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1 A new modelling approach to sediment bypass prediction
2 applied to the East Coast Basin, New Zealand

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8 **ABSTRACT**

9 Predicting when turbidity currents are erosional or depositional (i.e., leaving no depositional
10 record vs. leaving a deposit) remains challenging. Here we combined observations from
11 submarine channel morphology with a new sediment transport model to derive thresholds for
12 net erosional, equilibrium or net depositional flow and to predict how far turbidity currents can
13 transport different grain size classes down-channel. The approach was applied to the modern
14 Madden and Omakere channels, which traverse the Hikurangi subduction margin of the North
15 Island of New Zealand. A bathymetric dataset was used to establish the downstream change of
16 channel geometry. Taking account of centripetal and Coriolis forces, the flow super-elevation
17 method was used to estimate variations in flow velocity and concentration along the channels.
18 These parameters were used as model inputs in order to estimate the potential distribution of
19 sand in the system, assuming well sorted and poorly sorted sediment in suspension. The
20 predicted sand distribution maps deposited by poorly sorted flows in the channels show good
21 agreement with RMS amplitude mapping of the seafloor. These results confirm that thicker
22 flows, and those carrying well sorted suspensions can bypass sediment over lower slopes than
23 thinner flows and those carrying more poorly sorted suspensions. The net erosion and net
24 deposition thresholds derived from this study may help to guide and constrain predictions of

25 potential sediment bypass zones in seafloor and subsurface systems, and hence better constrain
26 the predicted *loci* of deposition.

27 INTRODUCTION

28 Deep-marine siliciclastic systems are volumetrically some of the most important sedimentary
29 environments on the surface of the earth (Covault and Graham, 2010; Meiburg and Kneller,
30 2010; Talling et al., 2015). Submarine gravity currents (e.g. turbidity currents) transport
31 sediment from shallow to deep-water, often developing complex depositional geometries (e.g.
32 Richards and Bowman, 1998; Wynn et al., 2002; Booth et al., 2003; Gardner et al., 2003;
33 Posamentier and Kolla, 2003; Deptuck et al., 2008; Ponce and Carmona, 2011; Dorrell et al.,
34 2015; Sychala et al., 2017). Whether suspended sediment of a particular grain size is either
35 transported up to the maximum flow runout distance, or deposited at any particular location
36 along the flow pathway plays a key role in: 1) the distribution of sediment across shelf-to-basin
37 slope profiles (Normark, 1978; Mutti and Normark, 1987; Prather et al., 1998; Wynn et al.,
38 2002; Hadler-Jacobsen et al., 2005; Carvajal and Steel, 2009; Pyles et al., 2011); 2) the
39 reservoir quality of turbidite sandstones through fractionation of different grain size classes
40 (Pyles and Jennette, 2009; Horseman et al., 2014; Marchand et al., 2015; Bell et al., 2018); and
41 3) the development of up-dip stratigraphic pinch-outs that trap hydrocarbon reservoirs (Straccia
42 and Prather, 2000; Carruth, 2003; Prather, 2003; Doré and Robbins, 2005; Milton-Worsell et
43 al., 2006; Horseman et al., 2014; Van der Merwe et al., 2014; Amy, 2019; Hansen et al., 2019).
44 However, determining whether a turbidity current transports or deposits sediment remains
45 challenging, despite recent work observing and monitoring turbidity currents (Vangriesheim et
46 al., 2009; Xu et al., 2014; Paull et al., 2018; Zhang et al., 2018).

47 Here we used a theoretical model for the threshold between net sediment erosion and
48 net sediment deposition of turbidity currents to determine the grain sizes that might be

49 transported or deposited along the Madden and Omakere slope channels of the East Coast Basin
50 (ECB), New Zealand. The submarine slope of the ECB represents an actively growing
51 subduction wedge (Nicol et al., 2007; Barnes et al., 2010), with a series of trench-slope basins
52 that are either supplied with sediment or starved, depending on the presence of slope channels
53 (McArthur et al., 2019). The flow properties of turbidity currents were calculated based on an
54 assumed relationship to the morphology of their confining channels. The thresholds between
55 erosion and deposition were calculated assuming flows carrying non-cohesive sediment of a
56 range of grain size classes and grain size distributions, accounting for the capacity and
57 competence of the flow, flow height and bulk sediment concentration. Furthermore, the results
58 from the model are validated by geophysical and petrophysical information. Here, we
59 demonstrate that the grain size distribution in the flow has a large impact on sediment transport
60 thresholds, therefore potentially controlling the sand distribution in the system.

61 **Terminology**

62 Despite its importance, there is no agreed definition of sediment bypass and bypassing flows
63 between disciplines that study both associated sediment transport processes and products. In
64 stratigraphic studies, *bypassing flow* or *bypass* have been broadly used for flows that partially
65 or completely transport their sediment load beyond a point of observation (e.g. Lowe, 1982;
66 Mutti and Normark, 1987; Amy et al., 2000; Cronin et al., 2005b; Kolla et al., 2007; Wynn et
67 al., 2007; Carvajal and Steel, 2009; Talling et al., 2012; Stevenson et al., 2013; Sylvester et al.,
68 2015). Furthermore, *bypassing flow* has also been used to refer to erosional flows despite the
69 fundamental differences (i.e. changes in flow capacity) between both flow types (e.g. Mutti
70 and Normark, 1987; Wynn et al., 2002; Hubbard et al., 2014; Stevenson et al., 2015; Lang et
71 al., 2017). In contrast, previous experimental and numerical studies have defined non-
72 depositional flows as equilibrium, self-sustaining or autosuspending flows (Bagnold, 1962;
73 Kneller, 2003; Sequeiros et al., 2009; Dorrell et al., 2018). These definitions describe a flow

74 state where there is a net balance between sediment erosion and sediment deposition, and
75 allows differentiation of non-equilibrium flow regimes (i.e., erosional versus depositional
76 flow).

77 In an attempt to make a clear differentiation of bypassing flows in both process and
78 stratigraphic contexts, the terms *bypassing flows*, *partially bypassing flows* and *depositional*
79 *flows* have been suggested by Stevenson et al. (2015). These definitions provide a useful
80 framework in stratigraphic terms to distinguish flows that transport their complete load from
81 those that leave a deposit. However, in terms of process, the definition of depositional flows
82 implies that there must be some bypassing fraction in suspension, which overlaps with the
83 definition for partially bypassing flows. Therefore, for the purpose of this paper, we use the
84 following terminology which applies to the suspension load transported by channelized
85 turbidity currents (Fig. 1).

86 *Erosional flows* are non-equilibrium, under-capacity flows which entrain sediment into
87 suspension and transport it beyond the point of observation. *Equilibrium flows* are flows at
88 capacity, sediment deposition is balanced with erosion and allows for transport of their
89 complete suspension load beyond the point of observation, resulting in absence both of a
90 depositional record and of erosional features (e.g. Stevenson et al., 2013). *Depositional flows*
91 are non-equilibrium flows that are over capacity and deposit a fraction of their suspended load
92 whilst the remainder is further transported downstream. Here, *sediment bypass* refers to the
93 process where sediment is transported beyond a point of observation by erosional flows,
94 equilibrium flows or depositional flows. The terms *deposition* and *bypass* preceded by the grain
95 size (e.g. sand deposition and silt bypass) will be used to differentiate the grain sizes that are
96 simultaneously deposited and bypassed at the point of observation.

97 **GEOLOGICAL SETTING**

98 The East Coast Basin (ECB) is a growing subduction wedge on and offshore of the eastern
99 margin of the North Island, New Zealand. The formation of the ECB is the result of collision
100 and oblique subduction of the Pacific Plate below the Australian Plate (Fig. 2) (Ballance, 1976;
101 Lewis and Bennett, 1985; Davey et al., 1986; Lewis and Pettinga, 1993; Nicol et al., 2007;
102 Bland et al., 2015; McArthur et al., 2019). The basin is limited to the east by the Hikurangi
103 subduction trench and to the west by the axial ranges of the North Island (Ballance, 1976;
104 Pettinga, 1982; Chanier and Ferriere, 1991; Nicol et al., 2007). The basin is dominated by NE-
105 SW striking thrust faulting, sub-parallel to the trench axis (Lewis and Pettinga, 1993) and is
106 divided into inner, mid and outer structural domains, each of which displays distinct
107 deformation styles (Fig. 2; McArthur et al., 2019).

108 Hikurangi Margin subduction initiated in the late Oligocene, at c. 25 Ma (Ballance,
109 1976; Chanier and Ferriere, 1991; Nicol et al., 2007; Reyners, 2013); this convergence has
110 created a series of elongate growth structures and trench-slope sub-basins, typically tens of
111 kilometers long by kilometers wide (Lewis and Pettinga, 1993; Barnes et al., 2010; Bailleul et
112 al., 2013; McArthur et al., 2019). Continued compression to the present day has resulted in
113 uplift and exhumation of the innermost trench-slope sub-basins, which are exposed at outcrop
114 (Bailleul et al., 2007), whilst the majority of the wedge remains submerged (Barnes et al.,
115 2010). Therefore, the basin predominantly experienced marine conditions during the Neogene,
116 with widespread deposition of mudstone and sandstone turbidites within bathymetric lows
117 (Bailleul et al., 2007; Burgreen-Chan et al., 2016), whilst submarine canyons and channels
118 incised ridges, acting as sediment conduits between sub-basins (McArthur and McCaffrey,
119 2019).

120 The present-day turbidite systems in the ECB dominantly transport very-fine sand and
121 silt (Barnes and Audru, 1999; Lewis and Pantin, 2002; Mountjoy et al., 2009; Wallace et al.,
122 2019). The sedimentation rates, character of sedimentary pathways and subsequent fill of sub-

123 basins varies throughout the subduction wedge, where a range of channels and submarine
124 canyons are observed (Mountjoy et al., 2009; Bailleul et al., 2013; McArthur and McCaffrey,
125 2019). Channels in the northern and southern parts of the wedge potentially connect with the
126 Hikurangi Channel (Fig. 2) (Mountjoy et al., 2009). Channels in the central zone terminate in
127 the mid-portion of the wedge delivering sediment to mid and outer trench-slope sub-basins
128 (McArthur and McCaffrey, 2019). This variation in sediment distribution systems has been
129 interpreted to result partly from the development of high-angle thrust faults forming steeper
130 ridges in the northern and southern zones of the basin (McArthur et al., 2019). Other sediment
131 transport processes in the basin include debris flows (Mountjoy and Micallef, 2012; McArthur
132 and McCaffrey, 2019) and contour currents (Carter et al., 1996, 2004; Bailey et al., 2020),
133 together with hemipelagic fallout.

134 **METHODS**

135 **Datasets and Morphological Analysis**

136 Bathymetric data of the East Coast Basin with a horizontal resolution of 100 m (provided by
137 the National Institute of Water and Atmospheric Research, New Zealand) were used to
138 calculate the channel thalweg, and channel dimensions in the Madden and Omakere channels
139 as illustrated in Figure 3. The calculations were conducted using the Topotoolbox program in
140 MATLAB, which enables landscape and drainage analysis of digital elevation models
141 (Schwanghart and Kuhn, 2010). The channel thalwegs were calculated from landscape
142 drainage and the bathymetric profiles were used to calculate the thalweg down-slope gradient
143 (S).

