



## **Correcting for Systematic Underestimation of Topographic Glacier Aerodynamic Roughness Values From Hintereisferner, Austria**

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Chambers JR, Smith MW, Smith T, Sailer R, Quincey DJ, Carrivick JL, Nicholson L, Mertes J, Stiperski I and James MR (2021) Correcting for Systematic Underestimation of Topographic Glacier Aerodynamic Roughness Values From Hintereisferner, Austria. Front. Earth Sci. 9:691195. doi: 10.3389/feart.2021.691195 Spatially-distributed values of glacier aerodynamic roughness ( $z_0$ ) are vital for robust estimates of turbulent energy fluxes and ice and snow melt. Microtopographic data allow rapid estimates of  $z_0$  over discrete plot-scale areas, but are sensitive to data scale and resolution. Here, we use an extensive multi-scale dataset from Hintereisferner, Austria, to develop a correction factor to derive  $z_0$  values from coarse resolution (up to 30 m) topographic data that are more commonly available over larger areas. Resulting  $z_0$  estimates are within an order of magnitude of previously validated, plotscale estimates and aerodynamic values. The method is developed and tested using plotscale microtopography data generated by structure from motion photogrammetry combined with glacier-scale data acquired by a permanent *in-situ* terrestrial laser scanner. Finally, we demonstrate the application of the method to a regional-scale digital elevation model acquired by airborne laser scanning. Our workflow opens up the possibility of including spatio-temporal variations of  $z_0$  within glacier surface energy balance models without the need for extensive additional field data collection.

Keywords: glacier,  $z_0$ , structure from motion, terrestrial laser scanning, aerodynamic roughness

## INTRODUCTION

The aerodynamic roughness length parameter ( $z_0$ ) is recognized as one of the key uncertainties in glacier surface energy balance (SEB) modeling (Cullen et al., 2007; Sicart et al., 2014; Litt et al., 2017; Fitzpatrick et al., 2019). It is defined as the topographically-controlled height above the surface at which horizontal wind velocity reaches zero (Stull, 1988; Garratt, 1992) and is typically derived from the direct observation of turbulent eddies through eddy covariance (Munro, 1989; Sicart et al., 2014; Fitzpatrick et al., 2019), or from extrapolation of log-linear fits of wind speed and air temperature profiles (Denby and Greuell, 2000; Brock et al., 2006; Quincey et al., 2017). Accurately quantifying  $z_0$  is essential for calculating and predicting glacier ablation because  $z_0$  is incorporated into calculations of the turbulent fluxes of sensible and latent heat between a surface and the adjacent atmosphere using the "bulk aerodynamic approach" (Hock, 2005). Using this approach, the sensible (*H*) and latent (*LE*) heat fluxes are defined as:

$$H = \rho_a c_a C_H u \left( T_a - T_s \right) \tag{1}$$

$$LE = \rho_a L_{\nu/s} C_E u \left( q_a - q_s \right) \tag{2}$$

where  $\rho_a$  is the density of air (kg m<sup>-3</sup>),  $c_a$  its specific heat capacity (J kg<sup>-1</sup> K<sup>-1</sup>) and  $L_{\nu/s}$  the latent heat of vaporization or sublimation. The bulk transfer coefficients  $C_H$  and  $C_E$  are derived from the logarithmic wind speed profile equation (assuming neutral stratification)

$$u(z) = \frac{1}{\kappa_0} u_* \ln\left(\frac{z}{z_0}\right) \tag{3}$$

in which wind speed (u, m s-1) is related to the logarithm of measurement height (z, m),  $\kappa_0$  is the dimensionless von Karman's constant, 0.4, and  $z_0$  is the surface roughness parameter. Substituting into **Eq. 1**, the bulk equation for the sensible heat flux becomes

$$H = \rho_a c_a \kappa_0 u_* \frac{T_a - T_s}{\ln\left(z/z_0\right)} \tag{4}$$

The turbulent fluxes commonly comprise ~35-50% of a glacier SEB and have an increasingly important role in cloudy and windy conditions (Giesen et al., 2014), when they can become the dominant source of energy over short timescales (daily and sub-daily), contributing >75% of melt energy (Fausto et al., 2016). In maritime climates their dominance increases (Anderson et al., 2010). Despite the importance of  $z_0$ , it is common for it to be generalized spatially across glacier surfaces and climatic zones, and through time (e.g., Lewis et al., 1998; Giesen et al., 2014; Bravo et al., 2017), at least partly because it is difficult to measure. Obtaining eddy covariance data or aerodynamic profiles from field measurements is challenging and provides only point-based  $z_0$ values; consequently, research has been driven toward estimation of spatially distributed  $z_0$  values from microtopography (Lettau, 1969; Munro, 1989), accelerated by the increasing availability of fine-resolution (sub-meter) and broad-scale topographic data.

Past work has identified that  $z_0$  is spatially and temporally dynamic, with topographic  $z_0$  values that have been validated against values from eddy covariance (EC) data or aerodynamic profiles (Brock et al., 2006; Irvine-Fynn et al., 2014; Miles et al., 2017; Quincey et al., 2017; Chambers et al., 2019; Fitzpatrick et al., 2019; Liu et al., 2020). In particular, the rapid data collection enabled by structure from motion photogrammetry (SfM; Smith et al., 2016b) and terrestrial laser scanning (TLS; Lemmens, 2011; Telling et al., 2017) has led to work focusing on development and validation of topographic methods, mostly concentrated on plotscale data (tens of meters). However, the acknowledged scale and resolution-dependency of topographically-estimated  $z_0$  (Rees and Arnold, 2006) complicates the application of these methods to glacier-scale distributed energy balance models, because coarser resolution data lead to substantial underestimates of  $z_0$  and an order of magnitude change in  $z_0$  can double the calculated turbulent fluxes (Munro, 1989).

Initial efforts to produce glacier-scale maps of  $z_0$  have been promising (Smith et al., 2016a; Fitzpatrick et al., 2019; Smith et al., 2020), but here we seek to further these attempts by producing more robust topographic estimates of  $z_0$  at the glacier scale. We employ a simple workflow to correct for systematic underestimation of  $z_0$  values derived from coarser resolution data and show that the workflow can be used to produce first order estimates of  $z_0$  across the surrounding region. Specifically, we first present a multi-scale analysis of topographic  $z_0$  from data collected during the 2018 ablation season on Hintereisferner, Austria, and identify power law relations between data resolution and derived topographic  $z_0$ . Second, through comparison with wind tower and EC data, we use these power laws to develop a correction factor for  $z_0$  estimates derived from coarse scale, widely available glacier surface topographic data, to bring them within one order of magnitude of likely true values, thus limiting the knock-on effects of over- or underestimation on the turbulent fluxes (c.f. Munro, 1989). Finally, we demonstrate the broader utility of the method by applying it to other nearby glacier surfaces covered by a freely available regional topographic dataset for Austria.

