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Jahn-Teller Distortions in Transition Metal Compounds, and their Importance in Functional Molecular and Inorganic Materials[†]

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This *tutorial review* discusses the structural and electronic consequences of the Jahn-Teller effect in transition metal complexes, focussing on copper(II) compounds which tend to be the most studied. The nature of a Jahn-Teller distortion in molecular complexes and extended lattices can be manipulated by application of pressure or temperature, by doping a molecule into a host lattice, or simply by molecular

¹⁰ design. Many of these results have been achieved using compounds with a *trans*-[CuX₄Y₂] coordination sphere, which seems to afford copper centres that are particularly sensitive to their environment. Jahn-Teller distortions lead to some unusual phenomena in molecular magnetism, and are important to the functionality of important classes of conducting and superconducting ceramics.

Introduction

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- ¹⁵ In 1937, Hermann Jahn and Edward Teller used symmetry arguments to demonstrate that a non-linear system in a degenerate energy state cannot be stable. It will spontaneously distort itself in some way so that its energy state will split, in such a way as to remove its degeneracy.¹ Such a distortion always has
- ²⁰ the effect of lowering the energy of the system to a small extent, and is thus energetically favourable (Scheme 1). This prediction has been borne out in practise in many areas of chemistry.²



Scheme 1. The two prototypical Jahn-Teller distortions of an octahedral ²⁵ copper(II) complex. Both distortions lower the energies of two electrons in the e_{g} subshell, while raising the energy of only one. Thus, there is a net reduction in electronic energy upon distortion of the complex. All Jahn-Teller distortions are driven by a comparable energy stabilisation.

Jahn-Teller distortions are most often associated with transition metal centres, either in molecular complexes or in ionic lattices. They are important, because Jahn-Teller effects strongly affect the electron distribution within a material or molecule, as well as ³⁵ the arrangement of their atoms. Thus, Jahn-Teller distortions of the metal ions in a ceramic material can have large consequences for its conductivity, magnetism or other physical properties,² examples of which are described below. Jahn-Teller distortions also occur in several types of organic radical,² including some ⁴⁰ alkyl and aryl radicals and ions,³ fulleride anions⁴ and ionised carbon nanotubes.⁵ Although they are phenomenologically distinct, closely related effects (Peierls distortions) are also seen in linear chain structures including solid lattices,⁶ conducting polymers⁶ and chains of atoms on surfaces.⁷

45 In transition metal centres, the splitting of the *d*-energy levels in a ligand field can often lead to degenerate electron configurations, that are subject to Jahn-Teller effects. For example, in octahedral symmetry d^1 , d^2 , d^4 (both spin-states), d^5 (low-spin), d^6 (highspin), d^7 (low-spin) and d^9 transition ions all have orbitally 50 degenerate electron configurations, that should be Jahn-Telleractive. In practise, ions with degenerate occupancy of the e_g subshell nearly always exhibit strong Jahn-Teller distortions, because of the M-L antibonding character of those orbitals. Thus, the Jahn-Teller effect is pronounced in six-coordinate complexes 55 of high-spin d^4 [chromium(II), manganese(III)], low-spin d^7 [cobalt(II)] and d^9 [copper(II)] ions. It is also important in less common oxidation levels of other transition ions with those delectron counts. The other *d*-electron configurations on the above list show degeneracy in the t_{2g} subshell, which makes a much 60 weaker, or zero, contribution to M-L bonding. In those cases, Jahn-Teller distortions are more subtle and can be so weak as to be unobservable using routine characterisation techniques (although not always - see below).⁸

Nearly any other coordination geometry can also be subject to a ⁶⁵ Jahn-Teller distortion in principle, as long as its *d*-electron count leads to a degenerate ground state (see below). A notable exception are linear complexes, which are not treated by the Jahn-Teller theorem.¹ While open-shell two-coordinate compounds

often show bent coordination geometries, these reflect the formation of additional agostic M–L bonding interactions rather than Jahn-Teller phenomena. 9

This article describes the measurement, manipulation and ⁵ consequences of the Jahn-Teller effect in transition metal chemistry. Six-coordinate copper(II) compounds are by far the most common, and the most studied, molecular Jahn-Teller system so they will feature prominently in the discussion.¹⁰⁻¹² Systems containing other ions and coordination geometries will ¹⁰ also be referred to where appropriate, however.²

Definitions of terms

A common point of confusion in copper chemistry in particular, is that not all copper(II) complexes are subject to the Jahn-Teller

- ¹⁵ effect. Complexes with octahedral and tetrahedral geometries are well known to be Jahn-Teller active,¹⁰ while it has recently been pointed out that the rare trigonal prismatic coordination geometry also leads to a strong Jahn-Teller distortion in *d*⁹ complexes.¹³ Notably, Jahn-Teller distortions in the two six-coordinate
- ²⁰ geometries are different. While octahedral copper(II) complexes exhibit the well-known *trans*-elongation along one L–Cu–L axis, trigonal prismatic examples instead undergo an elongation of two *cis* Cu–L bonds (Fig. 1). This can be understood as a consequence of splitting the degenerate e^{i} orbital pair (d_{xz} and
- ²⁵ d_{yz} , Scheme 2), which are oriented towards the donor groups in a trigonal prismatic ligand sphere.¹³ In tetrahedral examples, splitting of the t_2 orbital manifold is reflected in the bond angles about the metal ion, rather than its bond lengths, and leads to a flattening of the coordination sphere towards the square planar ³⁰ limit (termed a "tetragonal distortion").^{10.14}



Scheme 2. Crystal field splittings for the common coordination geometries of copper(II) complexes, grouped according to whether they display Jahn-Teller distortions. The splittings between energy levels are ³⁵ indicative only, and are not drawn to scale.

In contrast, five-coordinate and square planar copper(II) complexes are not Jahn-Teller active, because the singly occupied



Fig. 1 Crystal structures of copper(II) complexes with 2,2'-bipyridyl ligation, with arrows illustrating different Jahn-Teller distortions that occur (Å, °). Top, octahedral [Cu(bipy)₃]²⁺;¹⁶ middle, tetrahedral [Cu(bipy)₂]²⁺;¹⁷ bottom, a trigonal prismatic complex of a *tris*-bipyridyl 45 tripodal ligand.¹⁸ An undistorted [Cu(bipy)₂]²⁺ centre would show an inter-ligand dihedral angle of 90°.

d-orbital in those geometries is not degenerate (Scheme 2). In ⁵⁰ particular, the most common coordination geometry for copper(II) complexes is square-pyramidal, in which the apical Cu–L bond is longer than the basal ones by up to 0.5 Å. The lengthened apical bond is *not* a Jahn-Teller elongation in this case. Rather, it is simply a consequence of the double electron ⁵⁵ occupancy of the antibonding a_1 (d_{z^2}) orbital, and the single occupancy of b_1 ($d_{x^2-y^2}$), leading to increased antibonding electron density along the apical Cu–L axis.¹⁵ Second-order pseudo-Jahn-Teller effects (see below) can occur in five-coordinate complexes, however, leading to more subtle structural changes.

