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Wang, K, Ma, Y, Howlett, PR et al. (2 more authors) (2022) New Unsaturated Dynamic Porosity Hydromechanical Coupled Model and Experimental Validation. International Journal of Geomechanics, 22 (10). 04022171. ISSN 1532-3641

https://doi.org/10.1061/(asce)gm.1943-5622.0002545

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1	New unsaturated dynamic porosity hydro-mechanical coupled model and
2	experimental validation
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14	Abstract
15	Constitutive coupled modeling has developed rapidly in recent decades, with numerous new models
16	published. However, few models consider dynamic porosity, and experimental validation of such a
17	model remains a challenge due to multiple variables. In this study, a new constitutive model for
18	unsaturated soil with dynamic porosity was developed and validated using test data from two
19	experimental studies that yielded good results (relative average error AVRE = $0.8631$ - $1.3046$ , R <sup>2</sup> =
20	0.9028–0.9981). The sensitivity of the model to the four primary parameters was analyzed to investigate
21	the influence of model properties on the hydraulic and mechanical behavior. Results show that the
22	calculation of volumetric strain is most sensitive to Young's modulus (E), while the calculation of 1

specific water volume is most sensitive to permeability (k). Also, the sensitivity of the parameters changes with their value. Modeled results show that the porosity change significantly affects both hydraulic and mechanical behavior, even when soil undergoes relatively low deformation. Relative calculation error decreases notably after porosity change is considered (44.9% and 35.2% improvement in two different calculations). This study also finds that dynamic porosity affects the deformation energy of solids.

Keywords: Biot's theory; Mixture coupling theory; Unsaturated soil; Dynamic porosity; Model
 validation; Parameter sensitivity

31

# Introduction

32 Since Terzaghi developed the one-dimensional consolidation theory of saturated soil (Terzaghi 1943), 33 research into coupled hydro-mechanical models has been an important topic in many engineering 34 applications. Safety analysis for nuclear waste disposal must evaluate the long-term deformation of 35 barrier rock and potential environmental risk to the host rock and local groundwater (Félix et al. 1996, 36 Charlier et al. 2013, Schwartz 2018). The prediction of fluid inflow into tunnels as well as deformation 37 at dams requires an understanding of the coupled relationships between rock displacement, pore moisture 38 and water flow (Chen et al. 2016, Zhang et al. 2020). Hydro-mechanical coupling is required for studies 39 of rock slope deformation induced by precipitation infiltration (Masoudian et al. 2019) and research into 40 groundwater flow in shale rock and hydraulic fracturing during shale gas exploitation (Lisjak et al. 2017). 41 Many cases of coupled hydro-mechanical behavior occur in an unsaturated zone, which highlights the 42 importance of studies of coupled hydro-mechanical models in unsaturated media.

43 Two approaches can be used to derive most of the coupling theory of hydro-mechanical behavior: the

44	mechanics approach and the mixture approach. The mechanics approach is based primarily on the work
45	of Terzaghi (Terzaghi 1943) and Biot's theory (Biot 1941, Biot 1962). Many models have been
46	developed using this approach, including unsaturated hydro-mechanical models (Sanavia et al. 2002, Li
47	and Yang 2018), large-strain hydro-mechanical models (Meroi et al. 1995), multiphase flow models
48	(Edip et al. 2018), and coupled thermo-hydro-mechanical-chemical models (Li et al. 2006, Seetharam et
49	al. 2007). This approach has found to be useful when developing new models for specific applications
50	(e.g., introducing formulations of other fields to the constitutive model) and has supported many practical
51	areas of geomechanics and engineering. However, this approach lacks a systemic, self-developing
52	linkage between different fields, making it difficult to use to build a unified coupling theory for porous
53	media (Chen et al. 2016).
54	A different approach, mixture theory, originated from biological tissue engineering and was developed
55	by Truesdell (Truesdell 1962). Mixture theory is based on non-equilibrium thermodynamics but has yet
56	to describe the coupled interactions of different phases.
57	Recently, mixture coupling theory, which was previously called modified mixture theory, was developed
58	by Heidug and Wong (Heidug and Wong 1996), has been used to build more advanced coupled models
59	in geomechanics and geophysics (Chen and Hicks 2013a, Chen et al. 2013b, Chen et al. 2016, Ma et al.
60	2021). Instead of discriminating between solid and fluid phases, as the previous mixture theory does,
61	mixture coupling theory treats the fluid-infiltrated system as a single continuum material. Mixture
62	coupling theory builds a unified theory of multifield behaviors in soil and rocks using entropy production
63	as a single driving force.

64 Mixture coupling theory has demonstrated marked advantages in developing systemic multifield

84	Basic unsaturated hydro-mechanical coupled governing equations
83	Theory
82	nonnegligible porosity changes.
81	proposed model exhibits marked potential for characterizing hydro-mechanical coupling behavior under
80	and provides a benchmark for model calibration and evaluation, particularly for sensitivity analysis. The
79	coupling theory in the field of geotechnics, develops a new coupled model considering porosity change,
78	This study provides additional support for the development of the theoretical framework of mixture
77	in more detail.
76	porosity change. The deformation energy was studied to describe hydro-mechanical coupling behavior
75	different parameters is evaluated. The new model is compared with a previous model that ignores
74	and validated using experimental data, and the sensitivity of hydraulic and mechanical behaviors to
73	In this study, a new constitutive model for unsaturated soil considering dynamic porosity is developed
72	different parameters.
71	difficult to determine, and more effort is required to understand the sensitivity of the model results to
70	parameterization of models with material properties is currently inadequate. Some parameters are
69	geotechnology. Despite understanding the parameters of unsaturated hydro-mechanical coupling, the
68	experimental validation of mixture coupling theory, primarily due to its recent application in
67	hydro-mechanical-thermal model (Chen and Hicks 2013a, Chen et al. 2018). To date, there has been little
66	osmosis, and thermal osmotic processes were added into the hydro-mechanical-chemical model and the
65	coupling models in both unsaturated and saturated media. Chemical osmosis, including dual-chemical

