frac{1}{\sqrt{2\pi}^j} (G_1 \cdot [U * (G_2 \cdot [U * (\dots \\
 & \dots (G_2 \cdot [U * (G_2 \cdot [U * (G_2 \cdot (G_4 * U))])]) \dots]))](i\omega) \quad (3.16)
 \end{aligned}$$

Let  $x_0 = 0$  and introduce a basis  $\{E_j\}_{j \geq 0}$  of  $L^2[-\infty, \infty]$ .

If  $U(i\omega) = \sum_{j \geq 0} U_j E_j(i\omega)$  then (3.16) becomes

$$\begin{aligned}
 Y(i\omega) = & \sum_{k_1 \geq 0} U_{k_1} G_0(i\omega) E_{k_1}(i\omega) \\
 & + \sum_{j \geq 2} \sum_{k_1 \geq 0} \dots \sum_{k_j \geq 0} U_{k_1} \dots U_{k_j} \cdot \frac{1}{\sqrt{2\pi}^{j-1}} (G_1 \cdot [E_{k_j} * (G_2 \cdot [E_{k_{j-1}} \\
 & * (\dots E_{k_4} * (G_2 \cdot [E_{k_3} * (G_2 \cdot [E_{k_2} * (G_3 \cdot E_{k_1})])]) \dots]))](i\omega) \quad (3.17)
 \end{aligned}$$

Therefore,

$$Y_1 = \sum_{j \geq 1} \sum_{k_1 \geq 0} \dots \sum_{k_j \geq 0} V_{jk_1 \dots k_j} \cdot U_{k_1} \dots U_{k_j} \quad (3.18)$$

$$\text{where} \quad V_{1k_1} = \langle G_0(i\omega) E_{k_1}(i\omega), E_1(i\omega) \rangle \quad (3.19)$$

and

$$\begin{aligned}
 V_{jk_1 \dots k_j} = & \frac{1}{\sqrt{2\pi}^{j-1}} \cdot \langle (G_1 \cdot [E_{k_j} * (G_2 \cdot [E_{k_{j-1}} * (\dots \\
 & \dots E_{k_4} * (G_2 \cdot [E_{k_3} * (G_2 \cdot [E_{k_2} * (G_3 \cdot E_{k_1})])]) \dots \\
 & \dots]))](i\omega), E_1(i\omega) \rangle, \text{ for } j > 1 \quad (3.20)
 \end{aligned}$$

Hence, the diagram (3.5) induces the map  $s_{x_0} : \ell^2 \longrightarrow \ell^2$

given by

$$s_{x_0}((U_0, U_1, \dots))_1 = \sum_{j \geq 1} \sum_{k_1 \geq 0} \dots \sum_{k_j \geq 0} V_{jk_1 \dots k_j} \cdot U_{k_1} \dots U_{k_j} \quad (3.21)$$

#### 4. Realization theory:

In this section we shall consider the problem of the

realizability and the state space realization of an analytic map  $s : \mathcal{L}^2 \longrightarrow \mathcal{L}^2$  which defines a 'Generalized Frequency Response'. We shall present conditions under which such  $s$  is realized by a linear or a bilinear system.

#### 4.A: Linear system :

##### Theorem 2 :

A necessary and sufficient condition for a sequence of numbers  $\{G_{1j}\}_{1,j \geq 0}$  to be the 'Generalized Frequency Response' of a linear system with zero initial condition (with respect to a given basis  $\{E_k\}_{k \geq 0}$  of  $L^2[-\infty, \infty]$ ) is that there exists a strictly proper rational function  $G(i\omega)$  such that

$$\sum_{l \geq 0} G_{1l} \cdot E_l(i\omega) = G(i\omega) \cdot E_j(i\omega) \quad (4.1)$$

for all  $j \geq 0$ .  $G(i\omega)$  is then the Fourier transform of the impulse response of the linear system.

##### Proof :

immediate and shall be omitted.