⁶⁰ While Jahn-Teller elongations and compressions are both theoretically possible in octahedral complexes, in practise Jahn-Teller elongations are more common by far. This is particularly



Scheme 3. The pseudo-Jahn-Teller distortions commonly observed in heteroleptic six-coordinate copper(II) complexes. The splittings between energy levels are indicative only, and are not drawn to scale. The b_1 and b_2 levels in the $C_{2\nu}$ symmetry are not strictly degenerate, but will be very close in energy.

true in copper(II) chemistry, where genuine, experimentally verified Jahn-Teller compressions are almost unheard of.¹² This reflects a small additional stabilisation of $3d_{z^2}$ by mixing with the metal 4*s* orbital, which has the same a_{1g} symmetry as $3d_{z^2}$ in the

¹⁰ $O_{\rm h}$ and $D_{4\rm h}$ point groups. Jahn-Teller compressions are sometimes seen in complexes of other metal ions, however, notably high-spin manganese(III).¹⁹

A further complication, is that six-coordinate copper(II) complexes with cis- or trans-[CuX₂L₄] geometries can also

- ¹⁵ exhibit pronounced structural elongations along one of their degenerate L–Cu–X or L–Cu–L axes (Scheme 3). That would not be predicted from the arguments in the previous paragraph, since the unpaired *d*-electron does not occupy a degenerate *d* orbital in the C_{2v} or D_{4h} point groups associated with these geometries
- ²⁰ (Scheme 3). These distortions arise from vibronic coupling of the non- degenerate ground state with excited *d*-energy states, and are termed "pseudo-Jahn-Teller" effects.¹¹ Pseudo-Jahn-Teller distortions can also arise in five-coordinate copper(II) centres (see above), leading to more complicated structural changes.¹¹
- ²⁵ Although they are different in origin, Jahn-Teller- and pseudo-Jahn-Teller-distortions have similar consequences for the structural and electronic properties of a complex. They are therefore considered interchangeably in this article.
- A final grey area in the literature, is the description of ³⁰ coordination geometries of complexes with long axial or apical M...L contacts. This is a problem in copper(II) chemistry, since tetragonal copper(II) complexes with a square plane of ligand donors and one or two axial Cu...L interactions of 2.1-2.9 Å are very common.¹¹ Considering the covalent and van der Waals radii
- ³⁵ of copper (1.4 Å) and the common N/O/F donor atoms (1.35-1.5 Å),²⁰ an axial Cu–L distance of up to 2.4 Å might be considered as a genuine Cu–L bond; 2.4-2.8 Å represents a weaker, secondary Cu...L interaction, predominantly electrostatic in nature; and Cu...L > 2.8 Å is a van der Waals contact only.
- ⁴⁰ Corresponding Cu–L distances for P/S/Cl donors would be <2.8 Å (bond), 2.8-3.2 Å (secondary interaction) and >3.2 Å (van der Waals contact).

The above criteria represent a slightly arbitrary division of a continuum of bonding interactions. However, only the shortest ⁴⁵ interactions in the above list should probably be considered when

assigning a coordination geometry to a complex. For example, this has consequences for the tripodal ligand complex at the bottom of Fig. 1. Although described in the literature (and the Figure) as a Jahn-Teller distorted trigonal prismatic centre,^{13,18} ⁵⁰ the very long sixth Cu...N distance implies that a five-coordinate description may be equally valid. A multitude of other copper(II) crystal structures in the literature show similar ambiguities, too.

Dynamic and disordered Jahn-Teller distortions

⁵⁵ A straightforward example is provided by the salt [3-ClC₆H₄NH₃]₈CuCl₁₀,²¹ which contains centrosymmetric [CuCl₆]⁴⁻ octahedra with three unique Cu–Cl bonds (Fig. 2).²² These data show only minor changes on cooling to 156 K, and indicate a Jahn-Teller compressed copper(II) centre. This impression is ⁶⁰ reinforced by its powder EPR spectrum, which shows a pattern of *g*-values ($g_{\perp} = 2.20$, $g_{\parallel} = 2.04$).²³ At first glance, that is consistent with an unpaired electron in a d_{z^2} orbital on the copper ion, which would lead to $g_{\perp} > g_{\parallel} = 2.0023$,²⁴ and again implies a Jahn-Teller compressed structure (Scheme 1). However, the EPR g_{\parallel} value ⁶⁵ (2.04) is higher than would be expected for a $\{d_{z^2}\}^1$ configuration (2.00). Moreover, the crystallographic thermal ellipsoids of the chlorine atoms associated with the two longer Cu–Cl distances are significantly elongated, along the directions of the Cu–Cl bonds (Fig. 2). Both observations indicate that this interpretation ⁷⁰ may be erroneous. In fact, the copper ion adopts a Jahn-Teller



Fig. 2 Crystal structure of the $[CuCl_6]^{4-}$ ion in $[3-ClC_6H_4NH_3]_8CuCl_{10}$ at 295 K, showing the Cu–Cl bond lengths (Å). The thermal ellipsoids of Cl(2), Cl(3) and their symmetry equivalents indicate unresolved disorder, 75 involving expansion and/or contraction of those Cu–Cl bonds. Symmetry code (i): -x, -y, -z.²²

elongation, that is disordered between the Cl(2)-Cu-Cl (2^i) and Cl(3)-Cu-Cl (3^i) axes (Fig. 2).

Such ambiguities can be resolved by EXAFS, whose short timescale yields true bond distances that are uncontaminated by

- ⁵ disorder. In this case, an EXAFS study yielded a rhombic coordination sphere with Cu–Cl = 2.276(14), 2.377(18) and 2.83(5) Å at 10 K. The average of the two higher values, 2.60 Å, agrees perfectly with the two longer Cu–Cl distances from the crystallographic data above (Fig. 2).²⁵ The anomalous EPR data
- ¹⁰ reflect the fact that g_1 and g_2 are averaged in the spectrum because of the disorder. The "true" *g*-values of the copper centres were obtained by replacing 80 % of the copper in the material with diamagnetic cadmium(II) ions.²⁶ The remaining copper ions in [3-ClC₆H₄NH₃]₈Cd_{0.8}Cu_{0.2}Cl₁₀ are well-separated and
- ¹⁵ magnetically isolated, giving a higher resolution single-ion EPR spectrum with $g_1 = 2.30$, $g_2 = 2.09$, $g_3 = 2.04$. That maps perfectly onto the *g*-pattern expected for a Jahn-Teller-elongated, $\{d_{x^2-y^2}\}^1$ copper(II) centre $(g_1 > g_2 \ge g_3 > 2.0023^{24})$, while the true g_1 and g_2 values agree well with the averaged g_{\perp} value (2.20) from the ²⁰ initial study.²³