### DATA AND METHODS

#### Location

Field data were collected at Hintereisferner (46° 48′ N, 10° 47′ E) in the Austrian Alps (**Figure 1**), from 1 to 16 August 2018 during the Hintereisferner Experiment (HEFEX; Mott et al., 2020). Hintereisferner (HEF) is located high in the Rofenache catchment in the southern Otztal Alps, in the inner-alpine dry region (Strasser et al., 2018). The Rofenache catchment ranges from 1,891 to 3,772 m a.s.l, with an annual mean temperature of 2.5°C at 1,900 m a.s.l. (Strasser et al., 2018). Snow cover commonly persists from October through until June at elevations above 3,000 m a.s.l. Glaciers in the Otztal Alps, of which there are more than 50 (Abermann et al., 2009; Rastner et al., 2015), have been in retreat throughout the latter half of the 20th century and lost an average of 8.2% of their area between 1997 and 2006 (Abermann et al., 2009).

HEF is ~6 km long and ranges in elevation from ~3,740 m to the current (2018) terminus at 2,498 m. Its present-day area is around 6.22 km<sup>2</sup> but the glacier is receding rapidly, with an estimated reduction in area of 15% from 2001 to 2011, while the terminus retreated by around 390 m in the same decade (Klug et al., 2018). Mass balance records extend back to 1952/53 (Fischer, 2010) and HEF has been used as a type-site for gauging the overall health of Austrian glaciers and those of the wider European Alps, some of which, at current rates of retreat could be almost non-existent within a century (Vincent et al., 2017). Additional reconstructions suggest that HEF has been in near-constant retreat since the Little Ice Age, c.1855 (Greuell, 1992), with an increasing rate of mass loss as its tributary glaciers have become detached over the last two centuries (Fischer, 2010).

HEF has been the subject of a multitude of studies. Mass balance has been recorded using ablation stakes (Blümcke and Hess, 1899; Van De Wal et al., 1992; Kuhn et al., 1999) and numerical modelling (Escher-Vetter, 1985; Greuell, 1992; Schlosser, 1997; Fischer, 2010; Klug et al., 2018; Wijngaard et al., 2019). Remotely sensed data from airborne and satellite platforms have been used to study the glacier surface (Fritzmann



et al., 2011), its reflectance properties (Koelemeijer et al., 1993) and allowed it to be included in valuable regional glacier inventories that document the decline of ice masses in the Alps (Patzelt, 1980; Lambrecht and Kuhn, 2007).

## **Data Collection**

Data collection consisted of two main components: topographic surveys at multiple scales and meteorological data collection. Four plots were selected in the field (**Figure 2**), chosen to be as distinct in surface appearance as possible bearing in mind the safety risk associated with installing instrumentation and manually surveying very steep or heavily crevassed plots. Plot 1 (**Figure 2A**), the furthest down-glacier, was crosscut by supraglacial meltwater channels and was the most modified by

melt processes having been exposed the longest after snowline retreat. Plot 2, shown in **Figure 2B**, appeared smoother than Plot 1, with some low (<0.2 m) flow-parallel ridges, some perpendicular crevasse traces and small moulins. Plot 3 was centered on an area of streamlined, well defined longitudinal ridges with no discernable cross-glacier features (**Figure 2C**). Plot 4 was the smoothest visually, with only minor surface variability and some small meltwater channels (**Figure 2D**).

A wind tower was installed at each plot [see *Aerodynamic*  $z_0$  *Estimation* ( $z_{0WT}$ ,  $z_{0EC}$ )] during the period when topographic surveys were carried out. As fieldwork was completed near the peak of the ablation season, air temperatures were positive for the study duration (**Figure 3A**) with diurnal extremes of 3–4°C at night and 10–18°C during the day. The maximum temperature



was 19.8°C recorded on 5th August and the minimum was  $1.3^{\circ}$ C on 8th August. Mean wind speed for the data collection period was  $2.5 \text{ m s}^{-1}$ , although occasionally fluctuated above  $5 \text{ m s}^{-1}$ 

(Figure 3B). The maximum recorded was  $7.7 \text{ m s}^{-1}$  on 14th August, while the minimum was below the stall speed of the cup anemometers ( $<1 \text{ m s}^{-1}$ ) on multiple occasions. Wind



direction was recorded relative to the down-glacier direction and was predominantly around 165°, indicating the presence of downglacier flow characteristic of density driven katabatic winds (**Figure 3C**), which were also observed by Mott et al. (2020). Convective storms accounted for the majority of precipitation, with notable events in the afternoons of the 1st, 6th, 10th, and 13th August.

As part of HEFEX, four turbulence towers [TT1-4; see Aerodynamic  $z_0$  Estimation  $(z_{0WT}, z_{0EC})$ ] were present throughout the data collection period (Mott et al., 2020). TT4 was located within Plot 1, while TT1-3 were installed further down-glacier. Independent estimates of  $z_0$  from these towers was used for validation of topographic and wind tower  $z_0$ .

#### **Topographic Surveys**

Topographic  $z_0$  ( $z_{0DEM}$ ) was estimated using topographic data obtained via five methods covering a range of scales:

- 1. Small plots (10 m  $\times$  10 m): ground-based SfM using an Olympus OMD EM-10 camera mounted on a survey pole at ~6 m above ground level on 10, 11, and 13 August.  $z_0$  from this method is referred to as  $z_{0\text{Ground}}$
- 2. Large plots (~30 m  $\times$  70 m): ground-based SfM surveys encompassing [1], on 8, 10, 11, 12, and 13 August using the same camera as above but with the survey pole at ~9 m above the surface ( $z_{0Pole}$ )
- 3. Airborne plot surveys: SfM surveys of the same plots/dates as [2] using a DJI Phantom 3 uncrewed aerial vehicle (UAV) with gimbal-stabilized digital camera at ~30 m above the surface  $(z_{0UAV})$
- 4. Glacier-scale: the upper and lower glacier (see **Figure 1**) was surveyed on 3, 7, 12, and 16 August using a RIEGL VZ-6000 TLS situated on the true right of the valley, near the summit of "im Hinteren Eis", a vantage point (**Figure 1A**) from which most of the glacier surface can be seen ( $z_{\text{OTLS}}$ )
- 5. Regional-scale: Airborne Laser Scan (ALS) data, obtained from flights in 2001–2009 (between August and October) using ALTM3000 ALS, data freely available (Open Data Austria, 2020) ( $z_{0ALS}$ ).

