

Kinetics of the Gas-Phase Reactions of *syn*- and *anti*-CH₃CHOO Criegee Intermediate Conformers with SO₂ as a Function of Temperature and Pressure

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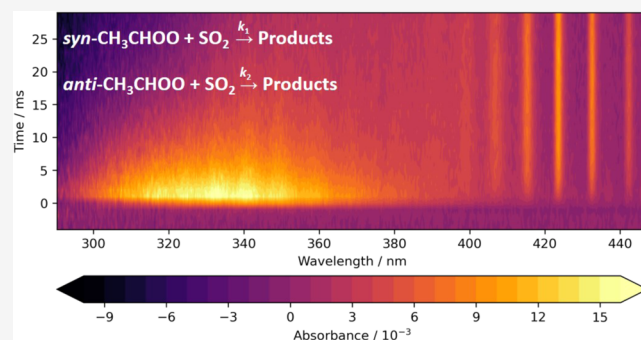
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ABSTRACT: Kinetics of reactions between SO₂ and CH₃CHOO Criegee intermediate conformers have been measured at temperatures between 242 and 353 K and pressures between 10 and 600 Torr using laser flash photolysis of CH₃CHI₂/O₂/N₂/SO₂ gas mixtures coupled with time-resolved broadband UV absorption spectroscopy. The kinetics of *syn*-CH₃CHOO + SO₂ are pressure-dependent and exhibit a negative temperature dependence, with the observed pressure dependence reconciling apparent discrepancies between previous measurements performed at ~298 K. Results indicate a rate coefficient of $(4.80 \pm 0.46) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ for the reaction of *syn*-CH₃CHOO with SO₂ at 298 K and 760 Torr. In contrast to the behavior of the *syn*-conformer, the kinetics of *anti*-CH₃CHOO + SO₂ display no significant dependence on temperature or pressure over the ranges investigated, with a mean rate coefficient of $(1.18 \pm 0.21) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$ over all conditions studied in this work. Results indicate that the reaction of *syn*-CH₃CHOO with SO₂ competes with unimolecular decomposition and reaction with water vapor in areas with high SO₂ concentration and low humidity, particularly at lower temperatures.



INTRODUCTION

The chemistry of Criegee intermediates (R₂COO) exerts potential impacts on air quality and climate through their involvement in atmospheric oxidation processes, and there has been considerable interest in the potential production of sulfate aerosols resulting from the reactions of Criegee intermediates with sulfur dioxide (SO₂) in the gas phase. Production of Criegee intermediates in the atmosphere occurs following the oxidation of unsaturated volatile organic compounds (VOCs) by ozone (O₃), with the Criegee intermediate initially produced with high internal energy.¹ Collisional stabilization of the nascent Criegee intermediate occurs in competition with unimolecular decomposition, leading to the production of stabilized Criegee intermediates (SCIs), which can participate in a range of processes, including reactions with SO₂.^{1,2}

For the reaction of the simplest SCI, CH₂OO, with SO₂, there is now general consensus regarding the kinetics at room temperature, with a current IUPAC recommendation of $(3.7^{+0.45}_{-0.40}) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ at 298 K.³ The kinetics of CH₂OO + SO₂ have been demonstrated to be independent of pressure under typical atmospheric conditions,^{4,5} with a negative temperature dependence⁶ and reaction products dominated by formaldehyde (HCHO)^{4,7} and sulfur trioxide (SO₃).^{8,9} Theory indicates that the reaction of an SCI with

SO₂ proceeds via the barrierless formation of a cyclic secondary ozonide (SOZ), followed by collisional stabilization of the SOZ or further rearrangement via one or more submerged barriers to form products including SO₃.^{10–15} In the case of CH₂OO + SO₂, there is negligible stabilization of the SOZ under atmospheric conditions,^{11,12} and contributions from two stereochemical pathways lead to the production of HCHO + SO₃ via submerged barriers.¹² For reactions of other SCIs with SO₂, there is greater potential for stabilization of the SOZ,¹¹ leading to the potential for pressure-dependent kinetics and product yields.¹¹

The atmospheric oxidation of SO₂ by CH₂OO is now expected to be limited, owing to competition with the rapid reaction of CH₂OO with water dimers,³ but several field studies have demonstrated that significant discrepancies between measured and modeled concentrations of gas-phase sulfuric acid (H₂SO₄) in the atmosphere remain, which potentially result from incomplete consideration of the

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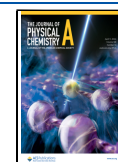


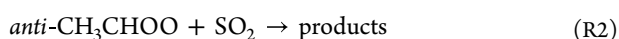
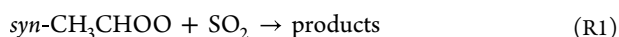
Table 1. Summary of Literature Results for k_1 and k_2 ^a

reference	method	photolysis λ /nm	T/K	p/Torr	bath gas	$[\text{CH}_3\text{CHI}_2]/10^{13}$ cm^{-3}	$[\text{SO}_2]/10^{13}$ cm^{-3}	$k_1/10^{-11}$ $\text{cm}^3 \text{s}^{-1}$	$k_2/10^{-10}$ $\text{cm}^3 \text{s}^{-1}$
Taatjes et al. ²⁵	LFP/PIMS	351	298	4	He		0.7–5	2.4 ± 0.3	0.67 ± 0.10
Smith et al. ²⁸	LFP/UV abs	248	295	7.5–500	N ₂	1300	155–600	2.0 ± 0.3	
Sheps et al. ²⁶	LFP/UV abs	266	293	10	He	1.5	0.8–4.8	2.9 ± 0.3	2.2 ± 0.2
Howes et al. ⁵	LFP/PIMS	248	295	1–2.5	N ₂	1–10	2–9	1.7 ± 0.3	
Zhou et al. ²⁷	LFP/LIF	248	298	10	Ar	1.9–10.4	0.3–2.2	2.5 ± 0.2	

^aLFP = laser flash photolysis, PIMS = photoionization mass spectrometry, UV abs = ultraviolet absorption.

