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# Quantitative Electrical Imaging of Three-Dimensional Moisture Flow in Cement-Based Materials

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## Abstract

The presence of moisture significantly affects the mechanical, hydraulic, chemical, electrical, and thermal properties of cement-based and other porous materials, and therefore, methods for detecting and quantifying the moisture ingress in these materials are needed. Recent research studies have shown that the ingress of moisture in porous materials can be *qualitatively* imaged with Electrical Impedance Tomography (EIT) – an imaging modality which uses electrical measurements from object’s surface to reconstruct the electrical conductivity distribution inside the object. The aim of this study is to investigate whether EIT could image the three-dimensional volumetric moisture content within cement-based materials *quantitatively*. For this aim, we apply the so-called absolute imaging scheme to the EIT image reconstruction, and use an experimentally developed model for converting the electrical conductivity distribution to volumetric moisture content. The results of the experimental studies support the feasibility of EIT for quantitative imaging of three-dimensional moisture flows in cement-based materials.

*Keywords:* Electrical Impedance Tomography (EIT), Imaging, Non-destructive testing, Unsaturated Moisture Flow, X-ray imaging, Cement-based materials.

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## 1. Introduction

The mechanical, hydraulic, chemical, thermal, and electrical properties of porous media are significantly affected by the moisture content within the media [1, 2, 3], and especially in cement-based materials, the majority of deterioration processes take place in the presence of moisture [4, 5, 6]. For these reasons, there has been a significant interest in developing methods of monitoring moisture flows in cement-based materials and other porous media. While some high-resolution imaging methods of monitoring moisture flow in porous media are available, e.g. X-ray imaging [7, 8, 9, 10],  $\gamma$ -rays [11, 12, 13], neutron imaging [14, 15, 16, 17, 18, 19, 20, 21, 22, 23], and nuclear magnetic resonance [24, 25, 26, 27, 28], these methods are often expensive, limited to testing small specimens, and/or have significant energy demand. Electrically-based methods, on the other hand, have often a limited spatial resolution but are generally inexpensive, non-invasive, and require comparatively lower energy [29, 30, 31, 32, 33, 34, 35]. One electrically-based method that has been recently shown to be particularly promising for monitoring moisture in porous media is Electrical Impedance Tomography (EIT) [36, 15, 37].

In EIT, the electrical conductivity distribution inside an object is reconstructed based on voltage measurements from the surface of the object. Previous studies have demonstrated the feasibility of EIT to *qualitatively* image unsaturated moisture flows in porous media. Such research includes: 1D and 2D unsaturated moisture flow in soil [38, 39, 40, 41] and cement-based material [42, 43, 37, 15], as well as 3D moisture flow in sand [44], sandstone [45], soil [46] and cement-based materials [36]. In all these works, (the spatial map of) the increase of the moisture content with respect to some initial state was inferred qualitatively

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## Nomenclature

$\alpha$	van Genuchten model fitting parameter	$H_s$	Hounsfield number for saturated material
$\bar{n}$	unit normal vector	$I$	tortuosity and pore connectivity parameter
$\ell$	arc length variable	$I_0$	intensity of incident X-ray
$\Gamma_e$	covariance matrix	$I_x$	intensity of transmitted X-ray
$\mu$	X-ray attenuation coefficient	$I_l$	total current of $l^{th}$ electrode
$\mu_{\text{air}}$	X-ray attenuation coefficient of air	$K$	unsaturated hydraulic conductivity
$\mu_{\text{water}}$	X-ray attenuation coefficient of water	$K_r$	relative hydraulic conductivity
$\Omega$	target volume	$K_s$	saturated hydraulic conductivity
$\partial\Omega$	target boundary	$L$	specimen length
$\rho$	electrical resistivity	$L_e$	Cholesky factor for the noise precision matrix
$\rho_w$	density of water	$L_\sigma$	Spatially weighted discrete differential operator
$\sigma$	electrical conductivity	$m$	van Genuchten model fitting parameter
$\sigma_{\text{max}}$	upper constraint of electrical conductivity	$m_w$	molecular weight of water
$\sigma_{\text{min}}$	lower constraint of electrical conductivity	$n$	van Genuchten model fitting parameter
$\Theta$	effective material saturation	$p(\sigma)$	regularization function
$\theta$	volumetric moisture content	$R$	electrical resistance
$\theta_i$	initial volumetric moisture content	$R_s$	universal gas constant
$\theta_r$	residual volumetric moisture content	$RH$	relative humidity
$\theta_s$	volumetric moisture content at saturation	$T$	temperature
$\hat{\sigma}$	conductivity estimate	$U$	EIT forward mapping
$\xi_l$	contact impedance of $l^{th}$ electrode	$u$	electric potential
$A$	specimen cross-sectional area	$U_l$	electrode potential
$d$	X-ray specimen width	$V$	measured electrode potentials
$e$	Gaussian distributed noise	$w/c$	water to cement ratio
$e_l$	$l^{th}$ electrode	$x, y, z$	spatial coordinates
$H$	Hounsfield number for unsaturated material		
$h$	capillary pressure		

20 from a relative increase in the electrical conductivity (or decrease in the electrical resistivity) distribution. The change in the electrical conductivity was estimated using a linear approximation in the solution of the non-linear EIT inverse problem, i.e., using the so-called *linearized difference imaging* scheme, resulting in a qualitative estimation of the conductivity change [15].

