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## Supplementary Information

## Porosity Percentage Calculation

The percentage of porosity in the samples was calculated using the $\mu \mathrm{CT}$ data. The total volume of the track above the substrate, with all pores filled, was measured. It is not possible to determine the volume of remelting into the substrate from the $\mu \mathrm{CT}$ scan, however a basic estimation has been calculated using a boundary box at the widest parts of the track, and a depth into the substrate at the deepest keyhole for each condition. This volume plus the track volume (above the substrate) were used as the estimation for total track volume. The volume porosity was then calculated as a percentage of this. This is a slight overestimation of total track volume, and thus underestimation of porosity percentage. The error was calculated by estimating a triangular section through the substrate, using the width of the boundary box and depth of the keyhole pores. This gave an error of $2-6 \%$ for the porosity percentage calculation.

## Supplementary Table 1

Linear energy density values from previous work and the present study. Hatch spacing has not been included in this calculation ( $L E D=\mathrm{P} / v t$ ) for direct comparison to the energy densities presented here.

| Laser Power, P <br> $\mathbf{( W )}$ | Scan Speed, <br> $\boldsymbol{v}$ <br> $\left(\mathbf{m m ~ s}^{-1}\right)$ | Powder layer <br> thickness, $\boldsymbol{t}$ <br> $(\mathbf{m m})$ | LED ( $\left.\mathbf{J ~ m m}^{-2}\right)$ | Reference |
| :---: | :---: | :---: | :---: | :---: |
| $175-400$ | $500-1100$ | 0.03 | 12 | $[1]$ |
| 275 | 750 | 0.05 | 7.3 | $[2]$ |
| 400 | $50-125$ | 0.25 | $13-32$ | $[3]$ |
| 100 | 560 | 0.04 | 4.5 | $[4]$ |
| $100-400$ | $400-1600$ | 0.05 | $5-10$ | $[5]$ |
| 42 | 200 | 0.03 | 7 | $[6]$ |
| 200 | $200-400$ | 0.1 | $5-10($ nominal $)$ <br> $3-9$ (adjusted) | Present <br> work |

Supplementary Figure 1: Graphs showing the change in layer 5 deposited powder thickness in (a) Sample A, (b) Sample C, and (c) Sample D.


Supplementary Figure 2: (a) Background corrected but unedited radiograph of Sample A layer 3 final morphology, with red highlighted region correlating to ( $b-\mathrm{d}$ ) in Figure 3. (b) Background, GC plate surface, and excess powder removed from the image to highlight surface shape. A section on this image was taken for Figure 3c.


Supplementary Figure 3: SEM image of Sample D, showing partially melted powder particles adhered to the track surface.


Supplementary Video 1: Sample A, Layer 1, showing ca. $50 \mu \mathrm{~m}$ build height, and laser keyhole ca. $150 \mu \mathrm{~m}$ below the powder surface.

Supplementary Video 2: Sample A, Layer 5, showing uneven track surface, and a large amount of powder and particle spatter.

Supplementary Video 3: Sample C, Layer 5, showing prominent surface undulations.
Supplementary Video 4: Sample D, Layer 5, showing uneven track surface and lack of fusion porosity between layers.

Supplementary Video 5: Sample B, Layer 5, showing large surface undulations and spatter, with a deep keyhole up to ca. $250 \mu \mathrm{~m}$ below the powder surface.

Supplementary Video 6: Sample B, Layer 5, background subtracted image to highlight powder layer thickness (dark grey), keyhole, and track undulations.

Supplementary Video 7: Sample C, Layer 1, showing ca. $100 \mu \mathrm{~m}$ track height at the ends of the track, with ca. $30 \mu \mathrm{~m}$ track height in the centre of the track. Powder spatter and porosity formation is visible. The keyhole depth was ca. 215
$\mu \mathrm{m}$ below the powder surface.
Supplementary Video 8: Sample C, Layer 1, cropped to ROI for keyhole porosity formation in the substrate.

Supplementary Video 9: Sample C, Layer 1, cropped to ROI, and locally averaged to reveal changing keyhole shape and size. Averaging method removes visible pores once they become stationary within the track.

Supplementary Video 10: Sample C Layer 2, full track, showing slight variations in track height and deep laser penetration, remelting all of layer 1.

Supplementary Video 11: Sample C, Layer 2, cropped to ROI for keyhole porosity formation in the substrate.

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# In situ radiographic and ex situ tomographic analysis of pore interactions during multilayer builds in laser powder bed fusion 

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#### Abstract

Porosity and high surface roughness can be detrimental to the mechanical performance of laser powder bed fusion (LPBF) additive manufactured components, potentially resulting in reduced component life. However, the link between powder layer thickness on pore formation and surface undulations in the LPBF parts remains unclear. In this paper, the influence of processing parameters on Ti-6AI-4V additive manufactured thin-wall components are investigated for multilayer builds, using a custom-built process replicator and in situ highspeed synchrotron X-ray imaging. In addition to the formation of initial keyhole pores, the results reveal three pore phenomena in multilayer builds resulting from keyhole melting: (i) healing of the previous layers pores via liquid filling during remelting; (ii) insufficient laser penetration depth to remelt and heal pores; and (iii) pores formed by keyholing which merge with existing pores, increasing the pore size. The results also show that the variation of powder layer thickness influences which pore formation mechanisms take place in multilayer builds. High-resolution X-ray computed tomography images reveal that clusters of pores form at the ends of tracks and when variations in the layer thickness and melt flow cause irregular


remelting and track height undulations. Extreme variations in height were found to lead to lack of fusion pores in the trough regions. It is hypothesised that the end of track pores were augmented by soluble gas which is partitioned into the melt pool and swept to track ends, supersaturating during end of track solidification and diffusing into pores increasing their size.

Keywords: Additive manufacturing, laser powder bed fusion, in situ X-ray imaging, Ti-6Al-4V, porosity

## 1. Introduction

Laser powder bed fusion (LPBF), also known as selective laser melting (SLM), is a rapidly evolving area of additive manufacturing (AM) technology using a layer by layer fabrication of 3D components from powder materials [1]. LPBF spreads layers of powder, which are then locally melted by a focussed laser beam according to a computer-generated programme. This process repeats until a full 3D part is produced [2]. The process is capable of producing components with complex geometries that cannot easily be made by conventional processing routes, e.g. casting [3]. The physics behind laser-powder interactions, e.g. the laser coupling to the metal surface and the melt pool dynamics, is very complex [4,5] and is related to both the material and process parameters.

