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1	A porous medium approach to the 3D modelling of an entire rotating packed bed
2	for post-combustion carbon capture
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6	Abstract: Rotating packed bed (RPB) technology shows great potential for post-
7	combustion capture. However, the capture process inside the full RPB is difficult to
8	simulate, due to the complexity of the process and the neglect of the CO ₂ capture in the
9	outer cavity zone. In this paper, a full 3D CFD model, including the packing and the
10	inner and outer cavity zones, has been established employing the Eulerian porous
11	medium method coupled with various sub-models. The CO2 capture performance in the
12	packing and outer cavity zones has been quantitatively analyzed under different
13	operating conditions. The simulation results show good agreement with the
14	experimental data, and the contribution of the outer cavity zone to the CO2 capture of
15	the RPB is in the range of 28%~42%. This work provides a new approach to efficiently
16	simulate the mass transfer process in the RPB.
17	Keywords: rotating packed bed, Eulerian method, 3D porous media model, CO2
18	absorption, outer cavity zone

20 Nomenclatures

a_S	specific area of the packing	$S_{h,i}$	heat of chemical reaction, $W\cdot$
	materials, $m^2 \cdot m^{-3}$		m ⁻³
Ae	effective interfacial area, $m^2 \cdot$	\vec{S}_{GL}	interfacial force between the
	m ⁻³		gas and liquid phases, $N \cdot m^{-3}$
С	concentration, kmol \cdot m ⁻³	t	time, s
C_{pG}	specific heat of gas phase, J.	T_i	temperature of phase <i>i</i> , K
	$kg^{-1} \cdot K^{-1}$		
C_D	drag coefficient	и	flow rate per unit area, $m \cdot s^{-1}$
D	diffusivity, $m^2 \cdot s^{-1}$	\vec{u}_i	velocity of phase <i>i</i> , $m \cdot s^{-1}$
Ε	enhancement factor	V	volume, m ³
E'_i	total energy, J	We	Weber number
f,f^*	drag function	x	radial length of a cell, m
f_e	ratio of wetted packing area to the	y_{i,CO_2} ,	CO ₂ mole fraction in the gas
	total packing surface area	$y_{o,CO_2},$	inlet and outlet and on the outer
		y_{op,CO_2}	packing surface, respectively
<i>Ē</i> disni	dispersion forces of phase i , N \cdot	Y	mass fraction of species
uisp,i	m ⁻³		
\vec{F}_{dragi}	interaction force between the	Greek	
urug,t	phase <i>i</i> and solids of the packing		
	material, N \cdot m ⁻³		
g	gravitational force, $m \cdot s^{-2}$	α	phase saturation
h	height of the packing, m	γ	packing void fraction (porosity)
h_j	enthalpy of the species, $J \cdot mol^{-1}$	β	dynamic contact angle, °
h_{GL}	heat transfer coefficient between	Е	volume fraction
	the gas and liquid phases, W \cdot		
	$m^{-2} \cdot K^{-1}$		
Н	Henry constant	δ	diffusion layer thickness, m
На	Hatta number	ω	angular velocity, rad \cdot s ⁻¹
Ī	mass diffusion flux of species, kg \cdot	κ_G	thermal conductivity of the gas
	$m^{-2} \cdot s^{-1}$		phase, $W \cdot m^{-1} \cdot K^{-1}$
k'_i	thermal conductivity, $W \cdot m^{-1} \cdot$	θ	angle of flow direction slop to
	K ⁻¹		the bed axis, °
k _i	mass transfer coefficient of phase	$ ho_i$	density of phase <i>i</i> , kg \cdot m ⁻³
	$i, \mathbf{m} \cdot \mathbf{s}^{-1}$		
<i>k</i> ₂	reaction rate constant, $m^3 \cdot$	Δho_{GL}	absolute value of the density
	$mol^{-1} \cdot s^{-1}$		difference between liquid and
			gas phases, kg \cdot m ⁻³
K_L	overall mass transfer coefficient,	μ	dynamic viscosity, $kg \cdot m^{-1} \cdot$
	$\mathbf{m} \cdot \mathbf{s}^{-1}$		s ⁻¹

K _{GL}	momentum exchange coefficient	μ_e	effective dynamic viscosity,
	between the gas and liquid phases		$kg \cdot m^{-1} \cdot s^{-1}$
$K_L a_e$	overall volumetric mass transfer	τ	particulate relaxation time, s
	coefficient, s ⁻¹		
M	molar mass, $g \cdot mol^{-1}$	$\overline{\overline{\tau}}$	stress tensor, N \cdot m ⁻²
N_{CO_2}	transfered CO ₂ between the gas	$ar{ar{ au}}_{eff,ij}$	effective shear tensor
	and liquid, $kg \cdot m^{-1} \cdot s^{-1}$		
Nu	Nusselt number	λ_{RT}	Rayleigh-Taylor instability
			wavelength
P	pressure, Pa	σ	surface tension, $N \cdot m^{-1}$
P_c	capillary pressure, Pa	ν	kinematic viscosity, $m^2 \cdot s^{-1}$
Pr	Prandtl number	Subscrip	ots
q	liquid initial velocity number	cell	mesh cell
Q	volume flow rate, $m^3 \cdot s^{-1}$	CO_2	carbon dioxide
$Q_{h,GL}$	transferred heat between the gas	G	gas phase
	and liquid phases, $W \cdot m^{-3}$		
r	radius or radial distance, m	i	=G, L
r _i	inner packing radius, m	j	species
r_o	outer packing radius, m	L	liquid phase
r_{CO_2}	CO ₂ reaction rate, $kg \cdot m^{-3} \cdot s^{-1}$	MEA	monoethanolamine
R	gas constant	0	outer cavity zone
R_{ij}	net rate of production of the	OW	computational cell next to the
	species j by the reaction, mol \cdot		casing wall
	$m^{-3} \cdot s^{-1}$		
Re	Reynolds number	Р	packing zone
Re [,]	relative Reynolds number	PW	computational cell next to the
			packing wall
$S_{m,i}$	mass transfer rate between gas and	S	solids phase for the packing
	liquid, kg \cdot m ⁻³ \cdot s ⁻¹		materials

21 **1. Introduction**

CO₂ emissions, as a kind of greenhouse gas, has become an urgent issue in the 21st century (Khan, 2017). In response to this important issue, the EU has a commitment for a net-zero greenhouse gas emissions target by 2050 (Garba, 2012), which aims to reach a balance between the amount of greenhouse gas emissions produced and the amount captured. Particularly, the total power plants contribute a huge amount of CO₂, which could contribute to 30% of the total CO₂ emissions. Therefore, reducing CO₂ emissions released by the power stations has become an emergent activity. 29 Currently, the amine-based post-combustion CO₂ capture (PCC) is the most common 30 and commercial way to capture CO₂ using the conventional packed bed (CPB), which 31 results in the CPB being large and costly. In addition, the high amine solvent flow rate 32 demanded by the large CPB would consume large heat energy for the rich solvent 33 regeneration. On the other hand, the rotating packed bed (RPB), one of the most 34 important applications in the area of process intensification, has the potential to enhance 35 the CO₂ capture performance and reduce the packed bed size by applying a high 36 centrifugal force (100-1000 times gravity) (Liu et al., 2019). However, the RPB is on 37 the early stage of the technological development for PCC, and the CO₂ capture 38 processes, including the hydrodynamics, heat and mass transfer, are still unclear inside 39 a RPB. Thus, it requires extensive research and design optimization.

40 In general, the RPB may be divided into three parts - inner cavity, packing and outer 41 cavity zones, see Figure 1. At the beginning, the liquid jets out from the liquid 42 distributor and passes through the inner cavity zone entering the packing region, then 43 the liquid contacts and reacts with the opposing gas flow on its way outwards through 44 the packing region. During this process, the liquid breaks into tiny droplets, or generates 45 thin films on the packing surface due to the strong interaction with the rotating mesh packing, which can enhance the interfacial area and micromixing (Yang et al., 2015b) 46 47 between gas and liquid. In addition, the liquid surface renewal rate in the packed bed is 48 improved, which also leads to a large increase in the CO₂ capture. In addition, the liquid leaves the packing with a large tangential velocity and enters the outer cavity zone in 49 50 the form of liquid droplets. Finally, the liquid droplets hit the cavity casing wall and the 51 liquid phase flows out from the liquid outlet tube under the influence of the gravitational 52 force. It should be noted that before the liquid leaves the outlet tube, the liquid has been 53 reacting with the gas phase all the time in the outer cavity zone. In addition, the space 54 ratio of the outer cavity zone to the packing region is large (Guo et al., 2019). As a 55 result, it can be concluded that the CO₂ capture in the RPB occurs not only in the 56 packing zone, but also in the outer cavity zone.

57 However, most investigations only focus on the packing region, the contribution of the 58 outer cavity zone for the CO₂ capture is rarely studied. In the experimental 59 investigations, it has been reported that the effective interfacial area of the outer cavity 60 zone could be up to 30% of the entire effective area (Guo et al., 2014; Sang et al., 61 2017a). Further, experiments showed that the contribution of the outer cavity zone to 62 the mass transfer was determined to take up approximately 13-25% of the total mass 63 transfer in the entire RPB (Sang et al., 2019; Yang et al., 2011). Therefore, the CO₂ 64 absorption in the outer cavity zone cannot be ignored.





66

Figure 1. Schematic diagram of a typical RPB for CO₂ capture.