144 Channel dimensions were measured from channel cross-sectional profiles as
145 represented in Figure 3. The cross-sections in each channel were taken perpendicular to the
146 channel thalweg approximately every 2 km downstream. In levee-bound channel sections, the

147 mean channel height (\bar{H}) was calculated from the maximum (H) and minimum (h) levee crest
148 heights, which represent the relief between the channel thalweg and the channel flank top (Fig.
149 3). In sections where the channels were fully or partially ridge-confined, the erosional rim was
150 used as a proxy for flow height and for the calculation of H , h and \bar{H} (Fig. 3). The maximum
151 channel width (W) represents the horizontal distance between the higher levee crest and the
152 lower levee crest, whereas the minimum channel width (w_m) is the horizontal distance between
153 the lower levee crest and the opposite channel wall. The mean channel width (\bar{W}) was
154 calculated from the maximum (W) and minimum (w_m) channel widths (Fig. 3). The width of
155 the channel floor (W_b) (Fig. 3) was measured perpendicular to the channel axis. The cross-
156 channel slope (γ) was calculated from the horizontal (W) and vertical difference ($H-h$) of the
157 levee crests or erosional rims, $\gamma = (H - h)/W$. The channel cross-sectional area (A) was
158 calculated using Eq. 1. The thalweg radius of curvature (r) (Fig. 3; cf. Keevil et al., 2007) was
159 measured at each channel bend.

$$160 \quad A = \frac{W_b + W_m}{2} h + \frac{W_m}{2} (H - h). \quad (1)$$

161 The canyon-confined portion of the Madden Channel (Fig. 4) was omitted in the analysis
162 because estimates of the width and vertical channel relief do not permit estimation of flow
163 properties where traversing flows are unlikely to overtop the erosional confinement (see
164 below).

165 A 3D pre-stacked depth-migrated seismic dataset (acquired at broadband frequency in
166 2017) provided by WesternGeco Multiclient was used to map and generate the seafloor dip
167 map using Schlumberger's Petrel© E&P software (Fig. 5), and the root mean square (RMS)
168 amplitude map of the shallow subsurface (<50 m). The dataset has an inline (NW-SE) and
169 crossline (SW-NE) spacing of 25 m and a vertical resolution of 6.7 m. The vertical resolution
170 was approximated through the calculation of the tuning thickness (Widess, 1973), using a

171 wavenumber of 0.037m^{-1} calculated from the instantaneous frequency attribute map of the
172 seafloor.

173 **Turbidity Current Modelling**

174 Downstream flow velocities and sediment concentration in the Madden and Omakere channels
175 were calculated using the cross-flow equation from Komar (1969), and the relationship
176 between the bed slope, densimetric Froude number and bed friction coefficient (Parker et al.
177 1987; Abad et al., 2011) that balances the gravitational force and drag force at the bed.

178 The cross-flow equation describes the balance of the Coriolis force, the centrifugal
179 force produced at channel bends and the pressure gradient in a turbidity current assuming a
180 bankfull flow (Komar, 1969), which can be used to reconstruct the flow properties of turbidity
181 currents (e.g. Bowen et al., 1984; Klauke et al., 1997; Pirmez and Imran, 2003; Stevenson et
182 al., 2018). In clockwise flows in the Southern Hemisphere, the centrifugal and Coriolis forces
183 are oriented towards the outer bank, causing preferential overspill and cross-sectional channel
184 relief asymmetry, where the left-hand side bank is higher (looking in a downstream direction)
185 (Cossu and Wells, 2010; Dorrell et al., 2013a); whereas in anticlockwise flows, the forces are
186 opposed. At high latitudes, in straight channel sections, cross-sectional channel asymmetry
187 arises in the absence of centrifugal forces, due to Coriolis force alone (Cossu et al., 2010). The
188 cross-channel slope (γ in Fig. 3) is then used as a proxy of the superelevation of the boundary
189 between the flow and the ambient fluid (Komar, 1969).

190 The momentum balance of the pressure gradient force, the Coriolis force and the
191 centrifugal force across the channel (Komar, 1969; Stacey and Bowen, 1988; Wells and Cossu,
192 2013) can be written as,

$$193 \quad gRC\gamma = \pm fU + \frac{U^2}{r}, \quad (2)$$

194 where $g = 9.81 \text{ m/s}^2$ gravity; $R = (\rho_s/\rho_f - 1)$ is the excess density (where $\rho_s = 2650 \text{ kg/m}^3$ is the
 195 density of quartz for the material in suspension, and $\rho_f = 1030 \text{ kg/m}^3$ is the density of the fluid);
 196 C , the bulk sediment concentration (vol./vol.); $\gamma = (H - h)/W$, the cross-channel slope
 197 (m/m); f , the Coriolis acceleration $f = 2\Omega \sin\theta$, where $-f$ for clockwise flows and $+f$ for
 198 anticlockwise flows, Ω is the Earth's rotation rate, and θ the latitude (values for $\theta = -40.5^\circ$ to -
 199 40.7° in the Madden Channel and -40.3° to -40.4° in the Omakere Channel); U , the downstream
 200 flow velocity in m/s; and r , the thalweg radius of curvature (m), in straight channel sections
 201 $r \rightarrow \infty$.

202 Eq. (2) can be rewritten in terms of the Froude number, Fr , (Cossu and Wells, 2010;
 203 Wells and Dorrell, 2020) where $Fr = U/\sqrt{gRC\bar{H}}$ (Parker et al. 1987),

$$204 \quad \gamma = Fr^2 \left(\frac{\pm f \bar{H}}{Fr \sqrt{gRC\bar{H}}} + \frac{\bar{H}}{r} \right), \quad (3)$$

205 where \bar{H} is the mean flow depth (m).

206 Further, using the model of Parker et al. (1987) which balances gravitational driving
 207 force with frictional drag at the bed, and through the entrainment of ambient water (e.g. Abad
 208 et al. 2011),

$$209 \quad S = \frac{C_d + e_w \left(1 + \frac{Ri}{2}\right)}{Ri}, \quad (4)$$

210 where S is the calculated down-channel slope in m/m; $C_d = 0.0025$ (Abad et al., 2011; Konsoer
 211 et al., 2013), the drag coefficient, which is considered constant in the downstream direction
 212 for the calculated flows; Ri , the bulk Richardson number (which scales inversely with Fr^2
 213 (Wells and Dorrell, 2020)) which is a measure of mixing of the flow-ambient fluid interface
 214 (Parker et al., 1987; Abad et al., 2011),

$$215 \quad Ri = \frac{gRC\bar{H}}{U^2} = \frac{1}{Fr^2}, \quad (5)$$

216 and e_w , the ambient water entrainment by mixing is parameterized using

$$217 \quad e_w = \frac{0.00153}{0.0204 + Ri}, \quad (6)$$

218 a relation empirically derived by Parker et al., (1987) for turbidity currents.

219 A non-linear least squares MATLAB solver was used to derive C and Fr using Eqs. (3)
220 to (6). Fr values were then used to calculate downstream flow velocities U in m/s using Eq. 7
221 (Parker et al., 1987; Abad et al., 2011). The shear velocity u^* was calculated in m/s via Eq. (8)
222 (Parker et al., 1987; Abad et al., 2011); the flow discharge Q in m³/s was calculated through
223 Eq. (9) and sediment discharge Q_s in kg/s with Eq. (10) using the derived sediment
224 concentration C :

$$225 \quad U = Fr \sqrt{g R C H}, \quad (7)$$

$$226 \quad u^* = \sqrt{C_d} U, \quad (8)$$

$$227 \quad Q = UA, \quad (9)$$

$$228 \quad Q_s = QC. \quad (10)$$

229 This analysis allows for improved estimations of bulk sediment concentrations for
230 turbidity currents compared to previous work, where analyses do not jointly solve for Coriolis,
231 gravitational and drag forces at straight channel sections and channel bends (Stevenson et al.,
232 2018). The mean flow depth and bulk sediment concentration derived from this analysis were
233 used as inputs into the sediment transport model to estimate the threshold between net sediment
234 erosion and net deposition of turbidity currents.

235 **Modelling of the Net Erosion and Net Deposition Threshold for Turbidity Currents**

236 A Flow-Power Flux-Balance type model (Dorrell et al., 2018) was used to model the threshold
237 of erosion and deposition for the turbidity current conditions estimated for the Madden and

238 Omakere channels. The model makes at-a-point predictions, and does not forward model the
 239 evolution of the currents nor the downdip transport of sediment in suspension (i.e. carried from
 240 source or eroded). In this analysis the average flow depths and sediment concentrations,
 241 calculated at each channel section from the cross-flow equation, were used to constrain the
 242 current hydrodynamics. The model incorporates the bulk capacity (maximum amount of
 243 sediment that can be transported in suspension by a turbulent flow), competence (particle class
 244 specific capacity) (Dorrell et al., 2013b), using a polydisperse description of sediment
 245 suspension. For each particle class the threshold between a net erosional and a depositional
 246 flow is given by

$$247 \frac{C_i^-}{C_m} E_i = C_i^+ w_{si} \forall i, (11)$$

248 where the total concentration of sediments at the bed is

$$249 \sum_{i=1}^N C_i^- = c_m, (12)$$

250 where C_i^- is the grain size class concentration in the active layer of the bed that freely exchanges
 251 material with material transported as suspended load (Dorrell et al., 2013b), C_m is the packing
 252 concentration, C_i^+ is the grain size class concentration at the bed and w_{si} the particle settling
 253 velocity for each grain size class. Further, the description of the vertical distribution of
 254 suspended sediment concentration was determined by the mass conservation equation

$$255 w_s C_i(Z) = -k_s \frac{dC_i(Z)}{dZ}, (13)$$

256 where $Z = z/L$ is the dimensionless flow depth, $k_s = ku^* L f(Z)$ is the eddy diffusivity which
 257 was assumed constant, therefore the flow length scale $L = h/6$, and the Rouse number is defined
 258 by $\beta = 6 \frac{W_{si}}{ku_i^*}$ (Dorrell and Hogg, 2012), $k=0.4$ is the von Kármán constant; thus, the diffusion
 259 profile is given by

260 $C_i(Z) = C_i^+ \left(\frac{1-e^{-\beta}}{\beta} \right). (14)$

261 The model incorporates a sediment entrainment function in which the power required
262 to move sediment and incorporate it into suspension is proportional to the depth-averaged flow
263 power:

264 $E_i = \varepsilon \rho (gR\bar{H})^{-1} \Delta u_i^{*3}, (15)$

265 where $\varepsilon = 13.2$ (Dorrell et al., 2018) is an empirical parameter describing entrainment
266 efficiency.

