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LETTER TO THE EDITOR

# X-ray properties of the white dwarf pulsar eRASSU J191213.9–441044★

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## ABSTRACT

We report X-ray observations of the newly discovered pulsating white dwarf eRASSU J191213.9–441044 with Spectrum Roentgen Gamma and eROSITA (SRG/eROSITA) and *XMM-Newton*. The new source was discovered during the first eROSITA all-sky survey at a flux level of  $f_X(0.2\text{--}2.3\text{ keV}) = 3.3 \times 10^{-13}\text{ erg cm}^{-2}\text{ s}^{-1}$  and found to be spatially coincident with a  $G = 17.1$  stellar *Gaia*-source at a distance of 237 pc. The flux dropped to about  $f_X = 1 \times 10^{-13}\text{ erg cm}^{-2}\text{ s}^{-1}$  during the three following eROSITA all-sky surveys and remained at this lower level during dedicated *XMM-Newton* observations performed in September 2022. With *XMM-Newton*, pulsations with a period of 319 s were found at X-ray and ultraviolet wavelengths occurring simultaneously in time, thus confirming the nature of eRASSU J191213.9–441044 as the second white-dwarf pulsar. The X-ray and UV-pulses correspond to broad optical pulses. Narrow optical pulses that occurred occasionally during simultaneous *XMM-Newton*/ULTRACAM observations have no X-ray counterpart. The orbital variability of the X-ray signal with a roughly sinusoidal shape was observed with a pulsed fraction of  $\sim 28\%$  and maximum emission at orbital phase  $\sim 0.25$ . The ultraviolet light curve peaks at around binary phase 0.45. The X-ray spectrum can be described with the sum of a power law spectrum and a thermal component with a mean X-ray luminosity of  $L_X(0.2\text{--}10\text{ keV}) = 1.4 \times 10^{30}\text{ erg s}^{-1}$ . The spectral and variability properties could indicate some residual accretion, in contrast to the case of the prototypical object AR Sco.

**Key words.** binaries: close – stars: individual: eRASSU J191213.9–441044 – X-rays: binaries

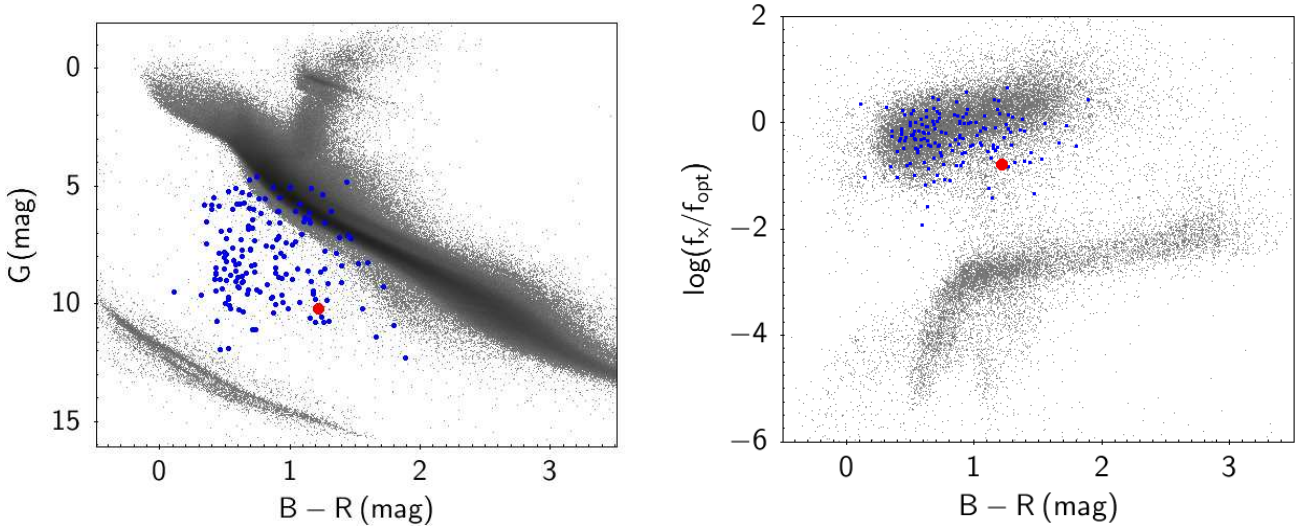
## 1. Introduction

AR Sco, identified in 2016 by Marsh et al. (2016), is a highly peculiar white dwarf main sequence (WDMS) binary, composed of a strongly magnetic white dwarf and an M-dwarf companion orbiting each other with a period of 3.56 h (Buckley et al. 2017). The spin of the WD is only 1.95 min and pulsed emission on the 1.97-min beat period has been observed from radio to X-ray wavelengths. AR Sco is thought to be powered by the spin-down of the white dwarf, and given the observed pulsed and highly polarimetric signal is said to be of synchrotron origin, the object has been deemed as the first white dwarf pulsar. Its X-ray properties were analysed by Takata et al. (2018, 2021), who found a pronounced variability of the X-ray brightness for both the orbital and the beat period between spin and orbital periods. The orbital-phase dependent maximum was observed at superior conjunction of the M star,

