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Investigating the microstructure of soft, microporous matter with synchrotron X-ray Tomography

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14	
15	Abstract
16	Soft porous matter is commonly encountered in artificial tissue applications, pharmaceuticals delivery
17	systems and in cosmetic and food products. These materials are typically opaque and tend to deform under
18	very small levels of shear; this makes the characterization of their microstructure very challenging,
19	particularly in the native state. Air-in-oil systems (oleofoams) are an emerging type of soft material with

21 dispersed in a liquid oil phase. Synchrotron radiation X-ray computed tomography (SR - XCT) is a non-invasive,

promising applications in cosmetics and foods, which contain air bubbles stabilized by Pickering fat crystals

20

22 non-destructive technique increasingly used to investigate multiphasic, porous materials, owing to its high

23 flux which enables sub-micron resolution and significant statistics at rapid acquisition speed. While the

penetration of high energy X-rays can provide high resolution images and allows the reconstruction of the
 3D structure of samples, the experimental setup and measuring parameters need to be carefully designed to
 avoid sample deformation or beam damage.

27 In this work, a robust methodology for investigating the 3D microstructure of soft, porous matter was 28 developed. Sample preparation and experimental setup were chosen to allow synchrotron tomographic 29 analysis of soft oleofoams with a low melting point (<30°C). In particular, the use of cryogenic conditions 30 (plunge-freeze in liquid nitrogen) provided stability against melting during the acquisition. Additionally, an 31 image processing workflow was designed for analysing the 3D microstructure of the samples using ImageJ. 32 Hence, the size and shape distribution of the air phase, as well as the thickness of the continuous gel phase 33 could be determined for samples with significantly different microstructures (fresh vs. heated). Furthermore, 34 the use of time-resolved X-ray radiography (XRR) allowed to study dynamic changes in the microstructure of 35 the samples during thermal destabilization, visualizing bubble coalescence and growth in optically opaque 36 foam samples with a sub-second timescale.

37 Introduction

Soft porous matter features in numerous contexts of scientific and commercial interests, ranging from 38 39 hydrogel-based scaffold for tissue regeneration (Sato et al., 2018; Kinoshita et al., 2020), porous 40 nanocellulose for delivery of pharmaceuticals (Sehaqui et al., 2010; Iftimi et al., 2019) gas marbles (Timounay 41 et al., 2017) to aqueous and non-aqueous foams used in food and personal care products (Fameau & Fujii, 2020; Hill & Eastoe, 2017; Luengo et al., 2021). These systems are comprised of a dispersed gaseous phase, 42 43 a liquid or semi-solid continuous phase and, in the case of foams, suspended stabilizing molecules, particles 44 or crystals that can also adsorb at the air phase boundary (Murray et al., 2011; Murray, 2020). The 3D 45 microstructure dictates the macroscopic properties of these materials, which in turn affects their 46 functionality, their stability against liquid drainage and liquid coalescence and, in the case of food, texture 47 and mouthfeel (Ciurzyńska & Lenart, 2016; Ellis et al., 2017; Herremans et al., 2013) Furthermore, the 48 relationship between raw ingredients, processing conditions and resulting microstructure in porous matter

is complex and still not fully understood. Nevertheless, it is extremely important for the design of novel
products with tuned properties, that new techniques and methodologies which clarify such relationship to
be developed (Lazidis et al., 2017).

52 Investigating the microstructure of soft porous matter is a challenging task, as these materials are usually 53 optically opaque, prone to deformation and often subject to melting at room temperatures (Ubbink et al., 54 2008; Murray et al., 2011). Standard optical microscopy is a readily available and commonplace technique 55 for characterization of materials; however, sample preparation is intrusive, as it requires the sample to be 56 placed between two glass covers, which can affect significantly the native structure of the specimen (Metilli 57 et al., 2020) and consequently, the measured bubble size and shape distribution. Furthermore, most 58 microscopy techniques provide only 2D, surface information on the sample microstructure, limiting the 59 accuracy of the measurement. Even confocal microscopy, which can probe the sample along the z-axis, 60 suffers from a limited field of view and accessible depth range, which hinders the collection of statistically 61 significant 3D data of macroscopic samples.

62 X-ray computed tomography (XCT) has been increasingly used to analyse the microstructure of soft and 63 porous materials, owing to its non-invasive and non-destructive approach (Barigou & Douaire, 2013). Soft 64 biological matter is usually weakly-absorbing with respect to X-rays, and hence a contrast agent may be 65 added during sample preparation to stain the continuous phase. Contrast agents such as sodium iodide (Nal) 66 increase the absorption of X-rays, which improves the contrast of different phases during acquisition. 67 However, the addition of the contrast agent might cause either excessive absorption of X-rays by the sample, 68 or alter its properties prior to imaging. The image contrast can also be enhanced using phase-contrast 69 imaging (PCI), which measures the X-ray refractivity in the sample, rather than the attenuation (Wang et al., 70 2018). Benchtop XCT instruments are becoming increasing accessible; however, they require lengthy 71 acquisition times, which make this technique unsuitable for temperature sensitive (e.g., low melting point) 72 or delicate samples, as prolonged beam exposure can cause heating, X-ray radiation damage and 73 deformation in the microstructure due to sample movement (Wang et al., 2018). Synchrotron radiation XCT 74 (SR-XCT), on the other hand, provides high flux X-ray sources with short exposure times and high signal-to-

noise ratio, reducing significantly the length of the experiment, as well as sub-micron voxel resolution.
Customizable sample environments, such as temperature and humidity control, and application of shear on
the sample allow dynamic microstructural investigation of specimens when subject to external stimuli (Rau
et al., 2017).

