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Von Neumann entropy and majorization

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Abstract

We consider the properties of the Shannon entropy for two probability distributions which stand in the relationship of majorization. Then we give a generalization of a theorem due to Uhlmann, extending it to infinite dimensional Hilbert spaces. Finally we show that for any quantum channel Φ , one has $S(\Phi(\rho)) = S(\rho)$ for all quantum states ρ if and only if there exists an isometric operator V such that $\Phi(\rho) = V\rho V^*$.

Keywords: Von Neumann entropy, majorization, quantum operation.

AMS Classification: 47L05, 47L90, 81R10

1 Introduction

In this paper we study aspects of the relation of majorization for classical probability distributions and for quantum mechanical density operators and connections with the action of quantum channels, utilising relevant properties of the Shannon entropy and von Neumann entropy. We obtain extensions of classic results from finite to infinite dimensional Hilbert spaces.

Let $\mathcal{B}(\mathcal{H})$ be the von Neumann algebra of all bounded linear operators on a separable Hilbert space \mathcal{H} over \mathbb{C} and $S(\mathcal{H})$ be the set of all *density operators* on \mathcal{H} . That is, $\rho \in S(\mathcal{H})$ if and only if $\rho \geq 0$ and $\text{tr}(\rho) = 1$. The elements of $S(\mathcal{H})$ are taken to represent quantum states in quantum physics, while the selfadjoint elements of $\mathcal{B}(\mathcal{H})$ represent (bounded) observables. The set $S(\mathcal{H})$ spans the Banach space $\mathcal{T}(\mathcal{H})$ of all trace class operators on \mathcal{H} .

We denote by $\mathcal{B}(\mathcal{H}, \mathcal{K})$ the space of all bounded linear operators from \mathcal{H} into another Hilbert space \mathcal{K} . An operator $V \in \mathcal{B}(\mathcal{H}, \mathcal{K})$ is called an isometry if $V^*V = I_{\mathcal{H}}$. In this case, $VV^* \in \mathcal{B}(\mathcal{K})$ is an orthogonal projection. For $x, y \in \mathcal{H}$, $x \otimes y$ denotes the (linear, rank-1) operator $z \mapsto \langle z, y \rangle x$ ($z \in \mathcal{H}$). Associated with each $\rho \in S(\mathcal{H})$ is the sequence $\lambda(\rho) \equiv (\lambda_1(\rho), \lambda_2(\rho), \dots, \lambda_n(\rho), \dots)$ of eigenvalues of ρ , arranged in non-increasing order. Thus $\lambda(\rho) \in c_0^*$, where c_0^* is the positive cone of sequences decreasing monotonically to 0, as defined in [7].

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Let $l^\infty(\mathbb{R})$ and $l_1^1(\mathbb{R}^+)$ denote the sets of all bounded real sequences and of all summable non-negative real sequences which have sum 1, respectively. For a vector $r \in l^\infty(\mathbb{R})$, we introduce $r^\downarrow = (r_1^\downarrow, \dots, r_n^\downarrow, \dots)$ as the vector whose elements are the elements of r re-ordered into non-increasing order. Adopting the definition of majorization given in [2],[7], for $r, s \in c_0^*$, we say that r is majorized by s , written as $r \prec s$, if

$$\sum_{i=1}^k r_i^\downarrow \leq \sum_{i=1}^k s_i^\downarrow, \text{ for } k = 1, 2, \dots \quad \text{and} \quad \sum_{i=1}^{\infty} r_i^\downarrow = \sum_{i=1}^{\infty} s_i^\downarrow.$$

Recently, the infinite dimensional Schur-Horn theorem and infinite majorization have received much attention. In [12], A. Neumann has given properties of infinite majorization in $l^\infty(\mathbb{R})$; V. Kaftal and G. Weiss [7] have obtained interesting results for infinite majorization in c_0^* . Arveson and Kadison [3] presented some other characterizations by different methods.

Some results relevant for our purposes are the following: if $r, s \in l_1^1(\mathbb{R}^+)$, then

$$r \prec s \iff r = Qs, \text{ with } Q_{ij} = |U_{ij}|^2 \text{ for some unitary } U \text{ [6, Theorem 1],}$$

and

$$r \prec s \iff r = Qs, \text{ for some orthostochastic matrix } Q \text{ [7, Corollary 6.1].}$$

Motivated by the above studies, we first consider the properties of the Shannon entropy for two elements in $l_1^1(\mathbb{R}^+)$ that satisfy majorization. Then we extend and study those properties for two density operators $\rho, \sigma \in S(\mathcal{H})$. Following the definition given for finite dimensional spaces (see [1]), we denote $\rho \prec \sigma$ for two operators $\rho, \sigma \in S(\mathcal{H})$ if $\lambda(\rho) \prec \lambda(\sigma)$.

Let $\mathcal{M}_n(\mathcal{B}(\mathcal{H}))$ be the von Neumann algebra of $n \times n$ matrices whose entries are in $\mathcal{B}(\mathcal{H})$, and let $\Phi : \mathcal{B}(\mathcal{H}) \rightarrow \mathcal{B}(\mathcal{H})$ be a linear map. Then Φ induces a map $\Phi_n : \mathcal{M}_n(\mathcal{B}(\mathcal{H})) \rightarrow \mathcal{M}_n(\mathcal{B}(\mathcal{H}))$ by the formula

$$\Phi_n((a_{i,j})) = (\Phi(a_{i,j})) \text{ for } (a_{i,j}) \in \mathcal{M}_n(\mathcal{B}(\mathcal{H})).$$