25 The aim of this study is to investigate whether EIT could be used for *quantitative* estimation of the three-dimensionally distributed water content in cement-based materials. In the experiments, moisture ingresses cement-based material containing fine aggregates. EIT measurements are collected during the moisture ingress, and used for reconstructing the 3D distribution of the electrical conductivity within the material. For the EIT reconstruction, the so-called *absolute imaging* scheme [47, 48, 49, 50] is adopted. The reconstructed electrical conductivity distributions are further converted to volumetric moisture content 30 distributions, by using an experimentally-developed model. The results of EIT are compared with X-ray

Computed Tomography (CT) and numerical simulations of unsaturated moisture flow. While the experiments are performed on cement-based materials, the methods developed in this work are also applicable to other porous materials.

## 2. Materials

### 35 2.1. General

For testing the feasibility of EIT for monitoring the moisture ingress in cement-based materials, two cylindrical mortar specimens were prepared. In the experiments, the moisture was added to the specimen from a PVC reservoir mounted on top of the specimen (Fig. 1). In one specimen, the reservoir was filled with water and in the second specimen, dilute solution of water and non-ionic radiocontrast agent (Iohexol, 40 120 mg Iodine/ml) was used. This was done to improve the contrast between the background and the regions of moisture ingress in X-ray CT imaging [36].

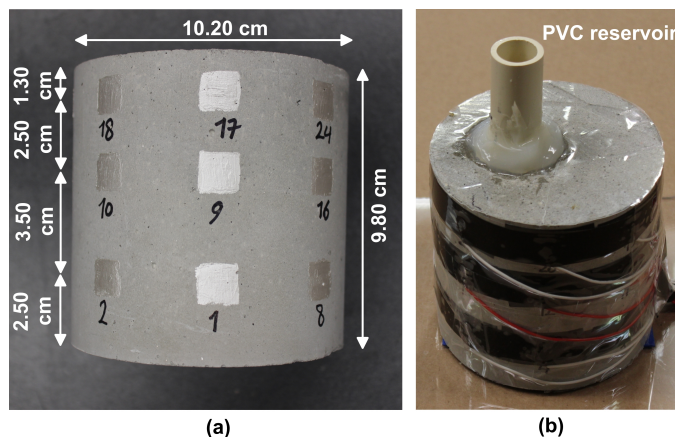


Figure 1: (a) Photograph of the specimen and locations of the electrodes, (b) prepared specimen.

### 2.2. Materials and sample preparation

The specimens used in this work were made of Portland cement mortar. The mortar was made with ordinary Type I Portland cement and fine aggregate consisting of natural river sand. The water-to-cement ratio (w/c) of the mortar was 0.60 and its total volumetric aggregate content was 40%. The high w/c and low aggregate content (high cement paste content) were used to accelerate the rate of moisture ingress and to reduce the test duration. The mortar mixture was prepared according to ASTM C192-06 [51]. A circular cylinder of diameter 10.20 cm and height 20.30 cm was cast from the mixture. The cylinder was demolded after 18 hours and cut in half to create the two specimens for the experiments.

After cutting, the specimen was moved to an oven at 50°C for 5 hours, and then sealed in two layers of plastic bags and placed inside an environmental chamber at 23°C for 21 days. This process was carried out to achieve a relatively uniform moisture content throughout the specimen. The uniform initial moisture content was not necessary for EIT measurements, but was considered to improve the contrast in X-ray CT images. Furthermore, the corroboration of the experimental results with the flow simulation required the knowledge of the initial moisture content. The initial moisture content was determined experimentally by completely drying separate specimens conditioned using the same procedure; the conditioning resulted in an initial moisture content of  $\theta_i = 0.07$ .

The electrodes used in this work were made of colloidal silver paint which was applied to the surface of the specimen using a small brush. The colloidal silver paint had a low resistivity and was fast drying [52]. A total of 24 square electrodes (with dimensions of 1.30 cm  $\times$  1.30 cm) were painted on the perimeter of the specimen. The electrodes were arranged in 3 rings of 8 equally spaced electrodes. Fig. 1a shows the specimen and the locations of the electrodes. After the silver paint had dried, an 18 gage wire was placed on the surface of the silver paint. Contact between the electrodes and wires was secured using electric tape and zip ties (Fig. 1b).

65 The water reservoirs were made of 5.0 cm long PVC pipes with internal diameters of 0.95 cm, and they were installed on the top of the specimens using fast-drying silicon-based caulking. The water reservoirs were positioned eccentrically 1.30 cm off-center. The purpose of this positioning was to induce asymmetric moisture ingress in the specimens and realize whether EIT reconstructions can capture the water source location and asymmetric flow within the specimens. After placing the water reservoirs, the cylinders were 70 wrapped with transparent tape to prevent evaporation of water from the specimens' surface during the test (Fig. 1b).

### 3. Electrical Impedance Tomography (EIT)

#### 3.1. Reconstruction methods in EIT

In Electrical Impedance Tomography (EIT), electric currents are injected into an object through a set of electrodes on the specimen's boundary. Corresponding to each current injection, a set of electrode potentials/potential differences are measured with respect to a common ground or between electrodes. Based on the measured potentials, the distribution of the electrical conductivity inside the object is reconstructed. A variety of methods for the image reconstruction in EIT have been developed [53]. Broadly, two imaging methods are commonly used in EIT: difference and absolute imaging.

80 In difference imaging, the temporal change of the conductivity distribution is estimated based on a difference between potential measurements corresponding to two states: before and after the change. The difference reconstructions are often relatively tolerant to systematic measurement and modeling errors, because the errors at least partly cancel out in the subtraction of the two data sets. However, in the conventional difference imaging scheme, the non-linear observation model of EIT (see below) is globally linearized; this can result severe bias in the reconstructed conductivity change, especially when the true conductivity change is large and/or if a good approximation for the initial conductivity is not available. In consequence, the linearized difference reconstructions are often only qualitative in nature [29, 50]. In [36], the linearized difference imaging was used for monitoring the water ingress in mortar; severe artifacts were observed, especially when the water ingress had caused a large change in the conductivity of the mortar with respect to the 85 initial state. The problems associated with the linearized difference imaging can be avoided to an extent by using the recently proposed non-linear difference imaging approach [47, 54]. In the present study, however, absolute EIT imaging is used – aiming at quantitative reconstructions of the conductivity distributions.