79 Bread dough can be considered a model soft porous material of great industrial relevance, in which the 80 microstructure is intimately linked to its rheology profile and ultimately taste perception (Jekle & Becker, 81 2011). Wheat flour dough displays relatively high viscosity, no creaming of air bubbles and good tolerance to 82 X-ray exposure, without requiring carefully designed imaging protocols. Early work reported by Babin and 83 colleagues demonstrated the efficacy of synchrotron XCT in characterizing the microstructure of bread dough 84 during its processing steps (*i.e.* proofing and baking) with a spatial resolution of 15 μ m (Babin et al., 2006). 85 Owing to the ability to track individual bubbles in the sample, the authors were able to model the growth of 86 gas cells developing in the dough, as well as their coalescence during proofing. The spatial resolution of the 87 instruments was a critical parameter in the detection of smaller air bubbles, as reported by more recent 88 publications by Trinh et al. (2013) and Koksel et al. (2016). Here the smaller voxel size (ca. 10 and 8.75 μm, 89 respectively) of the instrument resulted in a ten-fold increase in the bubble density detected in dough 90 samples compared with previous works; this allowed a more accurate description of the gas phase evolution 91 during dough mixing (Trinh et al., 2013) and of the time-dependent bubble disproportionation in non-yeasted 92 samples (Koksel et al., 2016). Sample density calculated from gravimetric methods is routinely used to 93 validate XCT results in these works (Campbell et al., 2001).

In the case of soft porous materials with higher susceptibility to shear and melting, additional steps are required during sample preparation and characterization. Ice cream is a complex multiphasic porous material made of a dispersed gas phase, ice crystals, and a continuous aqueous phase. The complex interplay between its ingredients and phases – which is still not fully understood – is crucial in ensuring high quality products and for the design of novel ice cream formulations (Bahram-Parvar, 2015). Over the last decade, XCT has been increasingly used to characterize ice cream. Pinzer et al. (2012) used a benchtop XCT scanner to investigate the microstructure of ice cream and its evolution upon thermal cycling. The instrument had a 101 nominal 6 µm voxel resolution, and sodium iodide was added to the ice cream recipe to enhance contrast 102 between the air and the ice crystals. To successfully segment the three main phases from the tomography 103 images, the authors developed an edge-preserving smoothing filter, based on the anisotropic diffusion 104 algorithm. The segmentation was then validated using the calculated ice fraction of the sample from 105 differential scanning calorimetry measurements.

106 A similar method was applied more recently on frozen sorbets by Masselot et al. (2021), using a 3D-printed 107 cold stage to keep a low sample temperature during the measurements. The voxel resolution of the scanner 108 used was 9 microns. This study focused mostly on measuring the size distribution of the air bubbles and the 109 ice crystals; CryoSEM was used to validate the XCT measurements and the two techniques showed values of 110 the same order of magnitude and similar range. In both this work and that of Pinzer et al. (2012), however, 111 the scanning time for one sample was between 10 and 15 minutes, which might not be suitable for other 112 types of soft porous materials that are particularly susceptible to X-ray damage. Furthermore, the authors 113 acknowledged that, due to low spatial resolution, air bubbles smaller than 20 - 15 microns were not 114 measured.

115 These limitations were overcome by the use of a synchrotron source, as presented by Guo and co-workers in 116 a series of recent publications (Guo et al., 2017; Guo et al., 2018; Mo et al., 2018). Here the authors focused 117 on the characterization of ice cream microstructure using SR-XCT, with a nominal voxel resolution of 0.8 μm 118 and fast acquisition times (in the order of minutes). The sample was initially maintained at low temperatures 119 (- 15°C). To improve further the segmentation of the different phases in ice cream – air, ice crystals and 120 unfrozen matrix – a novel computational approach was developed to reduce noise and improve intensity 121 homogeneity in the different phases. The effects of thermal cycling on individual air cells and ice crystals and 122 on the unfrozen matrix were investigated using a bespoke temperature-controlled stage and short 123 acquisition times for tomography scans (~ 2.5 minutes each). Nevertheless, there are still destabilization 124 mechanisms in soft matter that occur on shorter time scales than minutes (*i.e.*, seconds or milliseconds), such 125 as droplet or bubble coalescence, Ostwald ripening or disproportionation, or liquid film rupturing, which 126 require suitable time-resolved techniques to be captured and investigated.

127 Despite these recent advances in the use of SR-XCT in soft porous materials, at present many relevant 128 materials are still missing a suitable three-dimensional, close-to-native state methodology for the 129 characterization of their microstructure and its dynamics of destabilisation. Moreover, in light of the current 130 trends of products reformulation with more sustainable and biocompatible ingredients, soft material 131 characterization is of paramount importance (Cornwell, 2018; Manzocco et al., 2021; McClements, 2020). 132 Personal care products such as shampoos, which are surfactant-stabilized aqueous foams, owe their consumer appeal to their foamability and foam stability (Luengo et al., 2021); yet a description of their 133 134 aerated microstructure is currently lacking. Similarly, the fire suppression dynamics of firefighting foams are 135 highly correlated with their bubble size distribution and coarsening, which would benefit from a suitable 3D 136 characterization (Kennedy et al., 2015). Low-viscosity samples prone to movement may benefit from the use 137 of ultra-fast tomography techniques, which enable total acquisition times close to the second (Dittmann et 138 al., 2016).