If every Φ_n is a positive map, then Φ is called completely positive. Φ is said to be normal if Φ is continuous with respect to the ultraweak (σ -weak) topology. Normal completely positive contractive maps on $\mathcal{B}(\mathcal{H})$ were characterized in a theorem of Kraus [8, Theorem 3.3], which says that Φ is a normal completely positive map if and only if there exists a sequence $\{A_i\}_{i=1}^{\infty}$ in $\mathcal{B}(\mathcal{H})$ such that for all $X \in \mathcal{B}(\mathcal{H})$,

$$\Phi(X) = \sum_{i=1}^{\infty} A_i X A_i^* \quad \text{with} \quad \sum_{i=1}^{\infty} A_i A_i^* \leq I.$$

where the limits are defined in the strong operator topology. The sequence $\{A_i\}_{i=1}^{\infty}$, which is not unique, is also called a family of Kraus operators for Φ . In this case, Φ has a dual map Φ^\dagger defined by

$$\Phi^\dagger(X) = \sum_{i=1}^{\infty} A_i^* X A_i \quad \text{for } X \in \mathcal{T}(\mathcal{H}),$$

where the sum converges in the trace norm topology. It is easy to see that one has

$$|\mathrm{tr}[\Phi^\dagger(X)Y]| = |\mathrm{tr}[X\Phi(Y)]| \leq \|\Phi\| \|Y\| \mathrm{tr}(|X|), \quad X \in \mathcal{T}(\mathcal{H}), \quad Y \in \mathcal{B}(\mathcal{H}),$$

so $\Phi^\dagger(X) \in \mathcal{T}(\mathcal{H})$ and Φ^\dagger is well defined on $\mathcal{T}(\mathcal{H})$. In general, Φ^\dagger cannot be extended from $\mathcal{T}(\mathcal{H})$ into $\mathcal{B}(\mathcal{H})$. However, if $\sum_{i=1}^\infty A_i^* A_i \leq I$, then Φ^\dagger is well defined as a normal map on $\mathcal{B}(\mathcal{H})$. A normal completely positive map Φ which is trace preserving ($\Phi^\dagger(I) = I$, corresponding to $\mathrm{tr}(\Phi(X)) = \mathrm{tr}(X)$ for $X \in \mathcal{T}(\mathcal{H})$) is called a *quantum channel*. If a normal completely positive map satisfies $\Phi(I) \leq I$, then Φ is called a quantum operation [8, 9]. A quantum operation is unital if $\Phi(I) = I$, which is equivalent to $\sum_j A_j A_j^* = I$. A quantum operation is bi-stochastic if it is both trace-preserving and unital. In particular, Φ is said to be a mixed unitary operation if $\Phi(X) = \sum_{i=1}^n t_i U_i X U_i^*$, where $n < \infty$, the U_i are all unitary operators and $t_i > 0$, $\sum_{i=1}^n t_i = 1$.

The von Neumann entropy of a quantum state ρ is defined by the formula

$$S(\rho) \equiv -\mathrm{tr}(\rho \log(\rho)).$$

Here we follow the common practice of taking the logarithm to base two. In classical information theory, the Shannon entropy is defined by $H(p) = -\sum_i p_i \log(p_i)$, where $p = (p_1, p_2, \dots, p_n, \dots)$ is a probability distribution. If λ_i are the eigenvalues of ρ , then the von Neumann entropy can be re-expressed as

$$S(\rho) = -\sum_{i=1}^{\infty} \lambda_i \log(\lambda_i) = H(\lambda(\rho)),$$

where we use $0 \log 0 = 0$.

There are extensive recent studies of quantum operations which preserve the von Neumann entropy and the relative entropy of quantum states [5, 10, 15, 17]

Hardy, Littlewood and Pólya [4] showed for $\xi, \eta \in \mathbb{R}^n$,

$$\xi \prec \eta \iff \xi = Q\eta \text{ for some doubly stochastic matrix } Q.$$

In the quantum context, Uhlmann proved the following for any pair of density operators $\rho, \sigma \in S(\mathcal{H})$ acting in a finite dimensional Hilbert space \mathcal{H} :

$$\rho \prec \sigma \iff \rho = \Phi(\sigma), \text{ for some mixed unitary quantum operation } \Phi.$$

Uhlmann's theorem can be used to study the role of majorization in quantum mechanics.

Here we first consider the properties of the Shannon entropies of two probability distributions which obey majorization. Then we give a generalization of Uhlmann's theorem for infinite dimensional Hilbert spaces. Finally, we give a characterization of quantum channels that leave the von Neumann entropy invariant.

2 Shannon entropy of infinite probability distributions

The following lemma is a direct corollary of [13, Theorem 8.0.1] as the function $f(x) \equiv -x \log(x)$ is a concave function.

Lemma 2.1. *Let $a, b \in \mathbb{R}^n$ with $a_i, b_i \geq 0$, and $a \prec b$. Then $H(a) \geq H(b)$.*

This result extends to the infinite case as follows.

Theorem 2.2. *Let $a, b \in l_1^1(\mathbb{R}^+)$ and $a \prec b$. Then $H(a) \geq H(b)$.*

Proof. Suppose that $c = (c_1, c_2, \dots, c_n, \dots)$, where $c_i \geq 0$ and $\sum_{i=1}^{\infty} c_i \leq 1$. It is clear that for arbitrary n , we have

$$(c_1, c_2, \dots, c_n, 0, 0, 0, \dots) \prec \left(\sum_{i=1}^n c_i, 0, 0, 0, \dots \right),$$

so Lemma 2.1 implies $-\sum_{i=1}^n c_i \log c_i \geq -(\sum_{i=1}^n c_i) \log(\sum_{j=1}^n c_j)$. Letting $n \rightarrow \infty$, then by the continuity of the function $f(x) = -x \log x$ we get

$$-\sum_{i=1}^{\infty} c_i \log c_i \geq -\sum_{i=1}^{\infty} c_i \log \left(\sum_{j=1}^{\infty} c_j \right). \quad (2.1)$$

We assume first that b has only finitely many nonzero elements. Without loss of generality, we suppose

$$(a_1, a_2, \dots, a_n, \dots) = a \prec b = (b_1, b_2, \dots, b_m, 0, 0, \dots),$$

where $b_i > 0$, for $1 \leq i \leq m$. Thus there exists N such that $\sum_{i=N+1}^{\infty} a_i < \min\{b_1, b_2, \dots, b_m\}$, so

$$\left(a_1, a_2, \dots, a_N, \sum_{i=N+1}^{\infty} a_i, 0, 0, \dots \right) \prec (b_1, b_2, \dots, b_m, 0, 0, \dots),$$

which implies

$$-\sum_{i=1}^N a_i \log a_i - \sum_{i=N+1}^{\infty} a_i \log \left(\sum_{j=N+1}^{\infty} a_j \right) \geq -\sum_{i=1}^m b_i \log(b_i).$$

By inequality (2.1), we know $H(a) \geq H(b)$.