an observation used to locate the physical origin of the emission on the irradiated companion star.

AR Sco (and AR Sco-like objects) seem to play an important role in the evolution of close binaries and thereby contribute greatly to our understanding of how strong magnetic fields are generated in white dwarfs. According to Schreiber et al. (2021), the field is generated when the WD crystallizes, transforming a previously semi-detached non-magnetic cataclysmic binary (for a certain time interval) into a detached WDMS binary. With just one object of this kind known to date, it is difficult to find clear observational support for this scenario. The recent discovery of an AR Sco-twin, J191213.72–441045.1 (henceforth, J1912-44) by Pelisoli et al. (2023), with  $P_{\text{orb}} = 4.03\text{ h}$  and  $P_{\text{spin}} = 5.32\text{ min}$  has established AR Sco-like objects as a class, lending support to this evolutionary picture. Pelisoli et al. (2023) found pulsed emission from radio to X-ray wavelengths and pronounced polarimetric variability on the spin period. The spin-down has not yet been measured, hence, there is no constraint on the powering mechanism and the strength of the magnetic field.

\* Based on observations obtained with *XMM-Newton*, an ESA science mission with instruments and contributions directly funded by ESA Member States and NASA.



**Fig. 1.** Color-magnitude diagram of 161 CB candidates drawn from matching eRASS1 – *Gaia* sources (blue symbols) shown on the left. The background is made of  $\sim 230\,000$  unrelated *Gaia* sources with  $\pi/\sigma_\pi > 100$  ( $\pi$ : parallax). X-ray to optical flux ratio ( $\log f_x + \text{phot\_g\_mean\_mag}/2.5 + 4.86$ ) vs. optical colour of eRASS1 sources matching with *Gaia* within a radius of 1 arcsec, shown on the right. The lower branch on the right figure is populated by coronal emitters, the upper branch by AGN and accreting binaries. The position of J1912-44 is marked red in both diagrams and the 161 *Gaia* candidates in blue.

Pelisolli et al. (2023) described three kinds of pulses (flares): a short pulse dominating in radio and optical ( $\text{FWHM} < 10$  s) and a broader pulse seen at optical and X-ray wavelengths, both occurring with the same period, along with a flare with length  $\sim 1/4$  of a spin cycle that occurred only once at an interpulse phase.

Here, we describe the full body of SRG/eROSITA observations of J1912-44 obtained between 2020 and 2021 and we present an analysis of a pointed observation with *XMM-Newton* obtained in September 2022. Fortunately, it was possible to arrange for simultaneous observations with ULTRACAM at the NTT, which revealed some important insight about the timing properties of the source. Throughout this paper, we use the ephemeris  $\text{BMJD(TDB)} = 59\,846.050092 + 0.003696715 \times E$  to align our data in phase. This ephemeris puts the broad X-ray pulse at phase zero at the epoch of our observations. It corresponds to phase 0.715 of the spin ephemeris derived by Pelisolli et al. (2023).