Correlations between LPBF processing parameters and final part quality have been investigated through experimental $[6,7]$ and modelling methods $[8,9$. A number of studies have also investigated the transition from conduction melting to keyhole melting [10,11] as a function of process parameters. The volumetric energy density (VED) is defined by: $V E D=P / v h t$ [12], where $P$ is the laser power, $v$ is the laser scan speed, $t$ is the powder layer thickness, and $h$ is the hatch spacing. It is known that high energy density (i.e. high laser power coupled with low scanning speed) results in deep laser penetration (known as the keyhole mode) and often an increase in laser-induced features and spatter [13,14]. Very low energy density has been shown to produce wide, shallow melt pools, leading to discontinuous
tracks [15], interlayer porosity [16], and a heterogenous microstructure [17]. The microstructural evolution [18] and mechanical properties [19] of Ti-alloys produced by LPBF have also been investigated as a function of process parameters. However, most of the microstructural studies have been ex situ and do not reveal the underlying mechanisms controlling the development of the observed features.

A key area of interest in LPBF is the formation of internal porosity in the solidified structures. Lack of fusion [20], gas pores [21] and keyhole pores [10] are all known to occur in LPBF components. The former two pore types are also frequently found in electron beam melting AM $[22,23]$. Lack of fusion pores are often flat and elongated, arising from low energy density conditions where the prior track is not remelted. Gas pores are small and spherical and, for Ti-based alloys, may be due to the high solubility of hydrogen in the molten metal. Hydrogen gas can be absorbed from water vapour or contamination associated with the powder or the environmental chamber [24]. During cooling, the hydrogen solubility rapidly decreases and gas pores can nucleate in the melt [25,26]. The rapid cooling rates do not allow for diffusion and pore growth before solidification [27-29] and pores have been observed to be $10-300 \mu \mathrm{~m}$ in diameter. Pre-existing gas pores in the powder particles, arising from the atomisation manufacturing process, can also be a source of porosity in LPBF tracks. These pores can be entrained into the melt pool and can coalesce into larger pores [30,31].

The keyhole melting regime [32], although producing efficient energy transfer, can lead to the formation of keyhole pores [33] when a metal vapour/gas-filled bubble becomes trapped due to melt pool oscillations, unstable keyhole walls [34], or other complex hydrodynamic behaviour [35]. These bubbles can be composed of trapped inert chamber gas [4], metal vapour [10], or a combination of both. They are usually located near the bottom of the keyhole because rapid solidification prevents them from rising to the top of the melt [33,36].

A lack of understanding of the complex laser-matter interactions present during LPBF, including the formation of a dynamic melt pool, spatter (powder and droplet), metal vapour, plasma, and irregular powder entrainment [37], is hindering a more widespread uptake of LPBF technologies and thus further investigations are needed to help resolve this. Much of
the porosity analysis work has been carried out by ex situ studies, supported by computational models, to investigate formation hypotheses. In situ experiments with synchrotron X-ray radiography can provide critical information to substantiate these models and theories [38] by observing process phenomena such as spatter [39], melt pool flow [37,40], melt pool size [34,41], keyhole melting [36,42] and porosity formation [32,43-45].

In situ laser melting experiments on a solid substrate have determined a relationship between the laser power and scan velocity, and the transition from conduction to keyhole mode melting regimes, causing a change in the shape of the melt pool and keyhole [34]. In the case of LPBF, laser beam spot size and powder layer thickness can also affect the energy density transition to keyhole melting [10]. Deep, narrow cavities were described as the unstable keyhole zone, which led to the formation of keyhole pores [34]; the depth of these pores increased as energy density increased [42]. High-speed X-ray imaging has been used to observe the formation of these keyhole pores in situ during laser scanning [32,42,43,46].

Pore movement within the melt pool has also been observed during LPBF $[37,47]$ by in situ studies. Pores were shown to be swept with the Marangoni flow of the molten pool [37]. Pore shrinkage [42] and spheroidization [32] during solidification are other phenomena observed via X-ray imaging. Furthermore, oxidised powder has also been shown to significantly increase the internal porosity content [40], which is a key concern for reactive metals such as titanium and aluminium. These in situ studies have also been supplemented with porosity analysis via micro-computed tomography $(\mu \mathrm{CT})[16,30,48-50]$.

However, to date, most in situ synchrotron studies of the melt pool and pore formation in LPBF have involved the melting of a single layer of material. Experiments have been carried out on: a range of materials in overhang (melting onto powder) conditions [37,40,51]; a bare substrate without powder [34]; and on a substrate with a single layer of powder [32,36,41-44]. Since LPBF components are formed by many layers deposited on top of one another, an understanding of the multilayer process is essential. In the present study, we characterise multilayer builds in situ using high-speed synchrotron X-ray radiography, capturing the rapid dynamics of laser remelting, layer cohesion (or lack thereof), and changes in pore formation.

To conduct this work, a laser AM in situ and operando process replicator (ISOPR) [37] was developed to build multilayer tracks in situ during LPBF. The building of multilayer thin walls on a substrate was observed using four processing conditions (which gave a range of energy density values), and the influence of the powder layer thickness has also been investigated. The keyhole operating mode was selected as there is a need for a better understanding of porosity formation, and whether it can be controlled or minimised in this mode. Additionally, the track morphology and internal porosity were characterised ex situ by $\mu \mathrm{CT}$, to measure pore sizes and distributions, and better understand how changing the energy density can minimise pores and optimise track morphology.

## 2. Methods

### 2.1 Materials

Commercially pure (CP) titanium substrates with dimensions of $46 \mathrm{~mm} \times 17 \mathrm{~mm} \times 0.3 \mathrm{~mm}$ were used throughout this study. 0.3 mm substrate thickness was selected as the best compromise between realistic build conditions and sufficient X-ray transmission for good image quality. The melt pool width is less than the substrate thickness, so the influence of the walls is minimal, and the overall effects of sample dimensions on the powder bed is negligible [37].