## Plot-Scale Structure From Motion Surveys ( $z_{0Ground}$ , $z_{0Pole}$ , $z_{0UAV}$ )

SfM uses the principles of photogrammetry to digitally reconstruct surfaces or objects in 3D (Ullman, 1979; Brown and Lowe, 2005; Snavely et al., 2008) and has been used to obtain estimates of  $z_0$  on ice surfaces (Irvine-Fynn et al., 2014; Smith et al., 2016a; Chambers et al., 2019; Smith et al., 2020). We followed the same principles for each SfM survey, based on workflows and recommendations in James et al. (2017), O'Connor et al. (2017), completing surveys of Plot 1 on 10 and 13 August, Plot 2 on 8 August, Plot 3 on 11 August and Plot 4 on 12 August.

Camera specifications, camera calibration and survey area geometry (Table 1) were used to calculate the footprint of each image, which determined the distance between images required to achieve 60-80% overlap and the number of images needed for the survey area. Images were predominantly nadir and followed a regular grid pattern, with an additional ~10% of images taken obliquely (<20° off nadir). All survey plots were marked out using a regular grid of 9 ( $z_{0Ground}$ ) or 21 ( $z_{0Pole}$  and  $z_{0UAV}$ ) ground control points (GCPs), the locations of which were recorded using a Leica GS10 differential GPS, with a subcentimeter mean accuracy for each plot. Typical 3D root mean square (RMS) control point error was ±0.03 m and RMS reprojection error was 1.66 pixels (2.6  $\mu$ m). Uncertainty in  $z_0$ estimates associated with the SfM process is assumed to be negligible, following previous analysis in Chambers et al. (2019). Where  $z_{0DEM}$  values are given for any plot, the standard deviation of the relevant plot is also given as a proxy for any further uncertainties, and the same is included for  $z_{0TLS}$ values.

Data were processed in Agisoft Photoscan Professional Edition Version 1.4.0 using the following settings: high accuracy, generic preselection enabled for all methods, reference preselection enabled for UAV surveys (as the position of the UAV is recorded for each image), camera calibration parameters F, Cx, Cy, K1-3, and P1 and P2 included, high reconstruction quality and aggressive depth

Specification	Ground	Pole	UAV
Camera make and model	Olympus OM-D E-M10/Canon Powershot SX600 HS	Olympus OM-D E-M10	Sony EXMOR 1/2.3"
Camera Lens model	M. Zuiko Digital 14–42 mm/Built-in	M. Zuiko Digital 14–42 mm	FOV 94° 20 mm (35 mm format equivalent) f/2.8
Sensor size (mm)	17.3 × 13.0/6.2 × 4.6	17.3 × 13.0	6.16 × 4.62
Image size (pixels)	4,608 × 3,456	4,608 × 3,456	4,000 × 3,000
Pixel pitch (µm)	3.74/1.34	3.74	1.54
Height above surface	6 m	~9 m	~30 m
Focal length (mm)	~28/~25	14	4
GSD (mm)	0.8/0.3	2.4	11.6
Max. (mean) images per plot	71 (64)	583 (465)	343 (274)
Camera locations	360° survey (inward facing)	Regular grid/~20% of image total 360° survey	Regular grid/~20% of image total 360° survey
Camera angle	~20° off nadir	Nadir/~20% of image total <20° off nadir	Nadir/~20% of image total <20° off nadir
Camera trigger	Manual	Remote (smartphone app)	Continuous with 2 s interval

**TABLE 2** | List of grid resolutions and sliding neighborhood sizes used in multiscale analysis.

Source	Grid resolutions (m)	Sliding neighborhood sizes (m $\times$ m)
Z <sub>0Ground</sub>	0.005	0.1, 0.5, 1, 5
	0.01	0.1, 0.5, 1, 5
	0.05	0.5, 1, 5
Z <sub>0Pole</sub> /Z <sub>0UAV</sub>	0.01	0.5, 1
	0.05	0.5, 1, 5, 10
	0.1	0.5, 1, 5, 10
	0.5	5, 10
	1	5, 10
ZOTLS	5	5, 10
	10	50, 100, 150
	20	100, 200
	30	150, 300

filtering. Dense point clouds were imported into CloudCompare 2.10 (CloudCompare, 2020), where they were manually inspected/cleaned. Digital elevation models (DEMs) were constructed using linear interpolation with nearest non-empty neighbors, ensuring a regular grid shape with the top of the grid aligned with the direction of flow (roughly South-North). DEMs were produced at each of the grid resolutions shown in **Table 2**, using all of the neighborhood sizes also listed.

#### Glacier and Regional-Scale Surveys (z<sub>0TLS</sub>, z<sub>0ALS</sub>)

A permanent *in-situ* Riegl VZ-6000 terrestrial laser scanner (TLS) was used to survey the majority of the glacier ablation zone using a near-infrared wavelength (1,050 nm) suited to snow and ice surfaces (University of Innsbruck, 2020a). The TLS is housed in a climate-controlled container near the summit of "im Hinteren Eis" (46.79586° N, 10.78277° E, 3,244 m a.s.l.). The point acquisition rate was approx. 23,000 points per second. Horizontal and vertical spatial resolution was ~0.17 m at 1,000 m range, giving a theoretical density of 10 points per m<sup>2</sup> mid-glacier, and 2 points per m<sup>2</sup> at the head of the accumulation zone and near the terminus (University of Innsbruck, 2020b), due

to the beam angle (Carrivick et al., 2015). Validation of scans from this TLS suggests <0.15 m difference to ALS data and <0.1 m between TLS scans (University of Innsbruck, 2020b). Surveys were split into two sections, upper and lower glacier, and carried out on 3rd, 7th, 12th, and 16th August 2018. TLS DEM creation was carried out using CloudCompare 2.10 (as described in *Data Collection*).

Regional (gridded) elevation data were acquired for the entire Austrian Alps at 10 m resolution (Open Data Austria, 2020). This regional product was created by interpolation of 2.5 m airborne laser scanning (ALS) data obtained during flights (details in **Supplementary Table S1**) over the Ötztal Alps, Tyrol, from 2006 to 2012 (Bollmann et al., 2011; Fritzmann et al., 2011; Sailer et al., 2012; Fischer et al., 2015). The reported vertical accuracy on relatively flat and smooth surfaces is  $\pm 0.07$  m ( $\sigma = 0.07$  m), with an absolute standard error on slopes  $<37^{\circ}$  of  $\pm 0.04$  m (Bollmann et al., 2011; Sailer et al., 2011; Sailer et al., 2011; Sailer et al., 2013). ALTM 3,100 and Gemini ALS sensors were used in the Tyrol area, with an average density of 0.25 points m<sup>-1</sup>. Glacier outlines were taken from the Austrian Glacier Inventory 4 (Buckel et al., 2018; Buckel and Otto, 2018).