chemistry of SCIs other than CH₂OO.^{16–19} Measurements in a boreal forest in Finland led to the suggestion that reactions of SCIs with SO₂ may have been responsible for up to 50% of the H₂SO₄ observed in the gas phase,^{16,17} while observations in a rural location in Germany have indicated that SCI + SO₂ reactions could be responsible for up to 80% of H₂SO₄ produced at night.¹⁷ Similarly, field experiments in Texas, United States, have suggested that nighttime production of H₂SO₄ is dominated by SCI + SO₂ reactions, with potentially important contributions in the afternoon,¹⁸ and agreement between observations and model predictions for sulfate aerosol over the Southeast of the US has shown improvement when SCI + SO₂ chemistry is included in the model.²⁰ The role of SCI + SO₂ chemistry in the atmosphere has also been investigated using observations made on Corsica, in the Mediterranean, where it was found that SCI + SO₂ reactions could be responsible for 10% of the observed H₂SO₄ production during the day and 40% at night.¹⁹ Potential impacts of SCI reactions with SO₂ have also been reported in vehicle exhausts²¹ and power plant plumes.²² However, there have been few measurements of the kinetics of SO₂ oxidation by SCIs other than CH₂OO, and uncertainties in rate coefficients for SCI + SO₂ reactions have been highlighted as a constraint on our understanding of the atmospheric impacts of SCIs and the interpretation of field measurements of H₂SO₄.^{16–20,23}

The SCI CH₃CHOO exists in two conformers, *syn*-CH₃CHOO and *anti*-CH₃CHOO, which are separated by a significant barrier to interconversion ($\sim 160 \text{ kJ mol}^{-1}$)²⁴ and thus behave as distinct species under ambient conditions.^{25,26} The first direct measurements of the reaction kinetics of CH₃CHOO conformers, made using laser flash photolysis of CH₃CHI₂ in the presence of excess O₂ coupled with tunable synchrotron photoionization mass spectrometry (PIMS), demonstrated rapid reactions with SO₂ (R1 and R2).²⁵



The PIMS experiments were performed at 298 K and a total pressure of 4 Torr in He, giving $k_1 = (2.4 \pm 0.3) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ and $k_2 = (6.7 \pm 1.0) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$. Formation of SO₃ was observed, with a rate that suggested direct production from reactions of CH₃CHOO conformers with SO₂. Subsequent experiments using the PIMS technique at a fixed ionization energy gave a value for k_1 of $(1.7 \pm 0.3) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ at 295 K and total pressures between 1 and 2.5 Torr in N₂, with measurements indicating production of acetaldehyde (CH₃CHO) from R1 at a yield of (0.86 ± 0.11) at 2 Torr.⁵ The reaction of *syn*-CH₃CHOO with SO₂ has also been investigated by monitoring the kinetics of OH radical production from the decomposition of *syn*-CH₃CHOO

occurring in competition with R1, giving $k_1 = (2.5 \pm 0.2) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ at 298 K and 10 Torr in Ar.²⁷

Experiments using laser flash photolysis of CH₃CHI₂/O₂ mixtures with broadband UV absorption spectroscopy have also indicated that R1 and R2 are rapid.^{26,28} A rate coefficient of $(2.0 \pm 0.3) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ at 295 K and pressures between 7.5 and 500 Torr of N₂ was reported from experiments in which the conformer-specific contributions to the total absorbance were not resolved.²⁸ However, the result is expected to be dominated by *syn*-CH₃CHOO on the basis of results from the earlier PIMS experiments,²⁵ which indicated that *syn*-CH₃CHOO represents 90% of the total CH₃CHOO produced using the photolytic method. Conformer-specific measurements using broadband UV absorption spectroscopy have been achieved in experiments performed at 293 K and a total pressure of 10 Torr in He, giving $k_1 = (2.9 \pm 0.3) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ and $k_2 = (2.2 \pm 0.2) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$.²⁶ The conformer-specific UV experiments indicated that *syn*-CH₃CHOO is the dominant conformer produced,²⁶ in agreement with the earlier PIMS experiments,²⁵ although a lower yield of 70% was reported, which may result from the different experimental conditions or uncertainties in the UV absorption cross-sections, particularly for *anti*-CH₃CHOO.²⁶

There are discrepancies in the literature for values of k_1 and k_2 , but studies so far have all taken place at room temperature over a relatively narrow range of pressures (Table 1). Significant conformer dependence is shown for the reactivity of CH₃CHOO with SO₂,^{25,26} with studies also showing distinct conformer-dependent reactivity for reactions of asymmetric CIs with H₂O^{25,26,29} and acids,³⁰ as well as differences in their decomposition rates.³¹ CH₃CHOO is the simplest Criegee intermediate that exists as two conformers and can therefore be used as a prototype to characterize the reactions of the larger CIs, which requires rate coefficients to be well established across a range of conditions. In this work, we report the kinetics of R1 and R2 at temperatures between 242 and 353 K and pressures between 10 and 600 Torr determined using time-resolved broadband UV absorption spectroscopy.

EXPERIMENTAL SECTION

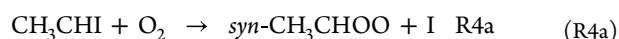
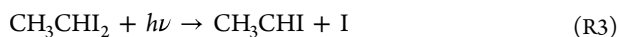
The kinetics of R1 and R2 were studied as a function of temperature and pressure using laser flash photolysis of CH₃CHI₂/O₂/N₂/SO₂ mixtures coupled with time-resolved broadband UV absorption spectroscopy. The experimental apparatus has been described in detail in previous work,^{6,31} and only a brief description is given here.

A dilute mixture of a known concentration of SO₂ (Sigma-Aldrich, 99.9%) was prepared manometrically in N₂ (BOC, 99.998%) and stored in a glass bulb before mixing in a gas manifold with N₂ (BOC, 99.998%) and O₂ (BOC, 99.5%) at known flow rates controlled by calibrated mass flow controllers (MKS Instruments). A known fraction of the total gas flow,

controlled by a needle valve, was then passed through a bubbler containing liquid CH_3CHI_2 (SynHet, 90%) held at a constant temperature in a water bath before being recombined with the rest of the gas flow and passed into a jacketed Pyrex reaction cell. Experiments were performed under pseudo-first-order conditions, with the concentrations of SO_2 in large excess over initial CH_3CHOO concentrations. Concentrations were varied in the range $[\text{CH}_3\text{CHI}_2] = (2.8\text{--}6.0) \times 10^{13} \text{ cm}^{-3}$, $[\text{O}_2] = (0.6\text{--}20) \times 10^{17} \text{ cm}^{-3}$, and $[\text{SO}_2] = (0.4\text{--}5.0) \times 10^{13} \text{ cm}^{-3}$, with typical initial $[\text{CH}_3\text{CHOO}]$ on the order of 10^{12} cm^{-3} .