For the general case of b , we note that for any positive integers s , the relation $a \prec b$ entails

$$(a_1, a_2, \dots, a_n, \dots) \prec \left(b_1, b_2, \dots, b_s, \sum_{i=s+1}^{\infty} b_i, 0, 0, \dots \right),$$

so the proof above yields

$$H(a) \geq -\sum_{i=1}^s b_i \log b_i - \sum_{i=s+1}^{\infty} b_i \log \left(\sum_{j=s+1}^{\infty} b_j \right).$$

Letting $s \rightarrow \infty$, we obtain $H(a) \geq H(b)$, as desired. \square

Lemma 2.3. Let $a, b \in \mathbb{R}^n$ with $a_i, b_i \geq 0$. If $a \prec b$ and $H(a) = H(b)$, then $a^\downarrow = b^\downarrow$.

Proof. Without loss of generality, we assume that

$$a^\downarrow = (a_1, a_2, \dots, a_m, 0 \dots 0) \prec b^\downarrow = (b_1, b_2, \dots, b_m, 0 \dots 0),$$

where $m \leq n$ and $a_i + b_i \neq 0$, for $1 \leq i \leq m$. For $0 < t < 1$, denote

$$c_t = (ta_1 + (1-t)b_1, ta_2 + (1-t)b_2, \dots, ta_m + (1-t)b_m).$$

It is easy to verify that $a \prec c_t \prec b$, so Lemma 2.1 implies

$$-\sum_{i=1}^m (ta_i + (1-t)b_i) \log(ta_i + (1-t)b_i) = H(a). \quad (2.2)$$

Taking the second derivative with respect to t on both sides of equation (2.2), we get

$$\sum_{i=1}^m \frac{(a_i - b_i)^2}{ta_i + (1-t)b_i} = 0,$$

which yields $a_i = b_i$ for $1 \leq i \leq m$, so $a^\downarrow = b^\downarrow$. \square

The following proposition shows that the Shannon entropy of an infinite probability distribution is strictly monotone with respect to the relation of majorization.

Proposition 2.4. Let $a, b \in l_1^1(\mathbb{R}^+)$ such that all $a_i b_i > 0$. If $a \prec b$ and $H(a) = H(b) < \infty$, then $a^\downarrow = b^\downarrow$.

Proof. Let

$$a^\downarrow = (a_1, a_2, \dots, a_n, \dots) \prec b^\downarrow = (b_1, b_2, \dots, b_n, \dots).$$

Then $a_1 \leq b_1$, so there exists k such that $a_1 \in (b_{k+1}, b_k]$, which implies $a_1 = tb_{k+1} + (1-t)b_k$, for some $0 \leq t \leq 1$. Denote

$$c = (b_1, b_2, \dots, tb_{k+1} + (1-t)b_k, (1-t)b_{k+1} + tb_k, b_{k+2}, \dots),$$

it is clear that $a \prec c \prec b$. Thus by Theorem 2.2 and the assumption, we have $H(a) = H(b) = H(c)$, so Lemma 2.3 implies $a_1 = tb_{k+1} + (1-t)b_k = b_k$. Let

$$\tilde{a} = (a_2, a_3, \dots, a_n, \dots) \text{ and } \tilde{b} = (b_1, b_2, \dots, b_{k-1}, b_{k+1}, \dots).$$

Thus $\tilde{a} \prec \tilde{b}$ and $H(\tilde{a}) = H(\tilde{b}) < \infty$, so by a similar proof, we conclude that there exists $k_1 > k$ such that $a_2 = b_{k_1}$. By mathematical induction, we get $a_{n+1} = b_{k_n}$, for some subsequence of b . However we have $\sum_{i=1}^\infty a_i = \sum_{i=1}^\infty b_i = 1$, and then the fact that all $a_i b_i > 0$ implies $a_{n+1} = b_{k_n} = b_{n+1}$, that is $a^\downarrow = b^\downarrow$. \square

Remark 2.5. The conditions that all $a_i b_i > 0$ and $H(a) < \infty$ are essential in Proposition 2.4. Indeed, it is obvious that both conditions $a \prec b$ and $H(a) = H(b) < \infty$ are not changed if we add some zeros for a . Furthermore, if $H(a) = H(b) = \infty$, we may replace a by $a' = (\frac{a_1}{2}, \frac{a_1}{2}, a_2, \dots)$. Then $a' \prec b$ and $H(a') = H(b) = \infty$. However, it is a contradiction that $a^\downarrow = b^\downarrow = a'^\downarrow$.

3 Von Neumann entropy of quantum states

Let us recall that two quantum states ρ_1 and ρ_2 are \mathcal{L}^1 -equivalent if and only if there is a sequence of unitary operators $\{U_i\}_{i=1}^\infty$ such that $\lim_{n \rightarrow \infty} \|\rho_1 - U_n \rho_2 U_n^*\|_1 = 0$.

Proposition 3.1. *Let $\rho_1, \rho_2 \in S(\mathcal{H})$. If $\rho_1 \prec \rho_2$ and $S(\rho_1) = S(\rho_2) < \infty$, then ρ_1 and ρ_2 are \mathcal{L}^1 -equivalent.*

Proof. By Proposition 2.4, all the non-zero spectral points (including multiplicity) of ρ_1 and ρ_2 are the same, which is equivalent to the fact that ρ_1 and ρ_2 are \mathcal{L}^1 -equivalent, by [3, Proposition 3.1]. \square

The following two results give extensions of Uhlmann's theorem in an infinite dimensional Hilbert space. We need some eigenvalue estimates of Weyl given in [16]. Let A be a compact operator with eigenvalues $\lambda_1 \geq \lambda_2 \geq \lambda_3 \geq \dots$ and \mathcal{P}_n be the set of all n -dimensional projections. Then $\max_{P \in \mathcal{P}_n} \text{tr}(AP) = \sum_{i=1}^n \lambda_i$.