67 In addition, the Volume of Fraction (VOF) method has widely used to computationally study the CO₂ capture in the RPBs. For instance, the micromixing and mass transfer 68 have been studied in a 2D packing region along with the Volume of Fraction (VOF) 69 70 method and user defined functions (UDF) (Guo et al., 2016; Yang et al., 2016). The 71 advantage of employing the VOF method is that it can reasonably, clearly and 72 accurately capture the interface between the gas and liquid, including the evolution of the liquid droplets in the RPB. Due to the multiscale nature of the liquid flow in an RPB, 73 74 the computational mesh should be fine enough to resolve the finest droplet and liquid 75 films, typically in millimeter and even micron meter scales, in addition to the packings, 76 within a meter scale RPB. When modelling in 3D this will require a prohibitively large number of computational cells and computational time. This is the reason why most of 77

78 the CFD simulations using CFD methods are for small RPB models and mostly in 2D. 79 If the packing in an RPB is treated as a porous medium using the Eulerian model, where the liquid flow is modelled as a continuum, there is no need to resolve the packing or 80 81 the liquid droplets. Instead, the interfacial area and the momentum transfer between the 82 phases are modelled using experimental correlations. Clearly, this method could not 83 produce a clear picture of the real discrete liquid droplet flow and its accuracy relies on 84 correlations employed for the resistance among the phases and the interfacial area. 85 Nevertheless, the Eulerian method could significantly reduce the simulation time and 86 with appropriate validations, it can potentially be utilized to address the significant 87 difficulties of modelling 3D and large-scale RPBs (Lu et al., 2019; Zhang et al., 2022a).

88 However, the Eulerian porous medium approach has so far seldom been employed for 89 the RPB because the momentum transfer between the phases has to be closed by closure 90 equations. In other words, a suite of mathematical formulations is essential for 91 accurately expressing the flow characteristics in the RPB, such as the sub-models for 92 the forces for the phases and the effective interfacial area between the phases (Fourati 93 et al., 2013; Iliuta et al., 2014). Since the characteristics of the multiphase flow are very 94 different when passing through the rotating packing and in the static cavity space, sub-95 models for the forces and the interfacial area should be separately introduced for the 96 packing region and the outer cavity zone, which is quite complex. For the packing 97 region, not only the porous resistance should be considered due to the wire mesh 98 packing, but also the dispersion force should be taken into account due to the strong 99 interaction between the liquid and the packing (Jiang et al., 2002; Xu et al., 2019; Zhang 100 et al., 2022b). In addition, it is complex to predict the contact area between gas and 101 liquid due to the non-uniform liquid distribution in the packing zone. The above force 102 and the effective mass transfer area models for the non-uniform multiphase flow in the 103 packing zone have been developed and validated in our previous paper (Zhang et al., 104 2022a). However, for the outer cavity zone, the liquid phase exists in the form of liquid 105 droplets in the cavity space and liquid film on the casing wall. As a result, different 106 contact area correlations for the liquid droplets and film should be proposed. In addition, 107 there is no existing force models for the outer cavity zone via employing the Eulerian 108 method. Therefore, the appropriate model and correlation for the outer cavity zone 109 should be developed for accurately analysing the hydrodynamics in this zone.

110 In addition to the momentum transfer between phases, the heat and mass transfer should 111 be modelled in the RPB. Many semi-empirical correlations for predicting the mass 112 transfer coefficients have been proposed in the literature based on different applications 113 and packing systems of the RPBs (Munjal et al., 1989), however, these correlations are 114 too general to be confidently employed for the CO₂-MEA absorption system employed 115 (Lu et al., 2019). A method that has been commonly used to model the chemically-116 enhanced mass transfer in the CPB, is to use an enhancement factor coupled with a 117 mass transfer theory, such as the two-film, penetration, or surface renewal theories. 118 Although penetration and surface renewal theories are more accurate in calculating the 119 mass transfer by introducing the contact time and surface renewal frequency, these two 120 parameters are difficult to measure in the RPB, especially under unsteady-state 121 conditions. Furthermore, the two-film reaction-enhanced mass transfer model has 122 already shown sufficient accuracy for expressing the CO₂ mass transfer among the 123 phases in many CPB investigations (Dashliborun et al., 2019; Kim et al., 2017; Pham 124 et al., 2015), which is introduced into the RPB model as presented in Sections 2.4 and 125 2.5.

The thermodynamics in the RPB should also be carefully considered because the liquid temperature and the chemical reaction rate could affect each other in this CO₂ chemisorption system. The heat released by the chemical reaction increases the liquid temperature; meanwhile, some of the heat would be carried away by the gas flow due to the temperature difference between the phases so that the gas temperature also increases. The increasing phase temperature would speed up the reaction rate until the whole CO₂ absorption system reaches a balance. In particular, the efficiency of the heat 133 transfer between the phases is determined by the diameter of the liquid droplet. The 134 smaller the droplet diameter, the more heat transfer takes place. The predicted heat and 135 mass transfer processes can be made accurate only when the liquid droplet diameters 136 in the packing and outer cavity zones are set properly.

137 In our previous work, the CO₂ capture performance in the RPB was explored in a 2D 138 packing region by the Eulerian porous medium method (Lu et al., 2019). However, the 139 predicted liquid outlet temperature and the overall volumetric mass transfer rate $(K_L a_e)$, 140 which was defined in Section 3.3, were lower than the experimental data because the 141 outer cavity zone was not taken into consideration at the time and the liquid droplet 142 diameter was set using the default value in the software. In this paper, for accurately 143 studying the CO₂ capture process within a RPB, a comprehensive 3D RPB model, 144 including the packing, inner and outer cavity zones, was built based on a practical pilot-145 scale RPB coupled with various sub-models. The $K_L a_e$ and liquid outlet temperature 146 under various operating conditions were analysed. Furthermore, the ratio of the CO₂ 147 capture performance in the packing zone and outer cavity zone was quantitatively 148 investigated. Therefore, this paper provides, for the first time, a new feasible approach 149 to effectively and accurately predict the mass transfer within the entire RPB.

150 **2. CFD** simulations

151 In this section, a 3D model has been built that incorporates the porous resistance, 152 dispersion force, two-film reaction enhancement mass transfer and heat transfer 153 models in order to achieve an accurate prediction of the flow dynamics, mass transfer 154 and heat transfer in the whole RPB. Since the resistance and dispersion models 155 developed for the packing zone have been illustrated in our previous paper (Zhang et 156 al., 2022a), only some of the most important models and equations, such as the 157 governing equations, mass and heat transfer models, and the correlations of the 158 effective interfacial area and liquid droplet diameter, are presented in this paper.

159 **2.1 Geometry of the RPB**

8

160 In order to establish and validate the 3D RPB model, the detailed information of the 161 experimental rig and the quality of the experimental results are required. Among all the published experimental works on RPBs, the Kolawole experimental data (Kolawole, 162 163 2019) has been selected as the comparison data. The reasons are: (i) the relative detailed 164 operating conditions and dimensions of the RPB have been offered, especially the size 165 of the packing and outer cavity zones. (ii) The CO₂ capture performance with different 166 operational conditions is analyzed by various indexes, including the overall mass 167 transfer coefficient, the liquid and gas phase temperatures at the outlets, etc., which 168 could help to validate the simulation results. And (iii) the experimental data has been 169 presented in the published paper (Lee et al., 2017) and previously utilized by Lu et al. 170 (2019) for verifying their models, thus meaning that the experimental data is reliable. 171 However, no experimental data related to the flow characteristics (i.e. liquid holdup) 172 has been shown in the thesis (Kolawole, 2019), thus, the flow dynamics obtained by 173 the experiment and simulation are not compared in this work.

174 The 3D RPB geometry has been reproduced in Figure 2 according to the experimental 175 rig (Kolawole, 2019). The inner diameter, outer diameter and height of the wire mesh 176 packing are 80, 300 and 20 mm, respectively, and it is made from stainless steel with 177 an expanded mesh grade 707. The void fraction and specific area of the packing are 178 0.801 and 663 m^2/m^3 , respectively. The diameter and the estimated height of the RPB 179 casing are 360 and 180 mm, respectively and the space ratio of the outer cavity zone to 180 the packing zone is about 11. There are two gas inlet tubes and one liquid inlet tube 181 with diameters of 60 and 14 mm, respectively. In order to make the mesh in the outer 182 cavity tetrahedral and further reduce the mesh number, the shape of the gas inlet tubes 183 and one liquid inlet tube is simplified to be a rectangle but with the same flow cross-184 sectional area. In addition, a gas outlet tube is outside the liquid inlet tube with a 185 diameter being 40 mm. A 2-arm liquid distributor is used, and its length and diameter 186 are 22.4 and 7 mm, respectively. In Figure 2, only two liquid nozzle outlet holes are presented. In order to save computational time, the liquid flows passing through the 187

liquid inlet tube and 2-arm liquid distributor has not been modelled. Initially, a liquid pressure outlet boundary was built at the bottom of the RPB to ensure that the liquid phase flows out from the RPB. However, after the simulation, we found that a large amount of the gas phase would flow out from this boundary at the same time. Thus, an elimination zone is employed at the bottom of the RPB to remove the liquid phase but to retain the gas phase, and the appropriate source equations for this zone can be found in (Zhang et al., 2022a).