The reaction cell was 100 cm in length and 3 cm in diameter and sealed with fused silica windows at each end. The temperature of the cell was controlled by flowing liquid from a recirculating thermostating unit (Huber Unistat 360) through the jacket surrounding the cell and calibrated by measuring the temperature of a flow of N_2 , under conditions identical to those used in kinetics experiments, at 5 cm increments along the length of the cell using a K-type thermocouple.^{6,31} Pressure in the cell was controlled by a rotary pump (EM2, Edwards) by throttling the exit to the reaction cell and measured by a capacitance manometer (MKS Instruments). The total flow rate through the cell was set to an equivalent of 1200 standard $\text{cm}^3 \text{ min}^{-1}$ (sccm) at 50 Torr and adjusted with pressure to maintain a constant residence time in the cell of ~ 2.6 s.

An excimer laser (KrF, Lambda-Physik CompEx 210) with output at $\lambda = 248$ nm and typical fluence of $30\text{--}40 \text{ mJ cm}^{-2}$ was aligned along the length of the reaction cell using a dichroic turning mirror (Edmund Optics) and used to initiate production of *syn*- and *anti*- CH_3CHOO in the cell via reactions R3 and R4a.



A delay generator (SRS DG535) was used to control the timing of the laser, which was operated with a pulse repetition frequency of 0.33 Hz to ensure that the gas mixture in the cell was replaced between each laser pulse.

Absorbing species in the cell were monitored by UV/visible radiation provided by a laser-driven light source (LDLS, Energetiq EQ-99X), which provided $\sim 10 \text{ mW cm}^{-2}$ of light with a near constant radiance from 200 to 800 nm. The LDLS output was collimated by an off-axis parabolic mirror (ThorLabs) and aligned through the reaction cell in a multipass arrangement consisting of ten mirrors (Knight Optical), each of 12 mm diameter, resulting in an effective path length of (595 ± 53) cm for the experiments described in this work, which was determined using the method in our earlier work.⁶ Light exiting the cell was passed through a sharp cut-on filter (248 nm RazorEdge ultrasteep long-pass edge filter) to reduce the impact of scattered 248 nm light and focused into a fiber optic via a fiber launcher (Elliot Scientific). Light exiting the fiber optic was directed through a 25 μm slit onto a diffraction grating with 600 grooves/mm and imaged onto a thermoelectrically cooled charge-coupled device (CCD) detector (FER-SCI-1024BRX, Princeton Instruments). Photocharge generated on the CCD was shifted from an illuminated region to a storage region shielded from incoming radiation at set time intervals throughout the reaction, with the experimental setup used in this work giving a spectral

resolution of ~ 1 nm and a temporal resolution between 70 and 100 μs . Intensity data were typically recorded for 500 photolysis shots and transferred to a PC for analysis.

RESULTS

Absorbance spectra were determined from measured intensity data and related to the concentration of each species present using the Beer–Lambert law (eq 1)

$$A_{\lambda,t} = \ln\left(\frac{I_{\lambda,0}}{I_{\lambda,t}}\right) = \sum_i \sigma_{i,\lambda} c_{i,t} l \quad (1)$$

where $A_{\lambda,t}$ is the total absorbance at wavelength λ and time t , $I_{\lambda,0}$ is the average pre-photolysis light intensity at wavelength λ , $I_{\lambda,t}$ is the post-photolysis light intensity at wavelength λ and time t , $\sigma_{i,\lambda}$ is absorption cross-section of species i at wavelength λ , $c_{i,t}$ is the concentration of species i at time t , and l is the effective path length, which has a value of (595 ± 53) cm for experiments reported in this work.

Figure 1 shows the typical absorbance measured following photolysis, which contains contributions from CH_3CHI_2 , *syn*-

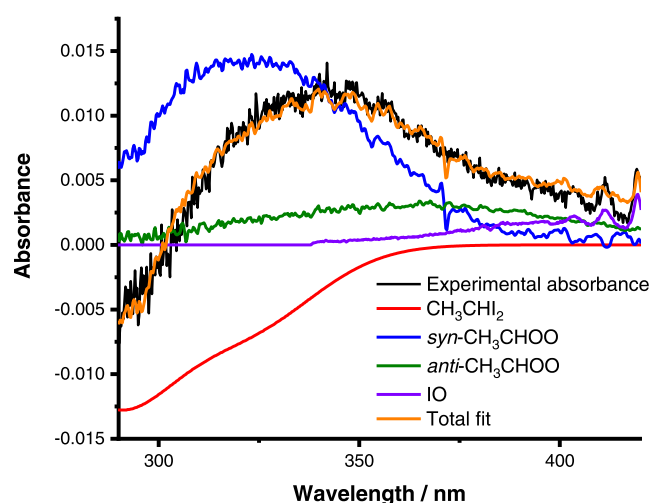


Figure 1. Observed absorbance (black), total fit (orange), and individual contributions from *syn*- CH_3CHOO ²⁶ (blue), *anti*- CH_3CHOO ²⁶ (green), CH_3CHI_2 ³² (red), and IO ³³ (purple) determined by fitting reference spectra to the observed absorbance using eq 1. For these data, $T = 353$ K, $p = 10$ Torr, $t = 1$ ms, $[\text{CH}_3\text{CHI}_2] = 6.0 \times 10^{13} \text{ cm}^{-3}$. The fit to the observed absorbance for these data gave $\Delta[\text{CH}_3\text{CHI}_2] = (5.6 \pm 0.07) \times 10^{12} \text{ cm}^{-3}$, $[\textit{syn}\text{-CH}_3\text{CHOO}] = (2.3 \pm 0.11) \times 10^{12} \text{ cm}^{-3}$, $[\textit{anti}\text{-CH}_3\text{CHOO}] = (7.6 \pm 0.27) \times 10^{11} \text{ cm}^{-3}$, and $[\text{IO}] = (5.5 \pm 0.14) \times 10^{11} \text{ cm}^{-3}$.

and *anti*- CH_3CHOO , and IO radicals produced by secondary chemistry within the system. Reference spectra for CH_3CHI_2 ,³² *syn*- and *anti*- CH_3CHOO ,²⁶ and IO ³³ were fit to the observed absorbance at each time point to determine the concentration of each species throughout the reaction. While absolute concentrations are reported here, it should be noted that uncertainties in the effective path length and absorption cross-sections do not contribute to uncertainties in measured kinetics for the pseudo-first-order conditions employed in this work.