85 This paper is based on the coupled unsaturated hydro-mechanical (UHM) model for soil, which is based

86 on mixture coupling theory (Chen and Hicks 2011), and provides a critical step forward by considering

87 dynamic porosity and analyzing deformation energy. To completely introduce the theory upon which this

study is based, the basic equation for the UHM coupled model is introduced as follows below.

#### 89 1. Mechanical equation

90 The final governing field equation for the solid phase can be written as (Chen and Hicks 2011):

91 
$$\left(K - \frac{2G}{3}\right)\frac{\partial^2 d_k}{\partial x_k \partial x_i} + G\left(\frac{\partial^2 d_i}{\partial x_j \partial x_j} + \frac{\partial^2 d_j}{\partial x_i \partial x_j}\right) - \zeta \nabla \left[\left(S_w + \frac{C_s}{v}p^w\right)p^w\right] = 0 \quad (1)$$

where G is the shear modulus; K is the bulk modulus; d is the displacement; ζ is Biot's consolidation
coefficient, which is defined as ζ = 1-(K/K<sub>s</sub>), where K<sub>s</sub> is the bulk modulus of the solid matrix); v
is the porosity; S<sub>w</sub> is the saturation of pore water; p<sup>w</sup> is the porewater pressure; and C<sub>s</sub> is the specific
moisture content, which is defined as C<sub>s</sub> = v ∂S<sub>w</sub>/∂p<sub>w</sub>) (Lewis et al. 1987).
Normally, C<sub>s</sub> is small and is assumed to be zero for the convenience of calculation; thus, Equation

97 (1) becomes:

98 
$$\left(K - \frac{2G}{3}\right)\frac{\partial^2 d_k}{\partial x_k \partial x_i} + G\left(\frac{\partial^2 d_i}{\partial x_j \partial x_j} + \frac{\partial^2 d_j}{\partial x_i \partial x_j}\right) - \zeta S_w \nabla p^w = 0$$
(2)

# 99 2. <u>Hydraulic equation</u>

100 The final governing field equation for the porewater can be written as (Chen and Hicks 2011):

101 
$$S_{w}\zeta \dot{d} - \frac{k_{rw}k}{v} (\nabla^{2} p^{w}) + (C_{s} + \upsilon \frac{S_{w}}{K_{w}}) \frac{\partial p^{w}}{\partial t} + S_{w}Q(S_{w} + \frac{C_{s}}{\upsilon} p^{w}) \frac{\partial p^{w}}{\partial t} = 0$$
(3)

102 where k is the permeability;  $k_{rw}$  is the relative permeability; v is the dynamic viscosity;  $K_w$  is the bulk

- 103 modulus of water; and Q is the void compressibility, which is defined as  $Q = (1/K_s)(\xi \upsilon)$ .
- 104 Normally,  $K_s$  is large compared to  $K_w$  (Alonso and Alcoverro 1999b), Q can be assumed to be zero,
- 105 and thus, Equation (3) becomes:

106 
$$S_{w}\zeta \nabla \cdot d - \frac{k_{rw}k}{v} \left(\nabla^{2} p^{w}\right) + S_{w} \upsilon q \frac{\partial p^{w}}{\partial t} = 0$$
(4)

107 where  $q = \frac{1}{K_w}$  is the poroelastic storage.

109 observations (Stormont and Daemen 1992) and is thus the primary focus of this study.

#### 110 Unsaturated hydro-mechanical coupled governing equations with dynamic porosity

- 111 In the absence of chemical influence, a porosity function can be obtained based on mixture coupling
- 112 theory (Chen and Hicks 2013a). This function will be incorporated in the coupled model:

113 
$$\dot{v} = \zeta \, \varepsilon_{ii} + Q \, p_{pore} \tag{5}$$

114 where  $\varepsilon_{ii}$  is the volumetric strain tensor; and  $p_{pore}$  is the pore pressure, which can be replaced by the

115 average pore pressure  $\overline{p}$ . If the porosity change induced by water pressure must be studied, Q should

116 be considered in both Equations (5) and (3).

117 According to previous research on the average pore pressure (Lewis and Schrefler 1982), because the air

118 pressure is assumed to be zero in this paper,  $\overline{p}$  can be defined as:

 $\overline{p} = S_w p^w \tag{6}$ 

120 According to Equations (5) and (6), with the assumption of  $\zeta = 1$ , the dynamic porosity can be expressed

121 as:

122 
$$\upsilon = \frac{\upsilon_0 + \varepsilon_{ii} + \frac{1}{K_s} (S_w p^w - S_{w0} p_0)}{1 + \frac{1}{K_s} (S_w p^w - S_{w0} p_0)}$$
(7)

123 where  $v_0$  is the initial porosity,  $p_0$  is the initial water pressure, and  $S_{w0}$  is the initial water saturation.