## Topographic $z_0$ Estimation and Correction Factor Development

 $z_0$  was estimated from topographic datasets using the DEM-based method of Smith et al. (2016b) alongside a sliding neighborhood operation, wherein an operation is applied to each cell of a grid with a specified number of surrounding cells forming the neighborhood. The center of the neighborhood then slides to next cell until the operation has been applied to the entire grid. For this study, the sliding neighborhood function within MATLAB<sup>®</sup> R2017a was used. This approach is similar to that used by Fitzpatrick et al. (2019) but differs in how each parameter is defined. Most DEM methods, including those of Smith et al. (2016b), Fitzpatrick et al. (2019) and that used here, are based on the Lettau (1969) equation

TABLE 3	Index of	correction	factors	for	each	arid	resolution.
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Grid resolution (m)	Neighborhood size (m $\times$ m)	Correction factor ( $log_{10}$ )		
0.005	0.1	0.22		
0.005	0.5	0.22		
0.005	1	0.22		
0.005	5	0.22		
0.01	0.1	0.32		
0.01	0.5	0.32		
0.01	1	0.32		
0.01	5	0.32		
0.05	0.5	0.56		
0.05	1	0.56		
0.05	5	0.56		
0.05	10	0.56		
0.1	0.5	0.66		
0.1	1	0.66		
0.1	5	0.66		
0.1	10	0.66		
0.5	5	0.9		
0.5	10	0.9		
1	5	1.01		
1	10	1.01		
5	25	1.25		
5	50	1.25		
5	100	1.25		
10	50	1.35		
10	100	1.35		
10	150	1.35		
20	100	1.45		
20	200	1.45		
30	150	1.51		
30	300	1.51		

$$z_0 = 0.5h^* \frac{s}{S_A} \tag{5}$$

where for each neighborhood, 0.5 is the average drag coefficient of roughness elements,  $h^*$  is their average vertical extent [here we used twice the standard deviation of elevations over the detrended mean plane, mm, following Munro (1989)], s represents average silhouette area (mm<sup>2</sup>) and SA is the surface area of the neighborhood in the horizontal plane (mm<sup>2</sup>). s was calculated as the sum of the heights (mm) of all cells which visible above their respective preceding cell, as seen from the prevailing downglacier wind direction, multiplied by cell width (mm). Detrending was performed on each neighborhood by removing the best-fit plane. TLS- and ALS-derived DEMs were additionally detrended by subtracting the moving mean calculated at 5 x grid resolution, which removed coarse scale topography but preserved finer scale topographic variability, i.e., perturbations of <1 m (Glenn et al., 2006; Smith, 2014). Resolution dependence was investigated by deriving z<sub>0DEM</sub> from grids with incrementally increased resolutions as listed in Table 2, while scale dependence was investigated by incrementally increasing neighborhood sizes.

Power law behavior displayed at Plot 1 (surveys from Day 10) was used to develop correction factors for each resolution. Correction factors were based on the quotient of modelled  $z_0$  for different grid resolutions and wind tower-derived  $z_0$  ( $z_{0WT}$ ), which was available from more locations across the glacier than  $z_0$  derived from turbulence towers ( $z_{0EC}$ ). Correction factors (CF) were calculated using:

$$CF(Res) = \frac{z_{0WT}}{a + b(Res)}$$
(6)

where *Res* is a given resolution, *a* is the intercept and *b* is the corresponding gradient of a linear model fitted to a plot of  $\log_{10}$  transformed grid resolution and  $z_{0DEM}$ . The model fitted to the Plot 1 Day 10 data had a goodness of fit of  $r^2 = 0.4$ , RMSE of 3 mm and p < 0.01, for 49 data points (47 degrees of freedom). The residuals of the model were normally distributed and displayed homoscedasticity when plot against predicted values. While some scatter remained in corrected  $z_{0DEM}$  values, the predictive performance of the linear model was superior to non-linear models which were also tested, which under-predicted  $z_{0DEM}$  at coarser grid resolutions and over-corrected it as a consequence.

Correction factors were calibrated using the data from Plot 1, then validated with Plots 2, 3, 4, and data from Plot 1 on Day 13, then TLS DEMs of Hintereisferner. Finally, they were used to project  $z_{0DEM}$  values for the ALS DEMs of other local glaciers. An index of correction factor values for grids with resolutions of 0.005 m up to 30 m is included in **Table 3**.

#### Aerodynamic $z_0$ Estimation ( $z_{0WT}$ , $z_{0EC}$ )

Two wind towers provided point-based profile  $z_0$  measurements ( $z_{0WT}$ ). Since the prevailing wind was down-glacier the wind towers were placed toward the down-glacier end of each of the 30 m × 70 m plots, thereby capturing part of the tower footprint within the plots (Fitzpatrick et al., 2019). One tower was placed at Plot 1 for the duration of the study and the second "roving" tower, with the same set up as the first, was erected at Plot 2 from 5 to 8 August, Plot 3 from 8 to 12 August and Plot 4 from 12 to 15 August.

Each tower comprised five NRG 40 cup anemometers (uncertainty  $\pm 0.14 \text{ m s}^{-1}$ ; starting at 0.3, 0.65, 1.22, 1.79, and 2.32 m above the surface, re-measured at each visit), one NRG 200P (uncertainty  $\pm 1.6^{\circ}$ ) wind vane and five shielded and passively-ventilated Extech RHT10 temperature and humidity loggers (uncertainty  $\pm 1^{\circ}$ C and  $\pm 3^{\circ}$ ), following previous installations and processing steps (Quincey et al., 2017; Chambers et al., 2019). Data were averaged over 15 min intervals and processed using an  $r^2$  filter of 0.99 for log-linear profile fits, a minimum wind speed threshold of 1 m s<sup>-1</sup> and a stationarity filter of 0.25°C m<sup>-1</sup>. A stability correction based on Monin-Obukhov (MO) similarity theory was applied to all profiles, as is common practice for glacier surfaces (Conway and Cullen, 2013; Stigter et al., 2017; Steiner et al., 2018).

**TABLE 4** Summary of  $z_{\text{OWT}}$  (mm) from wind towers. Mean  $z_{\text{OWT}}$  for the entire measurement duration for each plot is shown, including mean  $z_{\text{OWT}}$  corrected for stability, with standard deviation (mm) in square brackets. *n* is the number of profiles that fit MO similarity theory.