In absolute imaging, the conductivity distribution is reconstructed by solving the full non-linear inverse problem of EIT. The absolute reconstruction usually requires iterative solution. Furthermore, unlike in difference imaging, the reconstructions are generally severely intolerant to modeling errors, and hence, accurate (and computationally expensive) numerical models are needed for modeling the measurements. Consequently, the computational cost of absolute EIT imaging is often significantly higher than that of difference imaging.

100 In this work, the EIT measurements are modeled by the Complete Electrode Model (CEM), which is the most accurate forward model in EIT to date [55, 56]. The CEM consists of the partial differential equation

$$\nabla \cdot (\sigma \nabla u) = 0, \quad x \in \Omega \quad (1)$$

and the boundary conditions

$$u + \xi_l \sigma \frac{du}{d\bar{n}} = U_l, \quad x \in e_\ell, \ell = 1, \dots, L \quad (2)$$

$$\sigma \frac{du}{d\bar{n}} = 0, \quad x \in \partial\Omega \setminus \bigcup_{\ell=1}^L e_\ell \quad (3)$$

$$\int_{e_\ell} \sigma \frac{du}{d\bar{n}} dS = I_\ell, \quad \ell = 1, \dots, L \quad (4)$$

where  $\Omega$  is the target volume, and  $\partial\Omega$  is its boundary,  $\sigma$  is the electrical conductivity,  $u$  is the electric potential,  $\bar{n}$  is the outward unit normal,  $e_\ell$  is the  $l^{\text{th}}$  electrode, and  $\xi_l$ ,  $U_l$  and  $I_\ell$ , respectively, are the contact

impedance, electric potential and total current corresponding to  $e_l$ . Moreover, the current conservation law must be fulfilled

$$\sum_{l=1}^L I_l = 0 \quad (5)$$

and the potential reference level must be fixed, for example by writing

$$\sum_{l=1}^L U_l = 0. \quad (6)$$

The forward solution of EIT – computing the electrode potentials  $U_l$ , given the electrode currents  $I_l$  and the conductivity distribution  $\sigma$  – requires the solution of the CEM (1) - (6). In practice, the solution is approximated numerically, for example by using the Finite Element Method (FEM) [57, 58]. The finite element (FE) approximation was used also in the present work, leading to a numerical model  $U(\sigma)$  for the dependence between the electrode potentials and a discretized conductivity distribution  $\sigma$ . Further, assuming an additive Gaussian noise, the observation model for the EIT measurements becomes

$$V = U(\sigma) + e \quad (7)$$

where  $V$  is a vector consisting of the measured electrode potentials, and  $e$  is the Gaussian-distributed noise  $e \sim \mathcal{N}(0, \Gamma_e)$ . Here,  $\Gamma_e$  is the noise covariance matrix, which can usually be determined experimentally [59]. In principle, the reconstruction of the discretized conductivity distribution  $\sigma$  can be thought as a non-linear fitting problem spanned by model (7), i.e., finding a vector  $\sigma$  such that the computed voltages  $U(\sigma)$  fit to the voltage measurement data  $V$ . However, due to properties of the model  $U(\sigma)$ , this is mathematically an ill-posed inverse problem. This implies that classical solutions – such as the least-squares estimate for  $\sigma$  – are practically always unstable and non-unique [60]. In this work, we adopted the so-called Tikhonov regularization [57] to the solution. Furthermore, we introduced constraints for the conductivity values [47], and estimated the conductivity as

$$\hat{\sigma} = \arg \min_{\sigma_{\min} \leq \sigma \leq \sigma_{\max}} [\|L_e(V - U(\sigma))\|^2 + p(\sigma)] \quad (8)$$

where  $L_e$  is the Cholesky factor of the noise precision matrix  $\Gamma_e^{-1}$ , i.e.  $L_e^T L_e = \Gamma_e^{-1}$ , and  $p(\sigma)$  is a regularization function [61]. The function  $p(\sigma)$  is formulated such that it penalizes improbable features of  $\sigma$ . For the influence of  $p(\sigma)$  on the solution of the inverse problem, see e.g. [62, 63, 61]. In the present study,  $p(\sigma)$  was constructed to accommodate the so-called smoothness-promoting regularization, which is particularly applicable to diffusive processes [47], such as moisture flow in porous media. The smoothness-promoting function  $p(\sigma)$  can be written as  $p(\sigma) = \|L_\sigma \sigma\|^2$ , where  $L_\sigma$  is spatially weighted discrete differential operator [64, 60, 65].

The constraints ( $\sigma_{\min} \leq \sigma \leq \sigma_{\max}$ ) were selected based on experimentally determined physical limits for the electrical conductivity. First, the lower bound  $\sigma_{\min}$  was chosen to be equal to the conductivity corresponding to the initial moisture content  $\theta_i = 0.07$ . This choice was made based on the knowledge that the ingress of moisture does not decrease the conductivity of the mortar. Secondly, the upper bound  $\sigma_{\max}$  was chosen to be the electrical conductivity of the saturated material. The experiment for determining  $\sigma_{\min}$  and  $\sigma_{\max}$  using measurements of variably saturated mortar is reported below in Section 3.3.

The computations were performed using a MATLAB implementation which is an adaption of codes reported in [66, 67]. In the FE approximation of the forward model, a maximum element dimension of 1.75 mm was chosen, resulting in FE meshes with over 946,500 tetrahedral elements. For the Tikhonov regularized solution (8), Gauss-Newton method was used. The computations were performed utilizing parallel computing processes on eight quad-core processors using 32 Gb total memory and implementation of [68] to solve the systems of linear equations.