124 Equation (7) with Equations (4) and (2) build the general mathematical model for hydro-mechanical

125 coupling in unsaturated soil with dynamic porosity (UHM-P). This model considers porosity change due

126 to deformation (mechanical behavior) and pore pressure (hydraulic behavior). The UHM-P model

127 degrades to the UHM model if the porosity change is ignored.

#### 128 Deformation energy equation

The deformation energy W (i.e., the dual potential) describes the potential that controls the deformation (Chen and Hicks 2011). Its time derivative  $\dot{W}$  in soil with consideration of dynamic porosity change vis defined as:

132 
$$W = \left(\Psi - \upsilon \psi_{pore}\right) - \overline{p}\upsilon \tag{8}$$

133 
$$\dot{W} = tr(\mathbf{T}\dot{\mathbf{E}}) - \frac{1}{p}\upsilon \tag{9}$$

134 where  $\Psi = J\psi$  is the Helmholtz free energy density in the reference configuration of the total soil, J is

135 the Jacobian deformation gradient **F**,  $\psi_{pore}$  is the Helmholtz free energy density of the pore water, **T** is

136 the second Piola-Kirchhoff stress, and E is the Green strain.

137 Critically, the dynamic evolution of porosity influences deformation energy, which has remained unclear

138 in previous research. Substituting Equation (6) into (8) and (9) leads to:

139 
$$W = \left(\Psi - \upsilon \psi_{pore}\right) - S_w p \upsilon \tag{10}$$

140 
$$\dot{W} = tr(\mathbf{T}\dot{\mathbf{E}}) - (S_w p) \cdot \upsilon \tag{11}$$

The deformation energy is the energy obtained by deducting the energy of pore water pressure from the energy of the soil matrix, and its time derivative is a function of  $\mathbf{E}$ , p and v. With the small strain assumption, the Green Strain tensor  $E_{ij}$  and Piola-Kirchhoff stress  $T_{ij}$  can be replaced by the strain tensor  $\varepsilon_{ij}$  and Cauchy stress  $\sigma_{ij}$ . Therefore, equation (11) becomes:

145 
$$\dot{W} = tr\left(\begin{bmatrix}\varepsilon_{11} & \varepsilon_{12} & \varepsilon_{13}\\ \varepsilon_{21} & \varepsilon_{22} & \varepsilon_{23}\\ \varepsilon_{31} & \varepsilon_{32} & \varepsilon_{33}\end{bmatrix}\begin{bmatrix}\dot{\cdot} & \dot{\cdot} & \dot{\cdot}\\ \sigma_{11} & \sigma_{12} & \sigma_{13}\\ \dot{\cdot} & \dot{\cdot} & \dot{\cdot}\\ \sigma_{21} & \sigma_{22} & \sigma_{23}\\ \dot{\cdot} & \dot{\cdot} & \dot{\cdot}\\ \sigma_{31} & \sigma_{32} & \sigma_{33}\end{bmatrix}\right) - (S_w p) \dot{\cdot} \upsilon$$
(12)

#### Methodology

#### 147 Numerical model

146

148 The UHM model and UHM-P model are validated and analyzed in this study by considering the simple 149 but classic coupled hydraulic-mechanical problem of unsaturated "triaxial test" experiments. The 150 experimental data come from two published articles (Vassallo et al. 2007), and (Biglari et al. 2012). 151 In these experiments, soil samples were kept unsaturated throughout the entire period. After obtaining 152 some basic soil parameters, such as the plastic limit and initial water content, a sample of height H and 153 width W was mounted in the test apparatus (Fig. 1). The controlling water pressure on the lower boundary 154  $(P_w)$  and the air pressure on the upper boundary  $(P_a)$  are maintained to ensure that the suction  $S = P_a - P_w$ 155 remains constant. Concurrently, the mean net stress  $|\sigma| = P - P_a$  on the upper and lateral boundaries is also 156 a constant. This procedure is called the equalization stage, when the sample swells or compresses 157 according to its initial water content. The equalization stage ends when the pore water content and total 158 soil volume stabilize. Then, isotropic compression is performed, and this stage gradually increases the 159 mean net stress to a final value. A loading rate of 4 kPa/h is used by Vassallo to ensure fully drained and 160 constant suction conditions. The pore water content and volume change of the sample were measured 161 during these experiments. The 2D numerical model is shown in Fig. 1. 162 To solve the model equations, COMSOL Multiphysics software was used to build a numerical model. 163 Darcy's law was used for the hydraulic field, and the linear elastic materials model was used for the

164 mechanics field. The initial displacement of the sample (d<sub>0</sub>) was zero. The initial water pressure (P<sub>0</sub>) was

165	calculated from the initial water content using the Van Genuchten equation (Genuchten and Th. 1980).
166	The upper and right pressure boundaries for the mechanics field were pressure boundaries, and the stress
167	$\sigma$ was P-P <sub>a</sub> . The lower boundary was a 'roller' boundary, indicating no vertical deformation. The upper
168	and right boundaries for the hydraulic field were no flow boundaries; the lower boundary was a fixed
169	pressure boundary with a pressure value equal to -S; and the left boundary was a symmetric boundary.
170	All simulation settings and sample parameters used for model validation are shown in Table 1. The
171	Young's modulus, Poisson's ratio and poroelastic storage are reported in Alonso's studies (Alonso and
172	Alcoverro 1999a, Alonso and Alcoverro 1999b). The Van Genuchten equation parameters were obtained
173	by fitting the experimental water content and water pressure. The porosity used in the UHM model is a
174	constant value, which is the initial porosity value in the UHM-P model, as the porosity changes during
175	evolution.
176	Model validation method
177	The modeled change of specific water volume in the equalization stage by Vassallo and Biglari was

178 compared with the experimental change. The specific water volume  $v_w$  was calculated using the model

result: 179

$$v_w = G_s w = G_s S_w \nu \tag{13}$$

181 where  $G_s$  is the specific gravity of the soil solids and w is the water content.