Proposition 3.2. *Let $\rho_1, \rho_2 \in S(\mathcal{H})$. If ρ_1 is finite rank, then $\rho_1 \prec \rho_2$ if and only if there exists a mixed unitary operation Φ such that $\Phi(\rho_2) = \rho_1$.*

Proof. Necessity. If ρ_1 is finite rank, then denote $\lambda(\rho_1) = (\lambda_1(\rho_1), \lambda_2(\rho_1) \dots \lambda_m(\rho_1), 0, 0 \dots)$. As $\rho_1 \prec \rho_2$, we have that ρ_2 is also finite rank and $\lambda(\rho_2) = (\lambda_1(\rho_2), \lambda_2(\rho_2) \dots \lambda_n(\rho_2), 0, 0 \dots)$ where $n \leq m$. Using the spectral decomposition of the states ρ_1 and ρ_2 , we conclude that there exist two orthonormal bases $\{x_i\}_{i=1}^\infty$ and $\{y_i\}_{i=1}^\infty$ of \mathcal{H} such that

$$\rho_1 = \sum_{i=1}^m \lambda_i(\rho_1) x_i \otimes x_i \text{ and } \rho_2 = \sum_{i=1}^n \lambda_i(\rho_2) y_i \otimes y_i.$$

Let \mathcal{H}_1 be the subspace spanned by $\{x_i\}_{i=1}^m$ and U be the unitary operator defined by $Uy_i = x_i$, for $i = 1, 2, \dots$. Then $\rho_1|_{\mathcal{H}_1}$ and $U\rho_2 U^*|_{\mathcal{H}_1}$ can be represented by $m \times m$ matrices which have spectra $(\lambda_1(\rho_1), \lambda_2(\rho_1) \dots \lambda_m(\rho_1))$ and $(\lambda_1(\rho_2), \lambda_2(\rho_2) \dots \lambda_n(\rho_2), \dots, 0) \in \mathbb{R}^m$, respectively. Thus Uhlmann's theorem ([1] or [13, Theorem 4.1.1]) implies that there exists a mixed unitary operation $\tilde{\Phi}$ on $\mathcal{B}(\mathcal{H}_1)$, that is $\rho_1|_{\mathcal{H}_1} = \sum_{i=1}^m t_i V_i (U\rho_2 U^*)|_{\mathcal{H}_1} V_i^*$, where $V_i \in \mathcal{B}(\mathcal{H}_1)$ are unitary operators, $\sum_{i=1}^m t_i = 1$ and all $t_i > 0$. As $(U\rho_2 U^*)|_{\mathcal{H}_1^\perp} = 0$, we denote $U_i = \text{diag}(V_i, I_{\mathcal{H}_1^\perp})$, so $U_i \in \mathcal{B}(\mathcal{H})$ are unitary operators, and $\rho_1 = \sum_{i=1}^m t_i U_i (U\rho_2 U^*) U_i^*$.

Sufficiency. By Weyl's estimates [16], we get for $n = 1, 2, \dots$,

$$\begin{aligned} \sum_{i=1}^n \lambda_i(\rho_1) &= \max_{P \in \mathcal{P}_n} \{\text{tr}(\rho_1 P)\} = \max_{P \in \mathcal{P}_n} \left\{ \text{tr} \left(\sum_{i=1}^m t_i U_i \rho_2 U_i^* P \right) \right\} \\ &= \max_{P \in \mathcal{P}_n} \left\{ \sum_{i=1}^m t_i \text{tr}(\rho_2 U_i^* P U_i) \right\} \leq \sum_{i=1}^m t_i \max_{P \in \mathcal{P}_n} \{\text{tr}(\rho_2 U_i^* P U_i)\} \\ &= \max_{P \in \mathcal{P}_n} \{\text{tr}(\rho_2 P)\} = \sum_{i=1}^n \lambda_i(\rho_2). \end{aligned}$$

\square

Theorem 3.3. *Let $\rho_1, \rho_2 \in S(\mathcal{H})$. Then the following three conditions are equivalent:*

- (a) $\rho_1 \prec \rho_2$.
- (b) *There exists a sequence of mixed unitary operation Ψ_n and a bi-stochastic quantum operation Ψ on $S(\mathcal{H})$ such that $\lim_{n \rightarrow \infty} \|\Psi_n(\rho) - \Psi(\rho)\|_1 = 0$ for all $\rho \in S(\mathcal{H})$, and $\Psi(\rho_2) = \rho_1$.*
- (c) *There exists a bi-stochastic quantum operation Ψ such that $\Psi(\rho_2) = \rho_1$.*

Proof. (a) \Rightarrow (b). If $\rho_1 \prec \rho_2$, then $\lambda(\rho_1)^\downarrow \prec \lambda(\rho_2)^\downarrow$, it follows from [6, Theorem 1] that $\lambda(\rho_1)^\downarrow = Q\lambda(\rho_2)^\downarrow$, where Q is an infinite matrix satisfying $Q_{ij} = |u_{ij}|^2$ for the elements of a unitary matrix U . By virtue of the spectral decomposition of ρ_2 , there exists an orthonormal basis $\{y_i\}_{i=1}^\infty$ of \mathcal{H} such that $\rho_2 = \sum_{i=1}^\infty \lambda_i(\rho_2) y_i \otimes y_i$. Defining a unitary operator $\tilde{U} \in B(\mathcal{H})$ via $\langle \tilde{U} y_i, y_j \rangle = u_{ij}$ for $i, j = 1, 2, \dots$, we set $e_i = \tilde{U} y_i$ for $i = 1, 2, \dots$. It is easy to see that $\{e_i\}_{i=1}^\infty$ is an orthonormal basis of \mathcal{H} such that $\lambda_i(\rho_1) = \langle \rho_2 e_i, e_i \rangle$ for $i = 1, 2, \dots$. With respect to the basis $\{e_i\}_{i=1}^\infty$, then ρ_2 has an infinite matrix form given as follows:

$$\rho_2 = \begin{pmatrix} \lambda_1(\rho_1) & \lambda_{12} & \cdots & \lambda_{1n} & \cdots \\ \lambda_{21} & \lambda_2(\rho_1) & \cdots & \lambda_{2n} & \cdots \\ \vdots & \vdots & \ddots & \vdots & \vdots \\ \lambda_{n1} & \cdots & \cdots & \lambda_n(\rho_1) & \cdots \\ \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix},$$