### 3.2. EIT measurement strategy

The EIT measurements were carried out using an in-house developed EIT equipment described in [15]. In all measurements, an alternating current with 0.10 mA amplitude and 40.0 kHz frequency was used. This frequency was chosen based on Electrical Impedance Spectroscopy (EIS) measurements; at frequency 40 kHz, the imaginary component of the impedance was at the minimum. The accuracy of the potential measurements was  $\pm 1.0 \times 10^{-8}$  V. In the EIT measurements, current was injected between electrodes  $i$  and  $j$  where  $i, j = \{1, \dots, 24\}$  and  $i \neq j$ . Corresponding to each current injection, potential measurements were taken with respect to a common ground. 561 current injections and 4488 potential measurements were taken for each set of EIT measurements.

### 3.3. Experimental material model of electrical conductivity for variably-saturated mortar

To quantify the moisture content using EIT, we used an experimental model of  $\theta = \theta(\sigma)$  that relates the electrical conductivity  $\sigma$  to the volumetric moisture content  $\theta$ . For this aim, we first determined the relationship between the electrical conductivity and the equilibrium relative humidity  $RH$  (i.e.,  $\sigma = \sigma(RH)$ ). Secondly, we measured the relationship  $\theta = \theta(RH)$  between the volumetric moisture content and the equilibrium relative humidity, and finally, combined these two models to get the function  $\theta = \theta(\sigma)$ . Note that the dependence of  $\theta$  on  $\sigma$  (or the converse) in cement-based materials was shown to not exhibit hysteresis in [69].

To experimentally determine  $\sigma = \sigma(RH)$ , a total of six prismatic, 6.0 mm  $\times$  1.5 mm  $\times$  4.0 mm mortar specimens were cut from a mortar cylinder using a precision diamond saw. Stainless electrodes were mounted to the 1.5 mm  $\times$  4.0 mm surface of the specimens using conductive silver-based epoxy. The specimens were then suspended in an air-tight container containing saturated salt solutions of KCl, NaCl,  $Mg(NO_3)_2$ ,  $K_2CO_3$ , and  $MgCl_2$  which induce  $RH$  conditions of 84.0, 75.3, 52.8, 43.6, and 32.3%, respectively [70, 4]. Another specimen was submerged in water to reach saturation. After 5 days of conditioning, the electrical resistances of the specimens were measured in 24-hour intervals until no change in resistance was observed between three consecutive measurements. The electrical resistance  $R$  was then converted to conductivity using the relation  $\sigma^{-1} = R \frac{A}{L}$ , where  $A$  and  $L = 6.0$  mm are, respectively, the cross-sectional area and the length of the specimen.

The desorption isotherm,  $\theta = \theta(RH)$ , was experimentally obtained using an automated sorption analyzer [71, 72]. The function  $\theta = \theta(\sigma)$  was then obtained by a piecewise cubic interpolation [73] to the above mentioned experimental data. This interpolation method was selected to (i) create continuity between intermediate values of  $\theta$  and (ii) preserve the curve shape at high levels of moisture saturation. The experimentally obtained data and fitted curve are shown in Figure 2.

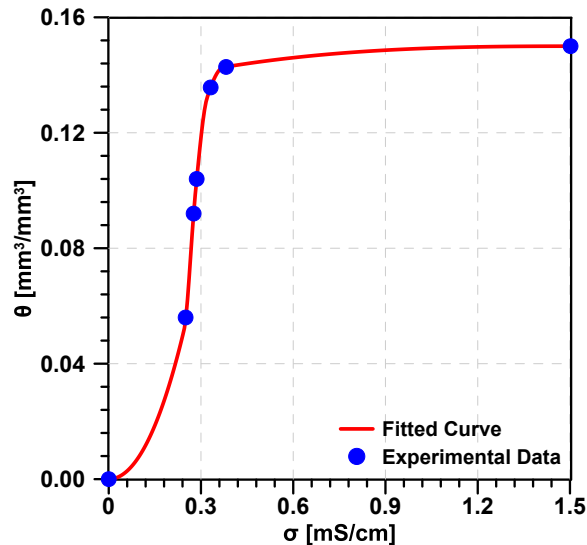


Figure 2: Fitted moisture retention curve for mortar using van Genuchten parameters

Using the initial moisture content  $\theta_i = 0.07$ , the moisture content at saturation  $\theta_s = 0.15$  and the above described model  $\theta = \theta(\sigma)$ , the minimum and maximum conductivity of the material were computed, resulting in values  $\sigma_{\min} = 0.28$  mS/cm and  $\sigma_{\max} = 1.5$  mS/cm. These values were used as the lower and upper bound in Eq. (8).

#### 4. X-ray CT imaging

To corroborate EIT reconstructions, X-ray Computed Tomography (CT) imaging was conducted using a full-scale medical CT scanner with a peak beam energy of 120 kV and intensity of 40-200 mA (Fig. 3). In the CT scanner, the X-ray tube rotates perpendicularly to the circular cross-section of the specimen. As the fan beam rotates around the specimen, detectors on the opposite side are activated and measurements are taken. Upon completing a full rotation, the X-ray absorption coefficient  $\mu(\text{mm}^{-1})$  is computed in voxels within the measured slice. X-ray radiation transmitted through the specimen is related to  $\mu$  via the Lambert-Beer law:

$$I_x = I_0 \exp\left(-\int \mu d\ell\right) \quad (9)$$

where  $I_x$  is the intensity of the transmitted X-ray passing through the specimen,  $I_0$  is the incident X-ray intensity, and  $\ell$  is the arc length variable along the line connecting the X-ray source and detector.