Gas atomised (GA) Ti-6Al-4V powder (supplied by Goodfellow, UK) with a particle size distribution of $15-45 \mu \mathrm{~m}$ was selected. The powder morphology was examined by scanning electron microscopy (SEM) in secondary electron imaging mode at 20 kV (JEOL JSM-6610LV, Tokyo). Figure 1a shows an essentially spherical morphology of powder, with few satellite particles. The particle size distribution was measured using a Mastersizer 2000 laser diffractometer (Malvern Instruments Ltd., UK). Figure 1b displays the cumulative size distribution plot for the powder where the $D_{10}, D_{50}$, and $D_{90}$ symbols represent the particle diameters for $10 \%, 50 \%$, and $90 \%$ of the cumulative volume (\%). The size distribution and the median $\left(D_{50}\right)$ particle size are representative of those typically used in LPBF.


Figure 1: (a) SEM image of Ti-6Al-4V gas atomised powder particles used in this study. (b) Cumulative powder particle size distribution with $D_{10}, D_{50}$, and $D_{90}$ labelled as determined by laser diffractometry.

### 2.2 In situ and operando process replicator

The experimental setup consists of a laser additive manufacturing in situ and operando process replicator (ISOPR) [37]; a small scale laser powder bed system, with the ability to produce multilayer parts. The ISOPR consists of: a continuous wave 200 W Yb-doped fibre laser, with a wavelength of 1030-1070 nm (SPI Lasers Ltd, UK); an environmental chamber (Figure 2a) containing a powder bed with $40-60 \%$ powder packing density; and a series of laser beam optics, namely a collimator, beam expander, and an $X-Y$ galvanometer scanner to control the laser line scanning, with an f-theta lens to focus the laser to a ca. $50 \mu \mathrm{~m}$ diameter spot $\left(4 \sigma_{x y}\right)$ [37]. The environmental chamber was evacuated and backfilled with argon gas at a constant flow rate of $4 \mathrm{Imin}^{-1}$ during experiments to maintain an inert atmosphere and prevent oxidation of the powder and molten pool.

The powder bed sample holder (Figure 2b and c) encases the substrate between two glassy carbon (GC) plates, which are used for their near transparency to X-rays. For each layer, the substrate is lowered by a pre-set layer thickness of $100 \mu \mathrm{~m}$, and the cavity between the GC plates is filled with Ti-6Al-4V powder using a vibrating gravity-fed powder hopper. A scraper behind the powder hopper ensured an evenly levelled powder surface on the deposited layer. This process is repeated for each subsequent layer in a build for 5 layers.

### 2.3 Experimental build conditions

Four processing conditions were selected, detailed in Table 1, henceforth referred to as Sample A, B, C, and D. Typical LPBF powder layer thicknesses range from $20-100 \mu \mathrm{~m}$ [52], however recent studies have investigated powder layers up to $250 \mu \mathrm{~m}$ for improved build rate [ 53,54$] .100 \mu \mathrm{~m}$ was selected for this study for the best radiography image quality within the standard operating range. The nominal linear energy density (LED) was calculated for each condition: $L E D=P / v t$ [17], where $P$ is the laser power, $v$ is the scan velocity, and $t$ is the powder layer thickness. LED is an adaptation of the volume energy density (VED) equation [12] where a single track was employed, and thus hatch spacing (h) was equal to 1. Supplementary Table 1 shows some examples of typical LPBF operating conditions and LED
values for comparison to the present study. The calculation shows that the energy density of this work is within the range used in typical industrial LPBF machines. VED or LED have been used in a number of studies [37,55-61] to correlate possible links between processing parameters, resultant energy densities, and build quality. However, VED and LED are not necessarily a reflection of absorbed energy density, due to powder ejection, and laser reflections within the keyhole [62] for example. The calculation also does not capture complex physics and is thus limited [55].

The ISOPR was employed to build a single line scan per layer for 5 layers onto a CP Ti substrate using a bi-directional scan strategy, alternating scan direction with each deposition layer. Layers 1, 2 and 5 were captured by synchrotron X-ray radiography. In the case of Sample A, a snapshot of the final track morphology was also recorded after layers 3 and 4.

### 2.4 In situ synchrotron X-ray radiography

In situ X-ray radiography experiments were carried out at the ID19 imaging beamline at the ESRF - The European Synchrotron, Grenoble, France [63]. The hard X-ray beamline uses a polychromatic beam, produced by two U32 undulators. The mean energy was ca. 30 keV . The attenuated X-ray beam was converted into visible light using a $200 \mu \mathrm{~m}$ thick LuAG:Ce scintillator (Ce-doped Lu_3AI_5O_12, Crytur, Czech Republic) and images were recorded with a FASTCAM SA-Z 2100K (Photron, USA) $4 x$ magnification, at $40,000 \mathrm{fps}$, an exposure time of $12.6 \mu \mathrm{~s}$ and an effective pixel size of $4.76 \mu \mathrm{~m}$. The field of view was 4.8 mm (width) $\times 2.4 \mathrm{~mm}$ (height).

### 2.5 Micro-computed Tomography ( $\mu \mathrm{CT}$ )

All samples were examined post-build by $\mu$ CT using a Nikon XTH225 (Nikon, Japan) to image and quantify internal porosity, Figure 2d; 3175 projections were collected, each with an exposure time of 1 s . The data was reconstructed using filtered-back projection and beamhardening algorithms embedded in CT Pro (Nikon), resulting in a voxel size of $2.7^{3} \mu \mathrm{~m}^{3}$ [51]. The image analysis was performed using Avizo 9.3 (ThermoFisher Scientific). A kernel of $3 \times 3 \times 3$ median filter was applied to remove noise, and a threshold applied to analyse internal
pores [50]. Any pore with a volume of $<27$ voxels was discounted from the analysis. Pore volume was converted from the number of voxels to equivalent diameter ( $D_{\text {eq }}$ ) using the equation:

$$
D_{e q}=\sqrt[3]{\frac{6 V}{\pi}}
$$

where $V$ is the volume of the pore [50] and hence 27 voxels is approximately a $D_{\text {eq }}$ of $10 \mu \mathrm{~m}$.
Pore sphericity, $\psi$, was calculated using the sphericity equation [64]:

$$
\psi=\frac{\pi^{\frac{1}{3}}\left(6 V_{p}\right)^{\frac{2}{3}}}{A_{p}}
$$

where $V_{p}$ is the volume of a pore, and $A_{p}$ is the surface area of a pore. Sphericity of pores with very small volume could exceed a value of 1 because the surface area measurements in Avizo 9.3 are based on chordal approximations, whereas volume measurements use the number of voxels and no approximations; this effect is more prevalent as voxel number reduced [65].