where $\lambda_{ij} = \overline{\lambda_{ji}}$. Denoting the sequence of rank-one projections $E_i \equiv e_i \otimes e_i$, then define $\Phi(\rho) = \sum_{i=1}^\infty E_i \rho E_i$, for $\rho \in S(\mathcal{H})$. Also for $n = 1, 2, \dots$, let Φ_n be the mixed unitary operation defined as $\Phi_n(\rho) = \frac{1}{n} \sum_{i=1}^n U^i \rho (U^i)^*$, where $U = \text{diag}(\omega, \omega^2 \cdots \omega^n, 1, 1 \cdots 1 \cdots)$ and $\omega = e^{\frac{2\pi i}{n}}$. By a direct calculation we obtain

$$\Phi_n(e_i \otimes e_j) = \frac{1}{n} \sum_{i=1}^n U^i (e_i \otimes e_j) (U^i)^* = \begin{cases} e_i \otimes e_i & \text{if } i = j \\ e_i \otimes e_j & \text{if } i, j > n-1 \\ 0 & \text{if } i \neq j \text{ and } i \leq n-1 \text{ or } j \leq n-1. \end{cases}$$

Thus $\Phi_n(\rho_2)$ becomes the following block diagonal matrix with respect to the subspace decomposition $\mathcal{H} = \bigvee_{i=1}^{n-1} \{e_i\} \oplus \bigvee_{i=n}^\infty \{e_i\}$:

$$\Phi_n(\rho_2) = \begin{pmatrix} \rho_{21} & 0 \\ 0 & \rho_{22} \end{pmatrix},$$

where

$$\rho_{21} = \text{diag}(\lambda_1(\rho_1), \lambda_2(\rho_1) \cdots \lambda_n(\rho_1)), \quad \rho_{22} = (I - P_n) \rho_2 (I - P_n)|_{(\bigvee_{i=n+1}^\infty \{e_i\})},$$

and P_n denotes the orthogonal projection on the subspace $\bigvee_{i=1}^{n-1} \{e_i\}$ spanned by $\{e_i\}_{i=1}^{n-1}$. Then

$$\|\Phi_n(\rho_2) - \Phi(\rho_2)\|_1 = 2 \text{tr}([\Phi_n(\rho_2) - \Phi(\rho_2)]^+) \leq 2 \text{tr}(\rho_{22}) \rightarrow 0 \text{ as } n \rightarrow \infty,$$

since $\text{tr}(\Phi_n(\rho_2)) = \text{tr}(\Phi(\rho_2))$ implies

$$\begin{aligned}\|\Phi_n(\rho_2) - \Phi(\rho_2)\|_1 &= \text{tr}(|\Phi_n(\rho_2) - \Phi(\rho_2)|) \\ &= \text{tr}([\Phi_n(\rho_2) - \Phi(\rho_2)]^+) + \text{tr}([\Phi_n(\rho_2) - \Phi(\rho_2)]^-) \\ &= 2 \text{tr}([\Phi_n(\rho_2) - \Phi(\rho_2)]^+); \end{aligned}$$

and $\Phi_n(\rho_2) - \Phi(\rho_2) \leq (I - P_n)\rho_2(I - P_n)$ yields $[\Phi_n(\rho_2) - \Phi(\rho_2)]^+ \leq P_+(I - P_n)\rho_2(I - P_n)P_+$, so

$$\begin{aligned}\text{tr}([\Phi_n(\rho_2) - \Phi(\rho_2)]^+) &\leq \text{tr}(P_+(I - P_n)\rho_2(I - P_n)P_+) \\ &\leq \text{tr}((I - P_n)\rho_2(I - P_n)) \\ &= \text{tr}(\rho_{22}), \end{aligned}$$

Here we use the notation A^+ , A^- for the positive and negative parts of the self-adjoint operator A , and P_+ is the orthogonal projection on the range of $[\Phi_n(\rho_2) - \Phi(\rho_2)]^+$. From the spectral decomposition of ρ_1 we have an orthonormal basis $\{f_i\}_{i=1}^\infty$ of \mathcal{H} such that $\rho_1 = \sum_{i=1}^\infty \lambda_i(\rho_1)f_i \otimes f_i$. Define a unitary operator V by $Vf_i = e_i$, for $i = 1, 2, \dots$, so $\rho_1 = V^*\Phi(\rho_2)V$. Denote $\Psi_n(\rho) = V^*\Phi_n(\rho)V$ and $\Psi(\rho) = V^*\Phi(\rho)V$ for all $\rho \in S(\mathcal{H})$. Then by a similar proof to the above, we obtain $\lim_{n \rightarrow \infty} \|\Psi_n(\rho) - \Psi(\rho)\|_1 = \lim_{n \rightarrow \infty} \|\Phi_n(\rho) - \Phi(\rho)\|_1 = 0$ for all $\rho \in S(\mathcal{H})$.

(b) \Rightarrow (c) is clear.

(c) \Rightarrow (a). Let $\rho_1 = \Psi(\rho_2) = \sum_{i=1}^\infty A_i \rho_2 A_i^*$. Applying Weyl's eigenvalue theorem, we get for $n = 1, 2, \dots$,

$$\sum_{i=1}^n \lambda_i(\rho_1) = \max_{P \in \mathcal{P}_n} \left\{ \text{tr} \left(\sum_{i=1}^\infty A_i \rho_2 A_i^* P \right) \right\} = \max_{P \in \mathcal{P}_n} \left\{ \text{tr} \left(\rho_2 \sum_{i=1}^\infty A_i^* P A_i \right) \right\}. \quad (3.3)$$

As $\sum_{i=1}^\infty A_i^* A_i = I$ and $\sum_{i=1}^\infty A_i A_i^* = I$, then $\sum_{i=1}^\infty A_i^* P A_i \leq I$ and

$$\text{tr} \left(\sum_{i=1}^\infty A_i^* P A_i \right) = \text{tr} \left(\sum_{i=1}^\infty A_i A_i^* P \right) = n.$$