The voxel values of the CT images,  $\mu - \mu_{\text{water}}$ , were normalized by  $\mu_{\text{water}} - \mu_{\text{air}}$  to get Hounsfield numbers  $H$  is defined by [74]:

$$H = 1000 \times \frac{\mu - \mu_{\text{water}}}{\mu_{\text{water}} - \mu_{\text{air}}}. \quad (10)$$

Further, the spatial distribution of the volumetric moisture content  $\theta$  was computed from the Hounsfield numbers [75]:

$$\theta = \frac{H_s - H}{1000}, \quad (11)$$

where  $H_s$  is the Hounsfield value of a saturated material. Here, the value of  $H_s$  was taken from H values corresponding to voxels directly below the water reservoir, where the material was assumed to be completely saturated [72, 8].

In this experiment, only the specimen where the absorbed water contained the radiocontrast agent (Iohexol) was imaged with X-ray CT. As pointed out in [36], the sensitivity of CT does not allow for monitoring the absorption of plain water in this medium and geometry. The specimen was scanned at times 1, 2, 4 and 22 hours after the start of the water ingress. 2D slices were taken through the horizontal and vertical cross sections in 0.6 mm slices. CT images were analyzed using digital imaging software [76].

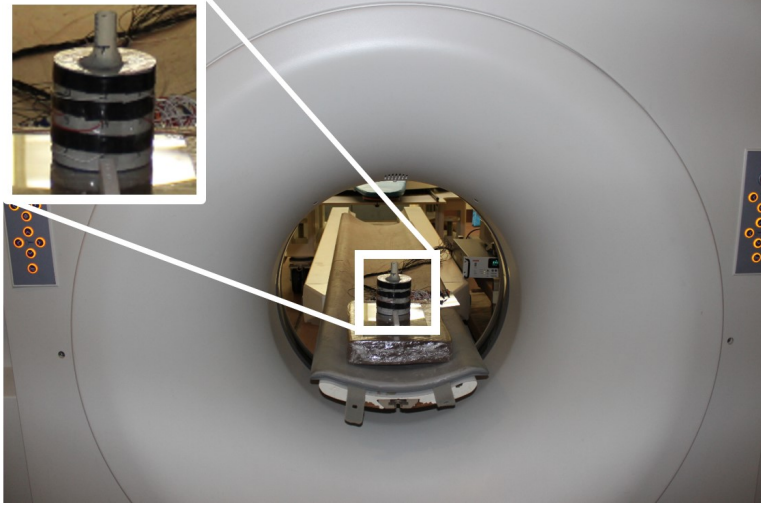


Figure 3: Full-scale medical CT scanner with a peak beam energy of 120 kV and intensity of 40-200 mA testing the specimen used in this study.

## 5. Numerical simulation of unsaturated moisture flow

### 5.1. General

200 Isothermal unsaturated moisture flow in porous media was simulated using Richards' Equation [77, 78]:

$$\frac{\partial \theta}{\partial t} = \frac{\partial}{\partial x} \left( K(h) \frac{\partial h}{\partial x} \right) + \frac{\partial}{\partial y} \left( K(h) \frac{\partial h}{\partial y} \right) + \frac{\partial}{\partial z} \left( K(h) \frac{\partial h}{\partial z} + 1 \right) \quad (12)$$

where  $K = K(h)$  (mm/hour) is the unsaturated hydraulic conductivity,  $\theta$  (mm<sup>3</sup>/mm<sup>3</sup>) is the volumetric moisture content,  $h$  (mm) is the pressure head, and  $x$ ,  $y$ , and  $z$  (mm) are spatial coordinates. In this work finite element approximation of Eq. (12) was used to simulate the moisture flow in mortar.

### 5.2. Material model

205 In general, the unsaturated hydraulic conductivity is expressed as the product of the relative hydraulic conductivity  $K_r$  and the saturated hydraulic conductivity  $K_s$ , such that  $K = K_r K_s$ . For unsaturated porous media,  $K_r$  is described by Mualem's equation [79]:

$$K_r = \Theta^I \left[ \frac{\int_0^\Theta \frac{1}{h(x)}, dx}{\int_0^1 \frac{1}{h(x)}, dx} \right]^2 \quad (13)$$

$$\Theta = \frac{\theta - \theta_r}{\theta_s - \theta_r} \quad (14)$$

210 where  $0 \leq \Theta \leq 1.0$  is the effective material saturation,  $\theta_r$  is the residual moisture content where  $\theta_r = 0$  is generally considered as an appropriate assumption for cement-based materials [72, 7],  $\theta_s$  is the saturated moisture content that is measured experimentally in this work ( $\theta_s = 0.15$ ),  $I$  is the tortuosity and pore connectivity parameter; in this work  $I = -9.0$  is taken from [80]. Discussion of  $I$  in cement-based materials is provided in [81] and [82]. In this work, saturated hydraulic conductivity  $K_s = 5.0 \times 10^{-3}$  (mm/hr) is obtained experimentally using Darcy's Law and the procedure in [70].