Plot	<i>z</i> <sub>0</sub> (mm)	n	Corrected z <sub>0</sub> (mm)	n
1	18.77 [13.88]	30	3.05 [1.24]	2
2	26.37 [10.99]	142	17.68 [7.61]	22
3	5.18 [5.95]	151	6.67 [5.29]	3
4	22.59 [16.82]	101	19.19 [0.00]	1

**TABLE 5** | Summary of  $z_{\text{OEC}}$  (mm) from EC towers. Mean  $z_{\text{OEC}}$  for the entire measurement duration for both levels at each tower is shown with standard deviation (mm) in square brackets, including overall mean. *n* is the number of values that fit the quality control criteria.

EC tower	Level 1		Level 2		Mean z <sub>0EC</sub> (mm)
	z <sub>oec</sub> (mm)	n	z <sub>0EC</sub> (mm)	n	
Π1	7.68 [14.48]	31	13.7 [22.3]	27	10.7
TT2	6.5 [14.19]	51	10.29 [14.3]	34	8.39
ТТЗ	7.62 [13.48]	24	8.46 [10.57]	22	8.04
TT4	2.29 [3.39]	76	1.84 [4.44]	71	2.07

Means and standard deviations for  $z_{0WT}$  are given in Table 4. The limited duration of data collection and other sources of uncertainty affect the results given by the wind tower method, which can end up yielding very few data points due to the prevalence of katabatic winds; in these conditions, where the wind speed maximum is close to the surface, profiles do not adhere to MO theory (Denby, 1999). The longest duration of wind tower data collection was available for Plot 1 (15 days), which was used to calibrate the parameters of the correction factor (Eq. 6). Despite the small number of profiles that fit MO theory once corrected for atmospheric stability, the resulting dataset is more representative of typical atmospheric conditions than the shorter duration of other Plots. Values of  $z_{0WT}$  from Plots 2, 3, and 4 were used for comparison with other z<sub>0DEM</sub> values, bearing in mind the inherent limitations of the method. Applying the MO stability correction ensured that, while the number of fitted profiles was smaller, profiles from katabatic conditions were not likely to be included erroneously.

The set-up of turbulence towers is described fully in Mott et al. (2020). On each tower, turbulence data were recorded by Campbell CSAT3 sonic anemometers at two levels, sampling at 20 Hz. Turbulent fluxes were calculated at 1 min intervals and averaged to 30 min, and  $z_{0EC}$  (**Table 5**) was derived following Fitzpatrick et al. (2019), with assumed neutral stratification (z/L > 0 and z/L < 0.2) and additional quality control filters in place for atmospheric stability, wind direction ( $150-250^{\circ}$ ), minimum wind speed ( $2 \text{ m s}^{-1}$ ), friction velocity ( $>0.1 \text{ m s}^{-1}$ ) and stationarity. TT4 was located next to the wind tower in Plot 1, with mean  $z_{0EC}$  for the entire study period providing an extra level of validation. Values of  $z_{0EC}$  from TT1-3 were used to provide independent validation, but did not overlap with any Plots or the TLS data.

## RESULTS

#### Correction Factor Development Scale and Resolution Relationships

Mean  $z_{0\text{DEM}}$  as an average of all topographic methods, was 1.6 mm ( $\sigma = 3.3$  mm), compared to 3.05 mm ( $\sigma = 1.24$  mm) for  $z_{0\text{WT}}$  and 2.07 mm ( $\sigma = 3.92$  mm) for  $z_{0\text{EC}}$ . Figure 4A shows that finer grid resolutions were associated with greater mean  $z_{0\text{DEM}}$  values for each resolution between 0.005 m (mean  $z_{0\text{DEM}} = 9.8$  mm,  $\sigma = 7.5$  mm) and 10 m (mean  $z_{0\text{DEM}} = 0.1$  mm,  $\sigma = 0.05$  mm). At coarser grid resolutions, i.e., 10, 20, and 30 m, mean  $z_{0\text{DEM}}$  increased by 0.1 –0.3 mm; these resolutions had fewer

associated data points and the <0.5 mm variation is considered inconsequential. Larger  $z_{0\text{DEM}}$  estimates were given at each resolution when larger neighborhood sizes were used, which increased the scale over which  $z_{0\text{DEM}}$  was calculated. The average increase in  $z_{0\text{DEM}}$  when neighborhood size was increased from the minimum was 0.7 mm ( $\sigma = 0.8$  mm). Exceptions to this observation included resolutions of 0.005, 0.01, 5, and 20 m, where in each case there was one instance of decreased  $z_{0\text{DEM}}$  with an increase of neighborhood size. Similar trends were observed for each of the other plots (**Figures 4C-F**).

#### **Correction Factor Calibration**

Calibration of correction factors was performed on the data from Plot 1 Day 10 (Figure 4A). Once grid resolution and  $z_{0DEM}$  were log<sub>10</sub> transformed, the fitted linear model (with 95% confidence intervals) had an intercept of –0.52  $\pm$  0.17 and slope of –0.34  $\pm$ 0.13 (RMSE = 3 mm; p < 0.01, 47 degrees of freedom). Correction factors increased with resolution from 0.22 at 0.005 m to 1.51 at 30 m. Application of the correction factors, as shown in Figure 4B, resulted in a shift in mean  $z_{0DEM}$  from 1.6 to 3.04 mm ( $\sigma$  = 3.1 mm). On average,  $z_{0DEM}$  was increased by 3.6 mm ( $\sigma$  = 3.1 mm). The mean difference between  $z_{0WT}$  and corrected  $z_{0DEM}$  was 2.2 mm ( $\sigma$  = 5.7 mm), and 3.15 mm ( $\sigma$  = 5.7 mm) between  $z_{0\text{ECT}}$  and corrected  $z_{0\text{DEM}}$ . Values of  $z_{0\text{DEM}}$  at the finest resolutions were slightly over-corrected, but the performance of the correction factors at the coarsest resolutions is the main focus, since it is desirable that similar corrections can be applied to coarse resolution data in other locations.