A survey of copper(II) complexes that are known to show Jahn-Teller disorder is given in the ESI[†]. Most of these involve polyatomic organic ligands, which complicates the interpretation of the crystal structures. The size and direction of the ligand

- ²⁵ donor atom thermal ellipsoids (Fig. 2) are not a good indicator of masked disorder in that case, since they contain contributions from every covalent bond those atoms are involved in (not just the Cu–L bond). That issue can be resolved by a TLS analysis, which deconvolutes the thermal ellipsoids to provide a measure ³⁰ of the atomic displacements along each bond in the molecule.^{12,27}
- TLS analyses are easily obtained using the program PLATON.²⁸ The Jahn-Teller disorder in [3-ClC₆H₄NH₃]₈CuCl₁₀ is static, and invariant with temperature. However, in many cases such disorder is dynamic and reflects fluxionality of the Jahn-Teller
- ³⁵ distortion in the molecule. This leads to temperature-dependent crystallographic and EPR behaviour, which show apparently "averaged" data at higher temperatures but where the true structure and *g*-values are slowly frozen out as the temperature is lowered.^{12,27} Analysis of the temperature dependence of Jahn-
- ⁴⁰ Teller disorder can yield the relative energies of the different orientations of the disorder axis, which depend on the steric properties of the crystal lattice as much as on the individual molecule (Fig. 3). The analysis in Fig. 3 was performed by modelling the metal/ligand vibrational potentials in isolated
- ⁴⁵ molecules,¹⁰ but more sophisticated treatments based on Monte Carlo simulations can reproduce cooperative Jahn-Teller fluxionality in an extended crystalline lattice.²⁹

Where Jahn-Teller disorder exists, placing a crystal under uniaxial stress can cause the Jahn-Teller elongation to order in an

- $_{50}$ orientation parallel to the applied tension. This has been demonstrated in salts containing the $\left[Cu(NO_2)_6\right]^{4-}$ ion. 31
- Such detailed studies of Jahn-Teller fluxionality have mostly been carried out using copper(II) complexes.^{10,27} However, comparable Jahn-Teller disorder has also been identified in six-
- ⁵⁵ coordinate complexes of other metal ions including vanadium(III),³² cobalt(II),³³ nickel(III),³⁴ rhodium(II),³⁵ and silver(II).³⁶



⁶⁰ **Fig. 3** Left: temperature dependence of the Cu–N bond lengths in two unique molecules in the crystal structure of $[CuTp_2]$ (Tp⁻ = hydrido-*tris*pyrazolylborate). The lines show the fit of the data to a model based on the electronic and vibrational energies of the copper coordination sphere. Right: the results of that calculation, plotted as the electronic and ⁶⁵ vibrational energies of the molecules in different Jahn-Teller states. The minima in the plots correspond to orientation of the Jahn-Teller distortion along the three different N–Cu–N axes in the molecules, with N(1)– Cu(1)–N(1ⁱⁱ) being the most stable. Symmetry code (ii): 1–*x*, 1–*y*, 1–*z*.³⁰

Manipulating Jahn-Teller distortions under pressure

⁷⁰ Most compounds of copper(II) behave similarly to those of other transition metal ions under hydrostatic pressure, in that a general contraction of M–L bonds is observed as the pressure is increased.³⁷⁻³⁹ This bond shortening tends to be more pronounced for the longer (weaker) Cu–L bonds in a Jahn-Teller elongated ⁷⁵ copper(II) centre, which can lead to a phase change under pressure,³⁹ or to the formation of new intermolecular Cu–L interactions as the atoms are forced together.⁴⁰⁻⁴² No more notable structural changes are usually observed.

However, two copper(II) compounds are known where application of hydrostatic pressure leads to rotation of the Jahn-Teller elongation axis about the metal ion, from one L–Cu–L axis to another. First to be discovered was the deuterated Tutton's salt,

[ND₄]₂[Cu(OD₂)₆][SO₄]₂. The Jahn-Teller elongation of the hexaaqua copper(II) centres is disordered between different O-Cu-O axes at above ca. 150 K, but becomes crystallographically ordered in one position below that temperature. Interestingly, 5 however, the orientation of the ordered elongation is different in the protio material [NH₄]₂[Cu(OH₂)₆][SO₄]₂, and deuterated $[ND_4]_2[Cu(OD_2)_6][SO_4]_2$.⁴³ The deuterated crystal reverts to the protio phase under pressure, so that the Jahn-Teller elongation axis switches from one O-Cu-O axis to another (Fig. 4). 43,44 The

- 10 critical pressure for this phase transition is ca. 500 bar (0.05 GPa[‡]) at room temperature, is temperature dependent, and shows a pressure hysteresis when cycled close to room temperature.⁴⁵ Hence, the high-pressure phase of $[ND_4]_2[Cu(OD_2)_6][SO_4]_2$ can also be obtained at 1 bar[‡], where it is metastable, if the crystal is
- 15 kept below 297 K when the pressure is released.⁴⁶ A pressureinduced Jahn-Teller reorientation also occurs in the chromium analogue $[ND_4]_2[Cr(OD_2)_6][SO_4]_2,^{47}$ but not in $M_2[Cu(OH_2)_6][SO_4]_2$ (M⁺ = NH₄⁺, K⁺) or Rb₂[Cr(OD₂)₆][SO₄]₂ which behave normally under pressure.44,46,47



Fig. 4 Neutron diffraction structure of the $[Cu(OD_2)_6]^{2+}$ ion in [ND₄]₂[Cu(OD₂)₆][SO₄]₂ at 15 K, with the Cu-O bond lengths for the deuterated and protio compounds at the same temperature, at 1 bar and [in square brackets] at 1.5 kbar[±] (Å).^{43,46} Symmetry code (i): -x, -y, -z.