## 2. eROSITA observations

### 2.1. eROSITA discovery and photometric properties

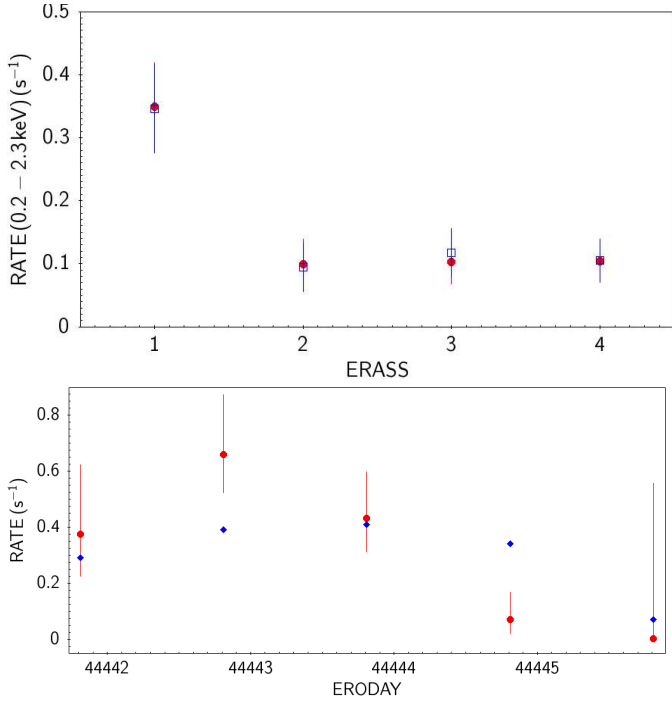
Observations by eROSITA on board SRG (Predehl et al. 2021; Sunyaev et al. 2021) of the field of J1912-44 were carried out four times during the all-sky surveys eRASS1 through eRASS4 in 2020 and 2021. The survey data were reduced with the eROSITA software eSASS version eSASSusers\_211214 (Brunner et al. 2022), using the latest calibration files. J1912-44 was detected as an X-ray source in each of the surveys individually. The data obtained during eRASS1 were used to search for accreting compact white dwarf binaries. A list of 161 high-confidence candidates with matching *Gaia*-sources was compiled using NWAY, a probabilistic cross-matching tool specially trained to identify cataclysmic binaries and related objects (Salvato et al. 2018, 2022). The location of all the candidates in an optical colour–magnitude and an X-ray to optical colour–colour diagram is shown in Fig. 1. Included among them is J1912-44 (*Gaia* designation *Gaia* DR3 6712706405280342784), shown as a red symbol. As many other cataclysmic binaries

(CBs) and candidates, the new WD pulsar is located between the ZAMS and the WD sequence in the CMD. Compared to the other 160 candidates, J1912-44 is relatively red and has a lower X-ray to optical flux ratio, but it is still a factor of 100 above coronal emitters with the same optical colour (we may also make a comparison with Figs. 15 and 16 in Schwope et al. 2022). An identification spectrum showing  $H\alpha$  in emission was obtained with the SAAO 1 m telescope in July 2022. The much larger body of optical data with higher S/N is described in Pelisolli et al. (2023).

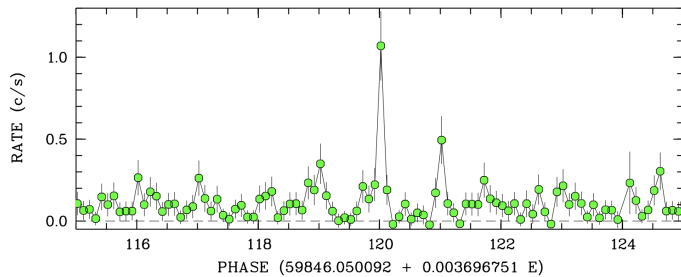
A joint detection on the stacked data from all four eRASS surveys (labeled eRASS:4) revealed the X-ray source at position  $\text{RA}_{2000} = 288.057345^\circ$  and  $\text{DEC}_{2000} = -44.17898^\circ$  with a statistical uncertainty of 1.8 arcsec. The X-ray position coincides within 1.2 arcsec with the position of the optical counterpart. The summed vignetting corrected exposure was 370.4 s (about 80 s, 75 s, 103 s, and 117 s in eRASS1, 2, 3, and 4, respectively) and revealed a total of 57 registered photons. The net exposure before correction was 740 s. The mean eRASS:4 rate was  $0.155 \pm 0.022 \text{ s}^{-1}$  in the energy band 0.2–2.3 keV that was used for the detection of the X-ray sources. The mean rate corresponds to a flux of  $F_X(0.2\text{--}2.3 \text{ keV}) = 1.44 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$ , if the standard ECF (energy conversion factor) is used. The mean rate was higher during eRASS1,  $0.35 \pm 0.07 \text{ s}^{-1}$ , compared to the following three scans where the rate was consistently lower at a rate of  $0.10 \pm 0.03 \text{ s}^{-1}$ . The fluxes in the 0.2–2.3 keV band were 3.2, 0.9, 1.1,  $1.0 \times 10^{-13} \text{ erg cm}^{-2} \text{ s}^{-1}$  for eRASS1, 2, 3, and 4, respectively. eROSITA variability is significant on long and short scales, namely, between different eROSITA surveys separated by half a year and, in eRASS1, between eRODays, as shown in Fig. 2. An eRODay corresponds to one full revolution of the spacecraft (14 400 s), while the length of a scan over a given position on sky is  $< 40$  s. An eRODay is close to but not identical to the orbital period of J1912-44.