139 Recently, air-in-oil systems have received significant attention due to their untapped potential in the field of 140 low-fat food products, oil-based cosmetics and pharmaceutical delivery systems (Heymans et al., 2017; 141 Fameau & Binks, 2021). Air-in-oil systems, also called oleofoams, consist of a continuous liquid oil phase in 142 which gas bubbles are stabilized by fat crystals. These materials exhibit a melting range close to body 143 temperature, they deform under small levels of shear and are subject to destabilization mechanisms such as 144 oil drainage and bubble coalescence (Heymans et al., 2018; Saha et al., 2020; Truong et al., 2019). For these 145 reasons the analysis of their microstructure is very challenging. While several studies have focused on the 146 relationship between crystal properties and resulting microstructure, the characterization of these materials 147 is usually carried out using optical or confocal microscopy. Up to the present, the native, three-dimensional 148 arrangement of the air bubbles in oleofoams has not been investigated due to the lack of suitable techniques 149 and methods to obtain meaningful parameters describing the microstructure.

150 In this paper, a novel methodology is proposed, which aims at investigating the microstructure of delicate 151 soft porous matter in a non-invasive fashion (*i.e.,* without the use of contrast agents), with a straightforward 152 cryogenic procedure to prevent deformation in the sample caused by melting, stage rotation or beam

153 damage. The method was demonstrated using cocoa butter-based oleofoams recently characterized in a 154 previous publication (Metilli et al., 2021). The effect of temperature on the microstructure of the specimens 155 was studied using SR-XCT and X-ray Radiography (SR-XRR), to track fast dynamic changes in the sample during 156 heating or cooling. The methodology presented in this work enables the extraction of several microstructure 157 descriptors, including the air volume fraction and its distribution within the sample, the size and shape 158 distribution of the gas phase and the thickness of the continuous phase. Furthermore, SR-XRR was also used 159 to dynamically monitor changes in the air phase heating of samples. While this method was demonstrated 160 with edible oleofoams, it is applicable to the analysis of similar types of sensitive, soft porous materials.

161 Materials and Methods

162 Sample preparation

163 The oleofoams investigated in this paper have been described and characterized in a previous publication 164 (Metilli et al., 2021). Briefly, mixtures of cocoa butter (CB) and high oleic sunflower oil in different weight 165 ratios were crystallized under shear in a lab-scale vessel to obtain an oleogel. For this work the samples were 166 obtained using two crystallization conditions: (1) samples containing 15% w/w CB and crystallized at a -0.10 167 °C/min nominal cooling rate (sample named "15S") and (2) samples containing 30% w/w CB and crystallized 168 at -0.75 °C/min (sample named "30F"). The oleogel was then aerated using a kitchen mixer in cycles of 5 169 minutes whipping and 10 minutes resting, for a total whipping time of 30 minutes. The whipping temperature was monitored during aeration and increased between 7 °C and 20 °C. In order to measure the air 170 171 incorporation, during the rest step the sample was weighed in triplicates using a cup of fixed volume (30 mL). 172 To ensure proper filling of the cup, the sample was added stepwise and set through percussion of the cup. 173 The oleofoam overrun (related to air incorporation) was calculated using Equation 1:

$$Overrun (\%) = \frac{(w_{oleogel} - w_{oleofoam})}{w_{oleofoam}} \times 100$$
 Eq. 1

where w_{oleogel} and w_{oleofoam} are the weight of the un-whipped oleogel and the weight of the oleofoam,
respectively. To calculate the overrun from the tomography data, the following equation was used (Equation
2):

$$Overrun (\%) = \frac{\varphi_{air}}{1 - \varphi_{air}} \times 100$$
 Eq. 2

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where ϕ_{air} is the air volume fraction of the measured Volume of Interest (VOI) of each sample. In this work, both oleofoam samples 15S and 30F were analysed with X-ray Tomography after 5 minutes of aeration and at the end of the aeration step. The overrun measured with the cup method was compared with the overrun calculated from the X-ray Tomography data. To verify any statistically significant difference between the two datasets, a t-test with a *p*-value of 0.05 was performed. The analysis was carried out on two repetitions.

183 Beamline setup

184 The experiments were carried out at beamline I13-2, Diamond Light Source (Didcot, UK), using a pink beam source with a mean energy of 27 keV ($\sigma_E = 5 \ keV$). The 2D projections for tomography and radiography were 185 acquired with a PCO edge 5.5 CMOS camera (2560 x 2160 pixels). The total optical magnification was set to 186 8x, with an effective pixel size of 0.8125 μ m. A small amount (*ca.* 1 mm³) of oleofoam sample was then gently 187 188 mounted on the top of a cut toothpick, minimizing deformation prior to the analysis. The toothpick was then 189 glued to a cryocap, and mounted on the tomography stage. A cryo-jet (Cryojet XL, Oxford Instruments, UK) 190 was installed to allow cooling and heating of samples on the beamline. A schematic of the experimental setup is provided in Figure 1. 191



Figure 1. Schematic of the tomography setup (a), oleofoam samples on cut toothpicks (b) and sample mounted on the rotational
 stage with the Cryojet temperature control.

Two experimental protocols were tested: in the first, samples were mounted on the rotational stage and imaged directly at room temperature without temperature control, with an exposure time of 10 ms and for a total acquisition time of 20 seconds. The second protocol involved flash-freezing samples by immersion in liquid nitrogen (-196 °C) prior to imaging and controlling their temperature during scanning using a cryogenic nitrogen jet (Figure 1c). The Cryojet temperature was set to -40°C, which allowed sample handling without the need of using cryogenic gloves. The exposure time for each X-ray projection was set to 100 ms, for a total acquisition time of 5 minutes. For both experimental protocols, the number of projections was set to 1001,

from an optimization range between 2000 and 500, while the optimal propagation distance was empirically
 determined and found to be 80 mm.