#### 4.B Bilinear systems :

Let  $e_k$ ,  $k \geq 0$  be the Laguerre functions defined by

$$e_k(t) = e^{-t/2} \sum_{m=0}^k \frac{(-1)^m}{m!} \binom{k}{m} t^m$$

They constitute a complete orthonormal basis for  $L^2[0, \infty]$ .

Consider  $E_k$  the Fourier transform of  $e_k$ , we have

$$E_k(i\omega) = \frac{1}{\sqrt{2\pi}} \frac{1}{\frac{1}{2} + i\omega} \left[ \frac{-\frac{1}{2} + i\omega}{\frac{1}{2} + i\omega} \right]^k$$

$\{E_k\}_{k \geq 0}$  is therefore a complete basis for  $L^2[-\infty, \infty]$ .

Remark 1 : The coefficients in the expansion of  $F(i\omega)$  with respect to the basis  $\{E_k\}_{k \geq 0}$  are the coefficients in the Taylor expansion of  $\bar{F}$

$$\bar{F}(z) = \sqrt{2\pi} \frac{1}{1-z} F\left(\frac{1}{2} \frac{1+z}{1-z}\right)$$

at  $z=0$ . So 
$$F_k = \frac{\sqrt{2\pi}}{k!} \left[ \frac{1}{1-z} F\left(\frac{1}{2}, \frac{1+z}{1-z}\right) \right]_{z=0}^{(k)}$$

where  $F(i\omega) = \sum_{k \geq 0} F_k \cdot E_k(i\omega)$ .

**Theorem 3 :**

A necessary and sufficient condition for a sequence of numbers  $\{V_{jk_1 \dots k_j l}\}_{j \geq 1, k_1, \dots, k_j, l \geq 0}$  to be the ' Generalized Frequency Response ' of a bilinear system with zero initial condition (with respect to the given basis  $\{E_k\}_{k \geq 0}$  of  $L^2[-\infty, \infty]$  ) is that there exist four matrices  $G_0(i\omega)$ ,  $G_1(i\omega)$ ,  $G_2(i\omega)$ ,  $G_3(i\omega)$  with dimensions respectively  $1 \times 1$ ,  $1 \times m$ ,  $m \times m$ ,  $m \times 1$  of strictly proper rational functions such that

$$(i) \sum_{l \geq 0} V_{lk_1 l} E_l(i\omega) = G_0(i\omega) \cdot E_{k_1}(i\omega) \tag{4.2}$$

$$(ii) \sum_{l \geq 0} V_{jk_1 \dots k_j l} E_l(i\omega) = \frac{1}{k_j! \dots k_2!} D_j^{k_j} \dots D_2^{k_2} \cdot [G_1(i\omega) \cdot G_2(i\omega - z_j) \cdot G_2(i\omega - z_j - z_{j-1}) \dots \dots G_2(i\omega - z_j - \dots - z_3) \cdot G_3(i\omega - z_j - \dots - z_2) \cdot E_{k_1}(i\omega - z_j - \dots - z_2) \cdot (-\frac{1}{2} + z_2)^{k_2} \dots (-\frac{1}{2} + z_j)^{k_j}] \Big|_{z_2 = \dots = z_j = -\frac{1}{2}}$$

for all  $k_1, \dots, k_j \geq 0$  ,  $j > 1$  (4.3)

where  $D_j^{k_j} = \partial^{k_j} / \partial z_j^{k_j}$  .

**Proof :**

(i) (4.2) can readily be obtained from (3.9) by premultiplying by  $E_1(i\omega)$  and summing over  $l \geq 0$ . Whereas (3.14) yields  $G_0(i\omega) = C(i\omega - A)^{-1} B$ .