### 2.2. eROSITA spectral analysis

During eRASS1 26 source photons were registered, whereas the 2nd, 3rd, and 4th surveys revealed 7, 12, and 12 photons,

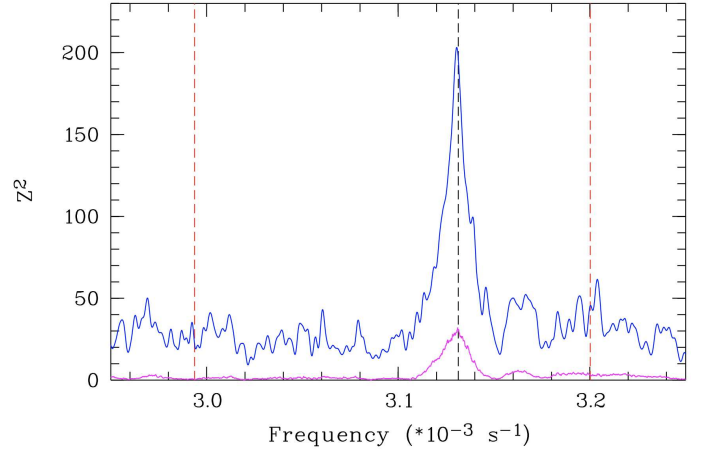


**Fig. 2.** X-ray variability of J1912-44 observed with eROSITA. Upper panel: long-term X-ray light curve displaying the mean rate per eRASS. Shown are mean count rates from stacked source detection followed by forced photometry to derive fluxes per eRASS (red symbols) and from source detection of the individual surveys (blue symbols). Lower panel: mean count rate (red symbols) and exposure time (factor 1/100, blue symbols) per eRODAY in eRASS1.

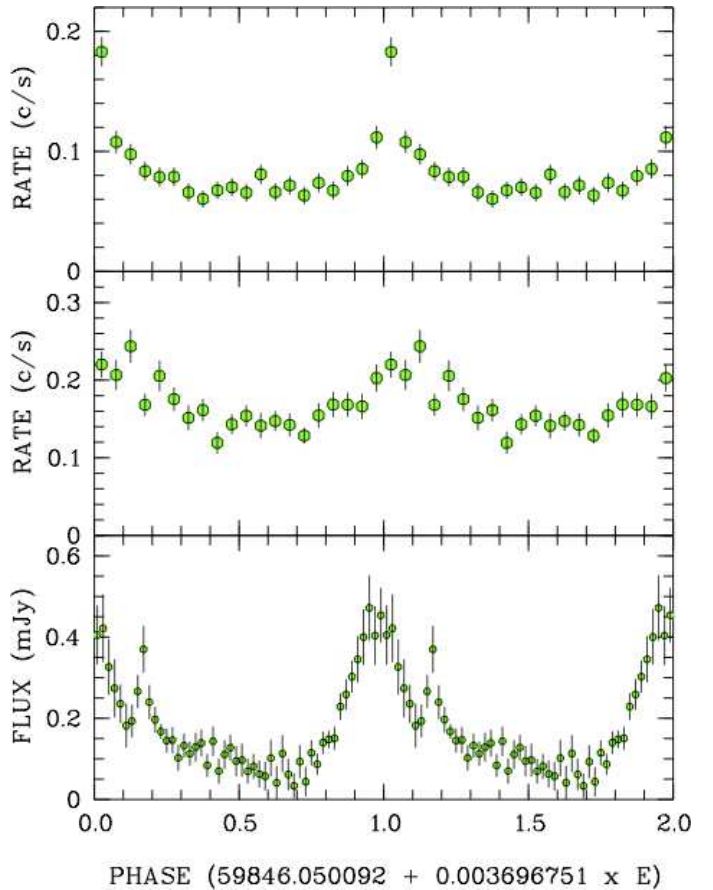


**Fig. 3.** Cutout of the sky-subtracted EPIC-pn light curve (orbital phase 333.226–333.444). Time bins are 32 s.

respectively. We attempted a rough spectral analysis for eRASS1. Source and background spectra together with the auxiliary files, the RMF and the ARF (redistribution matrix and effective area file) were extracted from the raw photon event file using standard eROSITA procedures. Then, XSPEC (version 12.12.0, Arnaud 1996) was used to model the observed spectrum, which was binned with one count per spectral bin. Two spectral models were tested, a thermal plasma emission model (Astrophysical Plasma Emission Code, i.e., the *apec* model in XSPEC terms, Smith et al. 2001) and a power-law model, each modified by some amount of cold interstellar matter (TBABS). Both have the capability to reflect the observed spectrum equally well. Given the low number counts, spectral parameters for either model are not well constrained. For a thermal plasma, the lower limit temperature is 1.7 keV at 90 percent confidence, whereas for a power law, the spectral index lies between 0.6 and 3.0, and the interstellar column density is smaller than  $N_{\text{H}} < 1 \times 10^{21} \text{ cm}^{-2}$ .