215 In order to calculate the integrals in Eq. (13), the effective material saturation is expressed as a function of the pressure head,  $\Theta = \Theta(h)$  using the van Genuchten model [83, 84]:

$$\Theta = \frac{1}{[1 + (\alpha h)^n]^m}, \quad m = 1 - \frac{1}{n} \quad (15)$$

where  $\alpha$  ( $\text{mm}^{-1}$ ) and  $n(-)$  are fitting parameters. For cement-based materials, instead of expressing water retention as  $\Theta = \Theta(h)$ , the material sorption isotherm is generally obtained experimentally as  $\Theta = \Theta(RH)$ , where  $RH$  denotes relative humidity. In this work, the desorption isotherm was experimentally measured using an automated sorption analyzer [71, 72] (see Section 3.3) For simplicity we neglected the effect of hysteresis. To convert the experimentally obtained isotherm  $\Theta(RH)$  to the moisture retention curve as  $\Theta(h)$ , the Kelvin-Laplace Equation is used [7, 85]:

$$h = \frac{\rho_w R_s T}{m_w} \ln(RH) \quad (16)$$

where  $\rho_w$  ( $\text{g}/\text{mm}^3$ ) is the density of water,  $R_s$  ( $8.845 \text{ JK}^{-1}\text{mol}^{-1}$ ) is the universal gas constant, and  $m_w$  ( $\text{g}/\text{mol}$ ) is the molecular weight of water. By fitting the model in Eq. (15) to the water retention curve  $\Theta(h)$ , van Genuchten parameters  $\alpha = 2.63 \times 10^{-2}$  ( $\text{mm}^{-1}$ ) and  $n = 1.77$  were obtained. The experimentally obtained data points and the fitted van Genuchten model are shown in Figure 4.

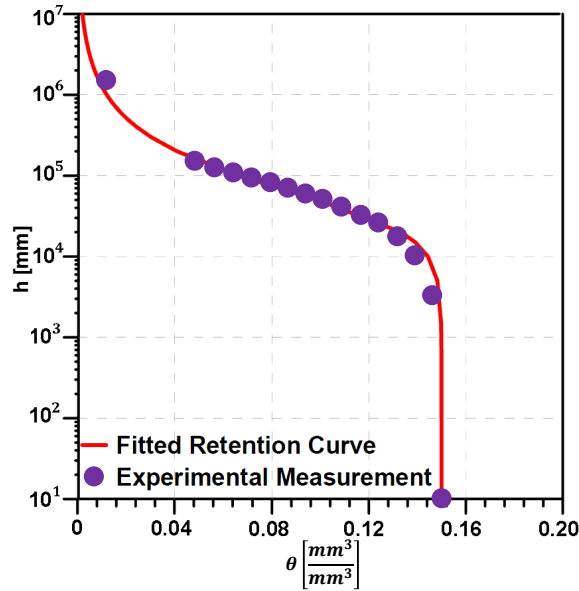


Figure 4: Fitted moisture retention curve for mortar using van Genuchten parameters

### 5.3. Simulation of moisture flow

The moisture flow in the specimen was modeled using a commercially available Finite Element Software HYDRUS 3D [86]. Zero-flux boundary conditions were considered at all surfaces except the surface from which water penetrated (water reservoir), which was modeled as a 9.5 mm diameter circle saturated throughout the simulation. Uniform initial moisture content,  $\theta_i = 0.07$  was selected based on the experiments described in Section 2.2. The Finite element mesh consisted of tetrahedral elements with a maximum dimension of 4 mm. The finite element model was solved in terms of moisture content.

## 6. Results and discussion

Figures 5 and 6 show the results of the experiment where the dilute solution of water and Iohexol was absorbed to a mortar specimen. The left column of Fig. 5 illustrates the 3D electrical conductivity distribution reconstructed based on EIT measurements after 1, 2, 4 and 22 hours of moisture penetration. In the images, the black rectangles mark the electrodes attached on the surface of the specimen. The series of EIT reconstructions clearly shows the ingress of moisture: After 1 hour of moisture ingress, a small volume directly under the water reservoir features increase of the electrical conductivity; in the subsequent time intervals, the volume with increased conductivity gets larger. Also the value of the conductivity under the

water reservoir increases over time, being close to conductivity at saturation,  $\sigma_s = 1.5$  mS/cm, after 22 hours of moisture ingress.

The EIT-based estimates for the 3D volumetric moisture content  $\theta$  at times 1, 2, 4 and 22 hours are shown in the right column of Fig. 5. These moisture content estimates were computed from the reconstructed 3D conductivity distributions  $\sigma$  using the experimentally developed model  $\theta = \theta(\sigma)$  described in Section 3.3. While the reconstructed conductivity distributions in the left column illustrate the ingress of moisture through the specimen qualitatively, the 3D volumetric moisture content estimates allow for quantifying the moisture ingress. For example, the reconstructions of  $\theta$  show that the moisture content is approximately 0.12 directly below the water reservoir after 1 hour of moisture ingress, and is nearly saturated ( $\theta_s = 0.15$ ) after 22 hours of moisture ingress, while the background moisture content remains at the initial value 0.07. Qualitatively, the spatial and temporal gradients of  $\theta$  look somewhat different from those of  $\sigma$ . This difference results from the non-linearity of the model  $\sigma = \sigma(\theta)$ .

The left column of Fig. 6 represents the EIT-based reconstructions of the 3D volumetric moisture content  $\theta$  on a vertical 2D plane crossing the center of the water reservoir and extending to the exterior of the specimen. The right column shows X-ray CT images at the respective times. In these images, the water reservoir appears in the center of the top surface, because the 2D slice crossing the center of the water reservoir is off-centric with respect to the specimen; this is due to the orientation of the sample during the CT scan (Fig. 3). Also the X-ray CT images show the increase of  $\theta$  directly below the water reservoir, and the value of  $\theta$  in this region is close to the saturation value 0.15 at times 1 h and 2 h. However, in the 4 h and 22 h CT images, the value of the moisture content has decreased under the water reservoir, indicating a low sensitivity of CT to the water content with this specimen size. Moreover, the X-ray CT images presented here detect only high levels of saturation (similarly to X-ray images in [7]). The CT images do not show the presence of low levels of saturation because of the large sample size used (10.20 cm  $\times$  20.30 cm cylinder). Some previous studies using X-ray imaging (e.g. [9] and [10]) showed high resolution images of unsaturated moisture flow in cement-based material; however, the specimen sizes were small (maximum dimension of 5.5 cm in cement-based specimen tested). These results demonstrate that EIT is capable of imaging the moisture flow in a specimen, the size of which exceeds the limit of a regular CT scanner.