For convenience, we denote $B = \sum_{i=1}^\infty A_i^* P A_i$ and $u_i = \langle B y_i, y_i \rangle$, where $\{y_i\}_{i=1}^\infty$ is an orthonormal basis of \mathcal{H} satisfying $\rho_2 = \sum_{i=1}^\infty \lambda_i(\rho_2) y_i \otimes y_i$. Thus $0 \leq u_i \leq 1$ and $\sum_{i=1}^\infty u_i = n$, so

$$\begin{aligned}\text{tr}(\rho_2 B) &= \sum_{i=1}^\infty \lambda_i(\rho_2) u_i = \sum_{i=1}^n \lambda_i(\rho_2) u_i + \sum_{i=n+1}^\infty \lambda_i(\rho_2) u_i \\ &\leq \sum_{i=1}^n \lambda_i(\rho_2) u_i + \lambda_{n+1}(\rho_2) \sum_{i=n+1}^\infty u_i \\ &\leq \sum_{i=1}^n \lambda_i(\rho_2) u_i + \lambda_{n+1}(\rho_2) \sum_{i=1}^n (1 - u_i) \\ &\leq \sum_{i=1}^n \lambda_i(\rho_2). \end{aligned} \quad (3.4)$$

By equation (3.3) and inequality (3.4), we conclude that for all $n = 1, 2, \dots$, one has the inequalities $\sum_{i=1}^n \lambda_i(\rho_1) \leq \sum_{i=1}^n \lambda_i(\rho_2)$, as desired. \square

Remark 3.4. In a finite dimensional Hilbert space, $\rho_1 \prec \rho_2$ is equivalent to $\rho_1 = \Phi(\rho_2)$, for some mixed unitary operation Φ . However, for an infinite dimensional Hilbert space, $\rho_1 \prec \rho_2$ does not imply $\rho_1 = \sum_{i=1}^{\infty} t_i U_i \rho_2 U_i^*$, where $\sum_{i=1}^{\infty} t_i = 1$, $t_i \geq 0$ and U_i are unitary operators for all i . Indeed, the condition $\rho_1 = \sum_{i=1}^{\infty} t_i U_i \rho_2 U_i^*$ yields $\dim(\ker(\rho_1)) \leq \dim(\ker(\rho_2))$. But we can supplement many zeros for $\lambda(\rho_1)$.

The following proposition was obtained in [10, 17] for the finite case, in which the condition of injectivity of ρ may be dropped.

Proposition 3.5. Let $\rho \in S(\mathcal{H})$ and Φ be a bi-stochastic quantum operation. If ρ is injective and $S(\rho) = S(\Phi(\rho)) < \infty$, then $\Phi(\rho) = U\rho U^*$ for a unitary operator U .

Proof. Suppose $\Phi(\rho) = \sum_{i=1}^{\infty} A_i \rho A_i^*$. Then by Theorem 3.3, we have $\Phi(\rho) \prec \rho$. We claim that if ρ is injective, then so is $\Phi(\rho)$. Thus, suppose that $\Phi(\rho)$ were not injective, then there is a vector $x \neq 0$ that satisfies $\Phi(\rho)x = 0$, so $\langle \sum_{i=1}^{\infty} A_i \rho A_i^* x, x \rangle = 0$, which yields $A_i^* x = 0$, for all i . Thus $x = \sum_{i=1}^{\infty} A_i A_i^* x = 0$, which is a contradiction. It follows from Proposition 2.4 that $\lambda(\rho)^\downarrow = \lambda(\Phi(\rho))^\downarrow$, so by the spectral decomposition theorem, we get $\Phi(\rho) = U\rho U^*$ for a unitary operator U . \square

Remark 3.6. In an infinite dimensional Hilbert space, the condition that ρ is injective may not be dropped.

In the following, we shall characterize the structure of a quantum channel that does not change the von Neumann entropy of any quantum state. In [11], Molnár and Szokol gave the structure of the map which preserves the relative entropy in a finite-dimensional Hilbert space. However, we were unable to find a publication that studied the structure of a quantum channel which preserves the von Neumann entropy. After a preprint version of the present paper was published as arXiv:1304.7442(v1), we received a sketch of a shorter proof of Theorem 3.9 from M.B. Ruskai based on the Stinespring representation theorem and some techniques from quantum information. Here we mainly use methods of operator theory and Kraus's theorem. The following two lemmas are needed. The set of all compact operators on \mathcal{H} is denoted $\mathcal{K}(\mathcal{H})$.

Lemma 3.7. (See [14].) Let $\rho \in S(\mathcal{H})$.

- (i) $S(\rho) \geq 0$, and $S(\rho) = 0$ if and only if ρ is a pure state.
- (ii) If $\dim(\mathcal{H}) = n$, then $S(\rho) \leq \log n$, and $S(\rho) = \log n$ if and only if $\rho = \frac{1}{n}I$.

Lemma 3.8. (See [9].) Let Φ be a quantum operation with a set of Kraus operators $K(\Phi) \equiv \{A_i \in \mathcal{B}(\mathcal{H}), i = 1, 2, \dots\}$ and such that $\Phi^\dagger(I) \leq I$. Then $\{B \in \mathcal{K}(\mathcal{H}) : \Phi(B) = B\} = \{B \in \mathcal{K}(\mathcal{H}) : \Phi^\dagger(B) = B\} \subseteq \mathcal{A}'$, where \mathcal{A}' is the commutant of $\mathcal{A} = \{A_i, A_i^* : A_i \in K(\Phi), i = 1, 2, \dots\}$.

Theorem 3.9. *Let Φ be a quantum channel on $\mathcal{B}(\mathcal{H})$. Then $S(\Phi(\rho)) = S(\rho)$ for all quantum states $\rho \in S(\mathcal{H})$ if and only if there exists an isometry $V \in \mathcal{B}(\mathcal{H})$ such that $\Phi(X) = VXV^*$ for all $X \in \mathcal{B}(\mathcal{H})$.*

Proof. Sufficiency is clear.

Necessity. For any unit vector x , we have $S(\Phi(x \otimes x)) = S(x \otimes x) = 0$, so $\Phi(x \otimes x)$ is a rank one orthogonal projection. Thus there exists a unit vector z such that $\Phi(x \otimes x) = z \otimes z$.