204 Synchrotron X-ray Radiography of heated samples

205 Selected oleofoam samples were subjected to controlled heating using the Cryojet. The evolution of their 206 microstructure was monitored using SR-XRR. The temperature profile was set as follows: equilibration at 293 207 K for 1 min, heating from 293 K to 300 K at 1 K/min, hold at 300 K for 5 minutes, and finally cooling from 300 208 K to 273 K at -6 K/min. The samples were imaged with X-ray Tomography before and after the thermal 209 treatment. Around 1500 2D radiographies of the oleofoam samples were collected during the heating profile, 210 with the aim of tracking dynamic changes in the microstructure due to temperature. The frames were 211 collected every 0.677 seconds. The radiography images were normalized with respect to the camera 212 background (dark field images) and the beam intensity distribution (flat field images), according to Equation 213 3:

$$I_{norm} = \frac{I_{raw} - I_{dark}}{I_{flat} - I_{dark}}$$
Eq. 3

where I_{norm} is the normalized pixel intensity of the image, I_{raw} is the pixel intensity of the sample image (projection) and I_{dark} and I_{flat} are averaged pixel intensities of 20 dark field and 20 flat field images, respectively. Due to the superimposition of bubbles in the 2D projection, a stack of difference images was produced by subtracting the pixel values between the *i*-th and the *i+1*-th frame, in order to visualize changes in the microstructure. The outline of the bubbles in the difference images was detected using the Image
 Processing Toolbox in MATLAB (Mathworks, USA), and the equivalent diameter and circularity of 10 bubbles
 was measured and compared with the 3D data obtained from X-ray Tomography.

To quantify the extent of microstructural changes in the sample during thermal treatment, the difference image stack was further analysed with Principal Component Analysis (PCA) in MATLAB (Mathworks, USA), using the *pca* function. The choice of PCA was justified due its ability to reduce redundancy in large datasets, and to detect changes in image sequences in an unsupervised fashion (Hussain et al., 2013; Celik, 2009), as bubble segmentation from the 2D difference images proved challenging to automate for analysis.

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227 Reconstruction and Image post-processing

The 2D projections were reconstructed into a tomography volume using the Savu processing pipeline (Wadeson & Basham, 2016) using the gridrec algorithm in TomoPy (Gürsoy et al., 2014). The reconstruction pipeline included the following steps: image normalization using dark and flat field images, correction of the ring artefacts (Vo et al., 2018) and a Paganin filter, which is used to restore the phase information generated by the inline phase contrast (Paganin et al., 2002). Finally, the reconstructed volume was obtained using the *Gridrec* reconstruction algorithm in the TomoPy software package (Dowd et al., 1999).

Image post-processing was applied to the reconstructed tomography volumes to obtain quantitative parameters describing the air phase and the continuous phase of oleofoams. To facilitate the computation burden, each tomography volume was divided into a number of Volumes of Interest (VOI) of approximately 500x500x500 µm³. A minimum of 5 randomly selected VOI were analysed and averaged for each sample. The stack of tomographic slices were processed using ImageJ 1.53 (National Institute of Health, USA) according to the following workflow: 3D median filtering, Otsu thresholding, 3D-Euclidean Distance Map Watershed (Legland et al., 2016).

The air cells were counted and measured using the BoneJ plugin for ImageJ (Doube et al., 2010) and excluding the objects on the edges of the VOI. The volume and surface areas were used to compute the equivalent diameter (D_{eq}) and the sphericity (Φ) of each air cell, according to the following equations (Eq. 4 and Eq. 5)

$$D_{eq} = \sqrt[3]{\frac{6V}{\pi}}$$
 Eq. 4

$$\Phi = \frac{\pi^{\frac{1}{3}}(6V)^{\frac{2}{3}}}{A}$$
 Eq. 5

244 D_{eq} was used to calculate the volume-weighted diameter (D[4,3]) of the air cells for a specific sample, using
245 Eq. 6

$$D[4,3] = \frac{\sum_{i=1}^{N} D_{eq} i^4}{\sum_{i=1}^{N} D_{eq} i^3}$$
 Eq. 6

246 Moreover, the BoneJ plugin measured the major, intermediate and minor axis ($a \ge b \ge c$, respectively) of each 247 air bubble, which were used to calculate two aspect ratios: the Elongation Index (EI = b/a) and the Flat Index 248 (FI = c/b). By plotting EI against FI, four shape classes were described: spheroids, oblate, prolate and blade 249 (Blott & Pye, 2008; Zhao & Wang, 2016), and the number of air bubbles belonging to each shape class 250 counted. Finally, BoneJ was also used to calculate the air volume fraction and overrun of each VOI, according 251 to Eq. 2. The thickness of the oleogel phase (*i.e.* the continuous phase) was also calculated using the same plugin, following the method described by Hildebrand & Rüegsegger (1997). The algorithm works by 252 253 inscribing spheres of maximal volume into the continuous phase structure and assigning each voxel the 254 diameter of the largest sphere it belongs to (Pinzer et al., 2012). In this publication, the distribution of the assigned voxels for each sample is represented, alongside with the volume-weighted average thickness value. 255

257 Results and Discussion

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- 258 Comparison of different imaging techniques
- 259 A comparison of different microscopy techniques for investigating the microstructure of oleofoams is
- 260 presented in Figure 2, using a 30F oleofoam as a sample.



Figure 2. Comparison of microscopy techniques for the characterization of the microstructure of oleofoams. Polarized light
 microscopy (a) confocal microscopy (b), CryoSEM (c) and one 2D slice taken from XCT (d). Figure (a) and (c) are adapted from Metilli
 et al. (2021).