195

196 **Figure 2.** Schematic diagram of the 3D RPB (1, gas outlet; 2, case; 3, gas inlet; 4,

197 liquid elimination zone; 5, liquid inlet; 6, outer cavity zone; 7, packing zone; 8, inner

cavity zone).

198

199 **2.2 Governing fluid flow equations**

200 The main assumptions of the RPB model are as follows:

- 201 (i) The wire mesh packed bed is a homogenous porous medium.
- 202 (ii) The gas is incompressible.
- 203 (iii) The pressure field is shared by the gas and liquid phases.
- 204 (iv) The liquid phase exists in the form of the film and droplets in the packing
 205 and outer cavity zone.
- 206(v)The mass transfer only happens in the packing and outer cavity zones207due to the small interfacial area and short contact time in the inner cavity208zone.
- 209 (vi) The heat transfer only occurs between the gas and liquid phases and the

casing wall is adiabatic.

210

213

The continuity equation utilized to calculate the overall mass conservation is given asfollows:

$$\frac{\partial}{\partial t}(\varepsilon_i \rho_i) - \nabla \cdot (\varepsilon_i \rho_i \vec{u}_i) = S_{m,i} \tag{1}$$

where ρ_i is the density of the *i*th phase (*i* = G for gas or L for liquid), *t* is the time, \vec{u}_i is the phase velocity, $S_{m,i}$ is the mass transfer rate between phases, and ε_i is the phase fraction, which is defined as follows:

217
$$\varepsilon_i = \alpha_i \gamma = \frac{V_i}{V_G + V_L} \gamma = \frac{V_i}{V_G + V_L + V_S}$$
(2)

where α_i is the phase saturation ($\alpha_i = \frac{V_i}{V_G + V_L}$), V_i is the volume of the *i*th phase, and γ and V_S are the porosity and solid volume of the packing, respectively, which are 1 and 0 for the outer cavity zone.

The momentum conservation equations for the packing and outer cavity zones are different since the drag, capillary and mechanical dispersion forces only exist in the packing region. The governing momentum equations for the whole RPB are as follows:

224
$$\frac{\partial}{\partial t}(\varepsilon_L \rho_G \vec{u}_L) + \nabla \cdot (\varepsilon_L \rho_L \vec{u}_L \vec{u}_L) = -\varepsilon_L \nabla P + \varepsilon_L \nabla P_c + \nabla \cdot (\varepsilon \bar{\tau}_L) + \varepsilon_L \rho_L \vec{g} - \vec{F}_{drag,L} + \vec{S}_{GL} + \vec{F}_{disp,L}(3)$$

225
$$\frac{\partial}{\partial t}(\varepsilon_G \rho_G \vec{u}_G) + \nabla \cdot (\varepsilon_G \rho_G \vec{u}_G \vec{u}_G) = -\varepsilon_G \nabla P + \nabla \cdot (\varepsilon \bar{\tau}_G) - \vec{F}_{drag,G} + \varepsilon_G \rho_G \vec{g} - \vec{S}_{GL} + \vec{F}_{disp,G}(4)$$

where *P* is the pressure, P_c is the capillary pressure, which is produced by the difference in the pressures across the fluid interface and only included in the liquid phase momentum equation, model details please refer to (Zhang et al., 2022a), $\bar{\tau}_i$ is the stress tensor, $\vec{F}_{drag,i}$ is the drag force between the phase and packing, \vec{S}_{GL} is the interfacial force between the gas and liquid, and $\vec{F}_{disp,i}$ is the mechanical dispersion force.

231 The species transport equation and energy equation are shown as follows:

232
$$\frac{\partial \left(\varepsilon_{i}\rho_{i}\vec{u}_{i}Y_{ij}\right)}{\partial t} + \nabla \cdot \left(\varepsilon_{i}\rho_{i}\vec{u}_{i}Y_{ij}\right) = -\nabla \cdot \left(\varepsilon_{i}\vec{J}_{ij}\right) + \varepsilon_{i}R_{ij}$$
(5)

233
$$\frac{\partial(\varepsilon_i\rho_iE_i)}{\partial t} + \nabla\cdot\left(\varepsilon_i\vec{u}_i(\rho_iE_i'+P)\right) = \varepsilon_i\nabla P\left(k_i'\nabla T_i - \sum h_j\vec{J}_{ij} + \left(\bar{\tau}_{eff,ij}\cdot\vec{u}_i\right)\right) + \varepsilon_iQ_{h,i} + \varepsilon_iS_{h,i}(6)$$

where Y_{ij} is the mass fraction of the species *j* in *i*th phase, such as the CO₂ in the gas phase or MEA in the liquid phase (*j*= species), R_{ij} is the production of the species *j* by the reaction, such as MEACOO⁻ (HOC₂H₄NHCOO⁻), \vec{J}_{ij} is the mass diffusion flux, E'_i is the total energy, k'_i is the thermal conductivity, h_j is the enthalpy of the species, $\bar{\tau}_{eff,ij}$ is the effective shear tensor, $Q_{h,i}$ is the transferred heat between the phases, see Section 2.6, and $S_{h,i}$ is the heat of chemical reaction, which may be obtained by $S_{h,i} = -\sum_j h_j R_{ij}$ (Mardani and Mahalegi, 2019).

241 **2.3 Effective interfacial area and forces**

Since the effective gas-liquid interfacial area and forces are quite different between the 242 243 packing zone and outer cavity zone, the effective interfacial area correlation and force 244 models for these two zones should be separately developed. In the packing region, 245 models for the interfacial, drag and dispersion (including capillary and mechanical dispersion) forces and the effective interfacial area for the non-uniform multiphase flow 246 247 in the RPB have been introduced in the previous work (Zhang et al., 2022a). It is worth 248 mentioning that the angle between the flow direction and the bed axis (θ) included in the porous resistance model is set as 10°, which is determined by a commonly used 249 250 specification of the wire meshes with the wire mesh diameter and the centre distance 251 between the wire mesh being 0.6 and 3.5 mm, respectively. Also, this value has been 252 employed in the work of Zhang et al. (2022c). Furthermore, when the dynamic contact angle (β) in the effective interfacial area model is set as 12°, the modelled fractional 253 254 effective interfacial area for the cases simulated in Section 3 is in a reasonable range of 255 0.29-0.52 (Luo et al., 2017; Yang et al., 2011) and the CO₂ capture coefficient matches

the experimental data well. In this section, only the effective interfacial area and force model for the outer cavity zone are discussed.

258 2.3.1 Effective interfacial area model

264

At present, no correlation of the effective interfacial area has been introduced for the outer cavity zone of the RPBs in the published papers. For a spherical bubble or droplet, the algebraic interfacial area concentration models are derived from the surface area to volume ratio $(A_p = \frac{\pi d_p^2}{\frac{1}{6}\pi d_p^3} = \frac{6}{d_p})$. When using the Eulerian multiphase model, a commonly used equation, as given in Eq. 7, could be utilized to estimate the effective

265
$$A_{e,0} = \frac{6\varepsilon_G \varepsilon_L}{d_{p,0}}$$
(7)

interfacial area, which has been built in Fluent ia-symmetric model.

where $A_{e,0}$ is the effective interfacial area in the outer cavity zone, $d_{p,0}$ is the average diameter of the liquid droplets in the outer cavity zone. The ia-symmetric model not only considers the gas and liquid volume fraction, but also takes into account the liquid diameter. In the experimental work of Sang et al. (2017b), they concluded that the liquid exists in the outer cavity zone mainly in the form of droplets and a correlation of the average droplet diameter in the outer cavity zone has been proposed as follows (Sang et al., 2017b):

273
$$d_{p,o} = 0.042We^{-0.272}Re^{0.068}(\frac{u_0}{\omega r_o})^{0.098}r_o$$
(8)

274
$$We = \frac{\rho\omega^2 r_o^3}{\sigma}; Re = \frac{\rho\omega r_o^2}{\mu}; q = \frac{u_0}{\omega r_o}; u_0 = \frac{Q_L}{2\pi r_i h}$$
(9)

where *We* is the Weber number, *Re* is the Reynolds number with outer packing radius as the characteristic linear dimension, ω is the angular velocity, Q_L is the liquid volumetric flow rate, σ is the liquid surface tension, μ is the liquid viscosity, and r_o , r_i and *h* are the outer radius, inner radius and height of the packing, respectively. Because the liquid on the casing wall exists in the form of the liquid film instead of the liquid droplet, Eq. (7)-(9), which are developed for the liquid droplets in the outer cavity space, cannot be used for estimating the contact area between the gas and the liquid film near the casing wall. If the casing wall surface is fully covered by the liquid film, then the effective interfacial area in the cell next to the casing wall ($a_{e,OW}$) is estimated as follows:

285
$$A_{e,OW} = \frac{A_{cell,OW}}{V_{cell,OW}} = \frac{1}{\Delta x_{OW}}$$
(10)

286 where $A_{cell,OW}$ and $V_{cell,OW}$ are the wall area surface and volume in the computational cell next to the casing wall, respectively, and Δx_{OW} is the radial length of the cell 287 288 normal to the casing wall. For the case studied in this paper, the radial length Δx_{OW} is 289 0.005 m, thus, the effective interfacial area near the casing wall is no more than 200 m^2/m^3 . According to Equation (10), it appears that the maximum effective interfacial 290 291 area $(A_{e,OW})$ depends on Δx_{OW} . Nevertheless, it should be noted that the unit of the interfacial area $(A_{e,OW})$ is m²/m³, and therefore it is more suitable to describe $A_{e,OW}$ 292 293 as the interfacial area concentration. The real physical interfacial area (m^2) could be 294 obtained by $A_{OW} = A_{e,OW}V_{cell,OW} = A_{cell,OW}$, which is independent on the first-layer 295 mesh size at the wall.