- ²⁵ The other example is *catena*-[CuF₂(OH₂)₂(μ -pyrazine)], which crystallises as a 1D coordination polymer, with the co-parallel chains being linked into three dimensions by O-H...F hydrogen bonding (Fig. 5).⁴⁸ The centrosymmetric copper ions have a trans, trans, trans disposition of fluoro, aqua and pyrazine N-
- 30 donor ligands, with a pronounced pseudo-Jahn-Teller elongation lying along the N-Cu-N axis at ambient pressure between 20 and 298 K. As the pressure is increased at 298 K, the compound undergoes two abrupt structural transitions at 0.9 GPa and 3.1 GPa. These were originally proposed to involve the consecutive 35 reorientation of the Jahn-Teller elongation along the O-Cu-O and
- F-Cu-F directions.49 A new, more detailed crystallographic study has partially borne
 - out that proposal, in that the transition at 0.9 GPa does indeed involve migration of the Jahn-Teller elongation axis to the O-
- 40 Cu-O axis, as described (Fig. 5).42 This is reflected in the material's magnetic response, which becomes more strongly antiferromagnetic in the intermediate-pressure phase. That is caused by reorientation of the copper $d_{x^2-y^2}$ magnetic orbital from being perpendicular to, to parallel with, the direction of the

⁴⁵ coordination polymer chain.⁴⁹ However, single crystal structures

above 3.3 GPa instead showed a new phase with a tripled unit cell, in which two out of every three coordination polymer chains become covalently linked through bridging fluoro ligands. Differences in the structural data for the highest pressure phase 50 between this study and the earlier one were noted, however, and have still to be reconciled at the time of writing.⁴²



Fig. 5 Crystal structure of *catena*-[CuF₂(OH₂)₂(μ -pyrazine)] at ambient pressure, at 150 K.48 Bond lengths are shown at ambient pressure, and [in 55 square brackets] at 1.2 GPa (Å).⁴² The O-H...F hydrogen bonds shown link the co-parallel chains in the lattice into a 3D network.

A comparable, temperature-dependent switching of a Jahn-Teller elongation axis, which leads to a magnetic transition in a family of copper(II) coordination polymers, is described below.

- 60 A different type of pressure-induced switching occurs in Rb₂CuCl₄·2H₂O, which contains layers of discrete trans- $[CuCl_4(OH_2)_2]^{2-}$ octahedra linked by O–H...Cl hydrogen bonds, with a pseudo-Jahn-Teller elongation along one Cl-Cu-Cl axis. Consistent with this, a copper EXAFS study showed three distinct $_{65}$ scatterers under ambient conditions at distances Cu–O = 2.00(5), and Cu-Cl = 2.25(5) and 2.85(5) Å. Application of hydrostatic pressure at first showed a typical gradual contraction of all these distances, before an abrupt transition at 16 GPa. Above this pressure, the Cu-O distance is unchanged at 1.96 Å, but all four
- 70 chlorine scatterers in the copper EXAFS experiment are now equivalent at Cu-Cl = 2.25 Å. This indicates suppression, rather than reorientation, of the pseudo-Jahn-Teller distortion in the complex under pressure.50

Comparable behaviour had been proposed for some layered 75 perovskites of type M_2CuCl_4 (M⁺ = an alkali metal or primary ammonium ion). These contain layers of *trans*- $[CuCl_2(\mu-Cl)_4]^{2-}$ octahedra, formed from pseudo-Jahn-Teller copper ions linked by alternating Cu...Cl–Cu bridges. Application of pressure leads to significant changes to their magnetic response, and vibrational

- ⁵ and UV/vis spectra, which has been interpreted as indicating a disappearance (quenching) of the Jahn-Teller distortion when $M^+ = Rb^+$ or $EtNH_3^+$.⁵¹ More recent structural studies under pressure, from diffraction or EXAFS methods, have not always been consistent with that conclusion however, and have shown that
- ¹⁰ canting of the $[CuCl_2(\mu-Cl)_4]^{2^-}$ layers can account for at least some of the magnetic and spectroscopic observations while retaining a pseudo-Jahn-Teller elongation at the copper centres. Currently it seems unlikely that Jahn-Teller quenching occurs in M₂CuCl₄ materials, at least at pressures below *ca*. 16 GPa.⁵¹
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Manipulating Jahn-Teller distortions by molecular design

- Some years ago, we showed that the EPR spectrum of $[Cu(L^{Mes})_2][CIO_4]_2$ (Mes = 2,4,6-trimethylphenyl; Scheme 4) is unprecedented for a six-coordinate copper(II) complex, and ²⁰ clearly indicates a $\{d_{z^2}\}^1$ electron configuration at the copper ion (Fig. 6).⁵² The more usual disordered Jahn-Teller-elongated $\{d_{x^2-y^2}\}^1$ structure could be ruled out for two reasons. First, the same
- spectrum (with minor differences) is observed in solution as well as in the solid state, ruling it out as a crystallographic artifact. ²⁵ Second, the observed values $g_{\perp} = 2.20$ and $g_{\parallel} = 2.01$ (in frozen MeCN solution) agree well with the $g_{\perp} > g_{\parallel} = 2.0023$ pattern expected for a $\{d_{z^2}\}^1$ radical.²⁴ In contrast, disordered Jahn-Teller elongated species with such an "inverse" pattern of *g*-values usually show higher-than-expected values of g_{\parallel} (see above).^{12,27}



Fig. 6 *Q*-band powder EPR spectra at 10 K of two complexes related to those in Scheme 4, with: a) structure A (without Jahn-Teller disorder); and b) structure B (Scheme 4).¹² The four-line splitting in each spectrum arises from hyperfine coupling to the Cu nucleus (63,65 Cu, $I = ^{3}/_{2}$).

³⁵ The $\{d_{z^2}\}^1$ electronic structure of $[Cu(L^{Mes})_2]^{2+}$ corresponds to structure B in Scheme 4, with four long Cu–N{pyrazole} distances.⁵² Importantly this is not a Jahn-Teller compression,

because the e_g *d*-orbitals are not degenerate in the idealised D_{2d} symmetry of these complexes.¹² Rather, it corresponds to a ⁴⁰ suppression of the expected pseudo-Jahn-Teller elongation (structure A, Scheme 4). Similarly unusual EPR spectra are shown by salts of $[Cu(L^{Pr})_2]^{2+}$ and $[Cu(L^{Ph})_2]^{2+}$ in the solid and solution phases,¹² but not by the parent complex $[Cu(bpp)_2]^{2+}$ which behaves like a typical Jahn-Teller elongated complex with ⁴⁵ this ligand type.⁵² Hence, the unusual electronic structures of $[Cu(L^R)_2]^{2+}$ (R \neq H) are probably sterically driven. The bulky distal substituents prevent any contraction of the Cu–N{pyrazole} bonds, forcing them all to be comparably long as observed.



⁵⁰ Scheme 4 The structures and ground-state configurations adopted by $[Cu(bp)_2]^{2+}$ (R = H; structure A) and its derivatives $[Cu(L^R)_2]^{2+}$ (R = *i*Pr, Ph or Mes; structure B). Short and long Cu–N bonds are indicated by solid and dashed lines, respectively].