267 The characteristic grain size classes (d_{50}) used to solve Eq. (11) to Eq. (15) range from
268 $\phi = 8$ to $\phi = -2$ (Wentworth scale; i.e., very fine silt to gravel). A log normal distribution derived
269 from empirical data (Dorrell et al., 2018) was used to model grain size distribution, where a
270 standard deviation $\sigma=0$ represents a monodisperse suspension, a standard deviation $\sigma = 0.5$ is
271 equivalent to a polydisperse suspension that is well sorted (Folk, 1966) and a standard deviation
272 $\sigma=2$ is equivalent to one that is poorly sorted (Folk, 1966). The slope gradient required to
273 maintain equilibrium conditions for a given grain size and grain size distribution was calculated
274 through Eq. (4) and presented in m/m and degrees. Then, the slope values in m/m calculated
275 for each down-channel section were plotted in the net erosion-deposition threshold plots (Fig.
276 8 and Fig. 9).

277 **RESULTS**

278 **Madden Canyon – Channel Morphology**

279 The Madden Canyon is located downstream of the Madden Sub-basin where it incises the
280 Madden Banks at water depths of ~ 1400 m. (Fig. 4). The Madden Canyon has been classified
281 as being detached from direct hinterland supply (McArthur and McCaffrey, 2019). The canyon

282 head exhibits a funnel shape measuring ~9 km wide and exhibiting a V-shape morphology (Fig.
283 4A and Fig. 4B, cross section A1).

284 Given the changes in channel morphology and flow characteristics, the following
285 division was established for the Madden Canyon - Channel:

286 1) The canyon-confined portion of the channel (not included in the turbidity current
287 modelling, see methods) initiates from the Madden Banks breach-point (-1500 m water
288 depth) to ~7 km downstream (-1800 m water depth) (Fig. 4A and Fig. 6A). A series of
289 crescentic bedforms stepping basinward, and the steepest gradients of up to 3° are found
290 in this section (Fig. 5).

291 2) The upper reaches of the channel (~35 km long) (cross section A2, Fig. 4B) comprise
292 an area with an erosional terrace on the left bank (looking downstream), characterized
293 by slide scars of mass failures from over steepened walls (Fig. 5), and a structurally
294 confined levee on the right bank (Fig. 5B). Sediment waves migrating in a SSW
295 direction perpendicular to the channel thalweg are present on the outer external levees
296 (cf. Hansen et al., 2015) and terraces have formed likely due to the inward collapse of
297 the external levees (inset in Fig. 5B). Knickpoints are observed in the canyon-confined
298 portion as the channel crosses the axis of a tectonic fold and at the beginning of the
299 upper reaches (Fig. 5B). The average channel height in the upper reaches is 380 m (Fig.
300 6B), the average channel width and area are 9.2 km and 1,899 km² respectively (Fig.
301 6C). Slope gradients in this section increase and decrease downstream, ranging from 2°
302 to 0.04° (Fig. 6 A).

303 3) The Madden Channel - Porangahau section (~10 km long) comprises the area where the
304 channel crosses the Porangahau Trough (cross-section A3, Fig. 4B). Sediment waves
305 are observed on the outer external levees of the right bank, migrating in a SSW direction

306 perpendicular to the channel thalweg. On the left bank, sediment waves are also
307 developed trending in a NNE direction. The levees have lower relief than those
308 observed upstream in the upper reaches. Scours are present downstream on the channel
309 floor (Fig. 5B). Channel heights (average channel height is 31 m), widths (average is
310 3.11 km) and area (57.7 km²) are lower compared to the channel upper reaches (Fig.
311 6B and Fig. 6B). Slope gradients are gentler, ranging from 0.02° to 0.04° (Fig. 6A).

312 4) The channel lower reaches (~20 km long) comprises the transect of the sub-basin
313 between the Porangahau Ridge and the start of the Akitio Trough (cross section A4 and
314 A5, Fig. 4B). In this area the channel is diverted by a growth ridge (cf. Clark and
315 Cartwright, 2009), changing briefly to a trough-axial flow direction and limiting the left
316 levee development (Fig. 5B). Sediment waves are developed in the outer right levee
317 where the channel is diverted by the ridge structure (Fig. 5B). Then, the channel
318 changes its course again, to resume a transverse orientation towards the Akitio Trough
319 (Fig. 5B). The average channel height in this section is 16 m, average width of 1.31 km
320 and area of 17.1 km² (Fig. 6B and Fig. 6C). The slope gradients in the lower reaches
321 also has recurring steep slopes (2°, 1.2°) followed by more gentle slopes as in the upper
322 reaches (Fig. 6A).

323 5) The Madden Channel terminates in the Akitio Trough which exhibits a uniform surface
324 (cross section A6, Fig. 4B) with local emplacement of mass failure deposits (Fig. 5B).
325 The lowest gradient of 0.01° is in this section (Fig. 6A).

326 **Modelled Turbidity Current Conditions in the Madden Channel**

327 The results from the turbidity current modelling in the Madden Channel are shown in Fig. 6.
328 The flow velocity and shear velocity in the upper reaches of the channel show the highest
329 values fluctuating from 8-12 m/s (Fig. 6D) and 0.4-0.6 m/s (Fig. 6F), respectively. A drop in

330 flow velocity to 3 m/s is observed at ~22 km decreasing gradually to 1.7 m/s in the last section
331 of the upper reaches, then to 0.7 m/s in the Porangahau section and finally there is an increase
332 up to 3.8 m/s in the lower reaches (Fig. 6D). The calculated velocities are within the ranges of
333 current velocities calculated from cable breaks and measured at submarine canyon systems
334 (0.02 m/s - 19 m/s) (Talling et al., 2013). *Fr* numbers from 2 to 1.2 downstream of the upper
335 reaches indicate super-critical flow conditions, which transition to sub-critical flow averaging
336 0.35. Flow deceleration and *Fr* values at this transition may indicate the development of an
337 internal hydraulic jump causing suspension and mixing near the bed which allows sediment
338 transport down-channel (Dorrell et al., 2016). The shear velocity decreases downstream to 0.15
339 m/s and 0.03 m/s in the upper reaches and Porangahau sections, respectively. Then, a slight
340 increase in flow velocity (up to 4 m/s) and shear velocity (up to 0.2 m/s) occurs at ~48 km
341 downstream in the lower reaches section. Supercritical flow conditions with *Fr* numbers from
342 1.9 to 1.6 prevail at the lower reaches as the channel passes through a steep area of structural
343 confinement.

344 The calculated bulk sediment concentrations exhibit less fluctuations in the upper
345 reaches and Porangahau sections with average values of 0.0074 v/v (0.7%) and 0.008 v/v
346 (0.8%) (Fig. 6G), respectively. A peak of 0.017 v/v (1.7%) occurs in the lower reaches at ~54
347 km downstream, the average sediment concentration in this section is of 0.009 v/v (0.9%). The
348 flow discharge and sediment discharge (Fig. 6H) display decreasing and increasing trends that
349 parallel those seen in the flow velocities and Froude number. The average flow discharge and
350 sediment discharge values in the upper reaches of the Madden Channel are 1.33×10^7 m³/s and
351 1.04×10^8 kg/s, respectively. In the Porangahau section, the average values reduce to 5.25×10^4
352 m³/s and 4.46×10^5 kg/s. The average flow discharge reduces to 5.14×10^4 m³/s in the lower
353 reaches but the sediment discharge depicts an increment to 5.53×10^5 kg/s, reflecting the higher

354 sediment concentrations shown in the same section in Fig. 6G. These calculated flow
355 conditions allow us to investigate for erosion, equilibrium or deposition along the channel
356 profile.

357 **Sediment Bypass Conditions in the Madden Channel**

358 *Well Sorted Flows*

359 The thresholds for equilibrium conditions of the modelled flows traversing each channel
360 section are shown from Fig. 8A to Fig. 8D. The contours represent the thresholds for well
361 sorted ($\sigma=0.5$), moderately sorted ($\sigma=1$) and poorly sorted ($\sigma=2$) sediment in suspension. For a
362 given slope value (S), the grain sizes (ϕ) above a given threshold represent sediment bypass
363 under erosional flow conditions and grain sizes below a given threshold represent deposition.
364 Note that the variation in particle size distribution (log-normal standard deviation), flow height
365 and down-slope channel gradient have an important effect on equilibrium.

366 The calculated flow of 380 m average height and 0.74% average sediment concentration
367 traversing the steep slopes (2° , 1.2°) in the upper reaches (Fig. 8A and Fig. 8E) can transport a
368 wider range of grain size classes (from very fine silt to medium sand) if they are suspended
369 within a well sorted flow ($\sigma =0.5$). However, for the same flow conditions and sorting, the
370 maximum grain size that can be maintained in suspension reduces to medium silt as the slope
371 gradients become gentler in the upper reaches (0.04° , 0.2°). As the flow reaches the Porangahau
372 section (Fig. 8B and Fig. 8E), the reduction in the average flow height to 31 m due to loss of
373 flow confinement in the Porangahau Trough, and gentle slopes of 0.04° causes a reduction of
374 the maximum grain size that can be transported in suspension to fine silt.

375 The steep slopes formed by the presence of the Porangahau Ridge at the start of the
376 channel lower reaches cause an increase in the driving force of the flow, consequently
377 increasing the flow velocity, shear velocity, bulk sediment concentration and discharge (Fig.

378 6). These changes in flow suggest that down-channel erosion might occur. Although the
379 sediment being eroded depends on the composition of the bed, we can estimate that the
380 maximum grain size class that can be transported in suspension is very fine sand, for the
381 calculated flow conditions, assuming a well sorted flow travelling through the lower reaches
382 (positions 6 and 7 in Fig. 8C and Fig. 8D, respectively). The channel terminates in the Akitio
383 Trough where flows become unconfined.

384 Although the flow conditions in the Akitio Trough could not be calculated in the
385 absence of channels, from the slope values in the trough (0.01°) it is estimated that most of the
386 grain size classes would become depositional except for the very fine silt (position 8, Fig. 8D).
387 Therefore, the maximum grain size that could be transported downstream by the modelled
388 flows, from the upper reaches of the channel into the Akitio Trough, assuming a well sorted
389 sediment in suspension, is very fine silt.

390 ***Poorly Sorted Flows***

391 The equilibrium thresholds for poorly sorted suspensions ($\sigma= 2$) occur at higher slope values
392 for all grain size classes compared to flows carrying well sorted suspensions ($\sigma= 0.5$). For the
393 same flow conditions, the maximum grain size class that can be suspended in the flow and
394 transported in the steep slopes (position 1, Fig. 8A and Fig. 8E) is very fine sand. However, it
395 reduces to coarse silt (position 2, Fig. 8A and Fig. 8E), fine silt (position 3, Fig. 8A and Fig.
396 8E) and very fine silt (position 4, Fig. 8A and Fig. 8E) as the slope gradient decreases
397 downstream.

398 As with the well sorted case, the equilibrium threshold for poorly sorted sediment in
399 suspension in the Porangahau Trough suggest that all grain size classes would start to become
400 depositional in this section (position 5, Fig. 8B), partly due to the low slope gradients and the

401 reduction in flow size. Slope gradients approximately above 0.1° would be required to keep the
402 very fine sand in suspension.