Figure 2 shows the concentration–time profiles for *syn*- and *anti*- CH_3CHOO in the presence of SO_2 , which were each fit according to a first-order kinetic loss (eq 2) convoluted with a Gaussian instrument response function (IRF) to describe the

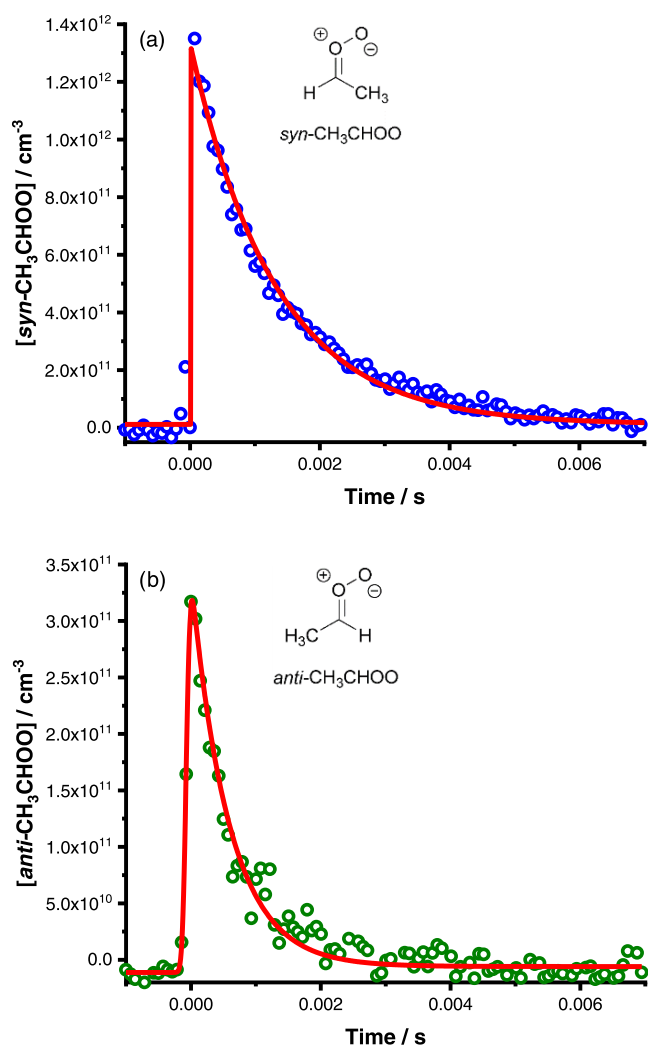


Figure 2. Observed concentration–time profiles for (a) *syn*-CH₃CHOO and (b) *anti*-CH₃CHOO. For these data, $T = 298$ K, $p = 50$ Torr, $[\text{SO}_2] = 1.1 \times 10^{13}$ cm⁻³, and $[\text{CH}_3\text{CHI}_2] = 2.8 \times 10^{13}$ cm⁻³. The fits to eq 2 (coupled with the instrument response function as detailed in the Supporting Information) (solid lines) gave an initial concentration of $(1.31 \pm 0.03) \times 10^{12}$ cm⁻³ and $k' = (765 \pm 15)$ s⁻¹ for *syn*-CH₃CHOO and an initial concentration of $(3.39 \pm 0.02) \times 10^{11}$ cm⁻³ and $k' = (2280 \pm 218)$ s⁻¹ for *anti*-CH₃CHOO. Instrument response parameters were $w = (2.99 \pm 0.10) \times 10^{-5}$ s and $t_c = -(4.80 \pm 0.09) \times 10^{-5}$ s for both conformers. Uncertainties are 1σ .

shifting of photocharge on the CCD detector (see the Supporting Information for further details).

$$C_t = C_0 \exp(-k't) \quad (2)$$

where C_t is the concentration of *syn*- or *anti*-CH₃CHOO at time t , C_0 is the initial concentration of the Criegee intermediate conformer, and k' is the rate coefficient describing the sum of first-order losses of the CH₃CHOO conformer and is given by $k' = k_x + k_1[\text{SO}_2]$ for *syn*-CH₃CHOO and $k' = k_x + k_2[\text{SO}_2]$ for *anti*-CH₃CHOO, where k_x represents losses of *syn*- or *anti*-CH₃CHOO via any reaction or process other than the reaction with SO₂. Unimolecular decomposition and bimolecular reactions with the CH₃CHI₂ precursor contribute significantly to k_x for both *syn*- and *anti*-CH₃CHOO,³¹ with potential additional contributions from reactions with iodine atoms, IO, or Criegee–Criegee chemistry as well as diffusion out of the probe region.

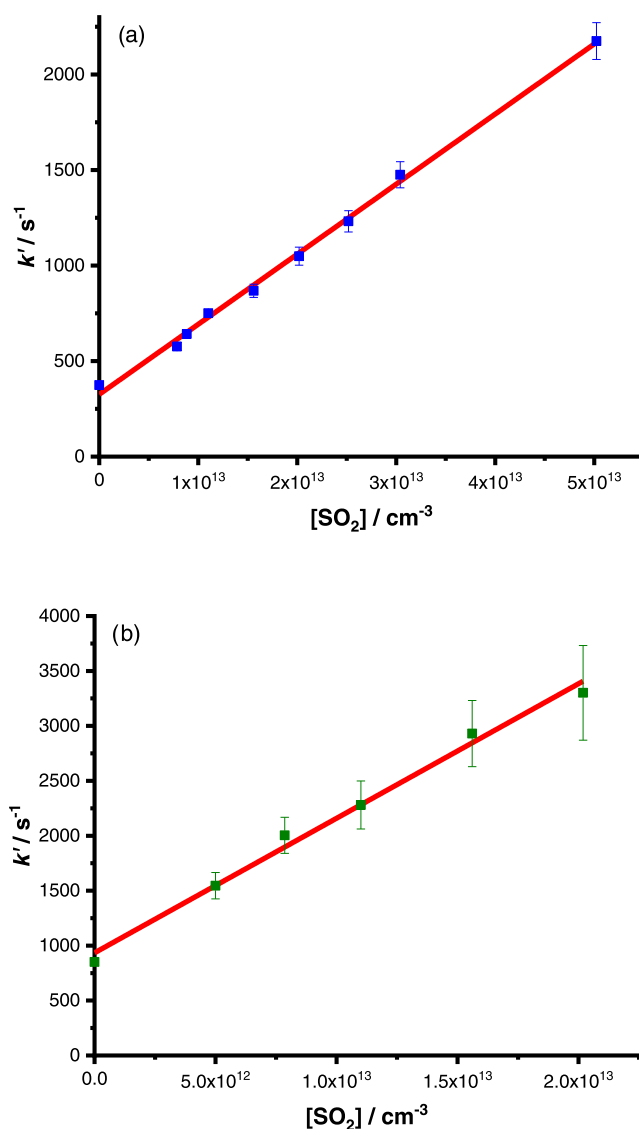


Figure 3. Dependence of k' on $[\text{SO}_2]$ at $T = 298$ K and $p = 50$ Torr for (a) *syn*-CH₃CHOO and (b) *anti*-CH₃CHOO. Fits to the data (solid lines) gave $k_1 = (3.67 \pm 0.07) \times 10^{-11}$ cm³ s⁻¹, with an intercept k_x of (326 ± 18) s⁻¹, and $k_2 = (1.22 \pm 0.10) \times 10^{-10}$ cm³ s⁻¹, with an intercept k_x of (934 ± 70) s⁻¹. Uncertainties are 1σ .