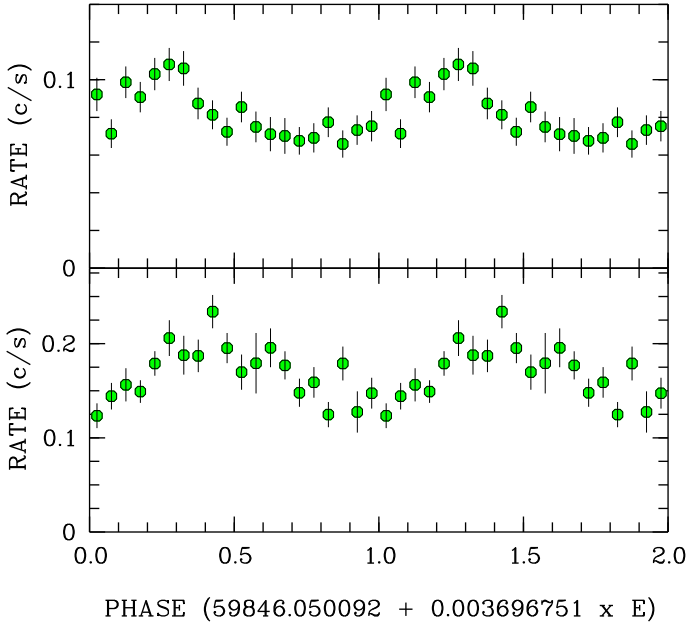
**Fig. 4.** Period search using Buccheri's method (blue curve) and using AOV (magenta). Dashed lines indicate the optically determined spin frequency  $\omega_{\text{spin}}$  (black), the beat frequency between spin and orbit  $\omega + \Omega$  and the sideband frequency  $\omega - 2\Omega$ .



**Fig. 5.** EPIC-pn X-ray, OM/UVM2, and ULTRACAM *g*-band light curves folded over the spin phase.

### 3. XMM-Newton observations

Following our DDT request, *XMM-Newton* observed the field of J1912-44 on Sep 17/18, 2022 for a total of 43 ks. The EPIC X-ray cameras were operated in full-frame mode and the optical monitor was used in fast-imaging mode (time resolution 0.5 s), with the UVM2 filter (effective wavelength 231 nm). The original data were reduced with the latest version of the *XMM-Newton*



**Fig. 6.** EPIC-pn X-ray and OM/UVM2 light curves folded over the orbital binary period.

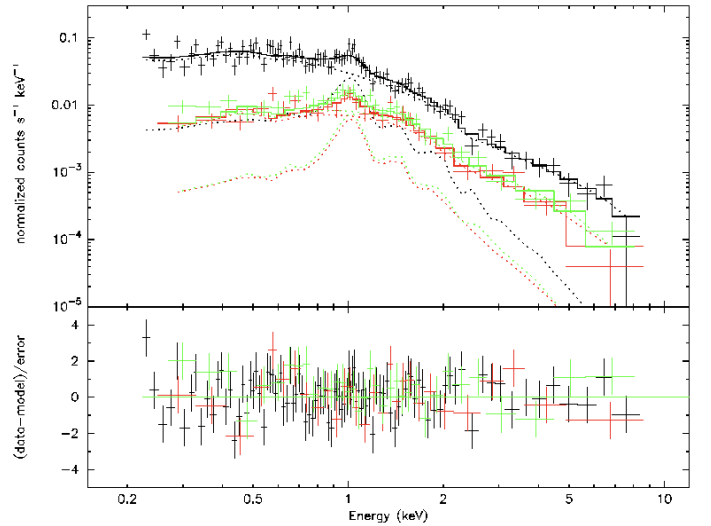
SAS (SAS 20.0) using the most recent calibration files. For the spectral analysis, the data were filtered for intervals of high background rates; the data loss was about 25% for EPIC-pn and almost negligible for MOS1 and MOS2. Times of arrival of individual photons were put on the TT time system and corrected to the solar system barycentre using SAS task `barycen`. Concentric annuli with radii of 30 and 60 arcsec were used to extract source and background photons and to correct the signal in the source region for background contamination. A total of 3430 photons from J1912-44 were collected with EPIC-pn, which gives a mean rate of  $0.0825 \text{ s}^{-1}$  (0.2–10 keV). Fortunately, about 50 min of simultaneous observations *XMM-Newton* and ULTRACAM at the NTT could be ascertained.

