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# Growth of crystalline $C_{60}$ by evaporation

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This application note describes the growth of crystalline thin films of  $C_{60}$ .

## 1 Introduction

$C_{60}$ , also known as Buckminsterfullerene, is the most common and stable of the fullerenes, comprising a near spherical cage of sixty carbon atoms.  $C_{60}$  has a wide range of applications in molecular electronics and spintronics, solar cells and molecular magnetism. [1–3] The growth of high quality, crystalline  $C_{60}$  films is vital for the production of hybrid meta-materials and devices. In this note, we outline the optimal growth parameters for the evaporation of highly crystalline  $C_{60}$  films from an effusion cell onto Pt substrates.

## 2 Growth

$C_{60}$  can be grown easily on various metallic and semi-conducting substrates.  $C_{60}$  grows poorly on oxides due to high surface tension, causing large clusters to form. [4] Epitaxial platinum with a (111) texture is an ideal metallic substrate for  $C_{60}$  growth. Pt (111) films were grown on  $Al_2O_3$  substrates (see application note for e-beam Pt growth). The substrate temperature was maintained at  $20^\circ C$ .  $C_{60}$  was evaporated from a graphite crucible in a single filament effusion cell. The effusion cell temperature was increased from standby at a rate of  $10^\circ C$  per minute to  $430^\circ C$ . The deposition rate was measured to be  $0.5 \text{ \AA/s}$  using a quartz balance. The rate varies during growth, reaching a maximum value of  $0.74 \text{ \AA/s}$ . The substrate was rotated at  $90^\circ/s$  during deposition. The background pressure was  $8 \times 10^{-10}$  mbar. After growth, the film was capped with a 15 nm thick film of Nb to protect it from oxidation.

## 3 Properties

Structural characterisation was obtained using X-ray reflectivity (XRR) and TEM. Figure 1 shows the XRR data taken of the bilayer film.  $C_{60}$  has a very low density of 1.65 g/cc. However, a structural peak is observable at  $11^\circ$ . Because of the low scattering length of  $C_{60}$ , this peak is only visible in samples with extremely high crystallinity. A GenX reflectivity fit was performed for this data. The low angle fit does not

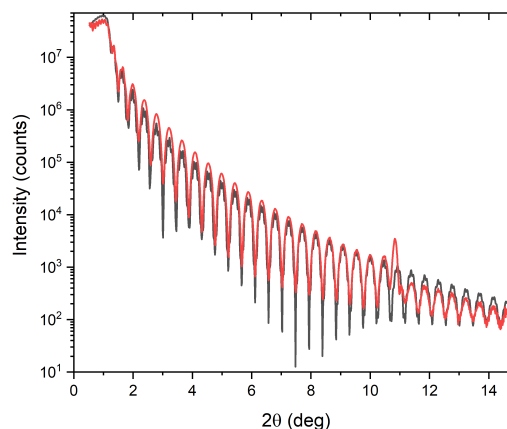


Figure 1: X-ray reflectivity data of a platinum  $C_{60}$  bilayer (red) and GenX fit (black). The structural  $C_{60}$  peak at  $11^\circ$  is not captured by the low angle fit.

capture the structural peak as this method does not simulate crystal structure.

A lamella,  $80 \pm 10 \text{ \AA}$  thick, was cut from the sample using a Dual-Beam Ga Ion FIB. This lamella was then measured using a Titan FEI TEM at 100 kV. The van der Waals lattice of the  $C_{60}$  is clearly visible with (111) vertical orientation. The interface between the Pt and  $C_{60}$  is atomically sharp. The lattice spacing of the  $C_{60}$  film is  $1.05 \pm 0.03$  nm.

The vibrational spectrum of  $C_{60}$  films was recorded

GenX fitting parameter	Value
Thickness - $C_{60}$ ( $\text{\AA}$ )	$975 \pm 0.7$
Density - $C_{60}$ (% of bulk)	$100 \pm 0.01$
RMS Roughness - $C_{60}$ ( $\text{\AA}$ )	$1.8 \pm 0.9$
Thickness - Pt ( $\text{\AA}$ )	$191 \pm 0.2$
Density - Pt (% of bulk)	$100 \pm 0.01$
RMS Roughness - Pt ( $\text{\AA}$ )	$1.7 \pm 0.1$

Table 1: Structural parameters obtained through the fitting of XRR data with GenX

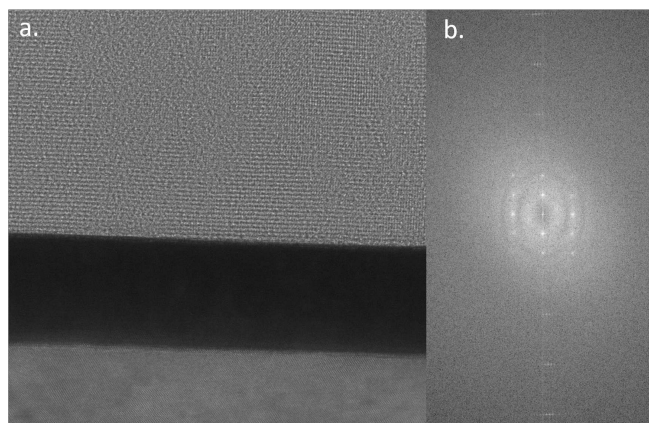


Figure 2: a. TEM obtained from the Pt/C<sub>60</sub> bilayer lamella, showing clear molecular layers and an atomically sharp interface. b. FFT of the C<sub>60</sub> layer, showing 3D crystal structure.

Mode	Raman Shift
Ag(2) - $\Delta\text{cm}^{-1}$	1470 $\pm$ 2
Hg(7) - $\Delta\text{cm}^{-1}$	1434 $\pm$ 5
Hg(8) - $\Delta\text{cm}^{-1}$	1575 $\pm$ 5

Table 2: Positions of the major Raman active modes of C<sub>60</sub>.

using a Horiba Raman Microscope with a 471 nm diode laser, Figure 3. The three highest intensity Raman active modes: Ag(2), Hg(7) and Hg(8) are used to determine the quality of the film. [5]

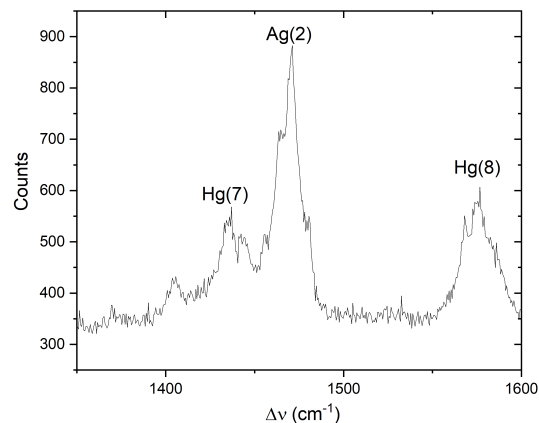


Figure 3: Raman spectrum of C<sub>60</sub> film. The three major Raman active modes are present and the characteristic mode of graphite at 1600  $\text{cm}^{-1}$  is absent. The splitting of the Ag(2) peak is characteristic of electron transfer across a metallic interface.

## References

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