Polarized light microscopy (PLM) (Figure 2a) images allowed visualization the air bubbles and the birefringent fat crystals, which were found to stabilize the air-oil interface in oleofoams (Metilli et al., 2021). The fat crystals were clearly visible in the continuous phase and bridging neighbouring bubbles. Non-spherical air cells were detected, with diameters between 10 and 100 µm. However, estimating a size distribution of the air bubbles with this method was challenging, as air bubbles were subject to severe deformation and coalescence during sample preparation. Confocal microscopy images (Figure 2b) provided better-resolved air bubbles, which also appeared non-spherical. Information on the fat crystals was not available, as the 272 fluorescent dye stained both the oil and the fat crystals. Despite the ability of CSLM to provide stacks of images in the z direction, lengthy acquisition times and a smaller field of view (FoV) compared with 273 274 tomography constituted a hindrance for measuring the bubbles' size and morphology, as the microscopy 275 images might not be representative of the whole 3D sample. Micrographs collected with CryoSEM (Figure 2c) 276 show a freshly cut surface from the bulk of the oleofoam, which is not accessible with other microscopy 277 techniques. The higher resolution of this technique and the cryogenic conditions allowed to visualize the 278 porous microstructure closer to its native state. Similarly to CSLM, the field of view of CryoSEM is limited 279 compared with tomography techniques, and the sample preparation is also lengthier, not to mention the 280 introduction of artefacts while imaging soft materials (Groves & Parker, 2013). Figure 2d shows a 2D tomography slice of an oleofoam sample, where the air bubbles appeared darker compared to the 281 282 continuous oil phase. Clusters of brighter pixels were visible and, by comparison with Figure 2a, this 283 suggested that XCT could locate the presence of fat crystals in oleofoams. However, further experiments with 284 higher spatial resolution and improved contrast are required to confirm this hypothesis.

285 Comparison of different imaging protocols for SR-XCT

286 The first imaging protocol – room temperature, 10 ms exposure time, for a 20 seconds total acquisition time - allowed to reconstruct an acceptable tomogram, with most of the air bubbles having a defined boundary 287 288 (Figure S1, Supporting Information 1). However, significant deformation on the edge of the sample occurred 289 during stage rotation, as highlighted with red circles in Figure S1. In some instances, the deformation was 290 excessive and prevented the reconstruction and the analysis of the 3D sample microstructure. Longer exposure times resulted in sample melting, again preventing reconstruction. Considering that CB oleofoams 291 292 display a melting point between 25 and 27°C (Metilli et al., 2021), enhanced stabilization by cooling and 293 controlling the sample temperature was sought to improve tomography acquisition procedure and the 294 quality of the images.

Figure 3 shows a tomography slice of an oleofoam sample (15S) acquired using the second imaging protocol, *i.e.* sample cooling with liquid nitrogen, temperature control with the Cryojet and 100 ms exposure time.



Figure 3. Tomography slice of a 15S oleofoam sample obtained with the second acquisition protocol (sample cooled with liquid nitrogen, maintained at -40°C and using 100 ms exposure time), reconstructed displaying the attenuation contrast (a). Same sample, displaying the phase-contrast mode (b). Zoomed areas showing crystal aggregates are displayed in the top right part of the image.

302 Because of the temperature control during the measurements and the cooling with liquid nitrogen applied 303 to the samples prior the experiment, no deformation was detected in the reconstructed tomogram. To 304 further improve the quality of the images, the phase-contrast mode, which is commonly used for weakly 305 absorbing specimen in XCT, was applied during reconstruction (Nielsen et al., 2016; Guo et al., 2017; Wang 306 et al., 2018). The Paganin-filtered image is shown in Figure 3b, displaying a more homogeneous pixel intensity 307 across the image compared with the attenuation-contrast mode (Figure 3a). The phase imaging mode also 308 resulted in a smoother contour for the brighter crystal aggregates visible in Figure 3a (insert, top right), leading to their pixel intensity being similar to the continuous phase and thus not distinguishable in the 309 310 reconstructed image. Hence, the improved signal-to-noise ratio induced by the Paganin filter was beneficial 311 for the characterization of the porous microstructure of the samples.

312 ImageJ Post-Processing

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On closer inspection of Figure 4a, it can be seen that bright pixels were present inside air cells – especially larger ones – which was mostly caused by ring artefacts in the tomography reconstruction. Hence, a segmentation of the air phase based on the greyscale value alone was not feasible.



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Figure 4. Image post-processing workflow developed in this methodology using ImageJ 1.53. Original reconstructed slice (a),
 binarized using Otsu threshold method (b), segmented with 3D Euclidean distance map watershed (c) and objects counted with the
 "Analyze Particles" function in BoneJ (d).

320 Therefore, an image post-processing workflow was then applied to enhance the quality of the images to 321 separate the air bubbles from the continuous phase. Application of 3D median filter and thresholding 322 resulted in a binarized image with several bubbles appearing connected (Figure 4b). The thresholding method 323 chosen was based on the Otsu algorithm, which returns for each image a threshold value that maximizes 324 inter-class variance; in other words, it divides the pixels into two classes, background and foreground (Russ, 325 2015). The validity of the method was assessed by visually comparing the thresholded images with the 326 starting greyscale analogues. The watershed function based on the Euclidean Distance Map was then 327 successfully applied to separate the connected air cells (Figure 4c). The counted objects in the slice are 328 displayed in Figure 4d, which excluded air cells partially on the edge of the VOI, to avoid underestimation of 329 the bubble size distribution. Recently, the use of machine learning-driven segmentation for image analysis 330 has become popular, and may be considered a promising alternative to more traditional workflows.

However, it does require a set of synthetic images to train the algorithm, which implies prior knowledge of
the sample microstructure (Ali et al., 2021).