The data obtained with the OM were reduced with the SAS task `omchain`, which generates background-subtracted light curves. Also the OM data were corrected to the solar system barycentre using the `barycen` task. Experiments were made using different time bins to create a light curve. Eventually bins of length 3.19390 s (corresponding to 100 phase bins of the measured periodicity) were used to then generate a phase-folded light curve with 20 phase bins (shown in Fig. 5). A fast mode image was created and inspected showing the source at the edge of the fast mode window. This may result in a loss of source photons of up to 50%.

### 3.1. X-ray and UV variability

A cutout of about 50 min of the sky-subtracted, vignetting, and dead time-corrected X-ray light curve is shown in Fig. 3 (here and in the following, we always use all photons in the 0.2–10 keV range). It shows apparent quasi-periodic variability with individual pulses separated by about 5 min. The chosen cutout for the figure contains the most prominent pulse during the whole observation, reaching a peak rate of about  $1.07 \text{ s}^{-1}$  in a 32 s time bin ( $1.69 \text{ s}^{-1}$  in a corresponding bin of 16 s in length).

We searched for periodic variations of the X-ray signal in EPIC-pn using two independent methods. Photons selected in a 30 arcsec region around the source were analysed using



**Fig. 7.** *XMM-Newton* spectra obtained with EPIC-pn, -MOS1, and -MOS2 with a power-law plus thermal model fit. Lower panel shows the error-normalized residuals (black: pn, red: MOS1, green: MOS2).

Buccheri’s  $Z_n^2$  statistic (Buccheri et al. 1983). The period range investigated was between 100 s and 450 s and the resulting periodogram showed one prominent peak at  $0.0031305 \text{ d}^{-1}$  corresponding to 319.43 s (Fig. 4). We also used the analysis of variance (AOV) implemented in MIDAS to search for periodic variability. The input used was the sky-subtracted, vignetting and dead time-corrected light curve with 1 s time bins. The AOV-periodogram, shown also in Fig. 4 (magenta line) has its maximum at a period of 319.35 s, the center of the distribution at 319.52 s.

Spin phase-folded X-ray, ultraviolet, and optical light curves (ULTRACAM *g*-band shown, the light curves of the other bands, *r* and *u*, are almost carbon copies of the *g*-band data) with 20 bins per cycle are shown in Fig. 5. X-ray and OM data show the same pattern, the broad pulse that appears superposed on a basal flux at about half the spectral flux. The ULTRACAM data additionally show the narrow pulse occurring a little less than 0.2 phase units after the main pulse. It is not detected at X-ray wavelengths. Whether the OM data show this feature is difficult to judge given the low count rate. The ephemeris used for folding,  $\text{MBJD(TDB)} = 59\,846.050092 + 0.003696751 \times E$ , locates the broad pulse in all wavebands at around  $\phi_{\text{spin}} = 0$  and, thus, suggest a common physical origin. A smooth interpolated continuum connecting the interpulse minima was subtracted from the ULTRACAM data before phase folding. The mean (basal) rate of the OM in the phase interval 0.3–0.7 (spin phase) was  $0.152 \pm 0.007 \text{ s}^{-1}$ , which corresponds<sup>1</sup> to  $3.34 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ \AA}^{-1}$ .

We searched for orbital phase-dependent variability in the X-ray and UV data from *XMM-Newton* by inspecting phase-folded light curves shown in Fig. 6 that were generated using the orbital ephemeris of Pelisoli et al. (2023). The phase-averaged X-ray light curve varies roughly sinusoidally with a maximum at around  $\phi_{\text{orb}} \sim 0.25$  and has a pulsed fraction  $p_f = (f_{\text{max}} - f_{\text{min}})/(f_{\text{max}} + f_{\text{min}})$  of 28%. The UV light curve is similar in shape, but the bright phase hump a little more extended (~60% of the cycle) and centered at a later phase, ~0.45. The UV light curve

<sup>1</sup> The conversion factor between counts and measured flux in the UVM2-band is  $2.2 \times 10^{-15} \text{ erg cm}^{-2} \text{ \AA}^{-1} \text{ ct}^{-1}$ , source <https://www.cosmos.esa.int/web/xmm-newton/sas-watchout-uvflux>

has the same shape and phasing as the phase-folded optical light curve from TESS (Pelisoli et al. 2023, their Fig. 1).