182 Note that the sample was divided into many elements in the model; thus, the surface average saturation 183 was used to calculate the  $v_w$  of the total sample.

184 For validation using results of both Vassallo and Biglari, the experimental data from one sample were

185 used for the model's calibration; however, experimental data from all samples were used for validation. 186 The error between the model results and experimental data was quantified using the following equations

187 during calibration and validation:

188 
$$AVRE = \sum_{i=1}^{n} |M_i/E_i|/n$$
 (14)

189
$$R = \frac{\sum_{i=1}^{n} (M_{i} - \overline{M}) (E_{i} - \overline{E})}{\sqrt{\sum_{i=1}^{n} (M_{i} - \overline{M})^{2} \cdot \sum_{i=1}^{n} (E_{i} - \overline{E})^{2}}}$$
(15)

190 where AVRE is the relative average error and R is the correlation coefficient. In Equations (14) and (15),  $M_i$  is the model value,  $E_i$  is the experimental value, *n* is the total number of data points,  $\overline{M}$  is the 191 average of the model data, and  $\overline{E}$  is the average of the experimental values. 192

#### 193 Sensitivity analysis method

194 To describe the effect of parameter variation on the change in model calculation more accurately, the 195 ratio of parameter variation (ROV) (Cheviron and Coquet 2009) was used. This variation relates the 196 relative variation in the results to analyze the parameter sensitivity. Four primary parameters were chosen: 197 permeability (k), poroelastic storage (q), Biot's coefficient ( $\zeta$ ) and Young's modulus (E). The average 198 specific water content, average volumetric strain for the hydraulic results and mechanical results were 199 chosen separately. The ROV was calculated using:

200 
$$ROV(t) = \frac{\left[c(t) - c_{ref}(t)\right] / \left|c_{ref}(t)\right|}{\left(p - p_{ref}\right) / p_{ref}} = \frac{out \operatorname{var}}{in \operatorname{var}}$$
(16)

**¬** //

The reference model is the model of sample 'E' that is validated by the experimental data of Biglari's  
results. 
$$c_{ref}(t)$$
 is the calculated value of the reference model at a specified time,  $p_{ref}$  is the parameter  
value of the reference model,  $c(t)$  is the calculated value of the variation model with the parameter  
changed and  $p$  is the parameter value of the variation model.

205 During each calculation of the sensitivity analysis, one parameter was changed, and the other parameters were held constant. Then, the sensitivity to that parameter was analyzed using *ROV*. The calculated value changes with time; thus, a maximum, average and minimum value are defined for each parameter as:

209  

$$ROV_{max} = \max |ROV(t)|$$

$$ROV_{min} = \min |ROV(t)|$$

$$ROV_{ave} = \sum_{t=0}^{t=T} ROV(t) / N$$
(17)

where  $ROV_{max}$  is the maximum effect of parameter change to the calculation value in the total period and is the maximum absolute value of ROV of one parameter during the calculation period;  $ROV_{min}$  is the minimum effect of parameter change and the ROV of minimum absolute value;  $ROV_{ave}$  is the average sensitivity of parameter and is the average value of ROV in the calculation period; N is the number of ROV; and T is the end time.

215 All of the input variations for parameters in the sensitivity analysis are shown in Table 2. These 216 parameters varied from -60% to 60% compared to their reference value. The variation ranges of  $\zeta$  and

217 *E* are based on previous studies by (Alonso and Alcoverro 1999b) and (Selvadurai et al. 2019).

218 Model validation results

219 In this section, the results of the UHM model are compared with the experimental data, and the result

220 error of each sample is analyzed. Comparisons of the experimental and modeled changes in specific

221 water volume  $\Delta v_w$  of four samples in researches of Vassallo and Biglari are shown in Fig. 2 and Fig. 3,

222 respectively.

223 The experimental data in Fig. 3 are smoother than those in Fig. 2 due to the finer measurement techniques 224 used by Biglari. Different responses of  $\Delta v_w$  can be observed in the eight samples. The water content rises

225 in all samples except for S100P400TG, samples F and G, which implies that the initial water pressures

226 of S100P400TG, samples F and G, were higher than the water pressure on the lower boundary (i.e., these 227 samples drain during the equalization period). For the other samples, the water pressure on the lower 228 boundary is higher than the initial water pressure; thus, the water flows to the upper part of the sample 229 from the lower boundary (i.e., there is a rising saturation in these wetting samples, and  $\Delta v_{u}$  remains 230 positive). 231 The rise of  $\Delta v_w$  slows and reaches zero after approximately 12,000 minutes for the samples in Vassallo's 232 study. The initial  $\Delta v_w$  error of S100P100TG is higher than that of the other three samples, which may be 233 due to inexact measurement of the initial water pressure in that sample. 234 The calculated error for each sample of Vassallo's study is shown in Table 3. The AVRE of all samples are low except for S100P100TG. R is larger than 0.9 in all samples, showing a good correspondence 235 236 between the experimental data and model results. 237 The rise in  $\Delta v_w$  also slows for samples in Biglari's study. Equalization is reached in approximately 40 h. 238 The permeability of samples in Biglari's study is larger than that in Vassallo's, and the water pressure 239 changes in Biglari's are also larger. These two factors promote faster pore water flow through samples in 240 Biglari's study, thus shortening equalization time. The model results are less accurate at D, showing 241 marked positive errors at both early and late stages, compared with the other samples. The calculated 242 error for each sample is shown in Table 4. The AVRE and R of all samples are generally better than 243 those in Vassallo's study and acceptable in numerical calculation. 244 Nearly all sample errors are positive throughout the calculation period. However, the errors between the experimental data and model results differ between samples that drain and wet. Most errors in the 245 246 modeling results of draining samples are higher in the early stage and lower in the late stage, while