In addition, within the packing bed, the method given by (Lu et al., 2019) has been employed to calculate the effective interfacial area in the cell next to the packing wall as follows:

299
$$A_{e,PW} = \frac{A_{PW}f_e + a_s f_e V_{cell,PW}}{V_{cell,PW}} = \left(\frac{1}{\Delta x_{PW}} + a_s\right) f_e \tag{11}$$

where $A_{cell,PW}$ and $V_{cell,PW}$ are the wall surface area and volume of the cell next to the packing wall, respectively, Δx_{PW} is the radial length of the cell normal to the packing wall, a_s is the specific area of the packing, f_e is the fractional effective interfacial area in the packing, which is the ratio of the wet area to the total packing area $(\frac{A_{e,P}}{a_s})$, and f_e of the packing walls is considered to be the same as that in the packing region.

305 2.3.2 Interfacial force model

The interfacial force is the only force for the liquid and gas phases in the outer cavity zone. It is due to the momentum exchange between the gas phase and liquid phase and this force is usually presented as an exchange coefficient in Ansys Fluent. There are many built-in models to calculate the interfacial exchange coefficient and a model, named universal drag laws (Ishii and Zuber, 1979), has been designed for the bubbleliquid and droplet-gas flows that could be utilized for modelling the interfacial exchange coefficient in the outer cavity zone, namely

313
$$K_{GL} = \frac{\varepsilon_G \varepsilon_L \rho_L f}{\tau_L}$$
(12)

314
$$\tau_L = \frac{\rho_L d_p^2}{18\mu_G} \tag{13}$$

$$f = \frac{C_D R e}{24} \tag{14}$$

316
$$Re' = \frac{\rho_G |\vec{u}_G - \vec{u}_L| d_p}{\mu_e}$$
(15)

317
$$C_{D} = \begin{cases} \frac{24}{Re'} & Re' < 1\\ \frac{24}{Re'} (1 + 0.1Re^{0.75}) & 1 \le Re' \le 1000\\ \frac{2}{3} \left(\frac{d_{p}}{\lambda_{RT}}\right) \left\{\frac{1 + 17.67f^{*6/7}}{18.67f^{*}}\right\}^{2} & Re' > 1000 \end{cases}$$
(16)

318
$$f^* = (1 - \varepsilon_L)^3; \ \mu_e = \frac{\mu_G}{(1 - \varepsilon_L)^{2.5}}; \ \lambda_{RT} = \left(\frac{\sigma}{g\Delta\rho_{GL}}\right)^{0.5}$$
(17)

319 where K_{GL} is the interfacial exchange coefficient, f and f^* are the drag functions, τ_L is 320 the particulate relaxation time, C_D is the drag coefficient, Re' is the relative Reynolds 321 number, μ_e is the effective viscosity, λ_{RT} is the Rayleigh-Taylor instability wavelength, 322 g is the gravity, and $\Delta \rho_{GL}$ is the absolute value of the density difference between liquid 323 and gas phases.

324 **2.4 Mass transfer rate model**

The two-film theory has by far been the most popular and useful theory for dealing with the CO₂ mass transfer among the phases (Gbadago et al., 2020; Kim et al., 2016, 2017; Pham et al., 2015). Generally, based on the two-film model and the Henry law, the overall mass transfer coefficient and enhancement factor are applied to model the CO₂ mass transfer between the gas phase and MEA solution, which can be expressed as follows:

331
$$N_{CO_2} = K_L A_e (C_{CO_2}^* - C_{CO_2})$$
(18)

332
$$\frac{1}{K_L} = \frac{RT}{H_{CO_2 - MEA}k_G} + \frac{1}{Ek_L}$$
(19)

where N_{CO_2} is the CO₂ mass transferred through the gas-liquid interface, K_L is the 333 334 overall mass transfer coefficient, k_G and k_L are the mass transfer coefficients in the gas and liquid phases, respectively, R is the gas constant, T is the temperature, H_{CO_2-MEA} is 335 336 the Henry constant, E is the enhancement factor, which is defined as the ratio of the 337 absorption rate with and without chemical reaction (Sebastia-Saez et al., 2015), and 338 $C^*_{CO_2}$ and C_{CO_2} are the CO₂ concentrations on the surface of the liquid and in the liquid 339 bulk flow, respectively. In particular, H_{CO_2-MEA} , E and $C^*_{CO_2}$ have been adequately 340 illustrated in many works (Borhani et al., 2018; Lu et al., 2019; Pham et al., 2015), so 341 that these factors are expressed in Table 1 along with some other parameters.

342 Since the CO₂ mass transfer resistance between the gas and liquid is dominated by the
343 liquid side, the CO₂ mass transfer resistance in the gas side is neglected. Therefore, Eq.
344 (19) may be simplified as the following equation:

 $\frac{1}{K_L} = \frac{1}{Ek_L} \tag{20}$

According to the film theory expression, the mass transfer coefficient is given asfollows:

$$k_L = \frac{D_{L,CO_2}}{\delta} \tag{21}$$

349 where D_{L,CO_2} is the diffusivity of the CO₂ in the liquid phase, which could be calculated

by the N₂O analogy method (Liu et al., 2006) and given in Table 1, and δ is the diffusion

- 351 layer thickness for the mass transfer.
- Guo et al. (1997) and Munjal et al. (1989) proposed the correlations for predicting the
- 353 film thickness for the RPBs, which are shown in Eq. (22) and (23), respectively,

$$\delta = 4.20 \times 10^8 \frac{Q_L}{2\pi r h} \frac{\nu_L}{\omega^2 r}$$
(22)

$$\delta = \left(3\left(\frac{Q_L}{2\pi r}\right)\frac{\nu_L}{r\omega^2}\right)^{\frac{1}{3}}$$
(23)

356 where *r* is the radial distance, v_L is the kinematic viscosity, and ω is the rotational 357 speed.

For the operating condition in the simulations, the enhancement factor (*E*) is approximately equal to the Hatta number (*Ha*) when the reaction is in the fast reaction regime ($5 \le Ha \ll E_i$) due to the high MEA concentration or the small CO₂ partial pressure given by Eq. (T-7) (Jiru and Eimer, 2013; Ying and Eimer, 2013). As a result,

362
$$K_L = Ek_L \approx Hak_L = \sqrt{\frac{k_2 D_{L,CO_2} C_{L,MEA}}{(k_L)^2}} k_L = \sqrt{k_2 D_{L,CO_2} C_{L,MEA}}, \text{ which means that the}$$

363 overall mass transfer rate is almost independent of the film layer thickness. In addition, 364 this conclusion is consistent with that drawn by Lu et al. (2019). Therefore, either of 365 these equations could be utilized in this paper and Eq. (22) is selected in this paper since 366 this equation is derived from the wire mesh (Guo et al., 1997), which is the experimental 367 packing material used in the reference paper (Kolawole, 2019).

368 **2.5 Chemical reaction rate**

The CO₂-MEA absorption system can be expressed by the zwitterion mechanism (Danckwerts, 1979), and the overall reaction between CO₂ and MEA is described as follows (Borhani et al., 2018):

372

$$2MEA + CO_2 \rightleftharpoons MEACOO^- + MEAH^+$$
 (24)

The above reaction could be separated into two-step reactions. According to the zwitterion mechanism, an intermediate product - the zwitterion ion (MEAH⁺COO⁻) is

375 generated in the CO₂-MEA absorption process (Ebadi Amooghin et al., 2018). Then,

 Table 1 Some of the correlations utilized in the current study.

Equation name	Equations	Eq. No.	Reference	
Diffusivity of the MEA	$D_{L,MEA} = \exp\left(-13.275 - \frac{2198.3}{T} - 0.078142C_{L,MEA}\right)$	T-1	(Snijder et al., 1993)	
	$D_{L,CO_2} = D_{L,N_2O} \frac{D_{w,CO_2}}{D_{w,N_2O}}$	T-2		
Diffusivity of CO ₂ in the MEA	$D_{w,CO_2} = 2.35 \times 10^{-6} \exp\left(-\frac{2119}{T}\right)$	T-3	(Liu et al.,	
solution	$D_{w,N_2O} = 5.07 \times 10^{-6} \exp\left(-\frac{2371}{T}\right)$	T-4	2006)	
	$D_{L,N_2O} = 5.07 + 0.865C_{MEA} + 0.278C_{MEA}^2 \exp\left(-\frac{2371 - 93.4C_{MEA}}{T}\right)$	T-5		
	$E = 1 + ((E_i - 1)^{-1.35} + (E_1 - 1)^{-1.35})^{-\frac{1}{1.35}}$	T-6	(Wellek et al., 1978)	
Enhancement	$E_i = 1 + \frac{D_{L,MEA}C_{L,MEA}}{2D_{L,CO_2}C_{L,CO_2}}$	T-7		
factor	$E_1 = \frac{Ha}{\tanh(Ha)}$	T-8	(Sebastia- Saez et al.,	
	$Ha = \sqrt{\frac{k_2 D_{L,CO_2} C_{L,MEA}}{(k_L)^2}}$	T-9	2015)	
CO ₂ saturation concentration	$C_{L,CO_2}^* = M_{CO_2} \frac{P_{CO_2}}{H_{CO_2 - L}}$	T-10	(Penttilä et al., 2011)	
Henry constant of the CO ₂	$H_{CO_2-L} = H_{N_2O-L} \left(\frac{H_{CO_2-H_2O}}{H_{N_2O-H_2O}}\right)$	T-11	(Penttilä et al., 2011)	