Let $y \perp x$ be another unit vector of \mathcal{H} . For convenience, denote

$$\Phi(x \otimes x) = x' \otimes x' \text{ and } \Phi(y \otimes y) = y' \otimes y'.$$

Setting $\rho_0 = \frac{1}{2}(x \otimes x + y \otimes y)$, we have

$$S\left(\frac{1}{2}(x' \otimes x' + y' \otimes y')\right) = S(\Phi(\rho_0)) = S(\rho_0) = 1. \quad (3.5)$$

In the following, we shall show $x' \perp y'$. Let $\mathcal{H}_0 \subseteq \mathcal{H}$ denote the two-dimensional space spanned by x' and y' . Then $\Phi(\rho_0)$ can be treated as an operator from \mathcal{H}_0 into \mathcal{H}_0 . Suppose $y' = \alpha x' + \sqrt{1 - |\alpha|^2} x'^\perp$, where $0 \leq |\alpha| \leq 1$ and x'^\perp is a unit vector in \mathcal{H}_0 orthogonal to x' . Then we have the following matrix forms for the operators $x' \otimes x'$ and $y' \otimes y'$, respectively:

$$x' \otimes x' = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}, \quad y' \otimes y' = \begin{pmatrix} |\alpha|^2 & \alpha\sqrt{1 - |\alpha|^2} \\ \bar{\alpha}\sqrt{1 - |\alpha|^2} & 1 - |\alpha|^2 \end{pmatrix},$$

where $\bar{\alpha}$ is the complex conjugate of α . Then

$$\frac{1}{2}(x' \otimes x' + y' \otimes y') = \frac{1}{2} \begin{pmatrix} 1 + |\alpha|^2 & \alpha\sqrt{1 - |\alpha|^2} \\ \bar{\alpha}\sqrt{1 - |\alpha|^2} & 1 - |\alpha|^2 \end{pmatrix}. \quad (3.6)$$

The characteristic polynomial of (3.6) is

$$\lambda^2 - \lambda + \frac{1 - |\alpha|^2}{4} = 0. \quad (3.7)$$

Furthermore, Lemma 3.7 (ii) and equation (3.5) imply that equation (3.7) has two equal roots $\lambda_1 = \lambda_2 = \frac{1}{2}$, which yields $\alpha = 0$. Thus, $x' \perp y'$ as required.

From the proof above, we know that the map Φ sends orthogonal pure states to orthogonal pure states. Let $\{e_i\}_{i=1}^\infty$ be an orthonormal basis of \mathcal{H} and P_n the orthogonal projection onto the subspace spanned by $\{e_i\}_{i=1}^n$. Then

$$\Phi(I) = \lim_{n \rightarrow \infty} \Phi(P_n) = \lim_{n \rightarrow \infty} \sum_{i=1}^n \Phi(e_i \otimes e_i),$$

so $\Phi(I)$ is an infinite dimensional orthogonal projection.

Let $\rho \in S(\mathcal{H})$ be injective and $\rho = \sum_{i=1}^\infty \lambda_i x_i \otimes x_i$, where $\{x_i\}_{i=1}^\infty$ is an orthonormal basis and λ_i are the eigenvalues of ρ . Then

$$\Phi(\rho) = \lim_{n \rightarrow \infty} \Phi\left(\sum_{i=1}^n \lambda_i x_i \otimes x_i\right) = \lim_{n \rightarrow \infty} \sum_{i=1}^n \lambda_i \Phi(x_i \otimes x_i) = \sum_{i=1}^\infty \lambda_i x'_i \otimes x'_i.$$

Denote $V_\rho x_i = x'_i$ for all i , then $\Phi(\rho) = V_\rho \rho V_\rho^*$.

By Kraus's theorem, we have $\Phi(\rho) = \sum_{i=1}^{\infty} A_i \rho A_i^*$, where $A_i \in \mathcal{B}(\mathcal{H})$, and $\sum_{i=1}^{\infty} A_i^* A_i = I$, as Φ is trace-preserving. Then

$$\sum_{i=1}^{\infty} A_i \rho A_i^* = V_\rho \rho V_\rho^*,$$

which yields

$$\sum_{i=1}^{\infty} V_\rho^* A_i \rho A_i^* V_\rho = \rho, \quad (3.8)$$

for all injective $\rho \in S(\mathcal{H})$. Furthermore, $\sum_{i=1}^{\infty} A_i A_i^* = \Phi(I) \leq I$, so

$$\sum_{i=1}^{\infty} V_\rho^* A_i A_i^* V_\rho \leq I \text{ and } \sum_{i=1}^{\infty} A_i^* V_\rho V_\rho^* A_i \leq \sum_{i=1}^{\infty} A_i^* A_i = I.$$

Then Lemma 3.8 and equation (3.8) imply

$$\rho \sum_{i=1}^{\infty} V_\rho^* A_i A_i^* V_\rho = \sum_{i=1}^{\infty} V_\rho^* A_i \rho A_i^* V_\rho = \rho,$$

and

$$\rho \sum_{i=1}^{\infty} A_i^* V_\rho V_\rho^* A_i = \sum_{i=1}^{\infty} A_i^* V_\rho \rho V_\rho^* A_i = \rho, \quad (3.9)$$

so

$$I = \sum_{i=1}^{\infty} A_i^* V_\rho V_\rho^* A_i,$$

which yields

$$\sum_{i=1}^{\infty} A_i^* V_\rho V_\rho^* A_i = I = \sum_{i=1}^{\infty} A_i^* A_i.$$

Thus

$$\sum_{i=1}^{\infty} A_i^* (I - V_\rho V_\rho^*) A_i = 0,$$

which implies $A_i^* (I - V_\rho V_\rho^*) = 0$, that is

$$A_i^* = A_i^* V_\rho V_\rho^*, \quad \text{for all injective } \rho \in S(\mathcal{H}). \quad (3.10)$$

Using equations (3.8) and (3.9), we get

$$\sum_i A_i^* V_\rho \left(\sum_j V_\rho^* A_j \rho A_j^* V_\rho \right) V_\rho^* A_i = \rho,$$

so equation (3.10) implies

$$\sum_{i,j} A_i^* A_j \rho A_j^* A_i = \rho, \quad (3.11)$$

which is equivalent to $\Phi^\dagger \circ \Phi(\rho) = \rho$, for all injective $\rho \in S(\mathcal{H})$.