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- Structure B can also be imposed on other related copper(II) complexes of linear tridentate ligands by steric hindrance (as above),⁵³ by ligand conformational strain,⁵³ and by tuning the basicities of the ligand donor atoms.⁵⁴ In the latter case, structure 60 B is obtained in the absence of steric hindrance, if the distal ligand-donor groups are less basic than the central one. In the latter regard, we have also unexpectedly observed that $[Cu(bpp)_2]^{2+}$ itself exhibits structure A as a pure compound, but transforms to structure B when doped into a host lattice, in $_{65}$ [M_{0.95}Cu_{0.05}(bpp)₂][BF₄]₂ (M²⁺ = Fe²⁺ or Zn²⁺).⁵⁵ The lower basicity of the pyrazole donor groups in $[Cu(bpp)_2]^{2+}$, compared to the central pyridyl donor, apparently makes structures A and B close in energy. The complex is therefore forced to adopt structure B in a rigid host lattice, that cannot undergo the local 70 structural rearrangements required to accommodate a dopant site with structure A. Comparable phenomena have been seen in a small number of inorganic materials containing copper(II) dopant ions, like K₂Cu₁Zn_{1-r}F₄ which also adopts a $\{d_{2}\}^{1}$ electronic structure when x is small.¹²
- ⁷⁵ There have been two recent observations of copper(II) complexes undergoing a temperature-dependent transition between structures A and B. In one example, a transition from structure B towards structure A was observed on cooling in a complex with a static Jahn-Teller structure, which was associated with the ⁸⁰ freezing out of anion disorder in the structure.⁵⁴ In the other, unusual EPR behaviour in a fluxional Jahn-Teller crystal implied that a fraction of the sample adopts a $\{d_{z2}\}^1$ configuration as the temperature is raised.⁵⁶ That situation is analogous to Fig. 3, but where one of the minima in the Jahn-Teller potential surface ⁸⁵ corresponds to an axial compression.

Structure B in molecular complexes has only been found in compounds with D_{2d} symmetry, like those in Scheme 4. Notably, these are stereochemically analogous to the *trans*-[CuCl₄(OH₂)₂]²⁻ ion in Rb₂CuCl₄·2H₂O, which undergoes Jahn-⁵ Teller quenching under pressure (see above).⁵⁰ Moreover, most of

the copper(II) sites in inorganic materials that adopt $\{d_{z^2}\}^1$ ground states also have *trans*-[CuX₄Y₂] coordination geometries.^{12,57} The electronic structure of copper(II) ions in D_{2d} symmetry seems to be unusually sensitive to their local environment.

Exceptions to the rule

As always, there are rare exceptions to the rules discussed in the previous sections. Jahn-Teller distortions in six-coordinate complexes with degeneracy in the t_{2g} orbital manifold are usually 15 small, because of the weak influence of those electrons on M–L

- σ -bonding.⁸ However, a few such complexes containing tridentate ligands related to terpyridine exhibit strong angular structural distortions away from the idealised D_{2d} symmetry (Fig. 7).^{58,59} Calculations have implied that this unusual
- ²⁰ stereochemistry is promoted by ligands imposing a narrow *cis*bite angle on the metal.⁶⁰ For example, the iron(II) compound in Fig. 7 has an intra-chelate *cis*-N–Fe–N angle of 72°, much lower than the ideal value of 90°. Other unexpectedly strong Jahn-Teller distortions of t_{2g} -degenerate ions are also occasionally seen in ²⁵ oxide materials.⁶¹

Cu-N1 = 1.963(2) Cu-N2 = 2.088(2) Cu-N3 = 2.072(2) Cu-N4 = 2.037(2) N4 C Cu-N5 = 2.286(2) Cu-N6 = 2.261(2) N1-Cu-N4 = 176.5(1) Inter-ligand dihedral angle = 88.4 Fe-N1 = 2.172(1) Fe-N2 = 2.218(1) Fe-N3 = 2.199(1) Fe-N4 = 2.172(1) Fe-N5 = 2.218(1) Fe-N6 = 2.199(1) N1-Fe-N4 = 154.2(1)Inter-ligand dihedral angle = 62.6

³⁰ **Fig. 7** Crystal structures of $[Cu(bpp)_2][BF_4]_2$ (top)^{52,62} and high-spin $[Fe(bpp)_2][PF_6]_2$ (bottom; bpp = 2,6-di{pyrazol-1-yl}pyridine),⁶⁰ showing their different Jahn-Teller distortions (Å, °). The copper structure is one of the three unique molecules in the low temperature polymorph of that compound.⁶²

Although the vast majority of square planar metal ions are lowspin, high-spin examples are known in a handful of iron oxide materials like SrFeO₂. The first high-spin square-planar complexes have recently been obtained, in salts of $[FeL_2]^{2-}$ (LH₂ ₄₀ = *meso*-3,4-dihydroxy-oxolane). A high-spin d⁶ ion with D_{4h}

35

symmetry is Jahn-Teller active. Although the crystallographic distortion in the compound is small, calculations imply that it has an important energetic role in stabilising the high-spin square planar site by reducing Fe–O covalency.⁶³

Jahn-Teller Effects in Molecular Magnetism

A bulk magnet retains its magnetisation when removed from a magnetic field. This requires strong coupling interactions between the electron spins in the material, whether they are ⁵⁰ localised electrons in an ionic or molecular solid, or itinerant electrons in a metal. Single molecule magnets (SMMs) are compounds whose individual molecules also retain magnetisation outside an applied field, albeit only at very low temperatures of <5 K.⁶⁴ They are of great current interest for quantum ⁵⁵ information storage applications,⁶⁵ and also for potential uses as magnetic refrigerants.⁶⁶

The best-known SMMs are polymetallic complexes with highspin magnetic ground states, arising from ferromagnetic coupling of the individual metal ions in the molecule. The effect also requires a strong magnetic anisotropy in the molecule; that is, its magnetic moment must be different in different directions. Manganese(III) complexes meet the latter criterion particularly well, because the Jahn-Teller elongation exhibited by sixcoordinate, high-spin d^4 manganese(III) ions lead to an anisotropic distribution of *d*-electron spins.