	$H_{CO_2 - H_2O} = \exp\left(145.369 - \frac{8172.355}{T} - 19.303\right)$	T-12	
	$H_{N_2O-H_2O} = \exp\left(158.245 - \frac{9048.596}{T} - 20.86\ln T - 0.00252\right)$	T-13	
	$H_{N_2O-MEA} = \exp\left(-9172.5 + \frac{39.598}{T}\right)$	T-14	
	$H_{N_2O-L} = H_{N_2O-H_2O}C_{L,H_2O} + H_{N_2O-MEA}C_{L,MEA} + 3524641.533(C_{L,H_2O}C_{L,MEA})^2 \left(1 - \frac{T}{324.718}\right) \exp(-13.219C_{L,MEA})$	T-15	
Viscosity of the MEA solution	$ \begin{aligned} \mu_{0.3} &= 0.3083 - 0.00262T + 7.4882 \times 10^{-6}T^2 - 7.17293 \times 10^9 T^3 \\ \mu_{0.9} &= 4.37711 - 0.03776T + 1.08945 \times 10^{-4}T^2 - 1.05031 \times 10^7 T^3 \end{aligned} $	T-16	(Amundsen et al., 2009)
	$\sigma_{H_20} = 0.18548(1 - \frac{T}{647.13})^{(2.717 + 3.554(\frac{T}{647.13}) + 2.047(\frac{T}{647.13})^2)}$	T-17	
Surface tension of the MEA solution	$\sigma_{MEA} = 0.09945(1 - \frac{T}{614.45})^{(1.067)}$	T-18	(Vázquez et al., 1997)
	$\sigma_L = \sigma_{MEA} + \frac{2.129(1 - y_{MEA})}{(1 + (2.129 - (1 - y_{MEA})))(\sigma_{H_2O} - \sigma_{MEA})}$	T-19	

this zwitterion transfers into carbamate (MEACOO⁻) via deprotonation by a base (MEA)
(Moftakhari Sharifzadeh et al., 2016), and the two-step reactions are as follows:

 $MEA + CO_2 \rightleftharpoons MEAH^+COO^-$ (25)

$$381 \qquad MEAH^+COO^- + MEA \rightarrow MEACOO^- + MEAH^+ \qquad (26)$$

Reaction (25) has a finite reaction rate and could be performed in Ansys Fluent through source term UDFs. Reaction (25) (i.e. carbamate formation) is the controlling step and presents the whole reaction kinetics (Sebastia-Saez et al., 2014; Sebastia-Saez et al., 2015). It can be regarded as a second-order irreversible reaction, and the reaction rate is expressed as follows (Luo et al., 2012):

387 $r_{CO_2} = -k_2[MEA][CO_2]$ (27)

388 where k_2 is the reaction rate constant. Different k_2 equations have been proposed for the CO₂-389 MEA absorption system based on various experimental conditions (Hikita et al., 1977; Luo et 390 al., 2015; Versteeg and van Swaaij, 1988; Ying and Eimer, 2013).

- Among them, the reaction rate reported by Versteeg et al. (1996) has been validated by Ying
 and Eimer (2013) and used in many studies (Dashliborun et al., 2019; Ying and Eimer, 2013),
 which is given as follows:
- 394

398

$$k_2 = 4.4 \times 10^{11} \exp\left(-\frac{5400}{T}\right) \tag{28}$$

395 2.6 Heat transfer rate

The amount of the transferred heat between the gas and liquid phases is a function of the temperature difference and the interfacial area:

 $Q_{h,GL} = h_{GL}A_{GL}(T_L - T_G)$ ⁽²⁹⁾

399 where $Q_{h,GL}$ is the transferred heat between the two phases, and h_{GL} is the heat transfer rate. 400 A commonly used model- the Hughmark model (Hughmark, 1967) was utilized to calculate 401 the heat transfer rate, namely

$$h_{GL} = \frac{\kappa_G N u_L}{d_p} \tag{30}$$

403
$$Nu_{L} = \begin{cases} 2.0 + 0.6Re_{L}^{\frac{1}{2}}Pr_{G}^{\frac{1}{3}} & 0 \le Re_{L} < 776.06, 0 \le Pr_{G} < 250\\ 2.0 + 0.27Re_{L}^{0.62}Pr_{G}^{\frac{1}{3}} & 776.06 \le Re_{L}, 0 \le Pr_{G} < 250 \end{cases}$$
(31)

404
$$Pr_G = \frac{C_{pG}\mu_G}{\kappa_G}$$
(32)

where κ_G is the thermal conductivity of the gas phase, Nu_L is the Nusselt number of the 405 liquid phase, Pr_G is the Prandtl number of the gas phase, and C_{pG} is the specific heat. The 406 heat transfer in the RPB is determined by the heat transfer rate, which is reversely proportional 407 to the liquid droplet diameter according to Eq. (30). The larger the liquid droplet diameter, the 408 409 lower is the heat transfer rate, which means that more heat could be retained in the liquid phase. 410 As a result, the liquid temperature changes and the CO₂ capture performance is changed. 411 Therefore, the liquid droplet diameter is a critical parameter for the CO₂ absorption, and this 412 parameter in the packing region and outer cavity zone should be carefully modelled. In 413 particular, the diameter of the liquid droplets in the outer cavity zone (Sang et al., 2017b) are 414 shown in Eq. (8) and the diameter in the packing region has been regressed as two 415 correlations based on the same set of the experimental data (Zhang, 1996), namely (Guo et 416 al., 1997; Yi et al., 2009):

417
$$d_{p,P} = 0.7284 \left(\frac{\sigma}{\omega^2 r \rho}\right)^{0.5}$$
(33)

418
$$d_{p,P} = 12.84 \left(\frac{\sigma}{\omega^2 r \rho}\right)^{0.630} u^{0.201}$$
(34)

419 where r is the radial coordinate of the packing from the centre, and u is the liquid flow rate 420 per unit area, which is only included in Eq. (34). Since the liquid diameter is influenced by 421 the liquid flow rate u, Eq. (34) could more accurately predict the diameter of the liquid 422 droplets in the packing region. Thus, Eq. (34) has been utilized in this work.

423 **2.7 CFD model setup**

Figure 3(a) shows the outside and inside mesh layouts of the 3D pilot-scale RPB model generated by the ANSYS Mesh. In addition to the packing region, the flow characteristics change evidently in the middle of the outer cavity zone and the inner cavity zone, the meshes 427 in these areas are finer compared with the top and bottom cavity zones. To reduce the number of cells and improve the calculation accuracy (Zhang et al., 2022b), the structured hexahedral 428 grids were generated in the whole RPB except in the inner cavity zone. Various numbers of the 429 430 grid cell were tested with total cells of 0.16 M, 0.26 M, 0.42 M and 0.76 M in order to obtain 431 a mesh independent solution, which is shown in Figure 3(b). According to Figure 3(b), until a grid of 0.42 M cells, the CO₂ capture rate and liquid outlet temperature reached stable. Finally, 432 433 the grid of 0.42 M cells, including 0.05 M cells in the packing zone and 0.28 M cells in the outer cavity zone, was applied. The minimum and largest mesh volumes are 1.4×10^{-10} and 434 3.3×10^{-7} m⁻³, respectively, and the average skewness and element quality are 0.11 and 0.85, 435 respectively. It should be noted that the mesh number used in this study is adequate only for 436 437 this case study. A solid mesh study may be required for other applications.



Figure 3. (a) Schematic of the mesh in the 3D model and (b) predicted CO₂ capture rate and
liquid outlet temperature using different grids.

The 3D transient simulations have been performed using the ANSYS Fluent 2021 R1 in a 441 double precision mode based on the High Performance Computing cluster in the University of 442 Sheffield. Various UDFs have been developed for defining the properties of the MEA solution, 443 implementing the extra forces in the momentum equations, and calculating the transferred CO₂ 444 in the mass equations, etc. The Phase Coupled SIMPLE method was applied, and the pressure 445 equations were discretized by the second-order scheme. The pressure-based method and the 446 absolute velocity formulation have been utilized. The time step was set as 5×10^{-4} s, and a 447 maximum of 35 iterations was employed per time step, and the convergence tolerance was 448

449 1×10^{-5} . The simulation case can be assumed to be the pseudo steady state when the CO₂ 450 concentration at the gas outlet and the liquid outlet temperature were reduced to within 1% in 451 ten seconds, and the governing equations' residuals were less than 5×10^{-5} .

For the RPB model, the realizable k- ε turbulence model is more suitable for implementation in the RPB than the standard k- ε model, and the reasons are illustrated in our previous paper (Zhang et al., 2022a). In addition, this turbulence model has been frequently used for the fluid flow simulations in RPBs (Wang et al., 2020; Wu et al., 2018; Yang et al., 2010), thus, this model has been employed in this study.