Rate coefficients k_1 and k_2 were determined from the dependence of k_1' and k_2' on $[\text{SO}_2]$, respectively, with the typical results shown in Figure 3. Potential impacts of second-order losses for the CH₃CHOO conformer through reactions such as CH₃CHOO + CH₃CHOO or CH₃CHOO + I were also investigated by fitting concentration–time profiles to mixed first- and second-order kinetic losses convoluted with the IRF. Results for k_1 and k_2 obtained from the mixed-order fits were within 5% of those obtained from the first-order fits. Further details are given in the Supporting Information. All results reported here were obtained from first-order fits (eq 2).

Figure 4 shows the results for k_1 , which are summarized in Table 2. At 298 K, the results demonstrate an increase in k_1 from $(3.02 \pm 0.32) \times 10^{-11}$ cm³ s⁻¹ at 10 Torr to $(4.66 \pm 0.52) \times 10^{-11}$ cm³ s⁻¹ at 600 Torr, where the uncertainties represent a combination of the statistical error and the systematic errors resulting from uncertainties in gas flow rates and in the concentration of SO₂, with results at other

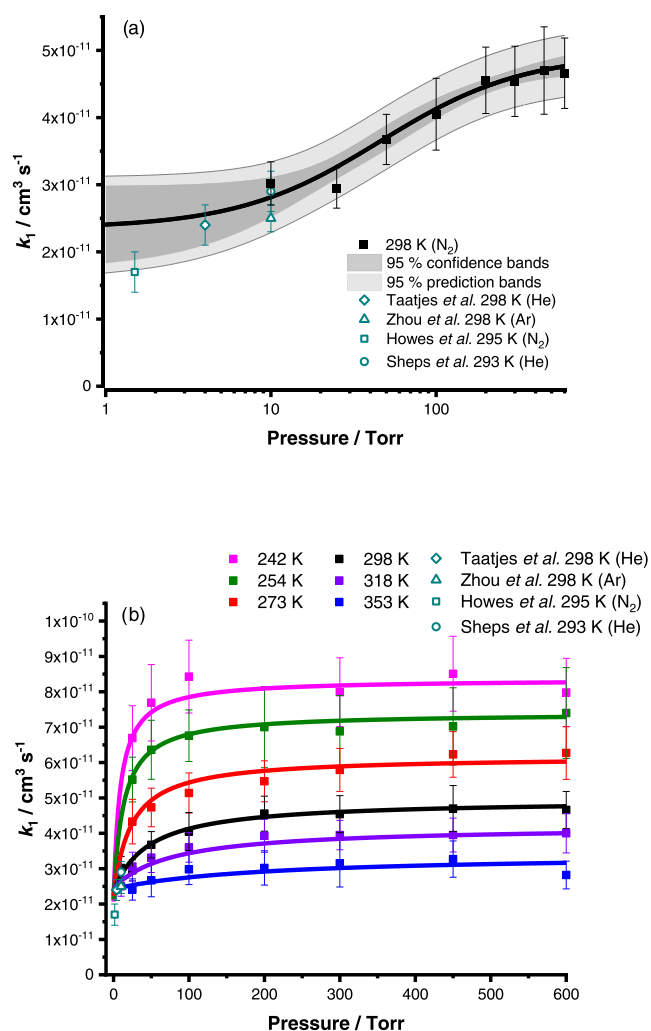


Figure 4. Effects of pressure on k_1 at (a) 298 K and (b) all temperatures studied in this work. Solid lines show the fits to k_1 using eqs 3–6 (which were performed globally using all data for k_1 measured in this work). Previous results reported for k_1 are also shown.^{5,25–27} Error bars represent a combination of the statistical error and the systematic errors resulting from uncertainties in gas flow rates and in the concentration of SO_2 .

temperatures also showing significant pressure dependence and overall negative temperature dependence. Equations 3–6, which describe a chemical activation mechanism with a non-zero rate coefficient at zero pressure,³⁴ were fit globally to results obtained in this work for k_1 over all temperatures and pressures to provide a parametrization for use in atmospheric models.

$$k_1 = k_{\text{int}} + \frac{(k_{\infty} - k_{\text{int}})k_0[M]}{(k_0[M]) + k_{\infty}} \quad (3)$$

where k_{int} represents the rate coefficient at zero pressure, k_0 is the low-pressure limiting rate coefficient, and k_{∞} is the high-pressure limiting rate coefficient, and these are given by eqs 4–6

$$k_{\text{int}} = A_{\text{int}} \times \left(\frac{T}{298} \right)^{n_{\text{int}}} \quad (4)$$

$$k_0 = A_0 \times \left(\frac{T}{298} \right)^{n_0} \quad (5)$$

$$k_{\infty} = A_{\infty} \times \left(\frac{T}{298} \right)^{n_{\infty}} \quad (6)$$

Fit results gave $A_{\text{int}} = (2.35 \pm 0.39) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$, $n_{\text{int}} = (0.61 \pm 0.79)$, $A_0 = (3.29 \pm 1.30) \times 10^{-29} \text{ cm}^3 \text{ s}^{-1}$, $n_0 = -(9.52 \pm 1.78)$, $A_{\infty} = (4.95 \pm 0.51) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$, and $n_{\infty} = -(2.52 \pm 0.29)$.

The fits of eqs 3–6 indicate a value for k_1 of $(4.80 \pm 0.46) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ at 298 K and 760 Torr. The pressure dependence observed in this work reconciles discrepancies between values for k_1 reported at room temperature in previous work^{5,25–27} at pressures below 10 Torr, as shown in Figure 4. While kinetics reported by Smith et al.²⁸ at 295 K over the pressure range of 7.5–500 Torr are in broad agreement with low pressure values for k_1 reported in this work and in previous work, Smith et al. were unable to distinguish between the *syn*- and *anti*-conformers, so the rate coefficient reported will contain contributions from the reactivity of both *syn*- CH_3CHOO and *anti*- CH_3CHOO .