333 Estimation of sample density

Tomography data were used to calculate the samples' bulk density and the values estimated were compared with the overrun values measured using the cup method described in the methodology section. Two oleofoam samples, 15S and 30F, were analysed after 5 minutes of aeration (Figure 5a and 5c), and after 30 minutes of aeration (Figure 5b and 5d). Table 1 contains the overrun of four individual VOI and its average, compared with the overrun measured experimentally.



Figure 5. Orthogonal projections of selected VOI, obtained from ImageJ 3D viewer plugin, of samples 15S 5Min (a), 15S 30Min (b), 30F 5Min (c) and 30F 30Min (d). Scale bar represents 250 μm.

	15S 5 Minutes		15S 30 Minutes	
	Sample 1	Sample 2	Sample 1	Sample 2
VOI 1	76.3	87.1	131.3	137.1
VOI 2	78.8	139.0	119.9	146.3
VOI 3	71.0	95.5	109.2	185.7
VOI 4	70.9	115.9	138.0	216.8
OR % XCT	73.7 ± 3.4 ^a	118.7 ± 24.0 ^ª	125.9 ± 15.7 ^ª	171.2 ± 26.0 ^ª

OR % Cup	76.3 ± 2.7ª	76.3 \pm 2.7 ^a 75.0 \pm 11.5 ^b		196.7 ± 10.2ª
	30F 5 Minutes		30F 30 I	Vinutes
	Sample 1	Sample 2	Sample 1	Sample 2
VOI 1	92.0	62.4	139.8	85.2
VOI 2	88.6	72.5	128.1	194.1
VOI 3	89.7	54.7	134.1	163.2
VOI 4	90.8	75.2	127.6	204.9
OR % XCT	89.7 ± 1.9ª	61.2 ± 11.3ª	135.4 ± 6.6ª	125.8 ± 31.9ª
OR % Cup	68.1 ± 19.4 ^b	70.8 ± 6.4^{a}	170.5 ± 17.3 ^b	133.5 ± 16.2ª

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Table 1. Calculated overrun from SR-XCT data for individual VOI from selected oleofoam samples, their average and the respective
 overrun measured with the cup method. Values in the same column labelled with different letters have a statistically significant
 difference (p = 0.05).

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347 The comparison of overrun estimation with the two different methods did not present a clear trend between 348 the two techniques, for both 15S and 30F oleofoams, after 5 minutes of aeration. While for some samples 349 the overrun measured with the cup method and SR-XCT were in agreement, in other repeats SR-XCT 350 overestimated the overrun by 60% or 20% (see Table 1, 15S 5 minutes sample 2 and 30F 5 minutes sample 1). After 30 minutes of aeration, on the other hand, there was a consistent overestimation by the cup method 351 352 of the overrun calculated by SR-XCT, for all samples. In particular, for sample 15S the overestimation was 353 between 15% and 43%, whereas for sample 30F was between 6% and 26%. The differences in estimated density with the two techniques were most likely due to the difference in volume being analysed, 30 mL for 354 the cup method and few cubic millimetres with SR-XCT. The overestimation by the cup method suggested 355 the presence of larger voids in the sample, which might result from filling the cup during weighing, or being 356 357 already present in the specimen, but remaining undetected by SR-XCT. Similarly, the presence of air bubbles 358 of comparable size with the VOI resulted in overrun values with large standard deviation in several samples 359 (Figure S2 in Supporting Information 2).

However, SR-XCT provides information on the distribution of the air bubbles and the homogeneity of overrun
 within samples that are not detectable with the cup measurement. From the observation of the different
 VOIs measured for each sample it is clear that sample 30F displayed, on average, more overrun homogeneity
 (smaller standard deviation) compared with sample 15S (larger standard deviation).

364 Density validation of XCT in bread dough research showed better agreement with gravimetric methods (Trinh 365 et al., 2013; Koksel et al., 2016); however, it should be noted that the sampled XCT volume was in the 366 magnitude of centimetres cube (as opposed to few millimetre cubes of this work) and that the method for 367 measuring dough density involves fluid displacement (Campbell et al., 2001), which is may not be applicable 368 to highly porous materials such as oleofoams. In other research fields, such as artificial bone scaffolds, cake 369 filters and steel pipes corrosion, the density of porous matter measured by XCT is routinely validated by gas 370 pycnometry or mercury injection porosimetry (Jones et al., 2007; Feng et al., 2020;; Wang et al., 2021). 371 However, these techniques are destructive and might not be suitable for soft materials that do not possess 372 a continuous pore network. The authors reported that differences in estimated density between porosimetry 373 and XCT are affected by the resolution of the tomographic scanner, and by the choice of thresholding 374 method. The air bubbles in oleofoams, in particular, usually cover the range between 10 and 100 µm in 375 diameter (Fameau & Saint-Jalmes, 2020; Heymans et al., 2017), hence are likely to be detected by the current SR-XCT setup. 376

377 The effect of thresholding on the resulting overrun for oleofoams samples was also explored. The Huang and 378 Wang (Huang & Wang, 1995) and Renyi's Entropy (Sahoo et al., 1997) methods were tested on sample 15S 379 after 30 minutes of aeration, and compared with the default Otsu method used in this work (Table 1). The 380 image segmentation and calculated overrun were similar between the Huang and Wang, and Otsu methods 381 $(134.7 \pm 7.2 \% \text{ vs. } 152.3 \pm 6.7)$, whereas the Renyi's Entropy produced binary images with excessive void 382 compared with the greyscale 2D image, and very high overrun values (241.5 ± 33.2) (see Supporting 383 Information 3, Table S1 and Figure S3). Therefore, the Otsu method was found to be the most reliable of the 384 ones tested.