12

#### 247 wetting samples exhibit an opposite trend.

Sensitivity analysis of model parameters 248 249 The sensitivity of permeability k, poroelastic storage q, Biot's coefficient  $\zeta$  and Young's modulus E 250 of the UHM model were assessed using the ratio of parameter variation (ROV). The parameter value of 251 the reference model comes from a previous study of clay and silt soil. Therefore, we investigated a 252 relatively tight range of -60% to 60%, where  $\zeta$  is -60% to -30%, to make specific statements about the 253 sensitivity of parameters of soft soil instead of covering general variation of parameters for all soils. The 254 variation range for all parameters is the same; thus, we can compare the sensitivity of each parameter 255 without the influence of parameter value because all the sensitivity of parameters is parameter value-256 dependent.

#### 257 Parameter sensitivity of mechanical behavior

258 The average volumetric strain of the sample using different input parameter values is shown in Fig. 4. In 259 all cases, the volumetric strain increases rapidly in the early stage, reaches its peak, and then decreases 260 slowly back to its initial value. The opposite effect of permeability and poroelastic storage on the 261 volumetric strain is clearly shown. The peak volumetric strain is delayed with decreasing k, and the 262 curves also become smoother. However, the decrease in q hastens the appearance of the peak volumetric 263 strain and the final equilibrium value, making the curves steeper. Also, the peak value of the volumetric 264 strain increases with increasing q, while the variation in k shows nearly no effect on the peak value of 265 the volumetric strain.

266 The variation in  $\zeta$  seems to have little effect on the volumetric strain. The change in the volumetric 267 strain using a variation of  $\zeta$  to the reference value is small, and the maximum difference is only

approximately  $2 \times 10^{-7}$ . The shape of the curve does not change with the variation of  $\zeta$ . However, the variation in Young's modulus *E* (the parameter of direct determination of the stress–strain relationship) exhibits a strong effect. Variation in *E* leads to different initial and end values of the volumetric strain, and the peak value changes from approximately  $-1.7 \times 10^{-4}$  to  $-6.7 \times 10^{-3}$ . A change in *E* also has no effects on the time to equilibrium of volumetric strain, with all curves of different E becoming horizontal lines after approximately 30 h.

274 The *ROV* for each parameter for average volumetric strain is shown in Table 5.

Table 5 shows that the average ROV of different parameters varies considerably within a range of -0.0034 to 2.5015. The higher |ROV| is, the more sensitive parameter is. A positive value means that the calculation results are positively correlated with parameters, while a negative value means a negative correlation.

279 The model appears to be more sensitive to E than all other parameters with respect to  $ROV_{ave}$ ,  $ROV_{max}$ 

and  $ROV_{min}$ . There is little change between  $ROV_{max}$  and  $ROV_{min}$  for E at one specific value, which shows

281 that the sensitivity of E is stable throughout the calculation period. Therefore, the strain or displacement

evolution does not affect the sensitivity of E, as observed in another study (Abdollahipour et al. 2019).

283 The peak value of the *ROV* for *E* appears at the initial time. The *ROV* for  $\zeta$  are small and steady, and

284 have negative effects on the results.

The *ROV* value for k and q change markedly throughout the calculation period. The *ROV* of k is positive in the early stage (i.e., before the peak of volumetric strain) and changes to a negative value after the peak. The *ROV* of q actually exhibits the opposite trend. As the parameter increases,  $ROV_{ave}$  of E

and k significantly decreases, indicating that these two parameters are more sensitive at low values.

Similarly, other research (Jin et al. 2016) has shown that displacement is more sensitive to these two parameters of low value in a saturated 3D model using Biot's equation. The  $ROV_{ave}$  of q and  $\zeta$  remain nearly stable during the variation, exhibiting strong symmetry.

#### 292 Parameter sensitivity of hydraulic behavior

293 The variation in specific water content  $\Delta v_w$  of the sample using different input parameter values is shown 294 in Fig. 5.  $\Delta v_w$  begins at 0 and rises to the same value, with the rates of  $\Delta v_w$  change decreasing with time. 295 The model suggests that k and q also have opposite influences on the specific water content. As k296 increases, the rate of increase of  $\Delta v_w$  increases and reaches its final value earlier. However, as q 297 increases, the rate of increase of  $\Delta v_w$  decreases, and the time to reach its final value lengthens. The ROV 298 of k and q change throughout the calculation period. The ROV of k is always positive but decreases to 299 nearly 0 at the end. The  $ROV_{min}$  of -60% variation of k is 0.0140 because in this case, it does not reach a 300 steady state at the end of the simulation. Conversely, the ROV of q remains negative in the model, which 301 is to be expected, as the increase in k promotes the water flow and shorten the time to water pressure 302 equalization according to Equation (4), while q has a negative effect on the water pressure change. 303 As with volumetric strain,  $\zeta$  has a weak influence on  $\Delta v_w$ , with little effect on the basic shape of the 304 curve. The influence of the change in E on  $\Delta v_w$  is also small; however, the pattern of change is unusual. 305 Close inspection of the curve between 30 h and 44 h shows that the value of  $E=2.428\times10^8$  is initially lower than that of  $E=1.670\times10^8$  but surpasses the latter after a few hours and then decreases again 5 h 306 307 later. In general, low values for E contribute to the highest value of  $\Delta v_w$ .