Figure 5 and Table 2 summarize the results obtained in this work for k_2 . In contrast to the results for k_1 , no significant dependence of k_2 on the temperature or pressure was observed. At 298 K, results gave a mean value for k_2 of $(1.15 \pm 0.16) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$ between 10 and 600 Torr, with results over all temperatures and pressures giving a mean value of $(1.18 \pm 0.21) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$. The effect of pressure on k_2 at each temperature is shown in Figure S3 in the Supporting Information.

The kinetics of R2 have been reported in two previous studies^{25,26} at room temperature. Taatjes et al.²⁵ performed experiments at 4 Torr using the PIMS technique and reported a value for k_2 of $(6.7 \pm 1.0) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$, while Sheps et al.²⁶ performed experiments at 10 Torr using cavity-enhanced UV absorption spectroscopy and reported a value for k_2 of $(2.2 \pm 0.2) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$. Differences between the studies reflect the challenges associated with measuring such rapid kinetics, with the lack of dependence of k_2 on temperature and pressure observed in this work potentially indicating that the kinetics for R2 are controlled by collision-limited or capture-limited kinetics. The difference in behavior between the *syn*- and *anti*-conformers is potentially influenced by lower steric hindrance for the *anti*-conformer, coupled with the higher ground state energy for *anti*- CH_3CHOO by $\sim 15 \text{ kJ mol}^{-1}$ ²⁴ compared to *syn*- CH_3CHOO and a higher dipole moment for *anti*- CH_3CHOO than *syn*- CH_3CHOO (5.53 D compared to 4.69 D, calculated at the B3LYP/AVTZ level of theory³⁵).

Figure 6 compares the experimental results for k_2 with estimated values using a collision model (eq 7) and a capture model (eq 8).

$$k_{\text{col}} = \pi(r_{\text{Cl}} + r_{\text{SO}_2})^2 \sqrt{\frac{8k_{\text{B}}T}{\pi\mu}} \quad (7)$$

where r_{Cl} and r_{SO_2} are the effective radii of *anti*- CH_3CHOO ³⁶ and SO_2 ,³⁷ respectively, k_{B} is the Boltzmann constant, T is the temperature, and μ is the reduced mass. The effective radius for *anti*- CH_3CHOO was assumed to be the same as that reported in the literature for *syn*- CH_3CHOO .³⁶

Table 2. Summary of Results for k_1 and k_2 ^a

T/K	p/Torr	[CH ₃ CH ₂]/10 ¹³ cm ⁻³	[SO ₂]/10 ¹³ cm ⁻³	$k_1/10^{-11}$ cm ³ s ⁻¹	$k_2/10^{-10}$ cm ³ s ⁻¹
242	25	3.0–4.1	0.6–3.1	6.69 ± 0.92	1.13 ± 0.13
	50			7.69 ± 1.08	1.17 ± 0.22
	100			8.43 ± 1.03	1.06 ± 0.17
	300			7.99 ± 0.97	1.06 ± 0.19
	450			8.51 ± 1.06	1.16 ± 0.25
	600			7.98 ± 0.97	
254	25	5.2–5.9	0.9–3.2	5.52 ± 0.63	0.90 ± 0.14
	50			6.36 ± 0.83	1.01 ± 0.20
	100			6.76 ± 0.73	1.00 ± 0.10
	200			7.00 ± 1.14	1.00 ± 0.18
	300			6.88 ± 1.02	
	450			7.02 ± 1.09	
273	25	3.4–4.3	0.6–2.7	4.33 ± 0.63	1.29 ± 0.14
	50			4.73 ± 0.55	1.50 ± 0.28
	100			5.14 ± 0.57	1.54 ± 0.17
	200			5.47 ± 0.58	1.56 ± 0.18
	300			5.79 ± 0.61	
	450			6.23 ± 0.64	1.32 ± 0.21
298	10	2.8–3.7	0.5–5.0	6.27 ± 0.74	
	25			3.02 ± 0.32	1.04 ± 0.13
	50			3.00 ± 0.30	1.33 ± 0.17
	100			3.67 ± 0.37	1.22 ± 0.16
	200			4.05 ± 0.54	1.19 ± 0.17
	300			4.55 ± 0.49	1.08 ± 0.14
318	10	4.1–5.3	0.7–4.1	4.54 ± 0.52	1.10 ± 0.20
	25			4.70 ± 0.65	1.26 ± 0.16
	50			4.66 ± 0.52	0.96 ± 0.16
	100			2.50 ± 0.48	1.22 ± 0.22
	200			3.06 ± 0.41	1.42 ± 0.15
	300			3.31 ± 0.42	1.20 ± 0.21
353	10	5.1–6.0	0.4–2.4	3.60 ± 0.42	1.05 ± 0.17
	25			3.93 ± 0.47	1.32 ± 0.22
	50			3.95 ± 0.42	1.41 ± 0.17
	100			3.95 ± 0.48	1.11 ± 0.20
	200			4.00 ± 0.56	1.08 ± 0.26
	300			2.41 ± 0.30	1.29 ± 0.20
	25			2.41 ± 0.30	1.09 ± 0.17
	50			2.67 ± 0.47	1.22 ± 0.25
	100			2.98 ± 0.43	1.10 ± 0.28
	200			3.02 ± 0.48	1.17 ± 0.19
	300			3.15 ± 0.67	1.13 ± 0.13
	450			3.27 ± 0.51	1.40 ± 0.25
	600			2.82 ± 0.39	0.88 ± 0.12

^aUncertainties represent a combination of the 1σ statistical error and the systematic errors resulting from uncertainties in gas flow rates and in the concentration of SO₂.

$$k_{\text{capt}} = C \sqrt{\frac{\pi}{\mu}} (D_{\text{Cl}} D_{\text{SO}_2})^{2/3} (k_B T)^{-1/6} \quad (8)$$

where C is a constant (4.08 for the case of isotropic capture)^{35,38} and D_{Cl} and D_{SO_2} are the dipole moments of *anti*-CH₃CHOO³⁵ and SO₂,³⁹ respectively.

The experimental results for k_2 obtained in this work are lower than the estimated rate coefficients using either the collision model or the capture model, with experimental values a factor of ~2 lower than those calculated from collision theory and a factor of ~6 lower than those calculated from capture theory. However, the calculated values do offer some insight

into the kinetics and suggest that R2 is close to the collision limit.