### 2.6 Image analysis

The open-source software Fiji version 1.52i [66] was used to analyse radiographs. 100 flatfield and 100 dark-field X-ray images were collected. The acquired radiographs were normalised using the flat-field correction (FFC) equation: $F F C=\frac{I-\text { Dark }_{\text {avg }}}{\text { Flatavg }^{\text {Darkavg }}}$, where $I$ is the raw image, Darkavg is the average of 100 dark-field images, and Flat $_{\text {avg }}$ is the average of 100 flat-field images [37]. This removed artefacts and noise variations inherited from the acquisition process. Local averaging over 50 neighbouring frames, 25 before and 25 after, was used to increase the contrast to reveal key features such as the laser keyhole, and remove stationary features; background subtraction using the first 50 static frames highlighted the deposited powder layer, shown in Figure 2e.


Figure 2: (a) Experimental build chamber with key components labelled. (b) Simplified schematic of the sample holder during in situ melting. (c) Schematic of the substrate and powder particles and melt track. (d) Schematic of $\mu C T$ sample scans. (e) Image processing methods.

| $\underset{\text { ID }}{\text { Sample }}$ | Laser power, $P$ <br> (W) | Scan velocity, $v$ ( $\mathrm{mm} \mathrm{s}^{-1}$ ) | Nominal powder layer thickness, $\boldsymbol{t}$ ( $\mu \mathrm{m}$ ) | $\begin{gathered} \text { Nominal } \\ \text { LED } \\ \left(\mathrm{Jmm}^{-2}\right) \end{gathered}$ | Average powder layer thickness, $\mathrm{t}_{\mathrm{m}}$, of the fifth layer ( $\mu \mathrm{m}$ ) | $\begin{aligned} & \text { Adjusted } \\ & \text { LED } \\ & \left(\mathrm{J} \mathrm{~mm}^{-2}\right) \end{aligned}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | $150 \pm 1$ | $200 \pm 1$ | 100 | 7.5 | $119 \pm 50$ | 6.3 |
| B | $200 \pm 1$ | $200 \pm 1$ | 100 | 10 | $115 \pm 45$ | 8.7 |
| C | $200 \pm 1$ | $300 \pm 2$ | 100 | 6.7 | $139 \pm 35$ | 4.8 |
| D | $200 \pm 1$ | $400 \pm 3$ | 100 | 5 | $163 \pm 50$ | 3.1 |

## 3. Results

Table 1 lists the LED values for the target processing conditions. The nominal $L E D$ was calculated using the desired powder layer thickness of $100 \mu \mathrm{~m}$. The average powder layer thickness measured, $t_{m}$, was obtained from the background subtracted radiograph images, with a threshold used to separate the powder layer. The powder thickness variation along the length of the track (shown later in Figure 4 f and Supplementary Figure 1) was averaged, and standard deviation calculated. $t_{m}$ was used instead of the nominal value $t$ to calculate an adjusted $L E D$, as $t_{m}$ was larger than $t$ for every condition. Previous studies have similarly observed thicker than intended powder deposition layers [67], however the disparity was due to the powder consolidation ratio rather than variations in track height as seen here. Typical LEDs used in previous literature are shown in Supplementary Table 1. The adjusted LED was thus lower than the nominal $L E D$ for every sample. This indicates that for multilayer builds, in addition to the aforementioned limitations of the LED calculation, the variation in powder thickness is another reason why LED is not necessarily a reliable method for quantifying absorbed laser energy density.

Initial observations of the radiographic data showed a large amount of powder spatter, as has commonly been seen in LPBF [39], caused by metal vapour jetting [68]. Powder entrainment into the jet can locally reduce the powder layer thickness ahead of the laser [69], and hence the laser beam can penetrate deeper into the previously deposited material. Due
to the nature of the sample holder, some excess powder was visible in the X-ray direction between the substrate and GC windows, along the entire length of the sample, and also on top of the GC windows. In some instances, this reduced the visibility of internal features such as pores. In multilayer builds, the track height was seen to vary along the length of the track, resulting in the large variation of powder layer thickness measured.

### 3.1 Multilayer track morphology

Figure 3 shows radiographs of a representative sample (Sample A) at various stages of the build. Background noise and excess powder has been removed in radiographs $(b-d)$ to highlight the track shape. Unedited final track radiographs are in Supplementary Figure 2 (layer 3), and Supplementary Videos $\mathbf{1}$ and 2 (layers 1 and 5 respectively). Layer 1 shows a uniform build height, largely free from undulations. Layer 3 shows a distinct undulating surface morphology along the length of the sample. Layer 5 shows an uneven surface but a reduction in amplitude of the undulations compared to layer 3.

The surface undulations significantly changed the thickness of subsequent powder layers, as the largest peak-to-trough distance was ca. $220 \mu \mathrm{~m}$ (Layer 3). Furthermore, local denudation, caused by recoil pressure and powder entrainment into metal vapor plumes [36,69], was more extensive where the powder layer was thinnest, (e.g. at the undulation peak) as vapour plumes are likely to eject a larger fraction of the available powder.

An uneven track surface also changed the depth of remelting of the previous layer along the track length. Each new powder layer was spread out with a level surface over the previous build. As the laser penetration depth was ca. $150 \mu \mathrm{~m}$ below the powder surface for Sample A, the peaks of the surface undulations - with less powder covering them - underwent deeper remelting than the troughs. Lack of fusion may have occurred when the laser penetration depth was less than the powder layer thickness.


Figure 3: Radiographs of final multilayer track morphologies in Sample A. (a) Full substrate length with added powder layer (100 $\mu \mathrm{m}$ ), and highlighted red region of interest for (b), (c) and (d) which show final melt track morphology in layers 1,3 and 5 respectively, with prominent surface undulations.

Figure 4 shows a time-series set of radiographs taken during the deposition of layer 5 in Sample B, with evidence of undulations in track height and powder layer. Figure 4a shows the large variation in powder layer thickness spread on top of the previous 4 solidified layers (yellow dotted outline). Figure 4b shows evidence of keyholing during the process; ca. 1.2 mm along the substrate, where the powder layer is around its thickest (ca. $200 \mu \mathrm{~m}$, Figure 4f). The purple dashed line in Figure 4c outlines the layer 5 track surface. The thickness added by the deposition of layer 5 is revealed by the difference between the dashed (purple) and dotted (red) lines. The keyhole extends ca. $350 \mu \mathrm{~m}$ below the track surface (Figure 4d and e). It is clear from these images that laser penetration extends well below the added powder layer and that significant laser remelting takes place. A large amount of powder spatter is seen behind the keyhole with a lesser amount in front of it. Dark and streaked features in Figure $\mathbf{4 e}$ indicate fast-moving spatter particles. The shape and depth of the keyhole are akin to the keyholes observed under similar energy density conditions for laser scanning with [34] and without [41] a powder layer, i.e. laser welding.