The properties of the gas and liquid phases are shown in (Xie, 2019). The gas phase consists of CO₂ and air, and the liquid phase is composed of CO₂, MEA (HOC₂H₄NH₂), MEAH+ (HOC₂H₄NH⁺₃), MEACOO⁻ (HOC₂H₄NHCOO⁻) and water. The correlations of the surface tension and viscosity for the MEA solution are present in Table 1 (Amundsen et al., 2009; Vázquez et al., 1997). In addition, the liquid density, thermal conductivity and specific heat of the two phases are evaluated by the mixing law based on every species in each phase. The operating conditions are shown in Table 2.

464

Table 2. The operational conditions of the experiments (Kolawole, 2019).

MEA concentration	Rotational speed	Gas flow	L/C maga ratio	
(wt%)	(rpm)	rate (kg/h)	L/O mass fatio	
30, 50, 70	600, 850, 1150	42	1.3-3.5	

465 In the experiment conducted by Kolawole (2019), all MEA solutions were preloaded to 0.1 mol CO₂/mol MEA before the MEA solution was fed into the RPB. The same MEA loading is 466 467 employed in the CFD model. The gas flow rate is fixed at 42 m³/h with a temperature of 40 °C according to the experimental settings and the gas outlet boundary is set as the pressure-outlet 468 469 with a zero gauge pressure. In terms of the reverse flow at the gas outlet boundary, the CO₂ 470 fraction on the inner packing surface is used as the CO₂ fraction in the gas outlet, which has 471 been performed iteratively. In addition, different L/G mass ratios, e.g. L/G = 2.7, 3.1, 3.5 for 472 30% MEA, are tested in the model. Accordingly, the liquid inlet velocity ranges from 0.20 to 473 0.54 m/s with a temperature fixed at 40 °C and the liquid disappears after entering the liquid elimination zone. The gravity force is set as 9.8 m/s² and in the -Y direction. The rotating speeds
employed in the experiment, including 600, 850 and 1150 rpm are used in the simulations. The
sliding model has been performed to achieve the motion of the packing. The wall boundaries
have been set as no slip and adiabatic.

478 **3. Results and discussion**

In order to validate the CFD model developed, 27 simulation cases have been compared with various experimental cases, and all the major variables have been critically analyzed. After that, the effects of various operational parameters on the mass transfer process have been investigated. In addition, the ratio of the CO₂ captured in the outer cavity zone to that in the full-scale RPB has been investigated.

484 **3.1 Validation and the distribution of variables**

For validating the developed 3D RPB model, the obtained CO₂ mole fraction in the gas outlet and the liquid temperature at the liquid outlet, which could reflect the mass and heat transfer performance, are compared with the experimental results in the Kolawole thesis (Kolawole, 2019). The full operating conditions are listed in Table 2. Although the flow dynamics may be compared via the liquid holdup, unfortunately, this parameter has not been measured in the experimental investigation.

491 Figure 4 presented a comparison of the CO₂ fractions in the gas outlet and the liquid outlet 492 temperatures obtained from the experiments and simulations under various operating 493 conditions, including the MEA concentration, rotational speed and the L/G mass ratio. From 494 Figure 4(a), it is observed that all relative deviations of the CO_2 mole fraction are within 15%. And the CO₂ fraction for the 50% MEA solution has a better agreement with the experimental 495 496 data, whose relative deviations are less than 6%. According to Figure 4(b), most of the relative 497 deviations of the liquid outlet temperature are less than 10%. Only a few points for the 70% 498 MEA solution with a lower rotational speed are larger than 10%. The difference between the 499 simulations and experimental results may mainly result from the difficulties and uncertainties in measuring the variables since the experiments have been conducted in a relatively small 500

501 reactor (Lu et al., 2019). In addition, the employed empirical correlations may not be suitable 502 for this RPB model since these correlations were derived from different RPB sizes, operating conditions, etc. Also, the phase properties may not be able to be accurately evaluated as the 503 504 temperature and CO₂ loading increase. Furthermore, the employed mesh grid with only 0.42 M cells could be a possible reason for the errors between the simulations and experiment 505 506 results. It can be found that the average absolute relative deviations in the CO₂ fraction and liquid outlet temperature are smaller than 7% and 6%, respectively, thus indicating that the 507 508 effective interfacial area, mass and heat transfer models could accurately describe the 509 thermodynamics and mass transfer processes within the RPB. Furthermore, the same sub-510 models, such as the hydrodynamics, including the interfacial, drag, capillary pressure and mechanical dispersion force models, have been verified in our previous paper (Zhang et al., 511 512 2022a) using a smaller RPB model. Therefore, it can be concluded that the developed 3D RPB 513 model gives a good match to the experimental results, and it can be employed to investigate 514 the CO₂ capture process within this RPB. It should be noted, although the overall CFD results agree with the experimental data for this RPB, further validation may be required when 515 applying the sub-models to other RPBs, in particular when different packings are employed. 516



518 Figure 4. Comparison of the (a) CO₂ mole fraction in the gas outlet; (b) liquid outlet
519 temperature (Kolawole, 2019).

517



521 hydrodynamics, thermodynamics, and reaction-enhancement mass transfer within the RPB. 522 Figure 5 presents the predicted variables for the 30% and 70% MEA solutions. Also, the cases 523 are performed with the rotational speed being 850 rpm and the L/G ratios being 3.1 and 1.5 for 524 30% and 70% MEA solutions, respectively. Since the vectors of the gas and liquid velocities 525 are similar for the 30% and 70% MEA solutions, only the vectors for the 30% MEA solution 526 in half of the RPB cut from the central plane is presented in Figures 5(a) and 5(b) in order to 527 observe them clearly. Figure 5(a) shows that the liquid flows tangentially followed the 528 rotational bed due to the porous resistance and its velocity gradually increases on the way outwards through the packing. Near the outer packing surface, the liquid velocity reaches a 529 530 maximum and then the liquid velocity suddenly reduces in the outer cavity zone due to the 531 interfacial force between the gas and liquid phases. In addition, the liquid also rotates with a 532 small velocity in the bottom outer cavity zone. From Figure 5(b), it is observed that the gas 533 phase that radially flows from the gas inlet, begins to rotate after entering the outer cavity zone 534 due to the effect of the rotational packing. Its motion is in rapid synchronization with the 535 rotating packing in the outer packing zone and it reacts with the MEA solution when passing 536 through the packing region. It is worth mentioning that although the vectors of the phase 537 velocity in Figures 5(a) and 5(b) are very similar, the radial flow directions of the gas and liquid 538 phases are opposite.

539 According to Figure 5(c), it is observed that the liquid holdup gradually reduces along the radial 540 direction in the packing zone and this is due to the increasing flow space and the increasing 541 liquid radial velocity due to the centrifugal force, which also can be seen in Figure 6(a). 542 Although the liquid holdup data is not available in the Kolawole's thesis, this phenomenon has 543 been accurately predicted by other simulation and experimental investigators (Burns et al., 544 2000; Lu et al., 2018). After flowing out from the packing region, the liquid radial velocity 545 suddenly increases without the restriction of the porous resistance, thus causing the liquid 546 fraction to become quite small in the outer cavity zone. Subsequently, the liquid droplets collide 547 on the casing wall and flow downwards under the influence of the gravitational force. Thus, 548 the liquid with a higher volume fraction may be observed on the casing wall surface. Also, it 549 can be found that the liquid holdup for the 30% MEA solution is clearly larger than that for the 550 70% MEA solutions as the ratio of the liquid flow rate for 30% to 70% MEA is 2.1. However, 551 the ratio of the liquid holdup in the packing regions is about 1.6, which is smaller than the ratio 552 of the liquid flow rate of 2.1. The first reason is that the liquid viscosity for the 70% MEA is 553 larger, which means that the 70% MEA solution is more likely to adhere to the wire mesh 554 surface (Xie et al., 2019). Furthermore, the liquid radial velocity increases and the liquid 555 residence time reduces when more of the liquid phase enters the packing zone, which allows 556 the 30% MEA solution to faster escape into the packing. These factors lead to a higher liquid 557 holdup for the 30% MEA solution and thus the ratio of the liquid holdup is not exactly equal to the ratio of the liquid flow rate. As a result, the effective interfacial area for the 30% MEA 558 559 solution is larger, see Figure 5(c), since more liquid is attached to the wire mesh packing surface, 560 or it is split into numerous small droplets in the packing zone (Liu et al., 2019; Yang et al., 561 2015a). This means that the contact area between the gas and liquid phases is larger, thus 562 indicating that the 30% MEA solution has more chance to interact with the CO₂.

563 Although the liquid holdup and contact area are larger for the 30% MEA solution resulting 564 from the larger liquid flow rate, the CO₂ fraction in the gas outlet is higher by comparing the 565 two illustrations in Figure 5(e). This is because the CO_2 capture performance is not only related 566 to the liquid holdup and contact area, but also it is affected by some other parameters, such as 567 the MEA concentration, liquid temperature and residence time. Clearly the MEA concentration 568 is the dominant factor in determining the CO₂ capture process. For the CO₂-MEA absorption system, the heat will be released when the CO₂ is captured in the MEA solution, thus leading 569 570 to an increasing liquid temperature. Simultaneously, some heat is taken away by the opposing 571 gas flow via the heat transfer and the rest of the heat remains in the liquid phase. This could be 572 the reason why the liquid temperature for the 70% MEA solution is higher than that for the 30% 573 MEA solution, see Figure 5(f), because more active MEA participates in the reaction and 574 releases more heat to the already smaller flow rate of the liquid phase (Vaewhongs et al., 2020). This phenomenon is more clearer in the top and bottom outer cavity zones, where the liquid 575 576 fraction is quite low and the liquid temperature could be even up to 89 °C. In return, the

577 chemical reaction is faster when the phase temperature is higher. In addition, the phase 578 temperature could also influence the physical properties, including the viscosity and density, 579 which further has an impact on the flow dynamics, such as the liquid holdup and residence 580 time. Thus, this indicates that the liquid holdup, CO₂ capture rate, and liquid temperature could 581 affect each other in this system. The predicted variables can be made accurate and stable only 582 when all the parameters related to the flow dynamics, mass and heat transfer are set properly 583 and these processes reach a balance within the whole RPB.