The reaction between *syn*-CH₃CHOO and SO₂ has been investigated using theoretical approaches, which indicate a barrierless reaction with a 98% yield of acetaldehyde (CH₃CHO) + SO₃ at 298 K and 200 Torr of He and a rate coefficient for CH₃CHO + SO₃ production of 4.49 × 10⁻¹¹ cm³ s⁻¹ at 298 K.¹³ However, the possible impacts of pressure were not fully discussed, and the reaction of *anti*-CH₃CHOO + SO₂ was not considered. The calculations¹³ predicted a positive temperature dependence for reactions of CH₂OO, *syn*-CH₃CHOO, and (CH₃)₂COO with SO₂, despite the reactions being barrierless, and this is in contrast to the

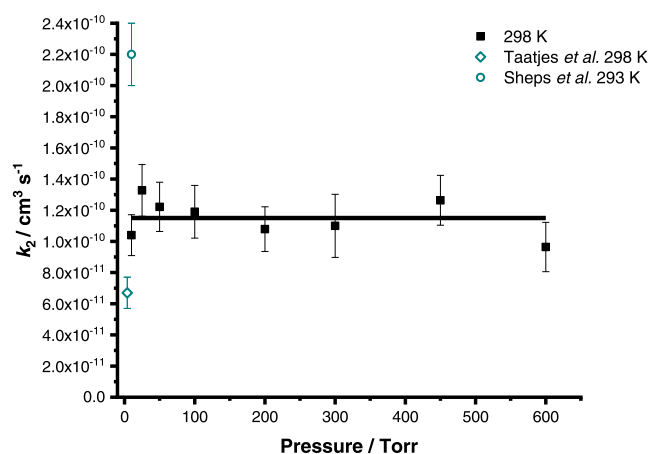


Figure 5. Effects of pressure on k_2 at 298 K. The solid line shows the mean value for k_2 at 298 K ($(1.15 \pm 0.16) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$). Previous results reported for k_2 are also shown. Error bars represent a combination of the statistical error and the systematic errors resulting from uncertainties in gas flow rates and in the concentration of SO_2 .

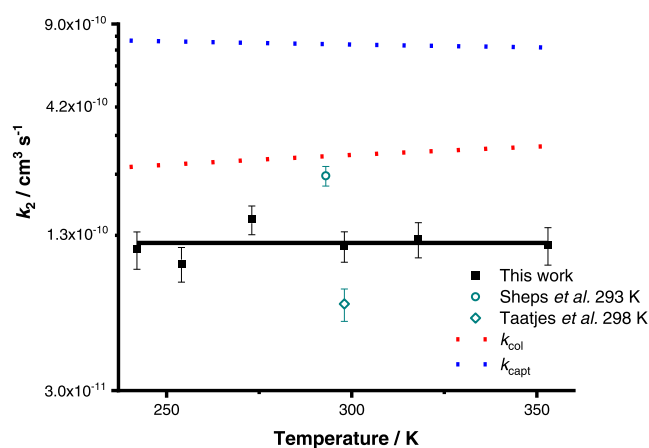


Figure 6. Mean values for k_2 determined at each temperature. The solid line represents that the mean value for k_2 over all conditions investigated in this work is $(1.18 \pm 0.21) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$. Previous results reported for k_2 and rate coefficients calculated using collision theory (k_{col} , red dashed line) and capture theory (k_{capt} , blue dashed line) are also shown. Error bars represent a combination of the statistical error and the systematic errors resulting from uncertainties in gas flow rates and in the concentration of SO_2 .

experimental results for $\text{syn-CH}_3\text{CHOO} + \text{SO}_2$ obtained in this work, our previous experiments for $\text{CH}_2\text{OO} + \text{SO}_2$,⁶ and experimental results for $(\text{CH}_3)_2\text{COO} + \text{SO}_2$ ⁴⁰ (see the Supporting Information for further details). Where potential impacts of pressure have been considered in detail in theoretical studies of $\text{SCI} + \text{SO}_2$ reactions, there is an agreement with the lack of observed pressure dependence in the kinetics for $\text{CH}_2\text{OO} + \text{SO}_2$ under atmospheric conditions,^{3,4,11,12} but there are differences in the predicted pressure dependence of the reaction between $(\text{CH}_3)_2\text{COO}$ and SO_2 .^{11,12} Vereecken et al. suggested that >80% of the SOZ formed by $(\text{CH}_3)_2\text{COO} + \text{SO}_2$ undergoes prompt decomposition to acetone ($\text{CH}_3\text{C}(\text{O})\text{CH}_3$) and SO_3 at 298 K and a pressure of 4 Torr, while >97% of the SOZ collisionally stabilizes at 298 K and 760 Torr, with the difference compared to $\text{CH}_2\text{OO} + \text{SO}_2$ attributed to the greater number of degrees of freedom in the SOZ formed via $(\text{CH}_3)_2\text{COO} + \text{SO}_2$, which

would also be relevant to the comparison between the SOZ formed via $\text{CH}_2\text{OO} + \text{SO}_2$ and those from reactions of CH_3CHOO conformers with SO_2 . However, Kuwata et al. calculated a different potential energy surface for the reaction between $(\text{CH}_3)_2\text{COO}$ and SO_2 compared to that reported by Vereecken et al., and thus a different mechanism for the reaction, with calculations predicting no significant collisional stabilization of the SOZ at 298 K and pressures below 10^4 Torr and SO_3 yields greater than 96% at 298 K and pressures from 1 to 760 Torr. Experimental measurements of the kinetics for $(\text{CH}_3)_2\text{COO} + \text{SO}_2$ have indicated significant pressure dependence and negative temperature dependence under atmospheric conditions,^{40–42} similar to the observations in this work for the reaction between $\text{syn-CH}_3\text{CHOO}$ and SO_2 . Differences between theoretical approaches and between experiments and theory indicate that the application of theory to the prediction of SCI kinetics remains a challenge.