For any unit vector x , we claim that there exists a sequence of injective states $\rho_n \in S(\mathcal{H})$ such that $\rho_n \rightarrow x \otimes x$ as $n \rightarrow \infty$ (in the weak-* topology). Indeed, let $\{f_i\}_{i=1}^\infty$ be an orthonormal basis of the orthogonal complement subspace of x , and $\rho_n = (1 - \frac{1}{n})x \otimes x + \frac{1}{n} \sum_{i=1}^\infty 2^{-i} f_i \otimes f_i$. Then $\lim_{n \rightarrow \infty} \text{tr}(|\rho_n - x \otimes x|) = 0$, so $\lim_{n \rightarrow \infty} \text{tr}[(\rho_n - x \otimes x)X] = 0$, for all $X \in \mathcal{B}(\mathcal{H})$. Thus

$$x \otimes x = \lim_{n \rightarrow \infty} \rho_n = \lim_{n \rightarrow \infty} \Phi^\dagger \circ \Phi(\rho_n) = \Phi^\dagger \circ \Phi(x \otimes x),$$

which says that

$$\sum_{i,j} A_i^* A_j (x \otimes x) A_j^* A_i = \Phi^\dagger \circ \Phi(x \otimes x) = x \otimes x$$

for all rank one projections $x \otimes x$. Then by Lemma 3.8 again, we have $A_i^* A_j (x \otimes x) = (x \otimes x) A_i^* A_j$, which implies that for all $i, j = 1, 2, \dots$, $A_j^* A_i = \lambda_{ji} I$ and $A_i^* A_i = \lambda_{ii} I$, so $\lambda_{ii} > 0$ and $\sum_i \lambda_{ii} = 1$. It is clear that

$$\begin{aligned} \begin{pmatrix} A_1^* \\ A_2^* \\ \vdots \\ A_n^* \\ \vdots \end{pmatrix} \begin{pmatrix} A_1 & A_2 & \cdots & A_n & \cdots \end{pmatrix} &= \begin{pmatrix} A_1^* A_1 & A_1^* A_2 & \cdots & A_1^* A_n & \cdots \\ A_2^* A_1 & A_2^* A_2 & \cdots & A_2^* A_n & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ A_n^* A_1 & A_n^* A_2 & \cdots & A_n^* A_n & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix} \\ &= \begin{pmatrix} \lambda_{11} I & \lambda_{12} I & \cdots & \lambda_{1n} I & \cdots \\ \lambda_{21} I & \lambda_{22} I & \cdots & \lambda_{2n} I & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ \lambda_{n1} I & \lambda_{n2} I & \cdots & \lambda_{nn} I & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix}. \end{aligned}$$

Denote

$$M = \begin{pmatrix} \lambda_{11} & \lambda_{12} & \cdots & \lambda_{1n} & \cdots \\ \lambda_{21} & \lambda_{22} & \cdots & \lambda_{2n} & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \\ \lambda_{n1} & \lambda_{n2} & \cdots & \lambda_{nn} & \cdots \\ \cdots & \cdots & \cdots & \cdots & \cdots \end{pmatrix},$$

then $\lambda_{ji} = \overline{\lambda_{ij}}$ and

$$|\lambda_{ij}|^2 \leq \lambda_{ii} \lambda_{jj}, \text{ for } 1 \leq i, j \quad (3.12)$$

since $M \geq 0$. Further equation (3.11) implies $\sum_{i,j} A_j^* A_i A_i^* A_j = I$, which yields

$$\sum_{i,j} |\lambda_{ij}|^2 = 1 = \left(\sum_{i=1}^\infty \lambda_{ii} \right)^2.$$

Then by a direct calculation, we get

$$\sum_{i \neq j} |\lambda_{ij}|^2 = \sum_{i \neq j} \lambda_{ii} \lambda_{jj},$$

so equation (3.12) implies

$$|\lambda_{ij}|^2 = \lambda_{ii} \lambda_{jj}, \text{ for } 1 \leq i, j \leq \infty.$$

Thus

$$A_j^* A_1 = \lambda_{j1} I = \sqrt{\lambda_{11} \lambda_{jj}} e^{i\theta_j} I,$$

for $j = 1, 2, \dots$. We denote $V_j = A_j / \sqrt{\lambda_{jj}}$, so V_j are isometric operators for $j = 1, 2, \dots$, which yields $V_j^* V_1 = e^{i\theta_j} I$. Then $V_j^* V_1 V_1^* = e^{i\theta_j} V_1^*$ and $V_j V_j^* V_1 = e^{i\theta_j} V_j$. Hence

$$V_1 V_1^* V_j V_j^* V_1 V_1^* = V_1 V_1^* \text{ and } V_j V_j^* V_1 V_1^* V_j V_j^* = V_j V_j^*,$$

so we get

$$V_1 V_1^* (I - V_j V_j^*) = 0 \text{ and } V_j V_j^* (I - V_1 V_1^*) = 0,$$

then $V_1 V_1^* = V_j V_j^*$, for $j = 1, 2, \dots$. Thus for all $X \in \mathcal{B}(\mathcal{H})$,

$$A_j X A_j^* = V_1 V_1^* A_j X A_j^* V_1 V_1^* = V_1 \frac{1}{\lambda_{11}} A_1^* A_j X A_j^* A_1 V_1^* = \lambda_{jj} V_1 X V_1^*,$$

which implies

$$\Phi(X) = A_1 X A_1^* + A_2 X A_2^* + \dots + A_n X A_n^* + \dots = V_1 X V_1^*.$$

□

The following result is an immediate consequence of Theorem 3.9.

Corollary 3.10. *Let Φ be a bi-stochastic quantum operation. Then $S(\Phi(\rho)) = S(\rho)$ for all quantum state $\rho \in S(\mathcal{H})$ if and only if there exists a unitary matrix U such that $\Phi(\rho) = U\rho U^*$.*

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