The most studied SMMs were also the first to be discovered, "Mn₁₂" the namely complex [Mn₁₂O₁₂(O₂CMe)₁₆(H₂O)₄]·2MeCO₂H·4H₂O and its derivatives (Fig. 8).⁶⁷ The acetate ligands can easily be replaced with other 70 carboxylates or different oxyanion ligands, leading to a large family of compounds that all show SMM behaviour. Their effectiveness as SMMs stems from their disk-like molecular structures, comprising a central [Mn^{IV}₄O₄] heterocubane surrounded by a belt of eight Mn^{III} ions (Fig. 8). The Jahn-Teller 75 elongations of the Mn^{III} centres are aligned approximately perpendicular to the plane of the disk, imparting a highly anisotropic magnetic structure to the molecule as a whole. Unusually, Mn₁₂ derivatives can adopt one of two different isomeric structures with different SMM properties. This is ⁸⁰ characterised by the effective energy barrier to the reorientation of their magnetisation in an alternating magnetic field (U_{eff}) , which is ca. 60 K in one isomer (the high-temperature or slowrelaxing form) and ca. 40 K in the other (the low-temperature or fast-relaxing form).⁶⁷ The isomers differ in the orientation of the ⁸⁵ Jahn-Teller elongation in one Mn^{III} ion, which is perpendicular to the plane of the disk in the slow-relaxing isomer, but lies within the disk in the fast-relaxing form (Fig. 8). This makes the magnetic structure of the fast-relaxing molecule less anisotropic, thus lowering $U_{\rm eff}$ and reducing its effectiveness as an SMM.

⁹⁰ In a small number of cases, both isomers of the same Mn₁₂ derivative can be obtained in pure form.⁶⁸ In one such example, [Mn₁₂O₁₂(O₂CCH₂*t*Bu)₁₆(H₂O)₄]·CH₂Cl₂·MeNO₂ (Fig. 8), placing crystals of the fast-relaxing isomer under hydrostatic pressure converts it into the slow-relaxing form. This conversion occurs ⁹⁵ gradually as the pressure is raised, from ambient to 1.4 GPa.⁶⁹

A very different phenomenon is shown by some copper(II) complexes of nitronyl nitroxide (NITR) radical ligands. Adducts of type $[Cu(hfac)_2 L^R]$ (hfacH = 1,1,1,5,5,5-hexafluoro-



Fig. 8 The fast-relaxing (top) and slow-relaxing (bottom) isomers in differently solvated crystals of $[Mn_{12}O_{12}(O_2CCH_2/Bu)_{16}(H_2O)_4]$. Methyl groups and H atoms have been omitted for clarity. Colour code: C, white; 5 Mn^{II} , green; Mn^{IV}, orange; O, red. The Jahn-Teller-elongated Mn^{III}–O bonds are emphasised, and indicated in pale green. The isomers differ in the orientation of the Jahn-Teller distortion in the highlighted manganese ion.⁶⁸

45

acetylacetone, and L^{R} is a 1-alkylpyrazol-4-yl-nitronyl nitroxide) ¹⁰ have been particularly well studied. These are 1D coordination polymers with bridging L^{R} ligands, which can be obtained in "head-to-tail" or "head-to-head" isomeric forms.



- The "head-to-head" isomers contain alternating [Cu(hfac)₂(pyrazole)₂] and [Cu(hfac)₂(NITR)₂] centres (Fig. 9). In the latter, the NITR radical donors coordinate to the copper axial sites, affording six-coordinate centres with a JahnTeller 20 elongation along the Cu–O{NITR} bonds at room temperature. This structure affords weak ferromagnetic coupling between the metal and ligand spins in the [Cu(hfac)₂(NITR)₂] units, leading to a magnetic moment of $\mu_{\rm eff} \approx 2.5$ BM per copper ion. Many of these materials undergo a magnetic transition on cooling, $_{25}$ associated with a reduction in $\mu_{\rm eff}$ to 1.8 BM per copper ion (Fig. 10).^{70,71} This reflects a temperature-dependent reorientation of the Jahn-Teller elongation in the [Cu(hfac)₂(NITR)₂] centres, to lie along one of the two unique equatorial O{hfac}-Cu-O{hfac} axes. This structural change in turn switches the Cu/NITR 30 coupling, from weakly ferromagnetic to strongly antiferromagnetic, leading to the observed reduction in μ_{eff} at low temperature (Fig. 10). Comparable magnetic switching related to a pressure-dependent Jahn-Teller rearrangement is seen in catena-[CuF₂(OH₂)₂(µ-pyrazine)] (see above).^{42,49}
- ³⁵ Although different in origin, these transitions resemble more conventional transition metal spin-crossover switching⁷² in exhibiting properties like thermal hysteresis⁷¹ and light-induced spin-state trapping.⁷³ Crystallographic order/ disorder transitions involving lattice solvent are important in propagating these
 ⁴⁰ cooperative phenomena through the bulk materials.⁷⁰ Although most work has been done on the [Cu(hfac)₂L^R] system, comparable magnetic transitions have also been observed in some copper(II) complexes of other nitroxyl radical derivatives.⁷⁴



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Fig. 9 Crystal structure of the head-to-head $[Cu(hfac)_2L^{Bu}]\cdot C_7H_8$ coordination polymer.⁷¹ The hfac⁻ ligands are de-emphasised for clarity, and H atoms and solvent molecules have been omitted. Colour code: C, white; Cu, green; F, cyan; N, blue; O, red.



Fig. 10 Temperature dependence of the effective magnetic moment of $[Cu(hfac)_2L^{Bu}]$ · $^{J}_{2}C_{7}H_{16}$, showing the Jahn-Teller-induced magnetic s transition near 150 K. Modified with permission from ref. 70. Copyright 2008, American Chemical Society.

Jahn-Teller Distortions in Conducting and Magnetic 10 Materials

- Superconductivity depends on the interaction between electron spins and lattice vibrations, which can promote strong pairing between electrons of opposite spin. If this coupling is strong enough, these electron pairs can be considered to be discrete 15 particles in their own right with zero overall spin, which are
- termed Cooper pairs. These Cooper pairs can migrate through a solid lattice much more readily than isolated electrons, giving rise to superconducting properties¶.⁷⁵
- The most important classes of high-temperature ²⁰ superconductors are cuprate materials, which can retain their superconducting properties at temperatures up to 130 K.^{76,77} A variety of such materials are known with varying elemental compositions and stoichiometries, but their lattices all contain CuO₂ sheets which are integral to their function (Fig. 11). These ²⁵ CuO₂ layers contain tetragonal copper(II) centres, with one or two
- longer Cu...O distances to axial oxide ion donors.⁷⁷ Dynamic Jahn-Teller distortions occur in the CuO₂ layers,⁷⁷ which are coupled to Cu–O vibrations (*c.f.* Fig. 3) and cause
- substantial changes to the local electron distribution (Schemes 1 ³⁰ and 3). This efficiently couples electronic structure to lattice vibrations, providing an alternative mechanism for electron pairing in cuprate superconductors that remains efficient at higher temperatures. Although the more detailed mechanism of superconductivity in these materials remains unclear, Jahn-Teller
- ³⁵ distortions in the cuprate layers are an important part of the process, and explain why these are still the most efficient superconductors known.⁷⁵

Another important class of magnetic materials are the manganites $[Ln_{1-x}A_x]_{n+1}Mn_nO_{3n+1}$ ($Ln^{3+} = a$ lanthanide ion, $A^{2+} = a$ group 2