308 The *ROV* of each parameter for the specific water content is shown in Table 6.

#### Analysis of the advantages and limitations of the new Model

The validation described in the previous chapter is based on the UHM model. In this chapter, the new model (UHM-P) is compared with the results of the UHM model, and the limitation of the new model is analyzed in detail.

#### 313 Influence of porosity change on hydraulic and mechanical behavior

314 As discussed before, during the equalization stage, the variation of the model from the experimental 315 results for draining and wetting samples appears to display different trends. Fig. 6a shows the variation 316 in calculation error of  $\Delta v_w$  for samples E and G using the UHM model. This figure clearly shows that 317 the error for sample E increases to approximately 0.002 in the first few hours and then decreases to 318 0.0001 at and beyond 20 h. However, for sample G, the error decreases to approximately -0.002 within 319 6 hours and then increases and remains at 0.002 after 20 h. These trends can be explained by the porosity change during the model period. A value of  $5 \times 10^{10}$  Pa for  $K_{s}$  (Alonso and Alcoverro 1999a) was 320 321 selected to calculate the theoretical true porosity. Fig. 6b shows the theoretical true porosity during the 322 equation stage. The porosities of both E and G are lower than the initial porosity of 0.539 due to the 323 confining pressure, while the UHM model assumes a constant value for porosity. This result may 324 overstate the true porosity during the experimental period, contributing to the positive errors in most of 325 the models (see Equation (13)). The evolution of porosity in E and G are different. As discussed earlier, 326 the water pressures in sample E increase during the experiments, while those in G decrease. Therefore, 327 the water pressure changes induced different changes in porosity. Water pressure shows a positive effect 328 during the late stage of E (the final porosity is higher than the earlier porosity), while the opposite impact 329 was found during the late stage of G.

330 The constant porosity value is replaced with the theoretical dynamic porosity obtained by Equation (7) 331 in the UHM-P model. The results of the equalization stage by the UHM-P model are shown in Fig. 7. 332 For both samples, the modified results are lower than the original model results. Modification of the 333 model reduces the error caused by the overvaluation of porosity during the experiment. As mentioned 334 above, the porosity change is relatively small in the equalization stage; however, the porosity change still 335 strongly influences the numerical calculation. Table 7 shows that the AVRE of the results improves 336 markedly for both samples E and G. The average relative error of E decreases from approx. 0.0985 to 337 0.0543 (a 44.9% improvement in relative error), while this error of G decreases from approx. 0.0962 to 338 0.0623 (a 35.2% improvement in relative error). However, the modification seems to have no marked 339 effect on R.

Fig. 8 shows the modeled results of  $log(p - p_a)$ : *V* during the isotropic compression stage of Vassallo's study, where *V* is the specific volume, which is calculated by:

342 
$$V = 1 + e = 1 + \frac{v}{1 - v}$$
 (1)

The modeled result of *V* remains zero if the porosity change is not considered in the model. The maximum porosity change reached 0.040 (9.2% less than the initial porosity) in S100CI and 0.036 (8.3% less than the initial porosity) in S200CI. The porosity change during the isotropic compression stage is much larger than that during the equalization stage. Therefore, porosity change is also important when studying the mechanical behavior of soil.

#### 348 Model limitations

The proposed model is based on mixture coupling theory, which has the advantage of engaging multiscale coupled processes, by analyzing the Helmholtz free energy difference between the pore water and the wetted matrix. Compared with the previous model, the proposed model considers the evolution of porosity and yields the dynamic porosity equation in the hydro-mechanical model. Several assumptions were made to simplify the discussion, including the following: the soil is homogeneous, isotropic, and linearly elastic during deformation; the coupling coefficient  $\zeta$ , Q and the elastic stiffness are all material-dependent,

or plastic-deformation problems. Additionally, because some soils may experience modulus changes
under different suctions (Sawangsuriya et al. 2009), the proposed model must be modified when dealing

which means that they remain constant in a specific material. Therefore, the model cannot manage large-

359 with such problems.

356

360 Another limitation is that all of the pore space is assumed to be isotropic in the hydraulic flow, which

361 may lead to large errors with fractured media. In this model, air is assumed to always remain at a fixed

- 362 pressure,  $p_{atm}$ , and the air is continuous. These assumptions ensure that no air flows in the porous media.
- 363 Although the air influence on liquid flow can be considered, the model cannot be used to solve a two-
- 364 phase flow problem in deformed porous media.

For some rocks, permeability change is more sensitive than porosity change (Raghavan and Chin 2004),
and the proposed model, which ignores permeability change, may face difficulty in such a problem.
Finally, this study focuses on experimental validation and porosity changes on a relatively small scale.
Water density variations in space and time are not considered in the model, which may be important in

369 field studies or in research focusing on intensive pressure changes.