403 The increase in flow, shear velocities and sediment flux at the start of the channel lower
404 reaches suggest downstream erosion (Fig. 6). When poorly sorted material is eroded (assuming
405 this characterizes the composition of the substrate), we can estimate that the maximum grain
406 size that can be kept in suspension is medium silt throughout the lower reaches (positions 6 to
407 8, Fig. 8C and Fig. 8D). Any sand-size grains eroded would become depositional, limiting
408 further transport into the Akitio Trough (Fig. 8D).

409 **Omakere Channel Morphology**

410 A series of gullies incising the shelf at ~70 km offshore Hawke Bay constitute conduits that
411 feed the Omakere Channel (Fig. 2). The Omakere channel initiates as a trough axial channel in
412 the Omakere Trough at ~1500 m water depth. It evolves into a transverse channel traversing
413 troughs and eroding the Paoanui and Porangahau ridges and terminates in the Akitio Trough at
414 ~2270 m water depth (Fig. 4 and Fig. 7A).

415 The following division is established for the Omakere Channel given its changes in
416 morphology and flow characteristics:

417 1) The upper reaches of the Omakere Channel (~35 km long) initiate in the Omakere
418 Trough and is confined by thrust-faulted NE-SW oriented ridges (Fig. 4) which have
419 limited the development of levees. The channel exhibits a wide u-shape morphology in
420 the first ~25 km (cross-section B1, Fig 4C) which evolves to a box-shape at the end of
421 the upper reaches (cross-section B2, Fig.4C) likely due to down-channel tectonic
422 confinement. Therefore, the channel area decreases from 552 km² to 57 km².
423 Knickpoints developed in the most confined section (Fig. 5B) before the channel is
424 diverted into the upper reaches sub-basin. The average channel height in this section is

425 147 m, and average width of 3.96 km. There is a gentle slope of 0.09° followed by an
426 increase up to 3° in the knickpoint area.

427 2) The Omakere Channel upper reaches sub-basin section (~15 km long) starts as the
428 channel is diverted to a NW-SE direction by the Omakere Ridge (Fig. 4). The channel
429 floor is smooth in this area. The average channel height decreases to 48 m due to loss
430 of ridge-confinement as it enters the sub-basin and the average channel width increases
431 to 6.14 km giving a wider U-shape morphology with subtle levee development (cross-
432 section B3, Fig. 4C). The slope gradient in this area is gentler (0.12°) and less variable
433 compared to the upper reaches.

434 3) The Omakere Channel-Paoanui section (~21 km long) starts where the channel course
435 is diverted into the Paoanui Trough and subsequently cross-cuts the Paoanui Ridge. The
436 channel widens in the trough giving a U-shape (cross-section B4, Fig. 4C) which
437 narrows downslope as the channel height and ridge confinement increases (cross-
438 section B5, Fig. 4C). The average channel height and area are 190 m and 783 km^2
439 respectively. The slope gradient in this area increase with respect to the upper reaches
440 sub-basin portion (Fig. 7A).

441 4) The lower reaches and the Akitio sections of the Omakere Channel (~15 km long)
442 comprise the area where the channel cross-cuts the northern portion of the Porangahau
443 ridge and other downstream ridges to terminate in the northern portion of the Akitio
444 Trough. A series of knickpoints are observed at the start of the lower reaches where the
445 slope gradient is steep (2° , Fig. 7A). The channel exhibits a box-shape as the channel
446 widens due to loss of ridge confinement (cross-section B6, Fig. 4C) and subsequently
447 becomes unconfined in the Akitio Trough (cross-section B7, Fig. 4C) where the slope
448 decreases to 0.008° (Fig. 7A).

449 **Modelled Turbidity Current Conditions in the Omakere Channel**

450 The flow velocity in the Omakere Channel decreases from 3.2 m/s average in the upper reaches
451 to 1.8 m/s average velocity in the upper reaches sub-basin and increase downslope to 6.3 m/s
452 in the Paoanui Trough, then to 12 m/s average in the lower reaches and drops to 0.6 m/s in the
453 Akitio Trough (Fig. 7D). As in the Madden Channel, the calculated velocity values for the
454 Omakere Channel are within the ranges of current velocities observed at slope canyon systems
455 (0.02 m/s - 19 m/s) (Talling et al., 2013). The shear velocity exhibits an average value of 0.09
456 m/s at the upper reaches sub-basin and peaks at 0.8 m/s in the lower reaches, followed by a
457 decrease to 0.5 m/s as the flow reaches the Akitio Trough (Fig. 7D). The flow is predominantly
458 super-critical with phases of sub-critical flow at the start of the upper reaches and in the sub-
459 basin, where slope gradients are gentler (Fig. 7E).

460 The calculated bulk sediment concentration of the flow suggests more diluted flows at the
461 upper reaches with average values of 0.004 v/v (0.4%), which increase down-channel in the
462 Paoanui channel section and peaks at the lower reaches ~0.023 v/v (2.3%) (Fig. 7G). As with
463 the flow and shear velocities, the concentration drops in the Akitio Trough to ~0.003 v/v
464 (0.3%). The flow and sediment discharge initiate with an average of 1.51×10^6 m³/s and
465 1.31×10^7 kg/s respectively in the upper reaches and increase down-channel to 3.98×10^6 m³/s
466 and 7.87×10^7 kg/s (Fig. 7H).

467 **Sediment Bypass Conditions in the Omakere Channel**

468 *Well Sorted Flows*

469 The thresholds for equilibrium conditions of the modelled flows traversing each channel
470 section are shown from Fig. 9A to Fig. 9D (see Well Sorted Flows section of the Madden
471 Channel on how to interpret the plots). The flow conditions calculated at the upper reaches
472 section, where slope gradients are gentle, show that the maximum grain size that can be kept
473 in suspension is coarse silt (0.09°, position 1, Fig. 9A). However, the maximum grain size

474 increases downstream over the steep slopes at ~25-31 km (positions 2, and 3, Fig. 9A) where
475 the erosional flows can transport up to fine sand.

476 The increase in sediment concentration suggests an erosional flow down-channel from
477 the upper reaches sub-basin (Fig. 9B) to the lower reaches (Fig. 9D), nevertheless, the different
478 grain sizes that can be suspended within such erosional flows varies at each channel section
479 due to the differences in flow height and variations in the slope gradient. As flows become less
480 confined and smaller in the upper reaches sub-basin, the thresholds widen and the maximum
481 grain size that can be maintained in suspension is coarse silt (position 4, Fig. 9B). The slope
482 gradient increases and maintains above 0.1° down-channel, flows increase in height in the
483 Paoanui Trough (positions 5 and 6, Fig. 9C) and the lower reaches (position 7, Fig. 9D),
484 therefore, the maximum grain size that can be suspended in these sections is fine sand.

485 As in the Madden Channel, the slope gradients in the Akitio Trough promote the
486 deposition of most grain size classes (position 8, Fig. 9D). Under well sorted conditions, the
487 maximum grain size that can be bypassed through the entire channel length and be deposited
488 in the Akitio Trough is coarse silt.

489 ***Poorly Sorted Flows***

490 The thresholds for equilibrium flow conditions assuming poorly sorted sediment in suspension
491 ($\sigma = 2$) show that the maximum grain size class that can be transported through the upper
492 reaches of the Omakere Channel is coarse silt (Fig. 9A) which reduces to very fine silt in the
493 upper reaches sub-basin (Fig. 9B). Nevertheless, flows evolving downstream could transport
494 larger grain sizes through erosional or equilibrium flows as they increase in height and
495 sediment concentration. In the Paoanui Trough, the maximum grain size that can be kept in
496 suspension is coarse silt (Fig. 9C) which increases to fine sand over the steep slopes in the

497 lower reaches (position 7, Fig. 9D). All grain sizes are calculated to be deposited as flows enter
498 the Akitio Trough (position 8, Fig. 9D).

499 **DISCUSSION**

500 **Controls on Sediment Bypass and Implications for Sand Distribution**

501 Results here suggest that in order to achieve very fine sand bypass (or coarser grain sizes)
502 through the ECB channels into outboard troughs, shallow flows (< 50 m) require steeper
503 gradients than thicker flows (> 140 m) under well sorted and poorly sorted conditions (Table
504 1). Furthermore, poorly sorted flows require steeper gradients than well sorted flows to achieve
505 very fine sand (or coarser) bypass, given that wider particle size distributions promote vertical
506 density stratification, hence, the magnitude of the shear stress in the flow must increase to
507 maintain sediment in suspension and reach equilibrium (Dorrell et al., 2018). Therefore, the
508 changes in the flow height, grain size distribution and slope gradient have an effect on the
509 equilibrium thresholds for sediment bypass and on the distribution of sand in the channels (Fig.
510 10).

511 Assuming a well sorted suspension of very fine sand and coarse silt, an equilibrium
512 flow develops in the upper reaches of the Madden Channel, enhanced by thick flows and steep
513 gradients. It may then become depositional as when reaching the Porangahau Trough, an area
514 of low confinement (hence the flow height reduces) and gentle slope gradients (Fig. 10A).
515 Therefore, very fine sand bypass into the lower reaches and Akitio Trough is limited and flows
516 rely on erosion to entrain very fine sand into suspension, if available, to continue sediment
517 transport downstream and deposit into the Akitio Trough (Fig. 10C). Similarly, the very fine
518 sand deposition in the Omakere Channel is more likely to occur in the Omakere Trough (upper
519 reaches section of the channel) where, although flows are more confined than in the Porangahau
520 section, it also constitutes an area of low gradient (Fig. 10A). As in the Madden Channel, flows

521 evolving downstream in the Omakere Channel must erode very fine sand in order to deposit it
522 in the upper reaches sub-basin and Akitio Trough (Fig. 10C).

523 Assuming a poorly sorted suspension of very fine sand and very fine silt to medium
524 silt, the wider grain size distribution promotes depositional flows with very fine sand deposition
525 at the uppermost reaches of the Madden Channel and the Omakere Channel, and silt bypass
526 down-channel (Fig. 10B). If erosional flows entrain very fine sand, the smaller magnitude of
527 the flows combined with a poorly sorted suspension promote very fine sand deposition in the
528 Madden and Omakere Channel sections as shown in Fig. 10D. Therefore, only the finer silt
529 fraction bypasses to the Akitio Trough.