(ii) From (3.20) we obtain

$$\sum_{l \geq 0} V_{jk_1 \dots k_j l} E_l(i\omega) = \frac{1}{\sqrt{2\pi}^{j-1}} (G_1 \cdot [E_{k_j} * (G_2 \cdot [E_{k_{j-1}} * (\dots \dots E_{k_4} * (G_2 \cdot [E_{k_3} * (G_2 \cdot [E_{k_2} * (G_3 \cdot E_{k_1})])])])])]) (i\omega)$$

Starting from the inner square bracket and proceeding outward,

we obtain, after a change of variable  $i\Omega = z$  and the use of residue theory,

$$- \{E_{k_2} * (G_3 \cdot E_{k_1})\}(i\omega) = \frac{\sqrt{2\pi}}{k_2!} D_2^{k_2} [G_3(i\omega - z_2) \cdot E_{k_1}(i\omega - z_2) \cdot (-\frac{1}{2} + z_2)^{k_2}] \Big|_{z_2 = -\frac{1}{2}}$$

$$- \{E_{k_3} * (G_2 \cdot [E_{k_2} * (G_3 \cdot E_{k_1})])\}(i\omega) = \int_{-\infty}^{\infty} E_{k_3}(i\Omega) G_2(i\omega - i\Omega) \frac{\sqrt{2\pi}}{k_2!} \cdot D_2^{k_2} [G_3(i\omega - i\Omega - z_2) \cdot E_{k_1}(i\omega - i\Omega - z_2) \cdot (-\frac{1}{2} + z_2)^{k_2}] \Big|_{z_2 = -\frac{1}{2}} d\Omega$$

$$= \frac{\sqrt{2\pi}}{k_2!} (-1) \int_{-i\omega}^{i\omega} E_{k_3}(z_3) G_2(i\omega - z_3) \cdot$$

$$D_2^{k_2} [G_3(i\omega - z_3 - z_2) \cdot E_{k_1}(i\omega - z_3 - z_2) \cdot (-\frac{1}{2} + z_2)^{k_2}] \Big|_{z_2 = -\frac{1}{2}} dz_3$$

$$= \frac{(\sqrt{2\pi})^2}{k_3! k_2!} D_3^{k_3} [G_2(i\omega - z_3) \cdot$$

$$D_2^{k_2} [G_3(i\omega - z_3 - z_2) \cdot E_{k_1}(i\omega - z_3 - z_2) \cdot (-\frac{1}{2} + z_2)^{k_2}] \Big|_{z_2 = -\frac{1}{2}} \cdot (-\frac{1}{2} + z_3)^{k_3}] \Big|_{z_3 = -\frac{1}{2}}$$

$$- \{E_{k_j} * (G_2 \cdot [E_{k_{j-1}} * (\dots E_{k_3} * G_2 \cdot [E_{k_2} * (G_3 \cdot E_{k_1})] \dots)])\}(i\omega) = \frac{\sqrt{2\pi}^{j-1}}{k_j! \dots k_2!} D_j^{k_j} [G_2(i\omega - z_j) \cdot D_{j-1}^{k_{j-1}} [G_2(i\omega - z_j - z_{j-1}) \dots \dots D_3^{k_3} [G_2(i\omega - z_j - \dots - z_3) \cdot D_2^{k_2} [G_3(i\omega - z_j - \dots - z_2) \cdot E_{k_1}(i\omega - z_j - \dots - z_2) \cdot (-\frac{1}{2} + z_2)^{k_2}] \cdot (-\frac{1}{2} + z_3)^{k_3}] \dots \cdot (-\frac{1}{2} + z_{j-1})^{k_{j-1}}] \cdot (-\frac{1}{2} + z_j)^{k_j}] \Big|_{z_1 = \dots = z_j = -\frac{1}{2}}$$

Hence, (4.3).

From (3.14) we get

$$G_0(i\omega) = C(i\omega I - A)^{-1} B \quad , \quad G_1(i\omega) = C(i\omega I - A)^{-1} N_1$$

$$G_2(i\omega) = N_2(i\omega I - A)^{-1} N_1 \quad , \quad G_3(i\omega) = N_2(i\omega I - A)^{-1} B \quad (4.6)$$