# 370 Influence of porosity change on deformation energy

371

372	considering porosity change. This figure shows the influence of porosity change on deformation energy.
373	The time derivative of deformation energy $W^*$ is positive in sample E but remains negative in sample G.
374	The $W$ of sample E decreased to approximately -0.2335 in the early stage and then increased to
375	approximately 0 at the end. While the $W$ of G exhibited the opposite trend compared with E, the peak
376	of $W$ appears at approximately 1.95 h in E and at approximately 0.58 h in G and is consistent with the
377	peak of water pressure, which is confirmed in Equation (12).
378	Similarly, the calculation error of $W$ is different in the two samples caused by the porosity change. The
379	$W^{\bullet}$ that considers porosity change is always higher than that that does not in sample E, while the $W^{\bullet}$
380	that considers porosity change is always lower than that does not in sample G. These results are due to
381	the different evolution of water pressure in samples E and G. Water pressure of sample E increased
382	during experiment, while the water pressure of sample G decreased, resulting in positive and negative
383	time derivative of water pressure in samples E and G, respectively. The calculation error index AVRE
384	of $W$ that does not consider porosity change in sample E is approximately 1.00051; in sample G, the
385	AVRE is approximately 1.00053.
386	The 2D contour maps of $W^{\bullet}$ of samples E and G are shown in Fig. 10 and Fig. 11. These two figures
387	show interesting evolutions of $W^{\bullet}$ at different subzones of the total sample. The $W^{\bullet}$ of the lower part is
388	markedly higher than that of the upper part in the beginning. Then, the high value subzone moves up
389	over time, and finally, the $W^{\cdot}$ of the upper part is higher than that of the lower part. This trend indicates

Fig. 9 shows the evolution of deformation energy at the center of samples E and G with and without

390 that the lower part of the sample deforms much earlier than the upper part in the experiment. In addition,

- 391 the absolute value of  $W^{\cdot}$  in the early stage is higher than that in the later stage, which means that the
- deformation is stronger in the early stage than in the later.
- 393 This evolution is similar in sample G. A higher absolute value of W is shown in the lower part at the
- 394 beginning, and a higher absolute value of the upper part is shown in the end. Additionally, a marked
- 395 decrease is observed in the absolute value of W in the total sample over time.
- 396 In this numerical analysis, no additional volume force was loaded on the sample except for the small
- 397 confined pressure in the lateral direction. The deformation is primarily due to the water pressure change,
- 398 which is caused by water absorption in sample E and water drainage in sample G. The calculation results
- of the first term in Equation (11),  $tr(T\dot{E})$  (range about  $4 \times 10^{-5}$  to  $3 \times 10^{-4}$ ), is much lower than the
- 400 second term  $(S_w p)^{\cdot} v$ , which explains the negative value of W in sample E and the positive value in
- 401 sample G.
- The deformation energy is useful in the analysis of the coupling behavior of soil. The lower boundary is the constant water pressure in this analysis. The water pressure change begins at the lower part of the samples, and the change in W and deformation also begins at the lower part. The decrease in the absolute value of  $W^{*}$  is caused by the decrease in the water pressure change rate in the samples. When the water pressure reached an equilibrium value,  $W^{*}$  reached nearly zero, and then the deformation energy Wfinally stabilized, leading to a new steady stage of the samples (i.e., no further deformation).
- 408 Conclusion

410 P) for soil was developed. The experimental data for two different unsaturated triaxial tests were used to

<sup>409</sup> In this study, a new coupled unsaturated hydromechanical model considering dynamic porosity ((UHM-

411 validate this model. Validation shows that the UHM-P model can basically describe the hydro-412 mechanical multifield behavior in unsaturated soil, when ignoring porosity changes. The average relative 413 errors of all the experimental samples are between 0.8631 and 1.3046, and all the correlation coefficients 414 are between 0.9028 and 0.9981. 415 Sensitivity analysis shows that Young's modulus (E) is the most sensitive parameter in the calculation 416 of volumetric strain, followed by permeability (k). The sensitivity of the model to poroelastic storage 417 (q) and Biot's coefficient  $(\zeta)$  seems to be stable over a range of parameter values, while the sensitivities 418 of k and E are higher at low values. The modeled hydraulic behavior is most sensitive to k, followed 419 by q. The sensitivity of  $\zeta$  is also stable over a range of parameter values, while the sensitivity of the 420 other three parameters is higher at low values. 421 Modeled results show that the UHM model cannot calculate the specific volume change. Even during 422 the relatively low deformation stage, the calculation results of the specific water volume are more 423 realistic when the dynamic porosity is considered, leading to 44.9% and 35.2% improvements in the 424 relative error in samples E and G, respectively. 425 The deformation energy results show that the relative calculation error when not considering porosity

426 change is approximately 0.00051 and 0.00053 in samples E and G, respectively.

#### 427 Data Availability Statement

- 428 Some or all data, models, or code that support the findings of this study are available from the
- 429 corresponding author upon reasonable request (Numerical model).