530 The results are compared to the RMS amplitude map of the seafloor (Fig. 10E), where
531 high RMS amplitude responses indicative of high impedance contrast, are interpreted as sand-
532 prone intervals, whereas low RMS amplitude responses indicate more homogeneous and finer-
533 grained sediments. The high RMS amplitude patterns indicate sandier deposits along the
534 Madden and Omakere channel thalweg (Fig. 10E), which are similar to the patterns projected
535 for poorly sorted flows that erode sand down-channel (Fig. 10D). In the Madden Channel,
536 higher RMS values along the channel thalweg are observed at the start of the upper reaches
537 (Fig. 10F), and in the lower reaches (Fig. 10G) whereas low RMS values predominate in the
538 Porangahau and Akitio Troughs. Furthermore, grain size analysis of drop cores in the Madden
539 system have suggested poorly sorted deposits from very fine silt to fine sand (McKeown,
540 2018). Drop core M1 located in the upper reaches (Fig. 10E and 10F) was records a series of
541 poorly sorted sediments, composed of interbedded very fine to coarse silt, silty fine sands and
542 very fine sand beds (McKeown, 2018). Drop core M2 located in the Porangahau Trough (Fig.
543 10E and 10F) also consists of poorly sorted sediments, but only very fine to coarse silt beds
544 were found here (McKeown, 2018). Sediments from drop core M3 located in the Akitio Trough
545 (Fig. 10E and 10G) are poorly sorted, very fine to medium silt and sandy silt beds. Sediments

546 of M1 and M2 are in agreement with the deposits predicted from the model in the scenario
547 shown in Fig. 10D and the RMS amplitude map, suggesting these deposits might have been
548 the product of poorly sorted depositional flows with flow depths and bulk sediment
549 concentrations comparable with the modelling. Samples were not available from the Omakere
550 Channel, however RMS amplitude mapping can be used to examine our findings. Higher
551 amplitude values are located in the upper reaches (Fig. 10H), the Paoanui Trough (Fig. 10I)
552 and the lower reaches (Fig. 10I), whereas lower amplitude values predominate in the upper
553 reaches sub-basin and northern Akitio Trough, which shows good agreement with the patterns
554 projected for poorly sorted deposits in Fig. 10D. Differences in the distribution of sand in the
555 system might be due to other processes such as contour currents or bedload sediment transport
556 which are not considered in the modelling.

557 The results show that flows can change from sand bypass to deposition depending on
558 the characteristics of the turbidity current and that the bypass slope is not unique. Furthermore,
559 the thresholds derived from this study might help constrain the conditions required to develop
560 equilibrium flows over submarine slopes more generally. If a system has an up-dip slope of
561 0.25° , it would constitute a bypass slope for very fine sand when flows are approximately
562 thicker than 140 m (Table 1) whereas, for the same slope gradient and flow thickness, sand
563 would be deposited if poorly sorted conditions prevail as they would require slope gradients
564 above 0.8° to sustain bypass (Table 1). Thick flows (>140 m) with well sorted sediment in
565 suspension might aid the development of detached systems and possibly upslope pinch-outs,
566 when up-dip slopes are above 0.1° . Poorly sorted flows transporting sand might be depositional
567 throughout the feeder channel, hence, either promoting the formation of attached systems or
568 starved basins where only the finest grain sizes can be bypassed to the basin as observed in the
569 ECB (Fig. 10).

570 The thresholds are dependent on the assumptions made on the modelling, therefore the
571 thresholds for erosion, equilibrium and deposition might change for a different set of flow
572 height and sediment concentration. However, the thresholds derived from these flow conditions
573 might help guide and constrain interpretations of zones associated with equilibrium or
574 depositional flows in study areas where the slope gradient, the grain size and grain size
575 distribution are known such as in surface or subsurface deep-marine systems. Furthermore this
576 approach might allow us to predict areas of slope accommodation that may be starved or well
577 supplied with sands and thus help constrain and predict sediment distributions.

578 **CONCLUSIONS**

579 We calculated equilibrium thresholds to estimate sediment bypass in turbidity currents
580 traversing the Madden and Omakere channels of the East Coast Basin using the cross-flow
581 equation and Flow-Power Flux-Balance type model. This methodology allowed us to
582 determine that sediment bypass in the channels is controlled by changes in the flow height
583 (calculated from the channel morphology), by sediment sorting in the flow and by changes in
584 the seafloor gradient. On this basis, we derived maps for the potential distribution of very fine
585 sand, where deposits estimated from poorly sorted flows show good agreement with the
586 observations from the RMS amplitude map of the seafloor. The conditions required to achieve
587 sand bypass derived from this study might serve to screen other flow pathways of the East
588 Coast Basin and could be applied to other margins.

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925 **FIGURE CAPTIONS**

926 Figure 1. Process terminology used in this study for a channelized turbidity current.

927 Figure 2. Study area. (A) The East Coast Basin is located on and offshore of the North Island,
928 New Zealand and is limited to the east by the Hikurangi Trough and to the west by the axial
929 ranges. (B) Offshore bathymetry map and bathymetric contours (500 m) of the East Coast
930 Basin (courtesy of the National Institute of Water and Atmospheric Research, New Zealand).
931 (C) Regional cross-section across the Hikurangi subduction complex (after Nicol et al.,
932 2007).

933 Figure 3. Diagram showing the morphometric parameters measured at channel bends, levee-
934 bound channel cross-sections and erosional channel cross-sections of the Madden and
935 Omakere channels.

936 Figure 4. (A) The bathymetry map and bathymetric contours (500 m) courtesy of NIWA,
937 show the Madden Channel and Omakere Channel systems. The channel thalweg line styles
938 differentiate the sub-channel sections described in this study. (B) Selected cross-sectional
939 profiles (with orientation looking downstream) of the Madden Channel highlighting the
940 down-channel evolution from the canyon-confined portion (transect A1) to the Akitio Trough
941 (transect A6). Transect location shown in A. (C) Selected cross-sectional profiles of the

942 Omakere Channel from the channel upper reaches (transect B1) to the northern Akitio
943 Trough (transect B7).

944 Figure 5. (A) Seafloor dip map of the Madden and Omakere channels derived from 3D
945 seismic data provided by WesternGeco Multiclient. (B) Interpreted seafloor dip map
946 showing morphological features of the channels resulted from tectonic and sedimentary
947 processes. The interpretation of the crests of thrust-faulted bathymetric ridges is derived from
948 Barnes et al. (2010).

949 Figure 6. Calculated dimensions and turbidity current conditions in the Madden Channel
950 from the upper reaches through the lower reaches section (sub-sections of the channel are
951 shown at the top of the plots). (A) Channel thalweg depth profile (solid line) and slope
952 gradient profile (dotted line) with values presented in degrees. (B) Average channel height
953 profile (solid line) and cross-channel slope values (triangles). (C) Average channel width
954 (solid line) and channel area (crosses). (D) Downstream turbidity current velocity profile. (E)
955 Froude number profile, dashed line at 1 indicates the threshold between sub-critical and
956 super-critical flow. Downstream (F) shear velocity, (G) bulk sediment concentration, (H)
957 flow discharge and sediment discharge profiles.

958 Figure 7. Calculated dimensions and turbidity current conditions in the Omakere Channel
959 from the upper reaches through the Akitio Trough section (sub-sections of the channel are
960 shown at the top of the plots). (A) Channel thalweg depth profile (solid line) and slope
961 gradient profile (dotted line) with values presented in degrees. (B) Average channel height
962 profile (solid line) and cross-channel slope values (triangles). (C) Average channel width
963 (solid line) and channel area (crosses). (D) Downstream turbidity current velocity profile. (E)
964 Froude number profile, dashed line at 1 indicates the threshold between sub-critical and

965 super-critical flow. Downstream (F) shear velocity, (G) bulk sediment concentration, (H)
966 flow discharge and sediment discharge profiles.

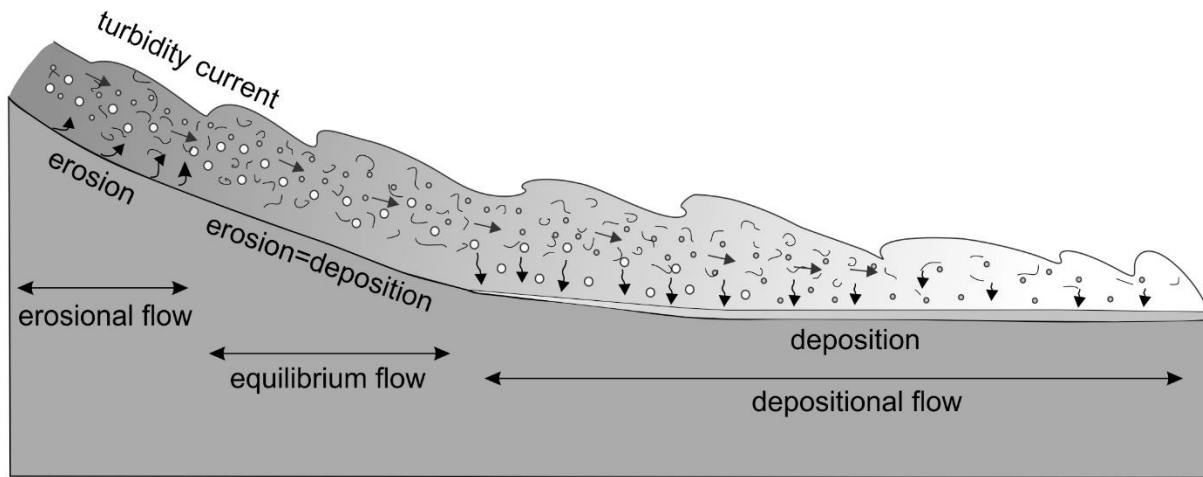
967 Figure 8. Plots showing the thresholds of equilibrium flow for the modelled turbidity current
968 traversing each section of the Madden Channel. The thresholds vary from well sorted ($\sigma=0.5$)
969 to poorly sorted ($\sigma=2$) sediment in suspension. For a given slope value (S), the grain sizes (ϕ)
970 above a given threshold represent sediment bypass under erosional flow conditions and grain
971 sizes below a given threshold represent deposition. The average flow height, \bar{H} , and average
972 bulk sediment concentration, C , used for the calculation of the thresholds are shown in each
973 plot for the (A) channel upper reaches (B) Porangahau Trough, (C) channel lower reaches
974 and (D) channel lower reaches and Akitio Trough. (E) Channel thalweg depth profile. The
975 numbers in black squares show the positions of the slope gradient values plotted in A-D.

976 Figure 9. Plots showing the thresholds of equilibrium flow for the modelled turbidity current
977 traversing the Omakere Channel. The thresholds vary from well sorted ($\sigma=0.5$) to poorly sorted
978 ($\sigma=2$) sediment in suspension. For a given slope value (S), the grain sizes (ϕ) above a given
979 threshold represent sediment bypass under erosional flow conditions and grain sizes below a
980 given threshold represent deposition. The average flow height, \bar{H} , and average bulk sediment
981 concentration, C , used for the calculation of the thresholds are shown in each plot for the (A)
982 channel upper reaches (B) upper reaches sub-basin, (C) Paoanui Trough and (D) channel lower
983 reaches and Akitio Trough. (E) Channel thalweg depth profile. The numbers in black squares
984 show the positions of the slope gradient values plotted in A-D.

985 Figure 10. (A-D) Potential sand distribution maps of the Madden and Omakere Channel
986 systems. The maps show the deposition from (A) a well sorted turbidity current transporting
987 very fine sand; (B) a poorly sorted turbidity current transporting very fine sand and very fine
988 sand to medium silt; (C) a well sorted turbidity current transporting very fine sand and coarse

989 silt, assuming downstream erosion of sand; (D) a poorly sorted turbidity current transporting
 990 very fine sand and very fine sand to medium silt assuming downstream erosion of sand. (E)
 991 Root Mean Square (RMS) amplitude map of the seafloor with drop core locations M1-M3.
 992 High RMS amplitude values are interpreted to represent deposits with higher sand content.
 993 Low RMS amplitude responses are interpreted to represent soft homogeneous finer-grained
 994 deposits. Enlarged areas showing the RMS amplitude distribution of the (F) Madden Channel
 995 upper reaches and Porangahau sections, (G) Madden Channel lower reaches, (H) Omakere
 996 upper reaches and (I) Omakere Channel Paoanui and lower reaches sections.