We examined whether the X-ray pulses show a dependency on the orbital phase. We divided the data set into five orbital phase intervals of same length and searched for the spin pulses in such reduced data sets individually. The pulsations are found at all phases with a tendency to show a larger pulsed fraction between orbital phases 0.0 and 0.4.

### 3.2. Spectral analysis

Mean X-ray spectra for the whole *XMM-Newton* observation of J1912-44 were extracted for the three EPIC cameras individually and then fitted jointly. The spectra were binned with a minimum of 25 counts per bin, so that a  $\chi^2$  minimization technique could be used to find the best model and to derive the parameter values. Initially a one-component model was tested, but neither a pure thermal emission (APEC) nor a single power law model yielded a satisfactory fit. The fit that reflects the data satisfactorily is a combination of the two emission models modified by a column of cold interstellar matter TBABS \* (APEC + POWL). The power law carries most of the flux (almost 90%), while the thermal model yields the lines around 1 keV to account for spectral residuals left by the power law component alone. The accepted best-fit parameters are listed Table 1, while the data, the model fit (with individual components shown), and the residuals are displayed in Fig. 7. The model chosen is the most simple, namely, the one with least amount of parameters to be determined. The prototype AR Sco was fitted with the superposition of several thermal spectra (Takata et al. 2018, 2021). This would be possible for J1912-44 as well. A fit with just two thermal components is not sufficient ( $\chi^2_{\nu} = 1.55$ ) but the sum of three thermal components with temperatures according to 0.2 keV, 1.1 keV, and 4 keV gives a satisfactory fit ( $\chi^2_{\nu} = 1.07$  for 154 d.o.f.).

Pelisoli et al. (2023, see their Fig. 3) compiled the spectral energy distribution (SED) from the radio to the X-ray regime using the X-ray spectral model derived here. Integrating this and using the *Gaia* distance, they find the bolometric luminosity to be  $\sim 10^{33}$  erg s $^{-1}$ . The X-ray luminosity of J1912-44, namely,  $L_X \sim 1.4 \times 10^{30}$  erg s $^{-1}$ , is a factor of 3 lower than that of AR Sco, while the fraction  $L_X/L_{\text{bol}}$  is only  $\sim 0.14\%$  for J1912-44. However, this fraction seems to be slightly higher (on the order of 1–2%) for AR Sco (Takata et al. 2021).

## 4. Results and discussion

We analysed the currently available X-ray data of J1912-44 from eROSITA and *XMM-Newton*. First, it is worth noting that J1912-44 was not detected by ROSAT, *Swift*, or by the *XMM-Newton* slew survey. J1912-44 shows X-ray variability on various time scales, the longest investigated being over half a year, namely, the time interval between individual eRASS scans. The mean brightness during eRASS1 with a flux of  $3.3 \times 10^{-13}$  erg cm $^{-2}$  s $^{-1}$  dropped to  $\sim 1 \times 10^{-13}$  erg cm $^{-2}$  s $^{-1}$  at all subsequent occasions. During eRASS1, J1912-44 was also variable on the eRODay scale (which roughly correspond to the orbital period). Pointed *XMM-Newton* observations revealed periodic variability of the X-ray and the UV signals on the spin period of 319.4 s. Phase-folded light curves on the spin period show a basal flux and the broad pulses seen also in the optical. Their maximum flux is a factor  $\sim 2$  higher than the basal flux. Phase-folded light curves on the orbital period have a pulsed fraction of 28% and a bright phase lasting about 0.5 phase units, and that is roughly centered on phase 0.25. The narrow pulse, which is prominent at radio and

**Table 1.** Parameters of best-fit models to the mean EPIC spectra.

Model: TBABS*(APEC + POWL)	
Parameters	
$N_{\text{H}}$ ( $10^{20}$ cm $^2$ )	$5 \pm 2$
$kT_{\text{APEC}}$ (keV)	$1.24^{+0.11}_{-0.10}$
Norm <sub>APEC</sub>	$(1.5 \pm 0.5) \times 10^{-5}$
$\kappa_{\text{PL}}$	$2.14 \pm 0.11$
Norm <sub>PL</sub>	$(3.11 \pm 0.38) \times 10^{-5}$
$\chi^2_{\nu}$ ( $\chi^2/\text{d.o.f.}$ )	1.12(173/154)
Observed fluxes ( $10^{-13}$ erg cm $^{-2}$ s $^{-1}$ )	
$F_{0.2-2.3}$	$1.05 \pm 0.04$
$F_{0.2-10}$	$1.66 \pm 0.09$
Unabsorbed fluxes, best-fit parameters ( $10^{-13}$ erg cm $^{-2}$ s $^{-1}$ )	
$F_{0.2-2.3}$	1.5
$F_{0.2-10}$	2.1
$F_{\text{bol}}$	12.3
$L_X$ (0.2–10 keV, $(D/237 \text{ pc})^2 10^{30}$ erg s $^{-1}$ )	1.4