21

#### 430 Acknowledgment

431 This study was supported by the National Key R&D Program of China (No. 2019YFC1805503), the

432 National Key R&D Program of China (No. 2018YFC1800905) and the Key Science and Technology

433 Projects of the Inner Mongolia Autonomous Region (2019ZD001).

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540

Fig. 1. 2D Numerical model domain and conditions





#### 544 S100P100TG, b for sample S100P200TG, c for sample S100P400TG and d for sample S100CI



547 Fig. 3. Comparison of experimental and model results for  $\Delta v_w$  in Biglari's test, a for Sample D, b







553

modulus







556 Fig. 5. Different responses of  $\Delta v_w$  to the variation in input parameters, a for poroelastic storage, b



for permeability, c for Biot's consolidation coefficient and d for Young's modulus





Fig. 6. a for the calculation error of the UHM model and b for the theoretical porosity



calculated by the UHM-P model during experimentation





porosity change, a for Sample E and b for Sample G

563







#### sample S100CI, b for sample S200CI



568 Fig. 9. Deformation energy evolution at the center of samples with and without considering

569

porosity change, a for sample E and b for sample G





571 Fig. 10. Deformation energy evolution of sample E with considering porosity change, a for the

profile at 0 h, b for 0.1 h, c for 5 h and d for 120 h





574 Fig. 11. Deformation energy evolution of sample G with considering porosity change, a for the

575

### profile at 0 h, b for 0.1 h, c for 5 h and d for 40 h



#### Table 1. Hydraulic and mechanical parameters and model setting

Biglari's test	Vassallo's test	Meaning	Parameters
2.428×10 <sup>8</sup> Pa	3×10 <sup>8</sup> Pa	Young's modulus	Ε
0.214	0.2	Poisson's ratio	$\theta$
$2.65 \times 10^3 \text{ kg} \cdot \text{m}^{-3}$	$2.748 \times 10^3 \text{ kg} \cdot \text{m}^{-3}$	Solid density	ρ
5.12 cm	$2.485 \times 10^4$ cm	Parameter of Van Genuchten equation	М
1.1459	1.4407	Parameter of Van Genuchten equation	n
8.5×10 <sup>-4</sup> Pa·s	$8.5 \times 10^{-4}$ Pa·s	Viscosity of water	V
0.539	0.441	Porosity	υ
1	1	Biot's coefficient	ζ
$1.3 \times 10^{-19} m^2$	$1 \times 10^{-20} m^2$	Permeability	k

5×10 <sup>-10</sup> Pa <sup>-1</sup>	5×10 <sup>-10</sup> Pa <sup>-1</sup>	Poroelastic storage	q
$1.2 \times 10^2$ h	$2 \times 10^4$ min	Calculation period	Т
0.1 h	10 min	Timestep	Step

# Table 2. Parameters used in the sensitivity analysis

60%	Tarameters			
 2.08×10 <sup>-19</sup>	1.69×10 <sup>-19</sup>	9.10×10 <sup>-20</sup>	5.20×10 <sup>-20</sup>	$k[m^2]$
8×10 <sup>-10</sup>	6.50×10 <sup>-10</sup>	3.50×10 <sup>-10</sup>	2.00×10 <sup>-10</sup>	$q[Pa^{-1}]$
-	-	0.7	0.4	ζ
3.885×10 <sup>8</sup>	3.156×10 <sup>8</sup>	$1.670 \times 10^{8}$	9.712×10 <sup>7</sup>	E[Pa]

578 Note: Only 0.4 and 0.7 were evaluated because Biot's coefficient is always <= 1.

579

577

Table 3. Error for  $\Delta v_w$  comparing the model to experimental data in Vassallo's test

S100CI	S100P400TG	S100P200TG	S100P100TG	Error estimation
1.0698	0.8631	1.1909	1.3046	AVRE
0.9289	0.9028	0.9981	0.9559	R

580

Table 4. Error for  $\Delta v_w$  comparing the model to experimental data in Biglari's test

G	F	Е	D	Error estimation
0.9038	0.9481	1.0985	1.1795	AVRE
0.9832	0.9975	0.9963	0.9959	R

Table 5. ROV of each parameter for average volumetric strain

ROV <sub>min</sub>	ROV <sub>max</sub>	ROVave	Variation (%)	Parameter
-0.0332	0.0212	-0.0126	-60	k
-0.0287	0.0176	-0.0077	-30	iv iv

-0.0210	0.0134	-0.0042	30	
-0.0188	0.0117	-0.0034	60	
-0.0174	0.0347	0.0056	-60	
-0.0158	0.0287	0.0054	-30	a
-0.0125	0.0217	0.0056	30	7
-0.0108	0.0193	0.0055	60	
-0.000024	0.00086	0.000094	-60	r
-0.000030	0.00087	0.000094	-30	4
2.4976	2.5087	2.5015	-60	
1.4242	1.4353	1.4293	-30	F
0.7665	0.7718	0.7694	30	Ŀ
0.6239	0.6266	0.6251	60	

# **Table 6.** *ROV* of each parameter for $\Delta v_w$

ROV <sub>min</sub>	ROV <sub>max</sub>	ROVave	Variation (%)	
0.0140	0.5912	0.2272	-60	
0.0030	0.5318	0.1465	-30	7
0.0002	0.4601	0.0864	30	ĸ
0.0002	0.4286	0.0729	60	
-2.3043	0.0002-	-0.1328	-60	
-1.3140	-0.0002	-0.1162	-30	a
-0.7074	-0.0014	-0.1090	30	Ч
-0.5748	-0.0029	-0.1050	60	
0	0.8823	0.0064	-60	
0	0.8839	0.0064	-30	ç

	-60	-0.0156	0	-1.9291
F	-30	-0.0124	0	-1.0990
L	30	-0.0104	0	-0.5907
	60	-0.0064	0	-0.4797

583	Table 7. Error estimation chan	ge of $\Delta v_{\rm m}$ in samples E	and G of Biglari's test with	th considering
		<b>8</b> W <b>I</b>		

porosity change				
G(UHM-P)	G(UHM)	E(UHM-P)	E(UHM)	Error estimation
0.9377	0.9038	1.0543	1.0985	AVRE
0.9831	0.9832	0.9963	0.9963	R