While the increase in overrun between 5 and 30 minutes of aeration was observed by both SR-XCT and the gravimetric method (*i.e.* the cup method), the latter technique provides a more reliable bulk density measurement due to the larger volume analysed. Nevertheless, SR-XCT provided essential information about the degree of aeration homogeneity in at the microscale, which directly affects the stability of the product during storage (Heymans et al., 2017; Fameau & Saint-Jalmes, 2020).

390 Effect of heating on oleofoam microstructure quantified by SR-XCT

After the image pre- and post-processing, SR-XCT data were used to estimate quantitative information about the air phase in oleofoams. Two samples with significantly different microstructure were analysed and compared. The first sample, 30F Fresh, was collected from the vessel after 30 minutes of aeration and imaged at -40°C with SR-XCT. The second sample, 30F Heated, was collected from the same batch of sample 30F Fresh, but was subjected to controlled heating, and then imaged with SR-XCT. The microstructure of the two samples is presented in Figure 6, as reconstructed tomographic slices.



397

Figure 6. Tomographic slices of a fresh 30F oleofoam sample (a) and a 30F Heated sample (b), obtained after holding the sample at
 300K (27°C) for 5 minutes.

The fresh sample (Figure 6a) contained mostly non-spherical small bubbles (average diameter < 50 μ m), with few larger ones (diameter *ca.* 100 μ m). The air phase was distributed homogeneously in the continuous oleogel phase, with domains containing both bubbles and a thin layer of oleogel in between. The heated sample (Figure 6b), on the other hand, presented fewer larger and rounder air bubbles, with diameters exceeding 300 μ m, along with a population of smaller bubbles (average diameter < 50 μ m). The oleogel phase comprised either very thin layers separating large bubbles, or areas where bubbles were not present at all.

- 406 The effect of heating on the microstructure of oleofoams was monitored with XRR, where 2D projections of
- 407 the sample from the side were collected during the temperature ramp.
- 408 The evolution of the microstructure of sample 30F Fresh to 30F Heated is shown in Figure 7.



409

<sup>Figure 7. XRR images of the 30F fresh oleofoam before heating (a) and at the end of the temperature ramp (f). Magnifications (b) to
(e) highlight the occurrence of a large air bubble during heating. Frame b) was taken after 7.21 minutes, c) after 8.29 minutes, d) after
8.45 minutes and e) after 12.37 minutes.</sup>

By inspecting Figure 7a, the porous microstructure of the sample was not straightforward to resolve, as 414 415 multiple layers of air bubbles were overlaid in the 2D projections due to the large field of view. However, 416 with increasing temperature (from Figure 7b to 7e) larger bubbles, approximately 300 µm in diameter, 417 appeared, as a result of the heating step. Moreover, by comparison of Figure 7a and 7f, it can be noticed that 418 the whole sample partially collapsed to a rounder structure, owing to the partial melting of the fat network in the continuous phase, as well as the rearrangement of the air bubbles driven by the surface tension forces 419 420 at the air-oil interface. In order to highlight the changes in the microstructure during heating, a stack of difference images was calculated from the stack of radiography frames by subtracting the pixel values of the 421

422 *i*-th frame from the pixel values of the *i*+1-th frame. An example of processed image can be found in Figure 423 S4 (Supporting Information 4), as well as the time-lapse sequence of the difference images (Figure S5, 424 Supporting Information 5). From these images, it can be seen that the air bubbles were subject first to a rapid 425 expansion (between one frame and the following, hence with a speed equal or faster to 0.667 seconds), 426 followed by a slower expansion that lasted over several frames (hence seconds). In particular, the 427 coalescence of two neighbouring air bubbles was captured during the experiment, showing that the newlyformed bubble relaxed "slowly" (i.e. over a few seconds in the following frames) to a slightly more spherical 428 429 shape. Analysis of ten difference images containing the contours of different air bubbles revealed that, during 430 heating, bubbles with an average equivalent diameter of $140.2 \pm 35.2 \,\mu$ m and an average circularity of 0.94 431 ± 0.02 appeared in the oleofoam microstructure. The values are clearly an estimate, as some larger bubbles 432 did not display a complete contour in the difference image to allow precise measurements of their size and 433 shape, and smaller bubbles did not exhibit enough contrast to be detected. To further characterize the effects 434 of heating on the sample, a principal component analysis (PCA) was performed on the stack of the difference 435 images collected during the heating profile. the first principal component (PC) score was plotted against time 436 and the temperature profile during heating (Figure S6, Supporting Information 6).

The PC score, which describes the changes in the pixel distribution in the stack of difference images, remained constant during most of heating ramp. A large variation in the PC score was observed in correspondence with the sample reaching 300 K (27 °C), which corresponded to the melting temperature of sample 30F. This variation reflected the occurrence of large air bubbles while the crystals melted, as the air phase was subject to coalescence without the stabilization of the Pickering crystals. Upon the start of the cooling ramp, the PC score returned to the baseline value with a steeper rate than the heating ramp, as the destabilisation of the foam microstructure was reduced at lower temperatures.

The microstructure characterization of some representative oleofoam samples is presented in Figure 8 - 10, displaying the three-dimensional rendering of a representative VOI each, the size distribution and the sphericity distribution, respectively. Table 2 contains the relevant parameters describing the size and shape of the air bubbles and the oleogel phase for each of the samples depicted in Figure 8.



Figure 8. 3D MATLAB renderings of selected VOIs of samples 15S Fresh (a), 15S Heated (b), 30F Fresh (c) and 30F Heated (d). The air bubbles are colour-coded based on their equivalent diameter, from smallest (blue) to largest (dark yellow).



454 Figure 9. Air bubbles' size distribution for samples 15S Fresh (a), 15S Heated (b), 30F Fresh (c) and 30F Heated (d), calculated with 455 MATLAB.