Figure 7 shows the results corresponding to the experiment where plain water was absorbed into a mortar specimen. The left column illustrates the volumetric moisture content  $\theta$  reconstructed using EIT, and the right column shows  $\theta$  computed by moisture flow simulation specified in Section 5.3. All images represent the vertical 2D cross-sections of the corresponding 3D distributions. The EIT reconstructions show a good correspondence with the moisture flow simulation. Indeed, the ingress of moisture predicted by the flow simulation is captured well by EIT, and the size and shape of the (nearly) saturated volume  $\theta \approx 0.15$  in EIT-based images are close to those in the simulation-based images. Also the shapes of intervals with lower saturation in EIT-based reconstructions of  $\theta$  compare relatively well with those in the simulations.

The addition of the radiocontrast agent (Iohexol) increases the viscosity and surface tension of water, and thus, the flow of the diluted solution in porous media is slower than the flow of plain water. The capability of EIT to distinguish between flow rates of these two liquids was shown in [36]. The comparison between the left columns of Figs. 6 and 7 further confirms this; especially, the volume of the saturated mortar is consistently larger in Fig. 7 than in the respective times in Fig. 6.

## 7. Summary and conclusions

In this paper, we investigated whether Electrical Impedance Tomography (EIT) could be used to quantitatively monitor 3D moisture flow in cement-based materials. For this aim, we conducted an experiment where mortar specimens were imaged with EIT during the ingress of moisture. The so-called absolute imaging scheme was used to reconstruct the 3D distributed electrical conductivity within the mortar. The reconstructed electrical conductivity distributions were further converted to volumetric moisture content distributions, by using an experimentally-developed model. One of the samples was simultaneously imaged by a regular X-ray CT scan. Moreover, the results of EIT were corroborated by a numerical simulation of unsaturated moisture flow.

EIT captured the moisture flow in mortar and provided quantitative information on volumetric moisture content. The EIT-based estimates for the moisture content compared well with numerical simulations of unsaturated moisture flow. In contrast to X-ray CT images, which showed only highly saturated material,

EIT reconstructions provided a good resolution at various levels of saturation. In the previous work [36], the difference imaging scheme was used to visualize moisture flow in cement-based materials, and significant artifacts were observed at the late stages of moisture flow. In the present work, however, the use of application-specific absolute imaging eliminated the artifacts in the reconstructions at late stages of moisture ingress, since, unlike difference imaging, absolute imaging is not based on the assumption of small changes in the electrical conductivity with regards to a reference state. Based on these findings, EIT shows promise for becoming a tool for imaging moisture flow in cement-based materials non-destructively and quantitatively. Moreover, the developed procedure for quantifying the moisture content using EIT is also applicable for other porous media.