- ⁴⁰ metal). These are derived from the LaMnO₃ perovskite structure but have differing ratios of octahedral manganese(III) and manganese(IV), depending on the level of divalent cation doping *x* and the oxide stoichiometry n.⁷⁹ Differing compositions lead to strong changes in the conductivity and magnetic behaviour of
- ⁴⁵ these compounds, which partly arise because of the Jahn-Teller character of the manganese oxidation states. Manganese(III) (high-spin d^4) is strongly Jahn-Teller active and has an

anisotropic distribution of electrons (there is one e_g electron, whose orbital occupancy is determined by the direction and form ⁵⁰ of the Jahn-Teller distortion). Thus, magnetic coupling between a manganese(III) site and its nearest neighbour metal ions depends on its Jahn-Teller state, and is strongly directional. In contrast, manganese (IV) (d^3) is not Jahn-Teller active, and has an isotropic spin distribution with one unpaired electron in each ⁵⁵ orbital of the t_{2g} subshell. Thus, varying the manganese(III):(IV) ratio in manganite prevoskites does not simply change their electron population, but has a much stronger effect on their electronic structure and magnetic ordering.⁷⁹



Fig. 11 Crystal structure of $La_{1.85}Sr_{0.15}CuO_4$, a superconductor with $T_c = 38$ K. Jahn-Teller disorder within the superconducting CuO₂ layers has ⁶⁵ been directly observed in this material.⁷⁸ The CuO₂ layers are highlighted, while the other atoms are de-emphasised for clarity. Colour code: Cu, green; La/Sr, cyan; O, red.

70 Manganite perovskites are particularly known as examplars of the colossal magnetoresistance phenomenon, where the electrical resistance of a material increases by orders of magnitude inside a magnetic field at high temperatures.⁸⁰ Although its detailed mechanism is still under study, colossal magnetoresistance 75 involves an interplay between three different types of transition, which can occur concurrently or independently in different materials. These are: magnetic ordering; charge ordering, delocalised between phases containing and localised manganese(III) and (IV) sites; and orbital ordering, involving a ⁸⁰ freezing out of Jahn-Teller disorder at the manganese(III) centres.81

Conclusions

Jahn-Teller phenomena in metal complexes continue to be discovered and developed, for new combinations of metal ion and ⁸⁵ coordination geometry.^{9,13,60,63,82} Temperature- and pressureinduced quenching or switching of Jahn-Teller distortions is now well-established in a limited number of copper(II) compounds,^{42-47,49,50,54,70-74} which can lead to functional materials with switchable magnetic moments or low-temperature magnetic ordering.^{49,70-74} The Jahn-Teller structure of copper(II) compounds with tridentate ligands can also be manipulated by molecular design, or by doping into an undistorted host lattice.⁵²⁻

- ^{5 55} Many of these studies have involved compounds with a *trans*-[CuX₄Y₂] coordination geometry, which seems to give rise to a particularly plastic electronic structure. Other copper(II) complexes with *trans*-[CuX₄Y₂] ligation may also display novel Jahn-Teller structures or switching properties.
- ¹⁰ Important classes of magnetic and superconducting ceramics also contain Jahn-Teller ions. The functions of these materials are intimately connected to their Jahn-Teller distortions, which provide an efficient mechanism for coupling their electronic structures to lattice vibrations.^{78,81} Comparable effects can also be
- ¹⁵ seen at the molecular level in some single molecule magnets, which depend on the strong magnetic anisotropy of the Jahn-Teller-distorted manganese(III) ion.⁶⁷⁻⁶⁹

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25 Notes and references

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³⁰ † Electronic Supplementary Information (ESI) available: a list of copper(II) complexes that have been shown to exhibit dynamic or static Jahn-Teller disorder in the solid state. This is an updated version of the list given in ref. 12. For ESI and crystallographic data in CIF or other electronic format see DOI: 10.1039/####.

 $\ddagger 1 \text{ bar} = 0.98 \text{ atm} = 0.1 \text{ MPa}.$

 \int Strictly speaking, the electronic configuration of structure A (Scheme 4) should be written as $\{d_{y^2-z^2}\}^1$ rather than $\{d_{x^2-y^2}\}^1$, because the structural 40 elongation occurs along the molecular *x*-axis, rather than the unique symmetry axis *z*.

¶ Quantum mechanical particles with integer spin, including s = 0, are bosons. Bosons have very different physics from half-integer spin

- ⁴⁵ particles (fermions), including individual electrons. In particular, bosons are described by delocalised wavefunctions, that are not centred on specific atoms or bonds. This allows Cooper pairs (which have effectively zero spin) to migrate freely through a lattice, without having to move between energy levels in a conduction band.
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Jahn-Teller Distortions in Transition Metal Compounds, and their Importance in Functional Molecular and Inorganic Materials

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Supporting Information

Table S1 Six-coordinate copper(II) compounds that have show unusual crystallographic coordination geometries, and the techniques used to demonstrate their true structure as a disordered (pseudo)-Jahn-Teller elongated octahedron. Only compounds for which these additional experiments have been done are included in the table.