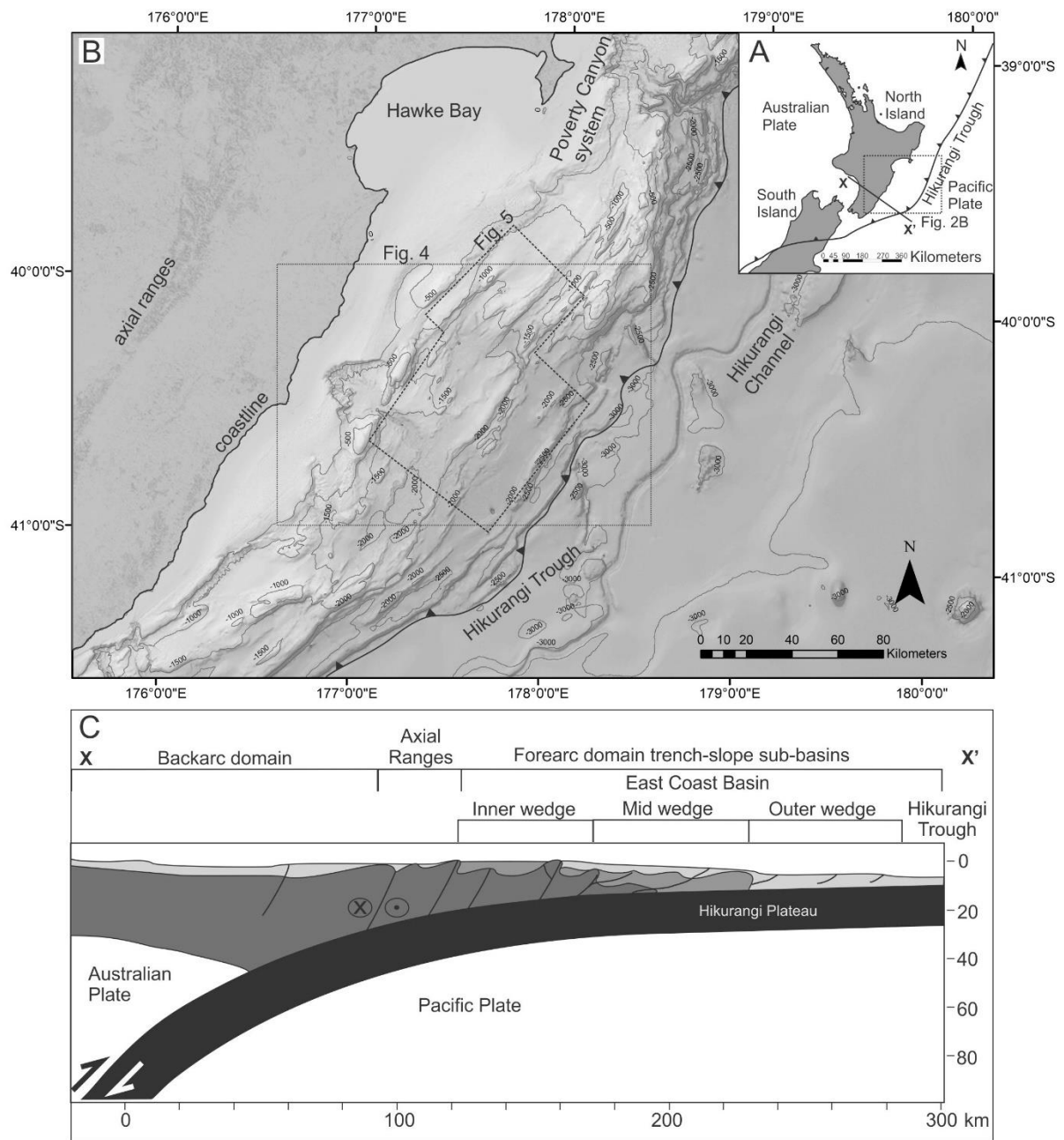
997 **FIGURE 1**



Process term	Definition
Erosional flow	A flow that erodes the substrate and entrains sediment into suspension, which might result in the development of erosional features.
Equilibrium flow	A flow where erosion is balanced with deposition, and transports its complete suspension load beyond the point of observation.
Depositional flow	A flow that deposits part of its suspension load and transports the remainder load downstream.

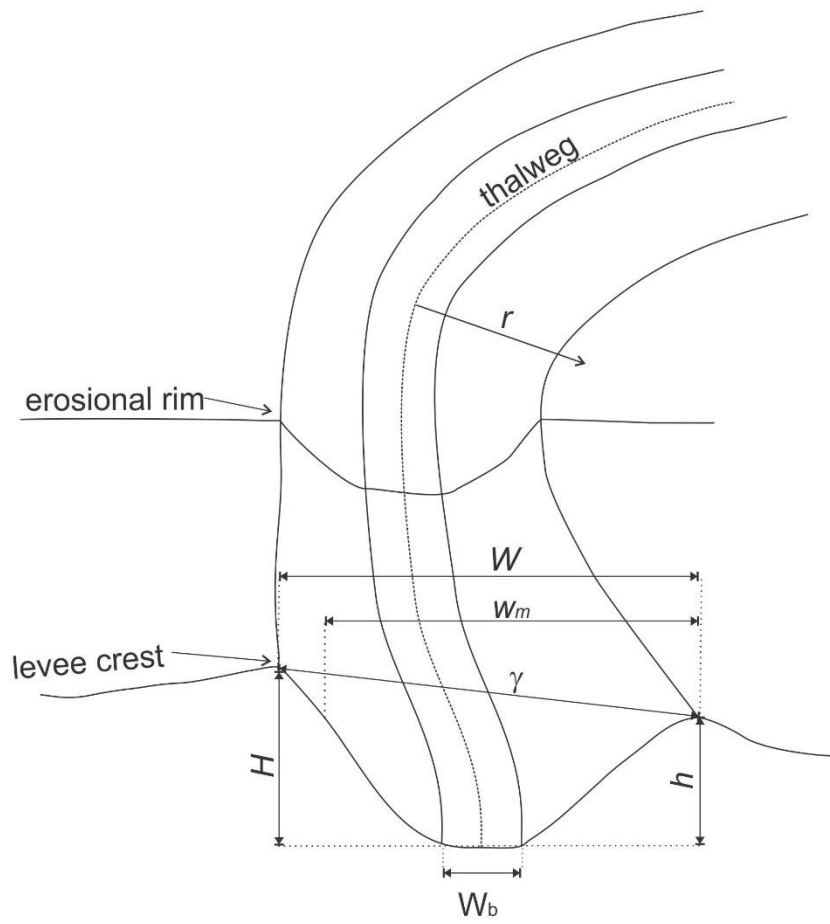
998

999 **FIGURE 2**



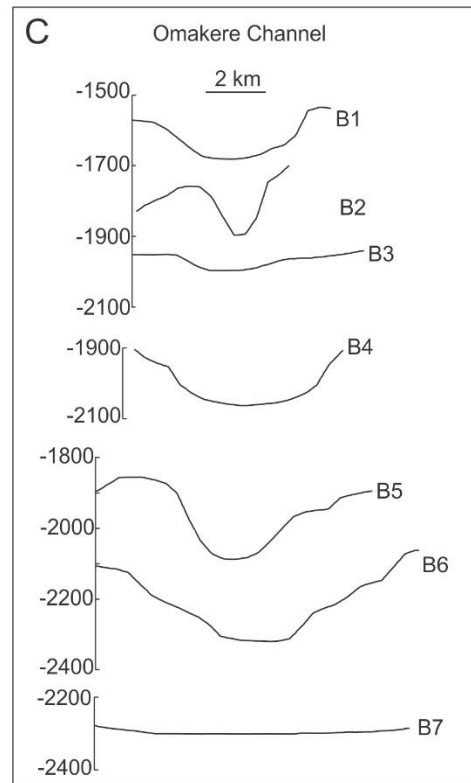
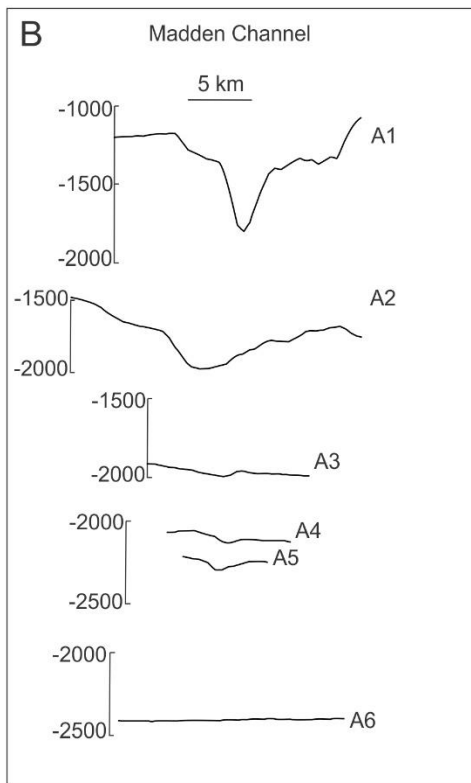
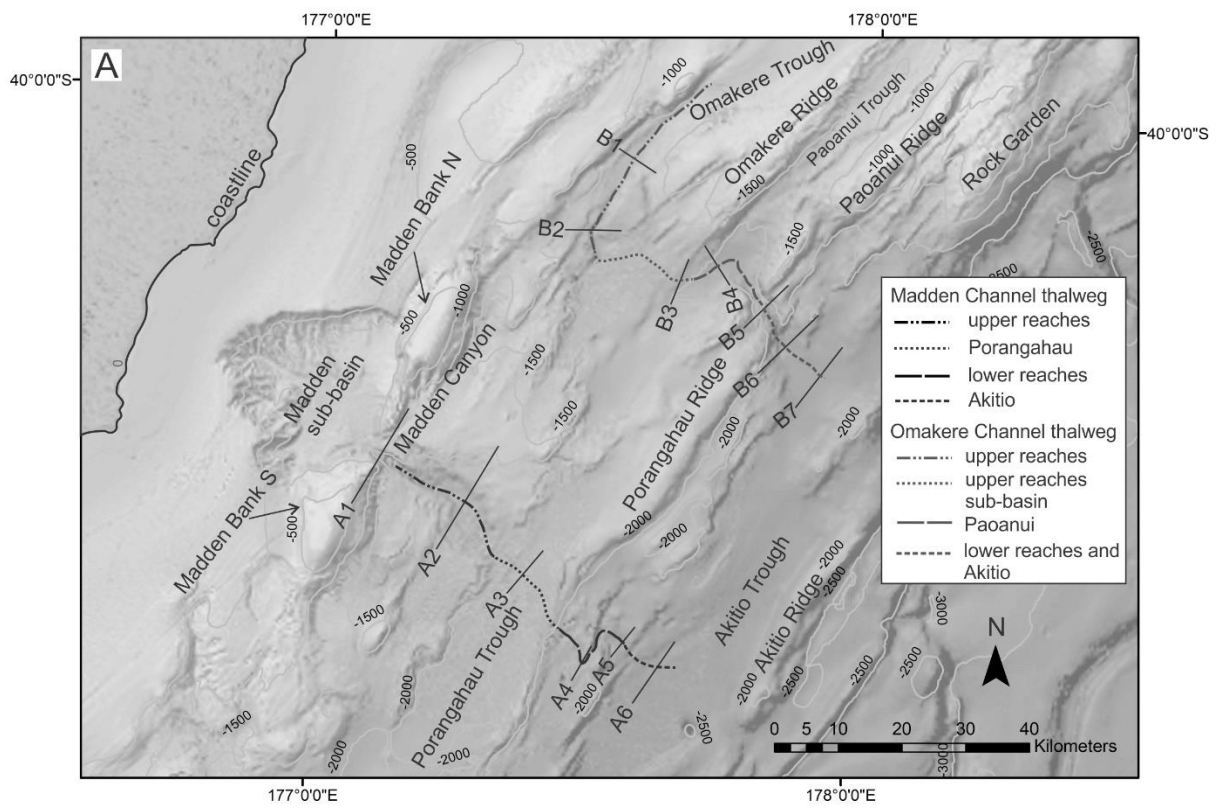
1000