Supplementary Videos 3 and 4 show the deposition of layer 5 in Samples C and D respectively, and also reveal prominent surface undulations and powder thickness variations. Notably, the video of Sample D shows significant porosity, characteristic of lack of fusion, at the end of the track.


Figure 4: (a)-(c) Background subtracted radiograph time series of Layer 5 melting in Sample B: (a) Track shape of prior 4 layers (yellow) with a deposited powder just before layer 5 is deposited. (b) ROI showing the formation of a deep keyhole depression during layer 5 melting (d-e). (c) Final track morphology of layer 5 (purple outline). (d) FFC radiograph of the laserinduced keyhole in layer 5. The red outline denotes the approximate position of the laser beam. (e) The keyhole shape is highlighted by locally averaging (d). (f) Graph showing the variation in powder layer thickness plot along the track length. Corresponding videos of the full track length melting with FFC and background subtraction (a-c) are shown in Supplementary Video 5 and 6 respectively. Powder thickness variation plots of Samples A, $C$ and $D$ are shown in Supplementary Figure 1.

### 3.2 Porosity analysis

$\mu \mathrm{CT}$ has been used to analyse pore shapes and sizes throughout all samples, to reveal overall trends in porosity. Pore formation mechanisms and interactions in multilayer builds are detailed in 3.3. Figure 5 shows $\mu \mathrm{CT}$ rendered images of the samples. The internal porosity has been highlighted by three colours, each representing a different $D_{e q}$ size range. Pores of $D_{\text {eq }}<25 \mu \mathrm{~m}$ are termed small, those with $25 \leq D_{e q}<45 \mu \mathrm{~m}$ are termed mid-sized, and pores with $D_{\text {eq }} \geq 45 \mu \mathrm{~m}$ are termed large. The surface roughness visible on the sides and top of each sample was caused by semi-melted powder particles which adhered to the track but were not fully consolidated into the melt pool before solidification, see example in Supplementary

Figure 3. Table 2 shows the total number of pores of different size fractions and volume porosity for each condition. The percentage porosity was calculated using volume measurements from the $\mu \mathrm{CT}$ data; the methodology is in Supplementary Information.

In Sample A, substrate remelting must have taken place during the first build layer as the 6 pores observed were in the substrate. The largest volume porosity was measured in Sample B, in which the majority of large pores were observed at either end of the sample. Figure 5 inset (a) illustrates the morphology of large irregularly shaped pores. The overall pore frequency histogram (below the $\mu \mathrm{CT}$ image for Sample $B$ and $C$ ) shows a higher number of pores at each end of the track, and pore frequencies tend to be higher below the peaks in the deposition profile.

Sample $C$ has about half the total volume porosity compared to Sample B, with reduced numbers of pores across all size ranges. The large pores are located towards one end of the track, where the solidified layer thickness is largest. Inset (b) shows a 'peanut' shaped large pore, in which two smaller pores appear to have coalesced. The pore frequency histogram reveals a correlation between the track height, and pore distribution, as more pores exist where the track is highest, which can be seen most clearly in Figure 5 Sample C. The laser penetration depth into solidified track is lower than into the powder layer, so track peaks have a very thin layer of powder and substantial track remelting, whereas trough regions require
the laser to penetrate predominantly into powder. This suggests that when the track height is higher, the laser penetration depth is not sufficient to remelt pores in the previous layer, and could increase the overall pore volume. The variation in laser penetration depth of layer 5 melting can be seen in Supplementary Videos 2, 3, and 5.

The volume porosity in Sample D was measured to be similar to that of Sample A. However, there are multiple regions showing poor layer cohesion, predominantly located at the end of the track. Inset (c) shows evidence of a lack of fusion pore, which due to the nature of single line scan tracks, propagated through the width of the sample. Although sample D has around $2 \%$ of the volume porosity compared with Sample B, this porosity analysis excludes surface connected pores such as inset (c), and is thus an underestimation of the total volume porosity, and care must be taken when comparing samples in this manner.

Table 2: Porosity quantification obtained from $\mu \mathrm{CT}$ data.

| Sample <br> ID | Number <br> of small <br> pores | Number <br> of mid- <br> sized <br> pores | Number <br> of large <br> pores | Total <br> number of <br> pores | Pore <br> Volume <br> $\left(\mu \mathrm{m}^{3}\right)^{*}$ | Percentage <br> Porosity <br> $(\%)^{*}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | 5 | 1 | 0 | 6 | $26 \times 10^{3}$ | 0.0032 |
| B | 54 | 57 | 13 | 124 | $2322 \times 10^{3}$ | 0.22 |
| C | 48 | 36 | 5 | 89 | $1126 \times 10^{3}$ | 0.15 |
| D | 4 | 2 | 0 | 6 | $41 \times 10^{3}$ | 0.0050 |

* The error range is from 2-6\%.


Figure 5: 3D rendered volumes from the reconstructed $\mu C T$ scans highlighting pore size, distribution, and pore location. Pore frequency histograms with a bin size of 0.2 mm along the track length of Samples B and C. Insets (a) and (b) show large irregular pores. Inset (c) shows an interlayer pore. For interpretation of the colours in this figure, the reader is referred to the online version of this article.

Figure 6 shows porosity measurements from the reconstructed $\mu \mathrm{CT}$ data, with pore volumes converted into $D_{\text {eq }}$ values. The number fraction of internal pores measured in each sample is given in Figure 6a. Sample A has few pores, almost $70 \%$ of which are in the 20 -
$25 \mu \mathrm{~m}$ size range. An even range of small and mid-sized pores are evident in Samples B and C. Sample $D$ has very few pores but a similar number fraction of pores across the small and mid-size ranges.