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	Crystallographic Cu site symmetry and apparent structural distortion from an ideal octahedral geometry	Ref.	Technique	Ref.
$\begin{array}{l} \alpha \text{-}M_2 Pb[Cu(NO_2)_6] \\ (M^+ = K^+, Rb^+, Cs^+, Tl^+)^a \end{array}$	$T_{\rm h}$, undistorted	1-3	TLS analysis VT powder and single crystal EPR EXAFS	4 5 6
$ \begin{split} \beta - M_2 Pb[Cu(NO_2)_6] \\ (M^+ = K^+, Rb^+, Cs^+, Tl^+)^a \end{split} $	D_{2h} , tetragonal compression	4, 7-9	TLS analysis VT powder and single crystal EPR Single crystal UV/vis EXAFS	4 5, 10-12 12 6
K_2CuF_4	D_{4h} , tetragonal compression	14	UV/vis and MCD measurements Susceptibility measurements Correction of incorrect space group	15 16 17
$\gamma\text{-}Cu_2(OH)_3Cl,$ $[Cu(\mu_3\text{-}OH)_6]$ site	C _i , tetragonal compression	18	Reformulation of compound as $Cu_3(Cu_{1-x}Zn_x)(OH)_6Cl_2 \ (x > 0.5)$	19
[NH4]2[Cu(OH2)6][SO4]2	<i>C</i> _i , small rhombic elongation	20	VT crystallography VT powder and single crystal EPR Comparative UV/vis and EPR spectra of other M ₂ [Cu(OH ₂) ₆][SO ₄] ₂ salts	21-24 21, 23 25, 26
$[Cu(pyO)_6]X_2 (X^- = BF_4^-, ClO_4^-, NO_3^-)^b$	S_6 , trigonal distortion	27-30	VT crystallography VT powder and single crystal EPR EXAFS	30 29, 31 32
$[Cu(en)_3]SO_4$	D_3 , trigonal distortion	34	VT powder and single crystal EPR EXAFS VT crystallography	35 6 36
[Cu(dien) ₂][NO ₃] ₂ ^c	C_1 , rhombical compression	37	VT crystallography and TLS analysis Single crystal UV/vis and EPR	38 39
[Cu(terpy)2]Br2.3H2O, high-T phased	C_1 , rhombic compression	42	VT powder and single crystal EPR	42, 43
[Cu(terpy)2][PF6]2, high-T phased	S ₄ , tetragonal compression	44	VT powder and single crystal EPR	42
$[Cu(terpy)_2][BF_4]_2^d$	C_1 , rhombic compression	45	TLS analysis and VT powder EPR	45
[Cu(1,1"-dioxoterpy) ₂][ClO ₄] ₂	S ₄ , tetragonal compression	47	TLS analysis	47
$[Cu(1-bpp)_2][BF_4]_2$	C_1 , small rhombic elongation	48	VT crystallography and TLS analysis VT powder and single crystal EPR	49-51 50
$\beta\text{-}[Cu(1\text{-}bpp^{Mes})_2][ClO_4]_2.2CH_3NO_2$	C_1 , small rhombic elongation	52	TLS analysis and VT powder EPR	52
$[Cu(\mu_3\text{-}btz)_6(Cu\{CNtBu\})_4]$	S ₄ , tetragonal compression	53	VT powder EPR	53
$[Cu(tach)_2][NO_3]_2^e$	C_{2h} , small tetragonal compression	54	VT powder and single crystal EPR	54
$[Cu(thch)_2][CH_3C_6H_4SO_3-4]_2$	C _i , rhombic compression	55	VT crystallography	55
$[H_3NC_6H_4Cl-3]_8[CuCl_6]Cl_4$	<i>C</i> _i , tetragonal compression	56	VT crystallography EXAFS VT powder and single crystal EPR Single crystal UV/vis EPR of Cu-doped Cd(II) analogue	57 58 59 59 60
$[Cu(ompa)_3][ClO_4]_2$	D_3 , trigonal distortion	61	VT powder EPR Single crystal EPR	62 63
trans-[Cu(O ₂ CCH ₂ OCH ₃) ₂ (OH ₂) ₂]	$C_{\rm i}$, rhombic compression	64	VT crystallography	65
$[Cu(HC\{PPh_2O\}_3)_2][ClO_4]_2$	C _i , tetragonal compression	66	EXAFS VT crystallography and powder EPR	67 68
cis-[Cu(vanillinate)(OH ₂) ₂]	C_{2v} , tetragonal compression	69	VT crystallography and TLS analysis	70
trans-[Cu(O2NO)2(tmbim)2]	C _i , rhombic compression	71	TLS analysis	71
$[Cu([9]aneS_3)_2][BF_4]_2$	C _i , small tetragonal elongation	72	VT crystallography VT powder EPR and powder UV/vis EPR of Cu-doped Fe(II) analogue	73 73, 74 73

Table S1 continued.

	Crystallographic Cu site symmetry and apparent structural distortion from an ideal octahedral geometry	Ref.	Technique	Ref.
$[Cu([9]aneN_3)_2][Cu(CN)_3].2H_2O^{f}$	C_1 , small rhombic elongation	752	VT crystallography VT powder and single crystal EPR Powder UV/vis	75 74, 75 74
$[Cu(MeTRI)_2][BF_4]_2$	D_3 , trigonal distortion	78	TLS analysis	78
[Cu(NMe ₃ sar)][ClO ₄] ₄	S_6 , trigonal distortion	79	VT powder EPR EPR of Cu-doped Zn analogue	79 79
<i>cis</i> -[Cu(ONO)(bipy) ₂]NO ₃ ^g	C_1 , small rhombic elongation	80	VT crystallography	80
cis-[Cu(hat) ₂ (OH ₂) ₂][ClO ₄] ₂			VT powder EPR	82
$[Cu(\mu-dpp)(OH_2)_2]_n[BF_4]_{2n}.2nH_2O^h$	C_1 , small rhombic elongation	83	VT crystallography	83
[CuTp ₂] (molecule B) ⁱ	<i>C</i> _i , small rhombic elongation	85	TLS analysis VT crystallography and powder EPR	85, 86 86
$[Cu(tpm)_2][NO_3]_2$	C_{3v} , trigonal distortion	87	EXAFS	87
$[Cu(tpp)_2]Br_2.8H_2O$	$C_{\rm i}$, tetragonal compression	88	VT powder and single crystal EPR	88

^aThe γ -phase of the same compounds exhibits a crystallographically ordered rhombically elongated octahedral Cu(II) centre.¹³ ^bThe solvated forms of these materials, [Cu(pyO)₆][NO₃]₂.2H₂O has a typical, static Jahn-Teller elongated geometry at 293 K.³³ ^cThe salts [Cu(dien)₂]Br₂⁴⁰ and [Cu(dien)₂]Cl[ClO₄]⁴¹ have a typical, static pseudo-Jahn-Teller elongated geometry at 293 K. ^dThe salt [Cu(terpy)₂][NO₃]₂ has a typical, static pseudo-Jahn-Teller elongated geometry at 293 K. ^dThe salt [Cu(terpy)₂][NO₃]₂ has a typical, static pseudo-Jahn-Teller elongated geometry at 298 K.⁴⁶ ^eThe salt [Cu(tach)₂][ClO₄]₂ has a typical, static Jahn-Teller elongated geometry at 293 K.⁵² ^fThe salts [Cu([9]aneN₃)₂][ClO₄]₂⁷⁶ and [Cu([9]aneN₃)₂]Cl₂⁷⁷ have a very similar, small rhombic elongated geometry at 293 K and so are also probably disordered although this was not studied. In contrast, [Cu([9]aneN₃)₂][CF₃SO₄]₂.H₂O has a static Jahn-Teller elongation in the crystal.⁷⁷ ^gSeveral other fluxional complexes of type [Cu(X)(bipy)₂]Y, [Cu(X)(phen)₂]Y or [Cu(X)(bipyam)₂]Y (X⁻ = NO₂⁻ or MeCO₂⁻; Y⁻ = NO₃⁻, BF₄⁻, ClO₄⁻, PF₆⁻ etc.) have also been studied by VT crystallography.^{81 h}The isomorphous ClO₄⁻ salt of the same compound also probably contains fluxional Cu centres, although this was not studied.^{84 i}There are two independent molecules in the crystal structure of this compound. The other molecule has a static Jahn-Teller elongated structure.

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