#### **Correction Factor Validation**

The linear model from Plot 1 Day 10 was applied to the data from the remaining Plots (Figures 4C-F), where the fit and predictive capability was checked. The data from Plot 1 Day 13 (Figure 4C) showed only minor differences to Plot 1 Day 10, including some small increases in  $z_{0DEM}$ , so the performance of the model was similar ( $r^2 = 0.4$ ). In Figure 4D, the data from Plot 2 (limited to UAV and Pole surveys) show a similar spread to those at Plot 1 and the mean corrected  $z_{0DEM}$  was 3.8 mm ( $\sigma$  = 2.9 mm). The model appeared to describe the data less well  $(r^2 = 0.1)$  due to relatively high scatter compared to the more extensive datasets at other Plots. At Plot 3 (Figure 4E) mean  $z_{0DEM}$  was corrected to 1.7 mm ( $\sigma$  = 3.6 mm), which may have been an under-correction compared to the other Plots arising from a smaller sample of  $z_{0TLS}$ data giving a weaker fit ( $r^2 = 0.3$ ). Plot 4 (Figure 4F) had the smallest  $z_{0TLS}$  values (minimum 0.01 mm) and the highest  $z_{0WT}$ values (19.2 mm). Here, mean  $z_{0DEM}$  was corrected to 2.1 mm ( $\sigma$  = 2.6 mm), with the model fitting the data as well as in Plot 1  $(r^2 = 0.4).$ 

### **Glacier-Scale Correction Factor Tests**

**Figure 5** shows corrected  $z_{0TLS}$  calculated from glacier-scale TLS data, produced from 10 m resolution input data and 10 m × 10 m neighborhoods. The lower glacier scan exhibited a smaller range of  $z_{0TLS}$  values (0.2–16.2 mm) than the upper (0.03–167.9 mm), with higher  $z_{0TLS}$  attributed to an area covered by a thin layer of supraglacial debris. The mean lower glacier  $z_{0TLS}$  was 1.7 mm



**FIGURE 4** [DEM grid resolution against  $z_{ODEM}$  on log-log scale. Each data point represents  $z_{ODEM}$  calculated using particular method, resolution and window size, extracted at the wind tower location within the Plot. (**A**) Plot 1 Day 10 uncorrected  $z_{ODEM}$ , along with power law model developed using Plot 1 Day 10 data. (**B**) Corrected  $z_{ODEM}$  for Plot 1 Day 10, where the correction factor for each DEM resolution is applied to raw input data. (**C**) Plot 1 Day 13 uncorrected  $z_{ODEM}$  with Plot 1 Day 10 model to demonstrate its applicability elsewhere. (**D**) Plot 2 uncorrected  $z_{ODEM}$ , again with model from Plot 1 Day 10 for illustration. (**E**) Plot 3 uncorrected  $z_{ODEM}$ . (**F**) Plot 4 uncorrected  $z_{ODEM}$ . The legends in (**A**) and (**B**) apply to all plots. The blue lines show  $z_{OWT}$  for each plot (see **Table 4**) and red line show  $z_{OEC}$  at Plot 1.

( $\sigma = 2.1$  mm). The  $z_{0TLS}$  values for the upper glacier covered four orders of magnitude, which reflected the presence of both heavy crevassing and icefalls alongside areas of smooth ice and those that were still snow-covered or recently exposed. The mean  $z_{0TLS}$  for the upper glacier was 0.9 mm ( $\sigma = 4.3$  mm), indicating that smoother surfaces were prevalent and that rougher surfaces

caused extreme high  $z_0$ TLS in a minority of cases. Inset panels (**Figure 5i–v**) show comparisons of the different  $z_{0DEM}$ ,  $z_{0WT}$  and  $z_{0EC}$  values for each Plot.

Generally, corrected  $z_{0\text{Ground}}$ ,  $z_{0\text{UAV}}$  and  $z_{0\text{Pole}}$  was smaller than  $z_{0\text{WT}}$ . Plot 1 Days 10 and 13 produced the most consistent  $z_{0\text{DEM}}$  values across all methods (**Figures 5i, ii**), unsurprisingly



**FIGURE 5** Upscaled and corrected  $z_{OTLS}$ . Data shown is  $\log_{10}$  transformed  $z_{OTLS}$  at 10 m resolution obtained using 10 m × 10 m neighborhoods. Inset graphs (i–v) show comparisons of  $z_{ODEM}$  from different sources. Ground, UAV and Pole data used 0.05 m resolution/0.5 m neighborhoods. TLS (A) used 5 m/50 m; TLS (B) 10 m/ 100 m; TLS (C) 20 m/200 m; TLS (D) 30 m/300 m. Error bars show the standard deviation of  $z_{ODEM}$  at all tested grid resolutions for each method. Inset (vi) shows the distribution of  $z_{OTLS}$  for the upper and lower glacier. The blue lines show  $z_{OWT}$  for each plot (see **Table 4**) and red lines show  $z_{OEC}$  at Plot 1.



FIGURE 6 | Corrected z<sub>OALS</sub> for HEF (A) with TLS area overlain. Corrected z<sub>OTLS</sub> shown in (B). Comparison of distributions for upper (i) and lower (ii) glacier shown in insets–ALS data in front of TLS data.



given the model was calibrated to Plot 1. All  $z_{0DEM}$  was lower than  $z_{0WT}$  at Plots 2, 3, and 4; however, the observed  $z_0WT$  values at Plots 2 and 4 were unexpectedly high (see *Robustness of*  $z_0$  *Estimates*). Mean correction error for available data at Plots 1 and 3 was 1.9 mm ( $\sigma = 5.2$  mm), 12 mm at Plot 2 and 16.1 mm at Plot 4. Plot 2 was not covered by TLS data so could not be compared fully. Plot 3 overlapped with the lower edge of the upper glacier scan, meaning  $z_{0TLS}$  was not available at resolutions of 20 and 30 m.

# Application to Regional Airborne Laser Scan Data

In order to demonstrate the potential for use of this upscaling approach at other locations without wind tower validation,  $z_{0DEM}$ 

was calculated and corrected for glaciers surrounding Hintereisferner that are >0.5 km<sup>2</sup> in area. Comparing corrected  $z_{0ALS}$  and  $z_{0TLS}$  values over Hintereisferner (Figures **6A,B**),  $r^2$  for the upper and lower glacier was 0.6 and 0.5, respectively (both p < 0.01) suggesting that the distribution and values matched well (**Supplementary Figure S1**). Due to the temporal misalignment between the datasets,  $z_{0ALS}$  was not expected to replicate  $z_{0TLS}$  exactly, yet the spatial patterns and the distributions of both upper and lower glacier are similar, and are included here for demonstration in Figures **6A,B** (insets). The sliding neighborhood operation allowed  $z_0$  to be estimated quickly for all glaciers at once from a single DEM of the region (Figure 7A). A correction factor of 1.35 was selected from the index (Table 3) according to resolution and neighborhood size (10 m resolution, 100 m × 100 m neighborhood size), and then applied to the uncorrected grid of  $\log_{10} z_{0ALS}$ .

Regionally, corrected  $z_{0ALS}$  was generally 0.1–10 mm (Figure 7A). Zones of  $z_0 < 0.1$  mm coincided with the highest elevation areas where persistent snow cover is likely, while zones of 10–100 mm covered areas of supraglacial debris, medial moraine (Figure 7C), light crevassing and rough ice. The roughest areas (>100 mm) matched heavily crevassed areas such as icefalls (Figure 7B), ice margins and where the glacier polygon overlapped with bedrock. Extreme values of  $z_{0ALS}$  (>> 100 mm) were coincident with areas of very high relief, such as icefalls, heavy crevassing and ice-margins where uncertainty in elevation data is greatest. Visual comparison of  $z_{0ALS}$  with inset images (Figures 7B–D) confirmed that expected spatial distributions of  $z_0$  can be derived from our workflow at the regional scale, as well as the glacier scale.