The pore volume fractions for each $D_{e q}$ bin size is shown in Figure $\mathbf{6 b}$, indicating their contribution to the total overall volume porosity. Although Sample $B$ and Sample $C$ have a low number of large pores, the plot shows that these contribute significantly to the overall volume porosity. The highest number of pores in Sample C was small pores, however, the contribution to volume porosity is low. The largest individual contribution being mid-sized pores; the keyhole pores at the base of each melt layer. Sample $D$ had a spike in pore volume from pores in the range $30-35 \mu \mathrm{~m}$ (even though there is no spike in the number of pores) indicating that fewer larger pores influence the volume porosity considerably more than many smaller sized pores. Sample $D$ had low overall pore volume because the methodology employed to analyse the $\mu \mathrm{CT}$ data did not quantify the surface connected pores, such as that shown in Figure 5 inset (c).

Figure 6c shows the pore sphericity measurements in a scatter plot. The small pores are typically close to spherical (sphericity value of 1 ). As the $D_{e q}$ value increases, the spread of the data also increases, showing pores to have a less uniform shape, with a minimum sphericity of ca. 0.7 for large pores. The pores shown in Figure 5 insets (a) and (b) are representative of pores with low sphericity values. Sample B had the largest range in sphericity values for all pore sizes of ca. 0.7 to 1.0 . In Sample C, $96 \%$ of pores had values $\geq 0.9$, with a minimum sphericity of 0.8 . All pores in Samples $A$ and $D$ had high sphericity $\geq 0.96$.


Figure 6: Measurements extracted from $\mu C T$ analysis. (a) Pore number fraction plot, with dividing lines for small, medium and large pore size terms. (b) Volume fraction plot for all pores in each sample. (c) Scatter plot of pore sphericity for all measured pores.

### 3.3 Keyhole melting and porosity formation observations

In the previous section, ex situ pore quantification has been reported. Through in situ X-ray radiography, it is possible to elucidate mechanisms by which the pores could have developed; described in this section. The melting of layer 1 in Sample C is shown in Figure 7, for a small region at the start of the melt track. Figure 7a shows a uniform powder layer of ca. $100 \mu \mathrm{~m}$ atop the substrate and Figure 7b-d show the evolution of the single layer melt track and keyhole pores in the substrate. As keyhole walls become unstable and collapse [34], pores form at the bottom of the keyhole. The insets in these figures $(\mathbf{b}-\mathbf{d})$ show locally averaged radiograph images of the keyhole shape and pore formation. Figure 7 e and f show that these pores are retained inside the melt track, and an uneven track surface is formed upon solidification of the first layer.

Layer 2 melting and pore interactions in Sample C are shown in Figure 8 for the same region of interest (ROI) as Figure 7. Details are shown in Supplementary Videos 10 (full track) and 11 (cropped ROI), for layer 2. Figure 8a shows the full track length of layer 2 after melting, with the ROI highlighted for ( $\mathbf{b} \mathbf{- k}$ ). Figure $\mathbf{8 b}$ shows a schematic (traced from radiographs) of layer 1 , with the powder deposited for the next layer, just prior to melting.

Figure $\mathbf{8 c} \mathbf{- h}$ show three different pore interactions that are observed to occur during multilayer laser melting and are described as follows: (i) the laser penetration depth of layer 2 melting is deep enough to interact with the solidified pore from the first layer, and the melt pool fills the void (Figure 8c-d, dotted yellow outline denotes a filled pore). (ii) The laser penetration depth is insufficient to fill the pre-existing pore with liquid metal (Figure 8e and $\mathbf{f}$ ). This can be due to laser processing parameters, or in the case of Sample C, changes in track height along the length of the sample affect the depth of remelting, as described in Figure 3. (iii) The laser keyhole is unstable [34] and produces pores, in the same way to those in layer 1 , and prior work $[42,69]$. These new voids may coalesce with existing pores to reduce interfacial energy in the melt pool, which increases their size. As the solidification rate is high,
these pores have not been able to rise to the surface and be removed before solidification (Figure 8g - $\mathbf{h}$ ).


Figure 7: Observation of porosity formation in Sample C, layer 1. (a) Radiograph of substrate and powder prior to melting. Time series schematics (traced from the radiographs) and corresponding radiograph insets in (b) to (d). (b) Melt track (blue) with laser keyhole in the substrate plate (grey). (c) Keyhole walls become unstable. (d) A new pore is formed at the keyhole. (e) Pore positions after layer 1 melting where the blue dotted line shows the approximate keyhole depth. (f) Radiograph showing the position of the keyhole pores. See keyhole pore formation in Supplementary Videos 7 (full track), 8 (cropped ROI), and 9 (locally averaged cropped $R O I$ ). Insets in $(\mathbf{b}-\boldsymbol{d})$ are made by locally averaged radiograph images to highlight keyhole shape. See Supplementary Videos 10 and 11.

In Figure 8h, the solidified layer 2 track is shown in blue, with solidified layer 1 outlined in the dashed red line. This shows that track remelting occurs and leads to a redistribution of material. Figure $8 \mathbf{i}$ is the final solidified layer 2 radiograph, with the $\mu \mathrm{CT}$ data as an overlay. It shows that the solidified pores observed remained in the sample. With high-resolution $\mu \mathrm{CT}$ scans ( $2.7 \mu \mathrm{~m} /$ voxel $)$ relative to the synchrotron X-ray imaging ( $4.76 \mu \mathrm{~m} /$ pixel $)$, we reveal additional pores in the samples which cannot be resolved in the X-ray radiograph. Pores produced in layers $3-5$ are also visible in the $\mu \mathrm{CT}$ overlay.


Figure 8: Schematic of porosity formation (traced from the radiographs) during deposition of the second layer in Sample C. (a) Radiograph of full track length after layer 2 melting, with the highlighted region of interest for (b-k). (b) Initial powder layer (100 $\mu \mathrm{m}$ ) on the substrate, prior to laser melting. ( $\mathbf{c}-\mathbf{g}$ ) Schematic of 3 laser interactions in multilayer melting. (h) Solidified layer 2 track; layer 1 outline shown in red. (i) Layer 2 final $X$-ray radiograph, with an overlay showing solidified pores in the sample from the $\mu C T$ data. For the radiographs from which $(b-h)$ are traced, please see Supplementary Videos 10 and 11.

## 4. Discussion

### 4.1 Morphological development in multilayer tracks

Periodic undulations similar to those found in the present in situ study have also been observed in recent studies of the LPBF process, e.g. in thin-wall structures [70,71], in single line scan tracks [11,61] and in cubes [72]. They have also been seen in laser welding [73] and been modelled computationally [59]. The undulations, also referred to as humping [11], have been attributed to Marangoni flow in the melt pool and surface tension effects [70], mainly the Plateau-Rayleigh instability, which describes the break-up of fluid into discontinuous elements [74]. It is denoted as balling [4] when the effects are sufficiently extreme for LPBF tracks to become discontinuous.