(c)



(d)

Figure 5. Vectors of the (a) liquid velocity; (b) gas velocity; contours of the (c) liquid holdup;
(d) effective interfacial area on the planes z = 0.01; (e) CO₂ fraction in the gas phase; and (f)
liquid temperature on the planes y=0 for 30% and 70 % MEA.

590 **3.2 Mass transfer in different zones in the RPB**

591 As we know, the CO₂ absorption within a RPB mainly occurs in two zones: packing and outer 592 cavity zones. Based on the newly developed 3D full-scale RPB model, where the CO₂ capture 593 process in the packing and outer cavity zones is separately analyzed by using the CFD method. 594 Taking the 50% MEA solution with a L/G ratio of 2.1 in Section 3.1 as an example, Figures 595 6(a), 6(b) and 6(c) show the circumferentially averaged liquid holdup and effective interfacial 596 area (per unit, m^2/m^3), liquid velocity, as well as the CO₂ mole fraction in the gas phase and the liquid temperature along the radial direction, respectively. For $r_i \leq r \leq r_o$, these 597 598 parameters are only circumferentially averaged in the packing region rather than including the 599 top and bottom cavity zones. When flowing into the inner packing zone, with a radial velocity 600 only, the liquid violently collides with the rotational packing and it is quickly dispersed, thus 601 generating numerous very small tiny droplets, or forming thin films on the wire mesh surface. 602 Due to the small flow space in the inner periphery of the packing, a large effective interfacial 603 area could be observed in Figure 6(a), which is defined as the "end-effect zone" - an area where 604 excellent micromixing occurs due to the strong interaction and liquid dispersion (Cortes Garcia 605 et al., 2017; Esmaeili et al., 2022). Although the phase surface renewal is fast and the effective 606 interfacial area is large in the end-effect zone, a sharp decrease in the CO₂ fraction is not 607 observed in Figure 6(c) due to the small local CO₂ fraction in the gas phase flow and the low 608 liquid temperature.

609 Soon after entering the bulk of the packing zone, the liquid starts to synchronize with the 610 rotating packing, and its tangential velocity almost coincides with the packing rotational 611 velocity in Figure 6(b). Consequently, the liquid dispersion is relatively weak, thus the effective 612 interfacial area decreases along with the radial position in Figure 6(a) (Zhang et al., 2017). 613 Since smaller droplets and thinner film are formed in this region, the liquid radial velocity 614 gradually decreases due to the larger interfacial and drag forces, which is shown in Figure 6(b). 615 In this packing region, most of the CO₂ is captured and a large amount of heat is released, and 616 this is due to the large interfacial surface area and the enhanced surface renewal of the phases

617 (Wang et al., 2021). Therefore, the CO₂ fraction in the gas phase reduces and the liquid
618 temperature increases on their way through the packing as shown in Figure 6(c).

619 After flowing out from the outer edge of the packing, the liquid droplets with large velocity 620 pass through the cavity space, causing a sudden decrease in the effective interfacial area and 621 liquid holdup in Figure 6(a). Without the interaction with the packing, the liquid flow direction is almost unchanged in the outer cavity space, thus causing the liquid radial velocity component 622 623 to increase and the tangential velocity component to decrease as shown in Figure 6(b). As a 624 result, the difference between the gas and liquid radial velocity increases, which leads to an 625 increase in the interfacial force between the two phases and a decrease in the total liquid 626 velocity. In addition, an increase in the liquid temperature can be observed in Figure 6(c) due 627 to the high liquid temperature in the top and bottom regions of the outer cavity zone (see Figure 628 5(f)). Then the liquid droplets collide on the inner casing wall, generating the liquid film and a 629 large number of splashing droplets (Sang et al., 2017b). These phenomena have been modelled 630 well with a steep increase in the liquid holdup and effective interfacial area, see Figure 6(a). 631 Figure 6(c) shows that the magnitude of the decreased CO₂ fraction is small in the outer cavity 632 zone due to the low average effective interfacial area.

In general, the outer cavity zone's volume is much larger than the packing region. For instance, the volume of the outer cavity zone to the packing region is about 11 for the present employed RPB model. However, most of the CO₂ is captured in the packing region instead of the outer cavity zone, which means that the larger the volume of the packing occupied in the fixed fullscale RPB, the more effective is the CO₂ absorption.

Based on the validation cases under various operating conditions shown in Figure 4, it is found that the outer cavity zone takes up 25%~40% of the total effective interfacial area and the effect of the outer cavity zone on the mass transfer could range from 28% to 42% for this case study. From the wide range of the contribution of the mass transfer and effective interfacial area in the outer cavity zone, it is known that the CO₂ capture process in the packing and outer cavity zones are quite different under various operating conditions. With the aim of providing some new, useful and important suggestions for RPB design and scaling up, the contribution of the
 CO₂ removal and effective interfacial area in different zones has been investigated, which is
 presented in the following section.



Figure 6. The circumferentially averaged (a) liquid holdup and effective interfacial area; (b)
liquid and rotating packing velocities and (c) CO₂ mole fraction in the gas phase and liquid
temperature along the radial direction.

651 **3.3 Effect of the operating parameters**

Among the three MEA concentrations investigated, the relative deviation of the CO₂ fraction
in the gas outlet for the 50% MEA solution is small, as discussed in Section 3.1. Thus, the 50%
MEA solution is selected as the baseline case for discussion, and its operating conditions are
listed in Table 3.

656

Table 3. The operational conditions for the baseline case.

MEA concentration	L/G	Rotational	Liquid inlet	Gas flow rate
(%)	ratio	speed (rpm)	temperature (°C)	(kg/s)
50	2.1	850	40	42

For evaluating the CO₂ absorption performance in the full-scale RPB, the overall volumetric mass transfer coefficient ($K_L a_e$), a parameter that may determine the rate at which a gaseous compound (CO₂) can transfer from the gas phase to the MEA solution, is introduced along with the liquid temperature at the liquid outlet. The equation for $K_L a_e$ is shown in Eq. (35). $K_L a_e$ becomes larger when the CO₂ fraction in the gas outlet (y_{CO_2out}) is lower, indicating a better CO₂ capture performance.

663
$$K_L a_e = \frac{Q_G}{\pi (r_o^2 - r_i^2) Z} \ln\left(\frac{y_{CO_2 in}}{y_{CO_2 out}}\right)$$
(35)

The ratios of the captured CO₂ (r_c) and the effective interfacial area (r_e) in the outer cavity zone to those in the whole RPB are examined in order to analyze the mass transfer in different zones. The expression for r_c and r_e are given as follows:

667
$$r_c = \frac{y_{i,CO_2} - y_{op,CO_2}}{y_{i,CO_2} - y_{o,CO_2}}$$
(36)

$$r_e = \frac{\int_O a_e dV}{\int_P a_e dV + \int_O a_e dV}$$
(37)

669 (i) Effect of the MEA concentration

It is known that the CO₂ capture efficiency increases when using the solution with a higher MEA concentration. In terms of the flow dynamics, the liquid phase can lead to early flooding or abnormal distribution in the packing of a CPB due to its large viscosity. However, one of the advantages of applying the RPB is that a higher MEA concentration solution with larger viscosity can fluently pass through the packing due to the higher gravitational environment.

675 Figure 7(a) illustrates the influence of the MEA concentration on the CO₂ absorption and 676 thermodynamics in the RPB in terms of $K_L a_e$ and liquid temperature at the liquid outlet. In this 677 case study, the MEA concentration varies from 30 to 70% while keeping all the other operating 678 conditions consistent with the base case. It can be observed that the magnitude of the increased 679 $K_L a_e$ and liquid outlet temperature significantly becomes larger when using a higher MEA 680 concentration solution since the increased active MEA available in the solution could remove more CO₂ from the gas phase and more reaction heat would be generated at the same time 681 682 (Qing et al., 2011). Although the solution with a higher MEA concentration favors for the CO₂ absorption, the greater trend for corrosion, degradation and foaming should be carefully 683 684 considered before the PCC industrial application.

Figure 7(b) illustrates the ratios of the CO₂ capture and effective interfacial area in the outer cavity zone to the whole RPB under the same operating conditions as Figure 7(a). On one hand, the effective interfacial area mainly depends on the liquid flow rate and rotational speed. On other hand, the increasing temperature in Figure 7(a) affects the phase properties, such as the density and viscosity. As a results, the r_e changes only slightly in Figure 7(b). On considering Figure 7(b), the liquid temperature approximately reaches the maximum after entering the outer cavity zone, which means that the average liquid temperature is relatively high in the outer cavity zone compared with that in the packing zone. The higher liquid temperature benefits the reaction rate. Therefore, the proportion of the CO₂ removed in the outer cavity zone (r_c) increases as the MEA concentration increases. From the above analysis, it is indicated that the higher MEA concentration could enhance the CO₂ capture in the full RPB, and more proportions of CO₂ are captured in the outer cavity zone.