**Notes.** The  $\kappa$  is the power law index.

optical wavelengths was not discovered in the 130 spin cycles covered by *XMM-Newton*. A broad interpulse flare, which could be an accretion event, seen in the optical and X-ray, would perhaps remained unnoticed without its prominent appearance in the X-ray data.

Pelisoli et al. (2023) explained the period of 319 s as the spin period of the white dwarf. The predicted beat period is 312.5 s, while (almost) no signal is seen in the periodograms at this period, neither in this work nor in Pelisoli et al. (2023). This is unexpected, because the emission lines, which are undoubtedly originating from the irradiated side of the donor star, show clear flux modulations that hint to a pronounced irradiation effect of the secondary star. These are expected to occur on the beat period. An observational confirmation is outstanding because the time coverage of the spectroscopic data was insufficient to distinguish between the beat and spin periods. The X-ray data show clear evidence for variability on the suspected spin period whereas there is at most marginal evidence for the beat period in the X-ray periodogram (Fig. 4).

A couple of questions remain. Narrow optical pulses during the joint ULTRACAM/*XMM-Newton* campaign are thought to be of synchrotron origin from the WD, but they have no X-ray counterpart. The narrow pulses are the only signal seen at radio wavelengths. The optical pulses were observed in cycle 331 at orbital phases 0.0526, 0.746, 0.967, and 0.1186 (with increasing flux) but were not observed at later phases, although the observations were continued through further six spin cycles up to orbital phase 0.26. Some parameters of the narrow flares are puzzling, in particular with regard to what is causing their phase dependence and what their genuine energy distribution is.

If the pulsed X-ray emission is due to synchrotron beaming, we must consider what the source of the basal flux is, for instance, whether it may be due to weak accretion. The measured X-ray luminosity is low, much lower than what a Roche lobe-filling cataclysmic variable at this orbital period would typically exhibit. An active secondary star might be the source of a wind that is captured by the strong magnetic field of the WD, as in the LARPs/PREPs (Schwope et al. 2002, 2009). This may lead to accretion-induced X-ray emission as in other LARPs (e.g., Vogel et al. 2007), although the accreted matter is guided to the footpoints of the magnetic fields and such an accretion geometry

typically leads to pulsed X-ray emission. An argument in favour of a scenario that involves accretion may be derived from the observed optical/X-ray flare. The basal flux might be composed by a number of unresolved miniflares. Flares or flare-like accretion might also be the cause of the observed variability between eRODays during eRASS1.

Orbital phase-dependent variability is observed at X-ray and UV wavelengths. It is important to take care to avoid overinterpreting the data, as they cover just three orbital cycles and phase-dependent behavior might be observed where it is not actually present. However, it appears as if the UV would be bright around phase 0.5. We consider whether this be the result of the irradiation of the donor star, which would then radiate at a level that corresponds to the flux difference between the mean flux at phases 0.0 and 0.5 (i.e.,  $\sim 1.7 \times 10^{-16} \text{ erg cm}^{-2} \text{ s}^{-1} \text{ \AA}^{-1}$ ). The X-ray light curve peaks earlier in orbital phase at around phase 0.25. It seems to display a slow rise and a faster decline from the maximum. The light curve asymmetry is reversed from left to right in comparison to AR Sco.

There are other dissimilarities that ought to be considered when making comparisons with AR Sco. In particular, for AR Sco, the X-ray and UV light curves appear similar, especially with regard to the phase of the minimum flux, which is the same. The curves in J1912-44 appear dissimilar. The spin phase-folded light curves in AR Sco are double-peaked, while those in J1912-44 are single-peaked. AR Sco seems to have a strong thermal component in its X-ray spectrum, whereas that of J1912-44 could be of non-thermal origin. The question of whether those differences can be explained simply by different geometries or whether they require different physical arguments can be addressed once data with higher signal-to-noise ratios are available.

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