A high length-to-width melt pool ratio would be expected to promote the PlateauRayleigh instability [11,61] and it has been reported that this increases the probability of such undulations developing [70]. It has also been observed that tracks produced under conditions of high LED, which increases melt pool length, produced more undulations [55,59,70]. However, in other studies, it has been found that decreasing the $L E D$ with a constant laser power led to a transition from a continuous track to surface undulations and to balling [55]. Evidently, there is a more complicated relationship between LED and formation of undulations, presumably because other process-related factors such as backward fluid flow in the molten pool [11] and recoil pressure [59] will play a role. In the present study, the range of $L E D$ values employed was not sufficiently great to provide further clarification. However, a notable feature from the in situ work is the observation in Sample A that the amplitude of undulations decreased as the number of layers increased, Figure 3. This could be related to the bidirectional melting strategy, but further research is necessary to confirm this.

The largest peak to trough distance observed was ca. $220 \mu \mathrm{~m}$ in Sample C, which has a substantial effect on layer remelting along the length of the track; as the keyhole depth was ca. $200 \mu \mathrm{~m}$ for these conditions in layer 5 (see Supplementary Video 5). This results in significant remelting of the previous layer at the wave peaks with little or none in the troughs.

The laser penetration depth itself also varied by $30-40 \mu \mathrm{~m}$, going deeper in the troughs where it is melting power rather than prior solid track. Despite this deeper penetration at the troughs, in some areas the prior layer was sufficiently deep that that lack of fusion pores or unmelted powder between tracks was observed, such as those present in for Sample D, Figure 5. The effect of variable layer remelting is to directly influence the pore populations in multilayer samples, and this will be explored further in the following sections.

### 4.2 Gas Porosity

For GA powders, it is possible that gas trapped within the feedstock powder, typically argon [75], as a result of the manufacturing route can be entrained into the pool during melting [76], as well as retained soluble gases such as hydrogen [77]. In LPBF, hydrogen can also from several other sources, such as the decomposition of water vapour on the surface of the powder, or water vapour in the environmental chamber [78]. This is most prevalent in materials such as aluminium alloys [25] where hydrogen solubility increases with increasing temperature. Although this is not the case in titanium alloys, hydrogen has still been shown to cause porosity in welds [79-81], where at the liquidus temperature hydrogen is twice as soluble in the liquid as the solid (partition coefficient of ca. 0.5 [82]). Hence, hydrogen is rejected into the melt pool by the advancing solid-liquid interface and the melt becomes supersaturated. Hydrogen and other soluble gases will be swept along in the melt pool until the end of the track. When the laser is turned off at the end of the track, these soluble gases can become highly supersaturated as the liquid pool shrinks, and will either nucleate new pores, or diffuse into pores formed by the keyhole, stabilising or increasing their size. This may also happen all along the track, but to a lesser extent as the supersaturation will be less. This mechanism will be discussed later in greater detail. It is also possible that some of the small spherical pores will contain argon gas which is entrained from the environmental chamber when the keyhole pores form $[4,33]$.

In some literature on the expected diameter of gas pores in additively manufactured materials, studies have classified any pores lower than $100 \mu \mathrm{~m}$ in diameter [29,83], or even between $100 \mu \mathrm{~m}$ and $300 \mu \mathrm{~m}$ [28] as gas pores. However, in the present study only those pores with $D_{\text {eq }}$ approximately $\leq 30 \mu \mathrm{~m}$ had sphericity values close to 1 which is regarded as a strong indicator of a gas pore. Larger pores tended to have lower sphericity values and those visible on the radiograph images, measured to be ca. $40-60 \mu \mathrm{~m}$ in size, clearly formed via a keyhole mechanism. Therefore, it can be concluded that size alone cannot be used to determine the mechanism by which pores form and that further in situ studies are required to clarify this.

### 4.3 Keyhole pore formation

Supplementary Video 6 and Figure 7 show keyhole pore formation over a 0.05 ms period. A keyhole is created, quickly distorted, and covered over by the unstable liquid metal above, resulting in a large entrapped pore. The final pore observed from this interaction is ca. $40 \mu \mathrm{~m}$; similar to the size found in industrial practice $[34,84]$. This formation mechanism has also been explored computationally $[4,33]$. Figure 5 shows that most irregular pores were found to be at the base of each melt track, and in the substrate for the initial layer. This supports their formation via the keyhole melting mechanism [10] and correlates with the stationary beam in situ experiments $[34,36]$ in which the closure of a keyhole was seen to take ca. 0.05 ms . The decrease in sphericity (from 1.0 to 0.7 ) as the pore size increases implies that keyhole melting and metal vaporisation effects produced unstable voids [33,34], which were then trapped during solidification. The pores with low sphericity values may either have become distorted when trapped between dendrites [85], or when two pores are trapped as they coalesce (see Figure 5 inset (b)). Large keyhole pores were most prevalent in Samples B and C and the number fraction of pores for each $D_{e q}$ bin size show a similar trend, Figure 6a.

### 4.4 Keyhole-pore interactions in remelted layers

Lack of fusion pores are potentially detrimental features in AM components as they are generally non-uniform in shape [86], can contain unmelted powder particles [87], and can act as crack initiation sites. Therefore, based on the observations in this study, sufficient laser penetration depth, possibly through the keyhole mode of heating, can assist to ensure adequate fusion to the previous layer. A phase-field model of multilayer scanning [88] supports this, as a higher percentage of lack of fusion pores were observed when energy density was decreased. The effects of track balling led to pores, which were not remelted in subsequent layer addition, however in this work keyhole pores were not studied. In the present study, such lack of fusion features became more prominent in Samples $C$ and $D$ as the scan velocity was increased, and the depth of keyhole formation was expected to decrease. In Sample C, the layers at the ends of the track were insufficiently fused together. In Sample D, there were large areas of unfused material at the end of the track, as well as a lack of fusion pore mid-track, Figure 5 inset (c). It can be estimated that the mid-track pore was ca. $60 \mu \mathrm{~m}$ long and the lack of fusion feature at the end of the track was ca. $450 \mu \mathrm{~m}$ long.