1001 **FIGURE 3**



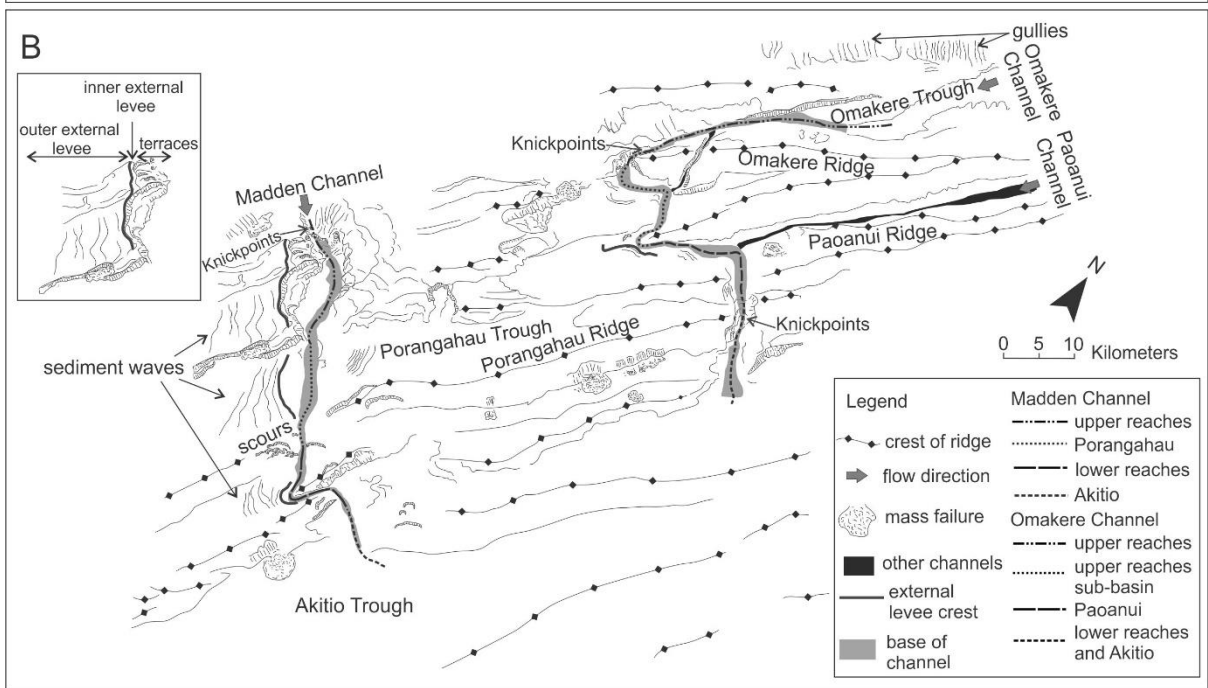
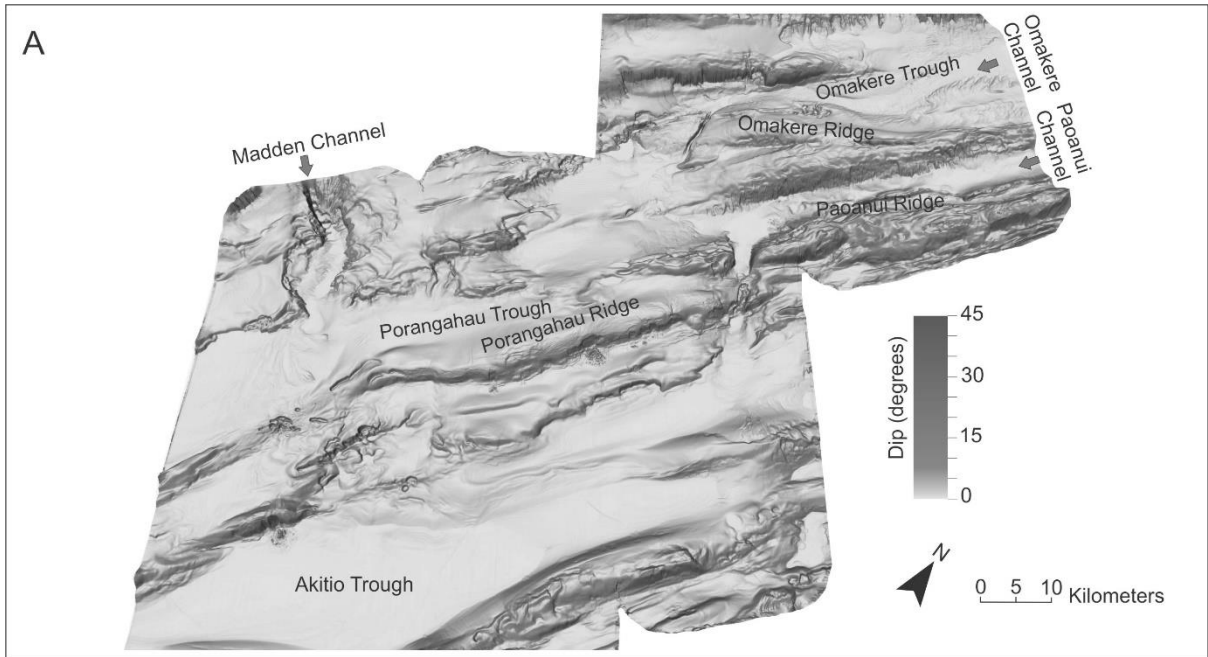
1002