■ ATMOSPHERIC IMPLICATIONS

The atmospheric impacts of CH_3CHOO conformer reactions with SO_2 depend on the competition with other CH_3CHOO conformer reactions, which are expected to be dominated by unimolecular decomposition for $\text{syn-CH}_3\text{CHOO}$ and reaction with water vapor for $\text{anti-CH}_3\text{CHOO}$.^{31,43–45} Figure 7 compares the pseudo-first-order losses for CH_3CHOO conformers through unimolecular decomposition and reactions with SO_2 and water vapor for a range of SO_2 and water vapor concentrations as a function of temperature at 760 Torr. Rate coefficients for unimolecular decomposition (k_{dec}) were taken from our recent work,³¹ and those for reactions with SO_2 (k_{SO_2}) were taken from those determined in this work. Rate coefficients for reactions with water vapor ($k_{\text{H}_2\text{O}}$ and $k_{(\text{H}_2\text{O})_2}$) were based on the upper limit for $\text{syn-CH}_3\text{CHOO} + \text{H}_2\text{O}$ reported by Sheps et al. at 298 K,²⁶ which forms the basis of the current IUPAC recommendation,³ and temperature-dependent measurements for $\text{anti-CH}_3\text{CHOO} + \text{H}_2\text{O}$ and $\text{anti-CH}_3\text{CHOO} + (\text{H}_2\text{O})_2$ reported by Lin et al.⁴⁵ Water dimer concentrations were calculated from the monomer concentration using equilibrium constants reported by Ruscic et al.⁴⁶ There are no current reports of rate coefficients or upper limits for a possible reaction of $\text{syn-CH}_3\text{CHOO}$ with water dimers, and it should be noted that current IUPAC recommendations³ for $\text{anti-CH}_3\text{CHOO}$ reactions with water vapor do not extend beyond 298 K, owing to uncertainties in temperature-dependent measurements, which will impact the analysis shown in Figure 7. For $\text{anti-CH}_3\text{CHOO}$, results show that the reaction with water vapor will dominate under all conditions relevant to the troposphere, but chamber studies employing high SO_2 concentrations and low humidity will need to consider the impact of R2. For $\text{syn-CH}_3\text{CHOO}$, the reaction with SO_2 will be competitive with other losses in the atmosphere in areas with high SO_2 concentrations and low humidity, particularly at low temperatures, contributing to the atmospheric oxidation of SO_2 . The pressure dependence of k_1 indicates that there may be significant collisional stabilization of the SOZ produced in the reaction between $\text{syn-CH}_3\text{CHOO}$ and SO_2 , potentially limiting the production of SO_3 and subsequently H_2SO_4 . However, the fate of the SOZ is uncertain, and even if there is significant stabilization of the SOZ, it may still contribute to the atmospheric production of H_2SO_4 through subsequent chemistry. A more detailed assessment of the atmospheric impacts of CI reactions with

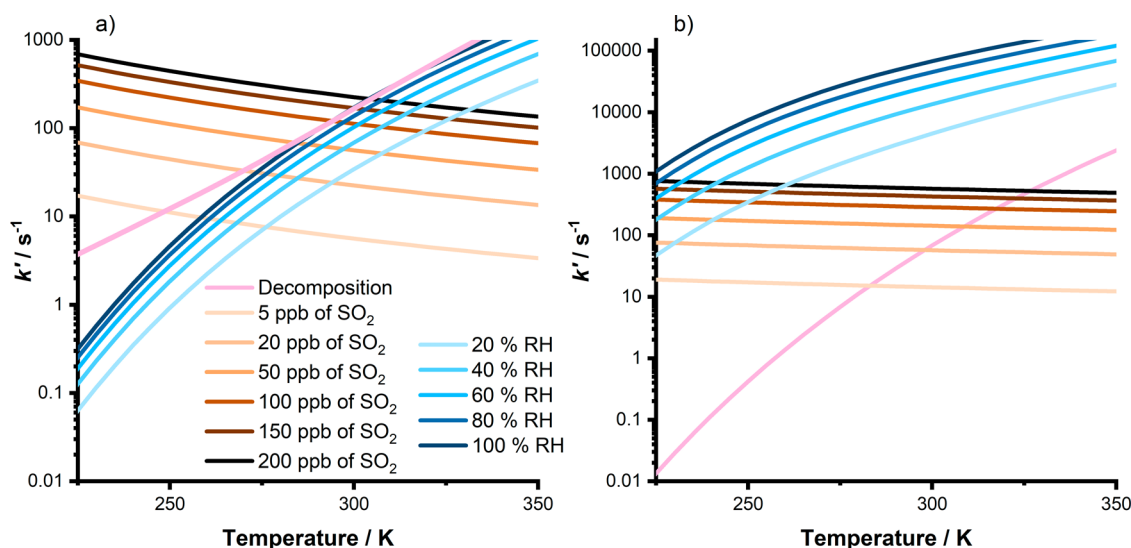


Figure 7. Pseudo-first-order losses of (a) *syn*-CH₃CHOO and (b) *anti*-CH₃CHOO through reaction with SO₂ (orange), reaction with water vapor (blue), and unimolecular decomposition (pink) at a total pressure of 760 Torr. Pseudo-first-order losses were calculated as described in the main text.

SO₂ would benefit from further experimental investigation of the nature and yields of the products, particularly as a function of pressure.

CONCLUSIONS

The kinetics of *syn*- and *anti*-CH₃CHOO reactions with SO₂ have been investigated in the temperature range from 242 to 353 K at pressures between 10 and 600 Torr using laser flash photolysis of CH₃CHI₂/O₂/N₂/SO₂ gas mixtures coupled with time-resolved broadband UV absorption spectroscopy.

Results for *syn*-CH₃CHOO + SO₂ show that the kinetics are pressure-dependent, with a negative dependence on temperature. The kinetics can be parametrized by a model that indicates a role for chemical activation, which gives a rate coefficient of $k_1 = (4.80 \pm 0.46) \times 10^{-11} \text{ cm}^3 \text{ s}^{-1}$ at 298 K and 760 Torr. The observed pressure dependence reconciles apparent discrepancies in previous measurements of *syn*-CH₃CHOO + SO₂ kinetics performed at ~ 298 K but at different pressures.

Kinetics of the reaction between *anti*-CH₃CHOO and SO₂ display no significant dependence on temperature or pressure over the ranges investigated. Results give a mean value for k_2 of $(1.15 \pm 0.16) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$ at 298 K and $(1.18 \pm 0.21) \times 10^{-10} \text{ cm}^3 \text{ s}^{-1}$ over all conditions studied in this work.

Comparisons with unimolecular decomposition kinetics of *syn*- and *anti*-CH₃CHOO and reactions with water vapor under typical atmospheric conditions indicate that the reaction with SO₂ will play an enhanced role in the removal of the *syn*-CH₃CHOO in areas of low humidity and at low temperatures and the removal of *anti*-CH₃CHOO is dominated by its reaction with water vapor under all conditions relevant to the troposphere.

ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acs.jpca.4c00199>.

Description of the instrument response function, comparison of first- and mixed-order analysis, effects of

pressure on k_2 across the studied temperature range, comparison to theoretical calculations, and summary of experimental data (PDF)

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Notes

The authors declare no competing financial interest.

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