## DISCUSSION

#### Robustness of $z_0$ Estimates

Several studies have demonstrated similarity between topographic and aerodynamic  $z_0$  over glaciers at the plot-scale (see Brock et al., 2006; Miles et al., 2017 for extensive compilations), lending confidence to our topographic estimates at finer resolutions. The values we obtained for  $10 \times 10$  m and  $30 \text{ m} \times 70 \text{ m}$  plots using ground-based, UAV and pole-based surveys fall well within expected bounds (same order of magnitude) of  $z_0$  for an ablating Alpine glacier (~0.1–10 mm; c.f. Munro, 1989; Brock et al., 2006), with the exception of the finest resolutions of 0.005 m. Power law behavior was observed in the distribution of 3D  $z_0$ DEM against DEM grid resolution (Figure 4), as also shown by Rees and Arnold (2006) for 2D transect  $z_0$  estimates. Where Rees and Arnold (2006) noted a difficulty in obtaining a dataset that spanned a continuous range of scales and resolutions, modern survey techniques have afforded us the opportunity to collect and analyze data at resolutions across five orders of magnitude and at scales across four orders.

The  $z_{0DEM}$  estimates within this study are also comparable to those given for different surfaces on Hintereisferner obtained using a slightly different topographic approach (Smith et al., 2020). A temporal model developed in Smith et al. (2020) noted that  $z_0$  was generally associated with the length of time that ice surface areas had been exposed by snow melt, as well as by physical factors such as surface gradient, which accounts for some of the highest  $z_0$  values. The  $z_{0TLS}$  values for rougher surfaces like crevasses and a rock pedestal were more extreme in this study compared to the plots used by Smith et al. (2020), which were up to ~20 mm, compared to 102 mm in this study. It is likely that our data incorporated some ice-marginal areas given the mismatch in resolutions between the TLS data and the glacier outline polygons derived from coarser resolution imagery. Other characteristic ice surfaces, such as the smooth/dirty ice, supraglacial channel and pressure ridge plots surveyed at the plot-scale in detail by Smith et al. (2020) compare well with the estimates of  $z_{0DEM}$  here (within <5 mm generally).

The sliding neighborhood algorithm used to obtain  $z_{0DEM}$ estimates in this study is similar to that described and used by Fitzpatrick et al. (2019), who used sliding neighborhoods to calculate a "drag parameter" ( $F_{\rm D \ local}$ ). Fitzpatrick et al. then attempted to account for fetch when calculating  $z_0$  for a point  $(z_{0 \text{ bloc}})$  by equally weighting all cells within an area then finding their sum. Despite some minor differences in the initial calculation of each parameter, the range of values found for  $F_{\rm D \ local}$  by Fitzpatrick et al. (2019) is the same as the range of corrected  $z_{0DEM}$  values found in this study (~10-5-~100 m). For further comparison, we also calculated  $z_{0 \text{ bloc}}$  for the  $z_{0 \text{ UAV}}$  grids of each Plot, by equally weighting  $z_0$  within all cells and finding their sum. We found that the given  $z_0$  bloc values from our Plots were at least an order of magnitude smaller than corrected  $z_{0UAV}$ values. It is worth noting that the area used for calculating  $z_0$  bloc here was only ~70 m  $\times$  30 m, compared to the ~2,000 m  $\times$ 2,000 m used by Fitzpatrick et al. (2019), so the smaller order of magnitude is likely to be a function of the smaller number of values within a plot. Additionally, we consider the sliding neighborhood method used here to be more practical when trying to model glacier-wide  $z_0$  for use in a distributed SEB model because it does not require topographic data from beyond the margins of the glacier; although fetch is not explicitly accounted for, doing so would require the incorporation of turbulence characteristics across different surfaces and especially near the margins of the glacier (Nicholson and Stiperski, 2020), which would detract from the intended simplicity of our approach.

A key area of uncertainty exists in the  $z_{0WT}$  values derived from Plots 2 (17.86 mm) and 4 (19.18 mm), which were both greater than expected considering field observations of the surface characteristics,  $z_{0WT}$  and  $z_{0EC}$  values at other plots and those from similar glaciers (e.g., Brock et al., 2006). Confidence in the value at Plot 4 is low because only one  $z_0$  value was obtained after stability correction; conversely, confidence is higher for Plot 2, where the 22 values given (Table 4) suggest either i) potential underestimation by topographic methods or ii) a stronger wind speed gradient likely due to a near-surface wind speed maximum under katabatic conditions (c.f. Mott et al., 2020). While  $z_{0WT}$  values were included for comparison, we acknowledge that the short duration of data collection at Plots 2, 3, and 4 (<4 days for each) reduces confidence in their results. It is typical to collect much longer time-series in order to ensure that the average aerodynamic properties of the surface are captured even after aggressive data filtering (Stull, 1988; Radić et al., 2017; Fitzpatrick et al., 2019). These uncertainties further highlight the merits of a topographic approach, which does not rely on a lengthy field campaign.

We are confident that the longer data collection period at Plot 1 (2 weeks) was sufficient for it to be used to calibrate the parameters of the correction factor and that, at worst, it indicated the likely order of magnitude of  $z_0$ . This is supported by topographic estimates of  $z_0$ , which were similar for Days 10 and 13, and by the mean  $z_{0EC}$  values for the same Plot which was less than 1 mm smaller and had a standard deviation that would still put  $z_0$  within the same order of magnitude. Other theoretical and methodological flaws in the retrieval of  $z_0$  from

wind profiles and EC data using MO similarity theory, such as the likely presence of katabatic winds (Mott et al., 2020) and collecting data over an ablating surface (Denby, 1999; Denby and Greuell, 2000; Litt et al., 2014; Radić et al., 2017), deterred us from seeking a perfect match between corrected  $z_{0DEM}$  and  $z_{0WT}$ .