1003 **FIGURE 4**



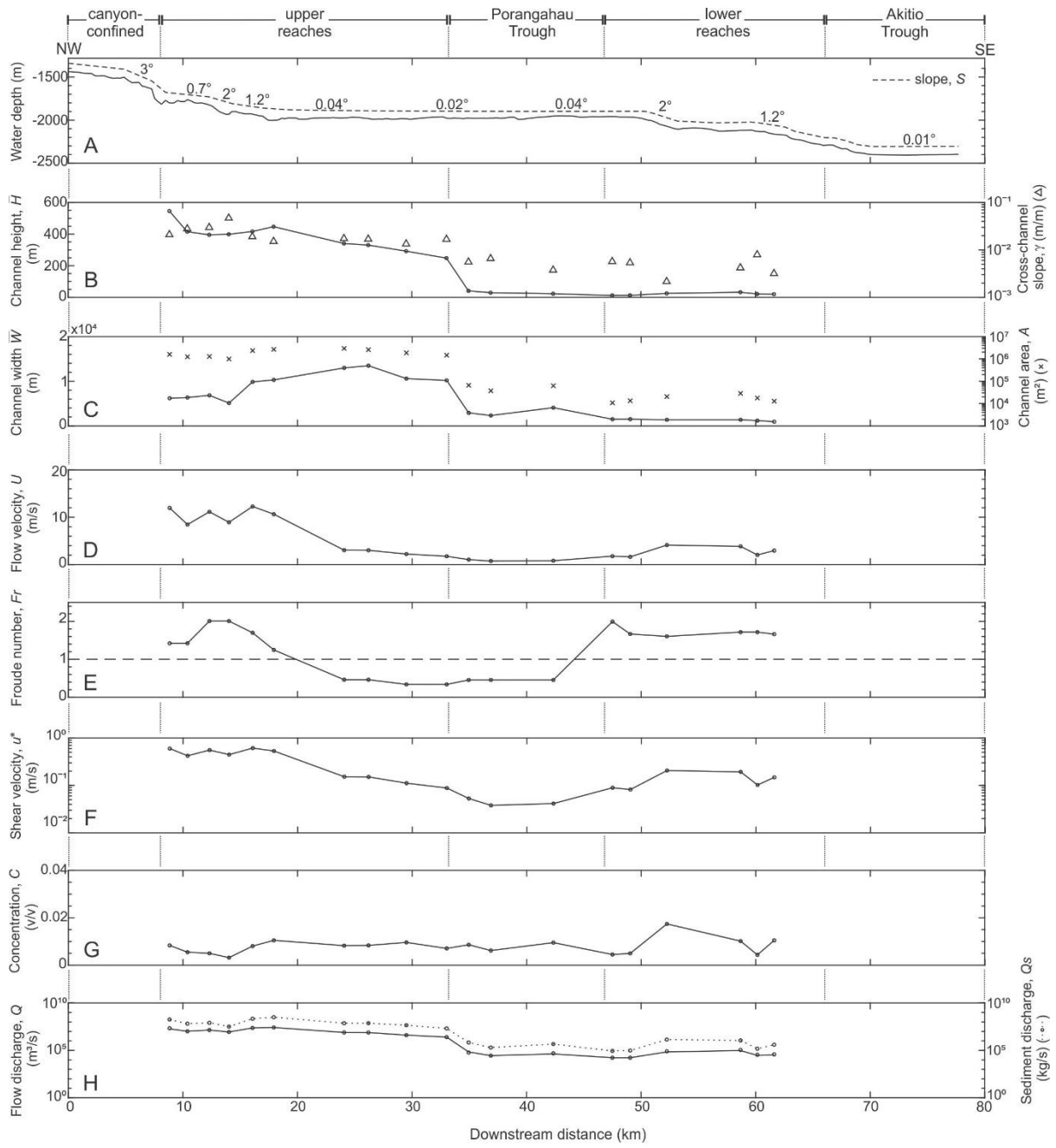
1004

1005 **FIGURE 5**



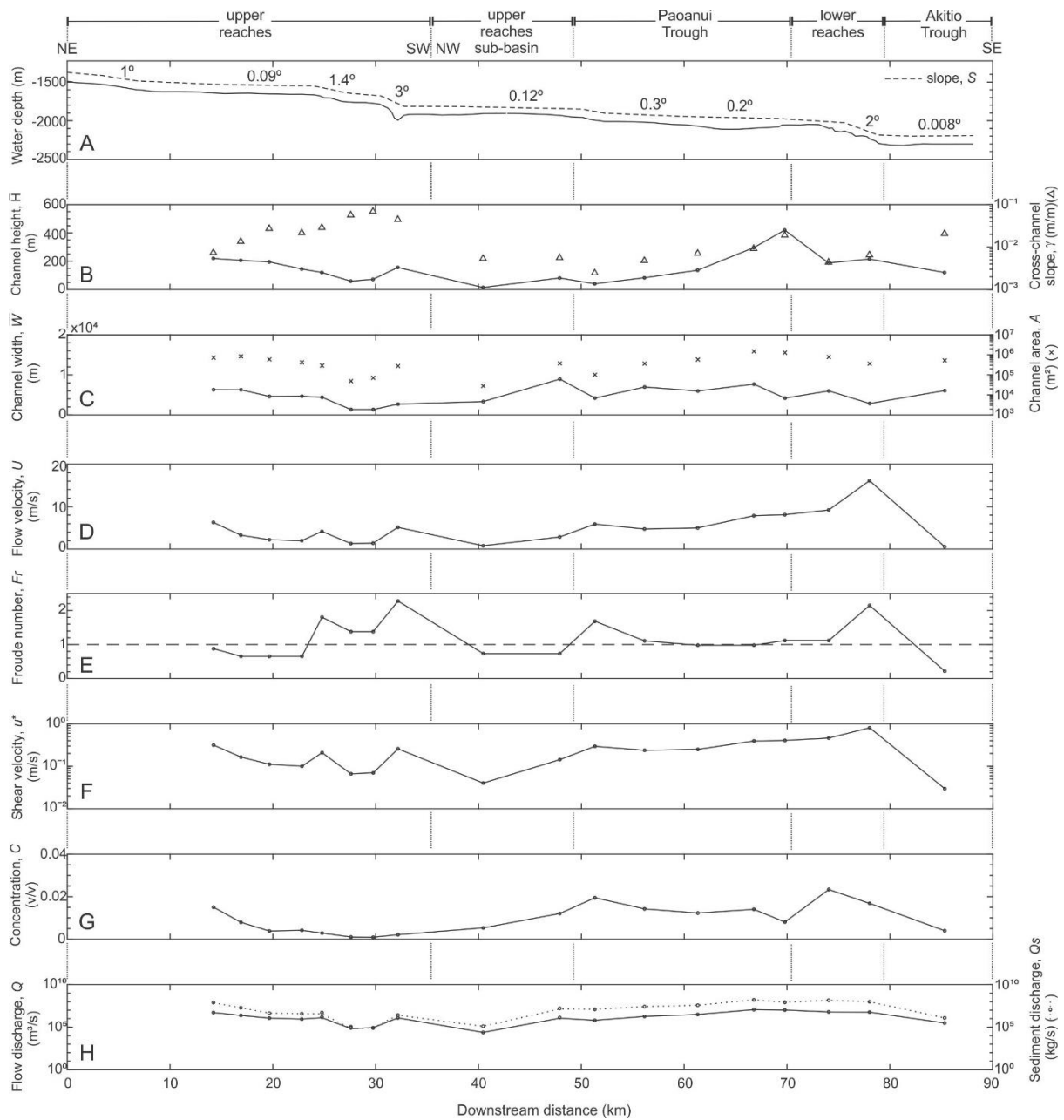
1006

1007 **FIGURE 6**



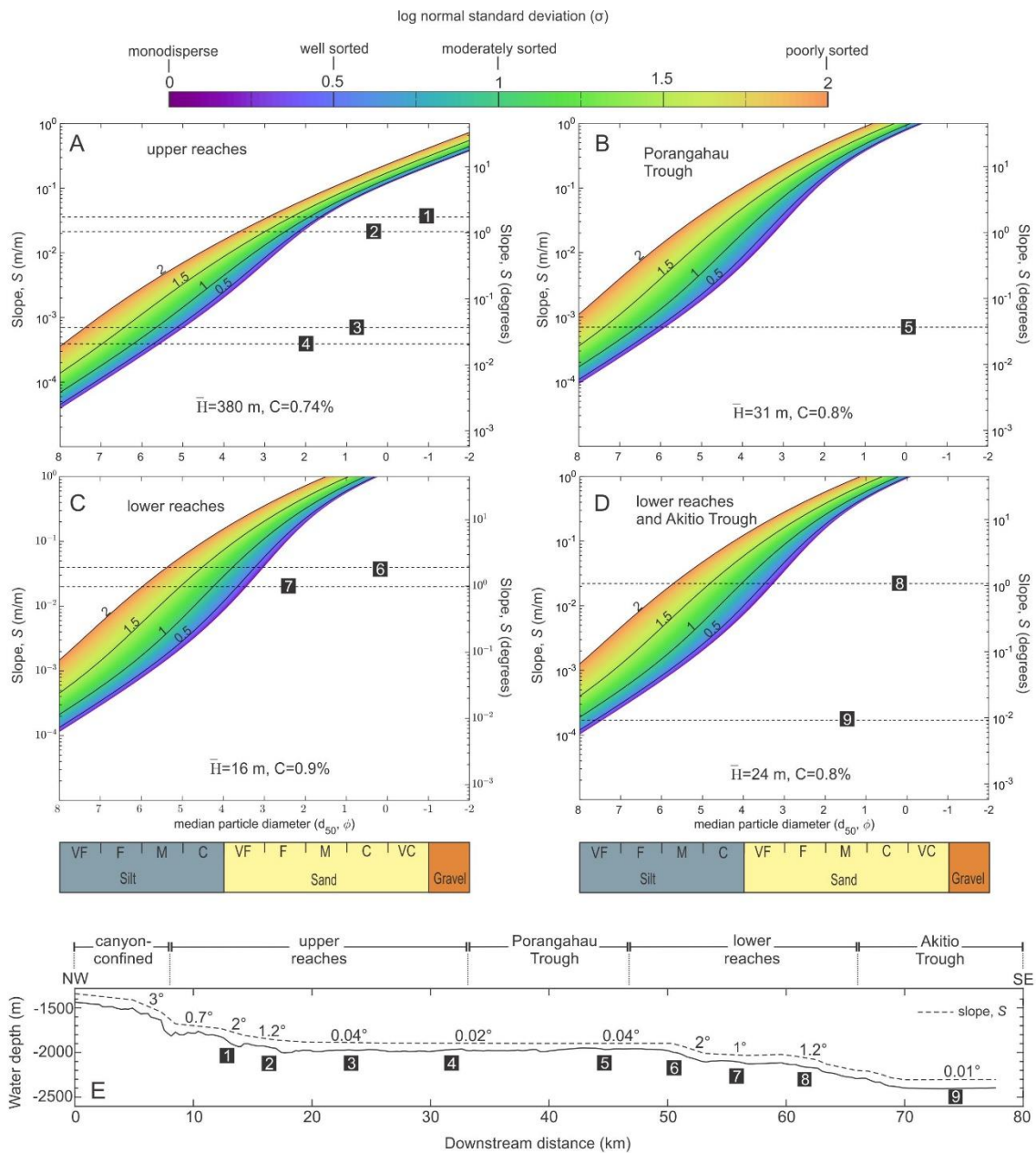
1008

1009 **FIGURE 7**



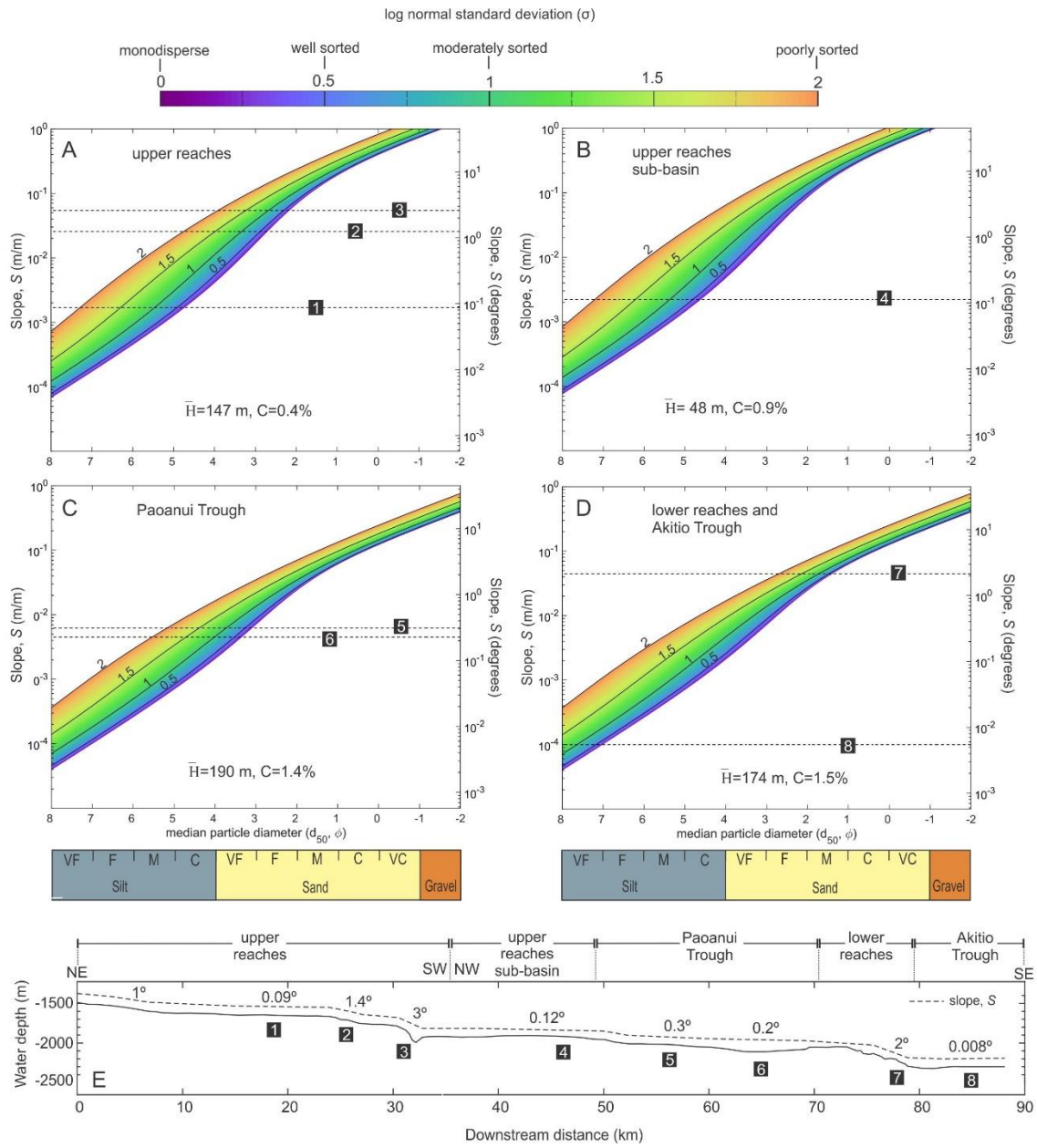
1010

1011 **FIGURE 8**



1012

1013 **FIGURE 9**



1014

1015 **FIGURE 10**