Whilst some remelting a previous layer is desirable, achieving this using keyholing may introduce keyhole pores at the interface between layers as shown in the in situ radiographs, and in the ex situ $\mu \mathrm{CT}$ reconstructions, Figure 5. However, the in situ radiographs, Figure 8, show that remelting can also be used advantageously, as large pores were eliminated from the track by a pore filling mechanism. Partial pore filling has also been observed previously [40]. Whether these pores were completely removed or redistributed as smaller pores [37] is unclear, as the spatial resolution of the radiographs is insufficient to confidently resolve pores with a diameter $<40 \mu \mathrm{~m}$. Remelting was also seen to increase the size of one pore, Figure $\mathbf{8 h}$, similar to that observed in overhang conditions [40], which is undesirable. The likelihood of this occurring could be related to the solidification rate, and whether pores have time to be filled by molten metal or rise to the surface of the melt before solidification. Therefore, careful control of layer remelting is an important factor to consider when selecting process conditions
in LPBF, as the multilayer in situ study has shown that complex interactions are taking place. The selection of suitable process parameters is further complicated by the variations in powder layer thickness, and consequently variable laser penetration to the previous layer. However, a more in-depth analysis with a larger number of deposited layers would be necessary to explore this. It must also be noted that this study only explored single melt track deposition, and that some of the findings may not be applicable to typical depositions used in LPBF employing multiple hatches and contours. However, this study is relevant to multilayer melting and can be used to validate models, and hypotheses can be applied to hatched samples.

### 4.5 Inhomogeneous pore distribution

The $\mu \mathrm{CT}$ results, Figure 5, clearly reveal that the large pores in Samples B and C form primarily at the ends of each melt track and that this is also true of mid-size pores in Sample C. Prior work has also observed a correlation between scan strategies and pore location [22,23,30,72]. The irregular shapes of such pores are highlighted in the insets to this figure and the time series radiographs in Figure 8 provide clear mechanistic information of the role of layer remelting in developing large irregular pores. Recently in situ methods [32] and computational methods [5] have been used to study point pore formation during turning, i.e. at the end of a track where the direction of travel changed. They concluded that pores form at laser turn points due to the formation and subsequent collapse of a deep keyhole. However, their powder-on-plate findings may not provide a full understanding of multilayer builds in LPBF.

Our results suggest that both gas and keyhole pores are combined to form larger pores during the melting and cooling of the track. The hypothesis is that firstly, the dissolved soluble gases are swept along in the melt pool concentrating at the end of the track, similar to the solute concentration that is observed in the Czochralski zone refining process [89]. Hydrogen is a solute element in the molten alloy, with a partition coefficient $k$ of ca. 0.5 [82]. During the transformation from liquid to solid, hydrogen solubility decreases (halving for titanium), and
thus partitioning at the interface will concentrate dissolved hydrogen in the melt pool, sweeping solute elements to the end of the track. Normally, keyhole pores containing superheated metal vapour and a small amount of argon gas shrink as the melt pool cools. However, at the end of the track, the soluble gas (hydrogen) is concentrated by the sweeping (or Marangoni convection), and even further so as the pool solidifies at track end. This super-saturated gas could diffuse into the keyhole pores, increasing or maintaining their size.

The 'peanut' shape formed as two pores joined together in inset (b) indicates that the solidification process happened so rapidly that the void could not reshape into a sphere to reduce surface tension and was frozen as an irregular pore [30]. It is also possible that complex fluid flow could distort pore shapes. In agreement with a prior work [90], porosity minimisation at the ends of the track or turning points could be achieved via a reduction laser power in these areas to prevent keyhole porosity forming. The present work supports this as Sample A and B only differed in laser power (150 W and 200 W , respectively), but Sample A had 99 \% less internal volume porosity. As can be seen from Supplementary Videos 2 and 5, the keyhole in Sample A is ca. $150 \mu \mathrm{~m}$ deep, while for Sample B the keyhole is ca. $250 \mu \mathrm{~m}$, generating many more keyhole pores. Previous work similarly shows higher porosity percentage in higher energy conditions [91]. In large LPBF components, directionality, hatch strategy, and the depth of layer remelting will influence pore location and size, and all need to be carefully considered.

## 5. Conclusions

This study investigates the track formation and internal porosity during laser powder bed fusion of Ti-6Al-4V single track, multilayer builds using high speed in situ synchrotron imaging. The samples were further examined by ex situ $\mu \mathrm{CT}$ to further support our findings. The following conclusions have been drawn:

1. Undulations in the track surface were observed, which led to different depths of remelting along the track length in subsequent track, and thus different pore interactions
in multilayer builds. The deposited powder layer thickness changes as track height changes, decreasing the uniformity of multilayer components through this inconsistent remelting. This is an important factor which needs further investigation and consideration when selecting processing conditions.
2. There interactions were identified to occur during multilayer laser melting: (i) pores are remelted and the void is filled with molten metal; (ii) the laser penetration depth is not sufficient to fill the pore and they remain in the track; (iii) the laser keyhole is unstable and produces new pores, which can join to existing pores, increasing them in size.
3. Keyhole pore formation was quantified, illustrating how pore size varies as a function of build parameters (laser velocity and power), ranging from $10-60 \mu \mathrm{~m}$, with those $<40 \mu \mathrm{~m}$ being detected ex situ with $\mu \mathrm{CT}$. In situ keyhole pore formation was found to occur in a process taking under 0.05 ms , forming pores with a size ca. $30-60 \mu \mathrm{~m}$, correlating well to prior in situ studies.
4. It is hypothesised that the formation of larger pores at the end of tracks is the result of the stabilising and growth of keyhole porosity by diffusion of supersaturated soluble gases (hydrogen) into these pores. This soluble gas is concentrated as the pool sweeps along the track due to the partition coefficient (k) of ca. 0.5 . for hydrogen in titanium. Upon cooling at the end of track, the solubility rapidly decreases, and hydrogen diffuses into nearby pores.

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## 7. Author contributions

PDL, LS, CLAL, and GB conceived the project. LS, YC and CLAL designed the experiments. All authors performed the beamtime, except GB and JG. LS performed $\mu \mathrm{CT}$, data analysis, and results interpretation. LS, CLAL, and PDL led the results interpretation and paper writing, with all authors contributing. JG performed the powder size distribution analysis.

## 8. Data Availability

Representative samples of the research data are given in the figures (and supplementary data - DOI if available). Due to their large size, other datasets generated and/or analysed during this study are available from the corresponding author on reasonable request.

## 9. Declaration of interest

The authors declare no competing financial interests.

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