698 **Figure 7.** The effect of the MEA concentration on (a) $K_L a_e$ and liquid outlet temperature and 699 (b) ratio of the contribution of CO₂ capture and effective interfacial area in the outer cavity 700 zone.

701 (ii) Effect of the liquid flow rate

697

For post-combustion capture using the CPB, it is reported that the liquid outlet temperature and
CO₂ capture level increase when the L/G ratio increases due to the larger amount of free amine
molecules and the higher effective interfacial area (Sønderby et al., 2013). However, a reduced
liquid outlet temperature is observed in the RPB from the experimental work (Kolawole, 2019).
Therefore, it is worth exploring the difference in the CO₂ capture processes under various L/G
ratios within the CPB and RPB.

Figure 8(a) presents the predicted impact of the L/G ratio on $K_L a_e$ and liquid outlet temperature

with the L/G ratio ranging from 1.5 to 2.7 while keeping the gas flow rate unchanged. As the

710 L/G ratio increases, more CO₂ is captured resulting from the larger amount of free amine

711 molecules and also the mass transfer resistance is reduced due to the enhanced gas-liquid 712 mixing (Luo et al., 2021) and liquid film refreshing. In addition, the increasing L/G ratio leads 713 to a higher effective interfacial area. These factors increase $K_L a_e$, but this increase tends to 714 slow down as L/G increases. Although increasing the number of tiny droplets that could be generated, the increase in the effective interfacial area is limited at the higher range of the liquid 715 716 flow rate (Wu et al., 2017). Similarly, although more heat is generated as the L/G ratio increases, 717 the heat generation is not high enough to increase the liquid temperature due to the large 718 sensible heat of the liquid phase. In addition, this may explain the decreasing liquid outlet 719 temperature in Figure 8(a), which is consistent with the Kolawole's experimental results 720 (Kolawole, 2019). This is different from that typically observed in a CPB because the liquid 721 flow rate in a CPB is the dominant factor that could significantly increase the interfacial area 722 leading to the significantly increased mass transfer and heat generation. From the above 723 analysis, the increased liquid flow rate, effective interfacial area and the decreased liquid 724 temperature together cause $K_L a_e$ to increase more slowly (Kuntz and Aroonwilas, 2009).

725 Figure 8(b) shows the ratios of the CO₂ capture and effective interfacial area in the outer cavity zone to the whole RPB for L/G ratio from 1.5 to 2.1. The r_c and r_e slightly change with 726 727 changing L/G ratios, but it is noted that a critical point is observed in Figure 8(b). As the L/G 728 ratio increases from 1.5 to 2.1, more liquid concentrates in the outer cavity zone and the temperature in this zone is relatively higher. As a result, both r_c and r_e increase. While the L/G 729 730 ratio continues to increase, the liquid turbulence is dominant in the packing region and the 731 liquid temperature continues to decrease in the outer cavity zone. As a result, the CO₂ that is captured in the packing increases, thus leading to a reduced value of r_c and r_e . 732



Figure 8. The effect of the L/G ratio on (a) $K_L a_e$ and liquid outlet temperature and (b) ratios of the contribution of CO₂ capture and effective interfacial area in the outer cavity zone.

736 (iii) Effect of the rotational speed

733

In our previous investigation (Zhang et al., 2022a), it was reported that the rotational speed has a considerable impact on the liquid holdup and liquid distribution, thus it will further affect the CO₂ capture in the RPB absorber. In addition, the rotational speed could directly affect the effective interfacial area and the reaction time between the liquid and gas phases. Thus, it is a key parameter for the PCC in a RPB.

Figure 9(a) presents the impact of the rotational speed on the $K_L a_e$, the liquid outlet 742 743 temperature of the entire RPB and the liquid temperature at the outer packing boundary. Figure 744 9(b) illustrates the ratios of the CO₂ captured and effective interfacial area in the outer cavity zone to the whole RPB under various rotational speeds. As the rotational speed increases, the 745 746 liquid holdup in the packing region reduces due to the stronger centrifugal force. However, more tiny liquid droplets are formed among the packing region, which results in an improved 747 effective interfacial area (Zhang et al., 2022a), although the magnitude of the improvement is 748 749 very limited. On the other hand, the liquid fraction in the outer cavity zone becomes larger 750 because the liquid accumulation in the outer cavity zone relies on the gravitational force to remove out of the RPB. In addition, more liquid is more likely to attach and cover the casing 751 wall when it flows out from the packing with a higher rotational speed, which leads to a 752 significant increase in the effective interfacial area in the outer cavity zone. Thus, the r_e in 753

Figure 9(b) gradually increases as the rotational speed increases. And the increasing effective interfacial area in both the packing and outer cavity zones makes the $K_L a_e$ increase in Figure 9(a).

757 From Figure 9(a), the liquid temperature flowing out the packing region (blue line) increases 758 with the rotational speed increasing and it is always lower than the liquid outlet temperature (red line). However, the liquid outlet temperature reduces when the rotational speed increases 759 760 from 600 to 850 rpm. The possible reason is that in the outer cavity zone, the liquid fraction 761 increases more significantly compared with the increase in the effective interfacial area when 762 the rotational speed increases from 600 to 850 rpm. And the reaction heat generated in this 763 zone at 850 rpm cannot significantly increase the temperature of the liquid with a large fraction. 764 Therefore, the liquid outlet temperature at 850 rpm slightly reduces from the value at 600 rpm. In addition, r_e increases as the rotational speed increases, thus, r_c is expected to increase 765 766 accordingly. However, r_c reduces in the range of 300 to 850 rpm as observed in Figure 9(b). 767 The possible reason is that the micromixing performance in the packing region improves 768 significantly at the lower rotational speed (Ouyang et al., 2019). As a result, the packing region 769 captures more CO₂ compared with the outer cavity zone when the rotational speed increases 770 from 300 to 850 rpm. Overall, the $K_L a_e$ and r_e increase with the increase in the rotational speed, 771 and r_c decreases first and then increases.



Figure 9. The effect of the rotational speed on the (a) $K_L a_e$ and liquid temperatures at liquid outlet and outer packing boundary and (b) ratios of the contribution of CO₂ capture and effective interfacial area in the outer cavity zone.

776 4. Conclusions

777 This paper has established a full 3D RPB model using the Eulerian porous medium approach, and the physical and chemical processes occurring in the packing and outer cavity zone of a 778 779 pilot scale RPB have been modelled. The overall CFD results agree with the experimental data. 780 Nonetheless, further validation may be required when applying the sub-models to other RPBs, 781 in particular when different packings are employed. The main findings of this paper are as 782 follows: (i) a new completed 3D Eulerian porous medium RPB model was established based 783 on a pilot scale RPB model. By using the Eulerian porous medium modelling method, the CO₂ 784 absorption performance within a whole RPB could be effectively investigated due to the 785 significantly reduced computational cost; (ii) the force, effective interfacial area, heat transfer, 786 mass transfer models were coupled with the Eulerian RPB model, thus, the hydrodynamics, 787 thermodynamics and mass transfer processes could be thoroughly analyzed; (iii) the effective 788 interfacial area, force and even liquid droplet diameter models were introduced separately for 789 the packing region and the outer cavity zone for accurately predicting the CO₂ capture 790 performance inside the full RPB, and the CO₂ absorption processes in these zones were 791 quantitatively evaluated under various operating conditions by the CFD method for the first 792 time.

The predicted simulation results are in good agreement with the experimental data by comparing the CO₂ fraction in the gas outlet and liquid outlet temperature. In addition, the end-effect zone is observed near the inner packing region where a large gas-liquid contact area and strong micromixing occur due to the liquid dispersion and strong interaction between the liquid and packing. However, the amount of the CO₂ transferred between the phases is small in the end-effect zone due to the small local CO₂ fraction in the gas phase and the low liquid temperature. 800 The outer cavity zone has an effect on the CO₂ capture inside the RPB. The contributions of 801 the outer cavity zone to the effective interfacial area and CO₂ capture are respectively in the ranges of 25%~40% and 28%~42% in the RPB investigated in this paper. However, of course, 802 803 these values may be different for different designs. Also, the simulation results show that the 804 CO₂ captured in the outer cavity zone becomes more as the MEA concentration increases, however, the MEA concentrations have little effect on the ratio of the effective interfacial area 805 806 in the outer cavity zone to the whole RPB. The contributions of the outer cavity zone to the 807 effective interfacial area and CO₂ capture change slightly with the L/G ratio increasing. With 808 the increase in the rotational speed, the ratio of the effective interfacial area in the outer cavity 809 zone to the whole RPB increases, however, the ratio of the CO₂ capture in the outer cavity 810 zone to the whole RPB first decreases slightly and then increases significantly.

811 The RPB model developed in this paper can successfully and effectively predict the CO₂ 812 capture process in the whole RPB, and this demonstrates the substantial potential of the model, 813 with further validation, to be used for process optimization and design of the large-scale RPB 814 for industrial PCC. There are two limitations of the CFD model proposed in this study. The first is that using the porous media approach, the details of the fluid flows are not resolved, and 815 816 therefore the characteristics of the formation of liquid droplets/films are unable to be revealed. 817 Secondly, the accuracy of the model is highly dependent on the applicability of the sub-models 818 employed and therefore, careful validation of the model should be considered, especially when 819 a very different packing is employed for the RPB.

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