457 Figure 30. Air bubbles' sphericity distribution for samples 15S Fresh (a), 15S Heated (b), 30F Fresh (c) and 30F Heated (d), calculated
458 with MATLAB.

Both fresh samples (15S and 30F) displayed a similar bell-shaped distribution, with an equivalent diameter of $32.10 \pm 10.3 \mu m$ and $34.2 \pm 14.1 \mu m$, respectively (Figure 9a,c). This value was in agreement, for sample 30F, with the microstructure shown in Figure 6a. At the same time, the sphericity distribution of the two fresh samples were similar, ranging from 0.70 to 0.95, with a mean value of 0.91 ± 0.06 (15S Fresh) and 0.89 ± 0.06 (30F Fresh) (Figure 10a, c). This large distribution reflected the non-spherical shape of the air bubbles,

also visible in the volume rendering of Figure 8a and Figure 8c. For sample 30F Fresh, this was confirmed by
the low volume fraction (44%) of spheroid-shaped bubbles; sample 15S Fresh, on the other hand, contained
a higher amount of spheroid-shaped bubbles (75%) (See Supporting Information 7, Figure S7) (Table2).

467 Upon heating, both 15S and 30F samples displayed a significant change in their microstructure, exhibiting a 468 bimodal size distribution with presence of two populations of bubbles: the first at ca. 27 µm for sample 15S 469 and at *ca*. 32 μ m for sample 30F, and the second with a maximum volume at *ca* 180 μ m and *ca*. 287 μ m, 470 respectively (Figure 9b, d). In fact, the D[4,3] value for heated samples increased to 107.7 \pm 67.1 μ m (15S 471 Heated) and 152.5 \pm 106.3 μ m (30F Heated), as the larger air bubbles contributed more significantly to the 472 distribution. The average diameter for sample 30F was also close with the value obtained from the 473 radiography images (140.2 \pm 35.2 μ m). The average sphericity increased for both samples, as well as the 474 volume fraction of spheroidal bubbles; in particular, for sample 30F, the variation in the shape of the air bubbles was more significant (44% vs. 98%), in agreement with the observed relaxation of the air bubbles 475 476 following the thermal treatment observed with XRR in Figure 7.

Furthermore, the normalized number density of bubbles decreased for both samples after heating: 47.1 ± 4.8 vs. 21.3 ± 3.0 for sample 15S and from 33.9 ± 1.9 to 13.6 ± 4.1 for sample 30F, explained by the occurrence of fewer, but larger bubbles due to coalescence after subjecting the sample to heating. This was visible also in Figure 6 by comparison of fresh and heated microstructure. The change in the microstructure was also reflected in the oleogel phase thickness, which increased from $11.5 \pm 5.2 \mu m$ to $20.0 \pm 9.9 \mu m$ (sample 15S) and from $8.8 \pm 3.6 \mu m$ to $15.1 \pm 5.8 \mu m$ (sample 30F). The change in the continuous gel phase was in agreement, for sample 30F, with the images in Figures 6a and 6b.

Table 2. Summary of the parameters describing the microstructure of sample 30F Fresh and 30F Heated, including the volume weighted mean equivalent diameter (D[4,3]), the sphericity, volume fraction of spheroidal bubbles, number of bubbles per VOI, and
 mean oleogel thickness.

	15S Fresh	15S Heated	30F Fresh	30F Heated
D[4,3] (μm)	32.0 ± 10.3	107.7 ± 67.1	34.2 ± 14.1	152.5 ± 106.3
Sphericity	0.91 ± 0.06	0.93 0.06	0.89 ± 0.06	0.94 ± 0.04

Vol. % of spheroidal bubbles (%)	75.7	89.5	44.0	98.6
Number of bubbles / 10 ⁶ * μm ³	47.1 ± 4.8	21.3 ± 3.0	33.9 ± 1.9	13.6 ± 4.1
Mean Oleogel Thickness (μm)	11.1 ± 5.2	20.0 ± 9.9	8.8 ± 3.6	15.1 ± 5.8

487

488 Conclusions

489 This work demonstrated the use of SR-XCT as a non-invasive technique for fast and accurate quantitative 490 investigation of the microstructure of thermally sensitive, soft porous matter specimens prone to 491 deformation with a straightforward, adaptable sample preparation and beamline setup. Sample stabilization 492 by means of plunge-freezing with liquid nitrogen prior to XCT analysis, followed by the use of a Cryojet 493 temperature control, enabled the collection of high-quality tomography data, suitable for the extraction of 494 quantitative information. Using propagation based phase-contrast mode allows to achieve the image quality 495 needed for segmentation without staining the samples with contrast agents. In particular, the phase-contrast 496 mode was applied to improve the quality of the reconstructed images.

497 The overrun of the samples calculated with XCT was compared with gravimetric measurements, which 498 highlighted the advantage of SR-XCT to study the aeration homogeneity of the specimen at the microscale. 499 A custom image processing workflow was developed to extract relevant descriptors of the porous 500 microstructure, such as bubble size distribution and morphology, together with the thickness of the 501 continuous phase. Furthermore, the use of time-resolved X-ray radiography enabled to track changes in the 502 microstructure of samples subject to external stimuli such as heating. While the method was demonstrated using edible oil-based foams, it is applicable to all porous soft matter that presents similar challenges in its 503 504 characterization.

505 Supporting Information

Supporting information contain additional figures, a time-lapse video of the radiography experiment, as
well as the scripts used in ImageJ and MATLAB for the current work.

508 Declaration of Interest

The authors declare that they have no known competing financial interests or personal relationships thatcould have appeared to influence the work reported in this paper.

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