## **Correction Factor Performance**

The correction factor was developed by calibrating topographic  $z_0$ from Plot 1 with  $z_{0WT}$  for the same plot, assuming  $z_{0DEM}$  to be representative of a portion of the footprint of the wind tower (c.f. Fitzpatrick et al., 2019). Our goal was to match the wind-towerderived  $z_0$  values to within an order of magnitude. Small differences in  $z_0$  will alter resultant turbulent flux values, but it is unreasonable to simply accept any one value as being absolute because  $z_0$  changes rapidly in both time and space over several orders of magnitude (Smith et al., 2016a; Smith et al., 2020). Furthermore, it is difficult to measure a "true" value of  $z_0$ , as even EC data (often seen as the benchmark) relies to an extent on modeling and associated embedded assumptions (Cullen et al., 2007; Radić et al., 2017). Therefore,  $z_{0WT}$  values were used as a calibration point to ensure  $z_{0DEM}$  was calibrated to within an order of magnitude of the true  $z_0$ . Of course, further aerodynamic data from this and other regions would strengthen the upscaling workflow developed here. However, this upscaling method provides a more reasonable estimate of distributed  $z_0$  than is achievable using aerodynamic methods by themselves or through the use of past topographic  $z_0$  techniques.

From visual inspection, and bearing in mind the absence of validation data beyond Hintereisferner itself, the correction factor appears to produce a realistic representation of spatial  $z_0$ variability at multiple scales (Figure 6) when resulting values are compared to past studies. Spatial patterns in the derived values are maintained and correspond to areas that appear smoother/rougher based on surface features identified in satellite imagery. Areas of dubiously high values (>102 mm) are coincident with areas where the DEM is least reliable, including heavily crevassed regions, icefalls and other areas of high relief (Sailer et al., 2014). While the distribution of  $z_0$ matches expectations from the perspective of surface roughness, it is worth noting that the footprint, or fetch, of each particular DEM cell is not considered. The fetch includes all aerodynamically important areas upwind of a point Kljun et al., 2015 which, over glaciers, has been observed to extend from ~100 to 200 m, depending on the wind speed (Fitzpatrick et al., 2019). This leads to the question of how the sharp transitions between areas of contrasting roughness that can be seen in Figures 5, 6 can be accounted for in a straightforward model such as that used in this study, or indeed if they can be at all, as the nature of turbulent air flow over glaciers is extremely complex when considering the influence of the valley sides, different surfaces and tributary valleys (Stiperski and Rotach, 2016; Mott et al., 2020; Nicholson and Stiperski, 2020). It may be possible to account for fetch in topographic models in the future, considering advances in the availability of fine-resolution fields of meteorological variables such as wind speed (Sauter and Galos, 2016) and capabilities for modelling their interactions with terrain at small (100 m) scales (Fiddes and Gruber, 2014; Peleg

et al., 2017). That being said, the topographic  $z_0$  methods employed here were developed only to use the properties of a surface to produce a value similar to an aerodynamically-derived value for the same point (Lettau, 1969; Munro, 1989); our workflow is intended to retain this rationale whilst indicating broader patterns. Additional complexities could be factored in with further work, depending on the application.

## Implications

Use of the correction factor presented here allowed a realistic  $z_0$ value to be derived from a relatively coarse resolution DEM of Hintereisferner, Austria, of the type typically available for whole glacierized regions of the cryosphere. Additional testing and ground-truthing, ideally with sonic anemometers, would allow further calibration of the power laws and correction factors, especially in areas with varying debris sizes and coverage (e.g., Miles et al., 2017; Nicholson and Stiperski, 2020) and in areas covered by snow, that were not sampled in our study. Our workflow is an advance in the use of topographic methods for estimating  $z_0$  that can be applied (with caution until further validated) to other glaciers without the need for additional field data collection. It provides a more robust estimate of  $z_0$  than is achievable by collecting 2D transects, or by implementing  $z_0$ values borrowed from the literature and we foresee that the prevailing practice of using a constant or assumed  $z_0$  value in both distributed and single-point SEB models could potentially be eliminated.

A worthwhile additional step will be to develop the workflow for spatially distributed  $z_0$  grids established here into a model that includes a temporal dimension. Fully distributed  $z_0$  could then be incorporated in a distributed surface energy balance (SEB) model, to analyze the effects that this new quantification of  $z_0$  has on the turbulent fluxes and resultant meltwater production. This development would be especially useful when combined with the more widespread availability of national ALS campaign datasets (such as that used herein for Austria), or regional DEM datasets such as the ArcticDEM, the Reference Elevation Model for Antarctica and the High Mountain Asia DEM (Shean, 2017; Porter et al., 2018; Howat et al., 2019) and increasing availability of multi-temporal DEMs from platforms such as WorldView 2, Pléiades and Geo-Eye-1 (Belart et al., 2017; Shean, 2017). A work-flow diagram and example code for calculating corrected  $z_0$  is provided in the Supplementary Information (Supplementary Figure S2) to facilitate replication of the method, with the caution that full testing across different sites exhibiting a variety of surface characteristics, and using topographic data from a range of different sources, is still required.

## CONCLUSION

We carried out a multi-scale investigation of topographic  $z_0$  over Hintereisferner, Austria. Data from SfM, TLS and ALS were employed to rapidly capture  $z_0$  at resolutions from sub-cm to 30 m.  $z_0$  values exhibited a power law behavior with data resolution. From this relationship we devised correction factors that adjusted topographic  $z_0$  so that it was within an order of magnitude of previously validated, more robust values from both finer resolutions and  $z_0$  calculated from wind profiles. While sensitivities and uncertainties in z<sub>0</sub> estimates persist due to scale/ resolution dependence and the simplification of aerodynamic processes, the method presented here has allowed  $z_0$  to be rapidly estimated at the glacier scale, capturing more detail and variability than is possible with point-scale or 2D techniques. We used the same upscaling method to demonstrate how topographic  $z_0$  can be estimated at the glacier scale across an entire region. Further analysis for sites within different climatic zones is required to calibrate the correction factors, but the workflow provides a robust foundation for obtaining spatially distributed  $z_0$  without the need for lengthy field campaigns with delicate meteorological equipment, by being compatible with regional elevation datasets that are increasingly becoming publicly available.

## DATA AVAILABILITY STATEMENT

The datasets supporting the conclusions of this article are available from the UK Polar Data Centre repository, DOI: https://doi.org/ 10.5285/57EB95F0-1650-4670-A1BB-8CCC90CB1C5F

## **AUTHOR CONTRIBUTIONS**

Fieldwork was carried out during the HEFEX campaign. JRC, TS and MS collected UAV, pole, ground and wind tower data. Turbulence tower data was collected by IS, LN and JM, and TLS data by RS. JRC performed data analysis and wrote the

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manuscript with supervision from MS and DQ. IS processed the turbulence data and provided manuscript feedback along with JLC, LN, JM, RS, and MJ.

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## SUPPLEMENTARY MATERIAL

The Supplementary Material for this article can be found online at: https://www.frontiersin.org/articles/10.3389/feart.2021.691195/full#supplementary-material

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**Conflict of Interest:** The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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