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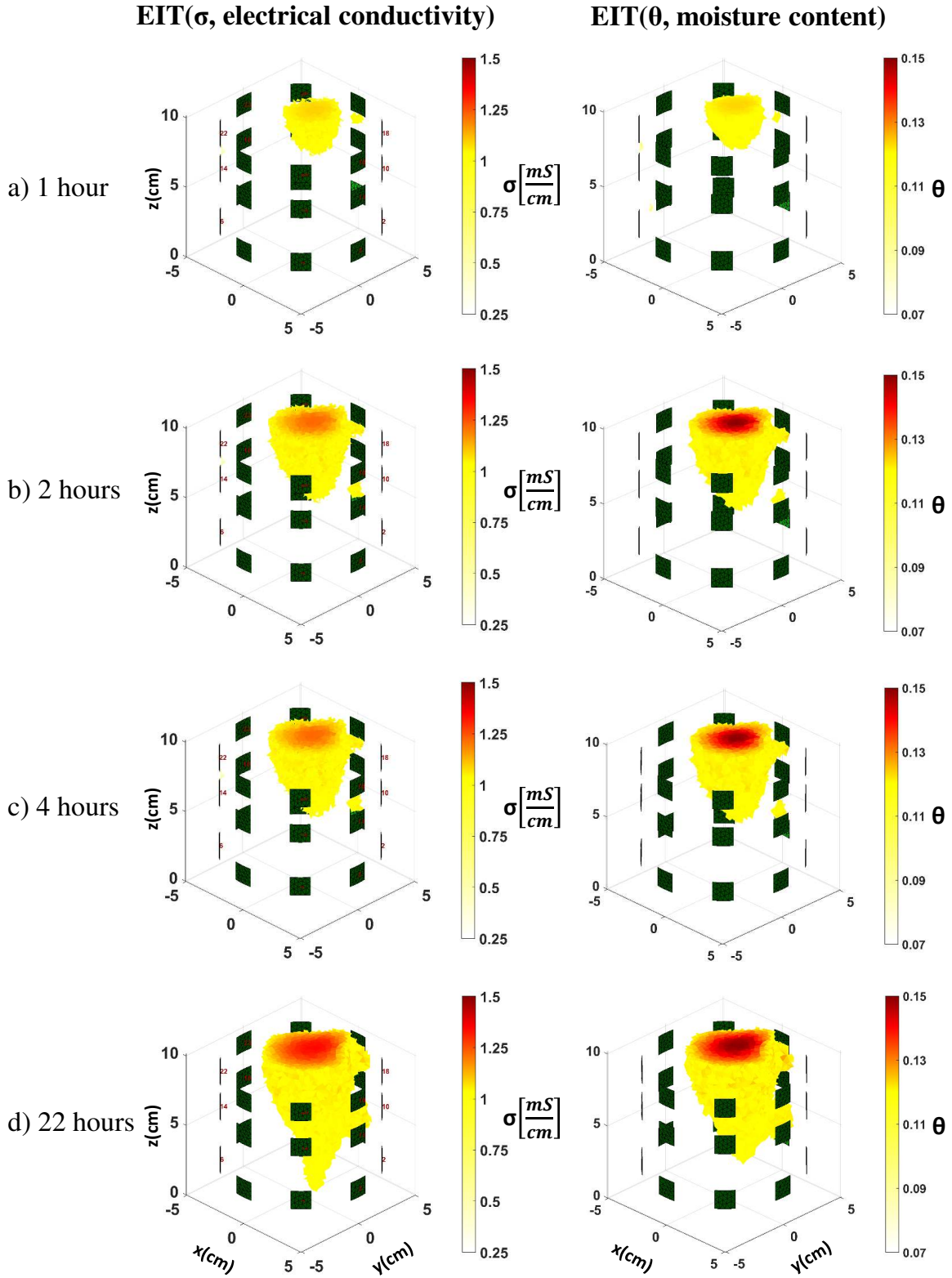
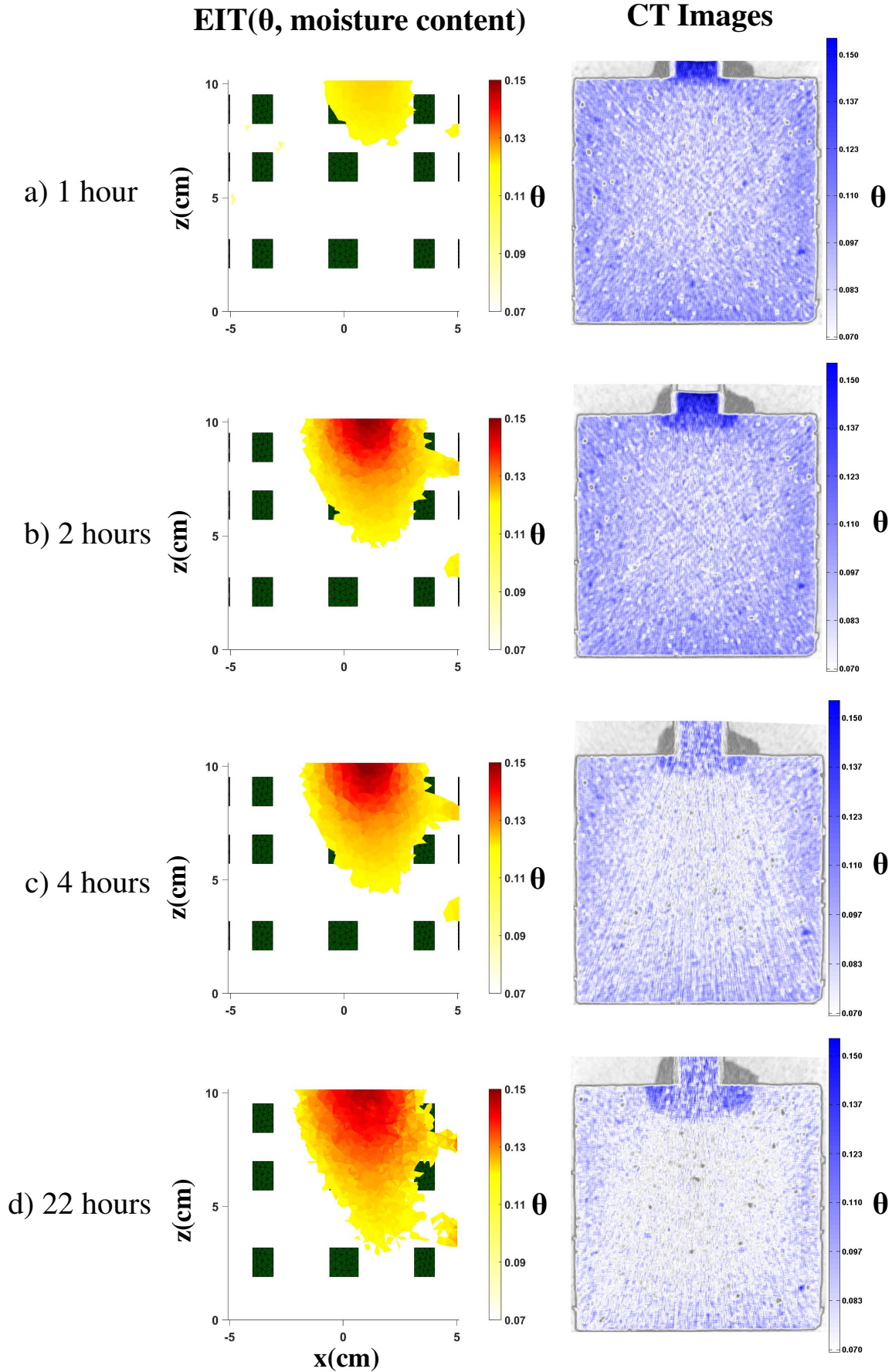


Figure 5: Results of the experiment with dilute solution of water and Iohexol; 3D EIT reconstructions depicting electrical conductivity ( $\sigma$ , left column) and volumetric moisture content ( $\theta$ , right column) for moisture ingress at: (a) 1 hour, (b) 2 hours, (c) 4 hours (d) 22 hours.



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 Figure 6: Results of the experiment with dilute solution of water and Iohexol; 2D slices of the EIT-based 3D reconstructions of the volumetric moisture content (left column) and CT-based images of the volumetric moisture content (right column) at: (a) 1 hour, (b) 2 hours, (c) 4 hours (d) 22 hours

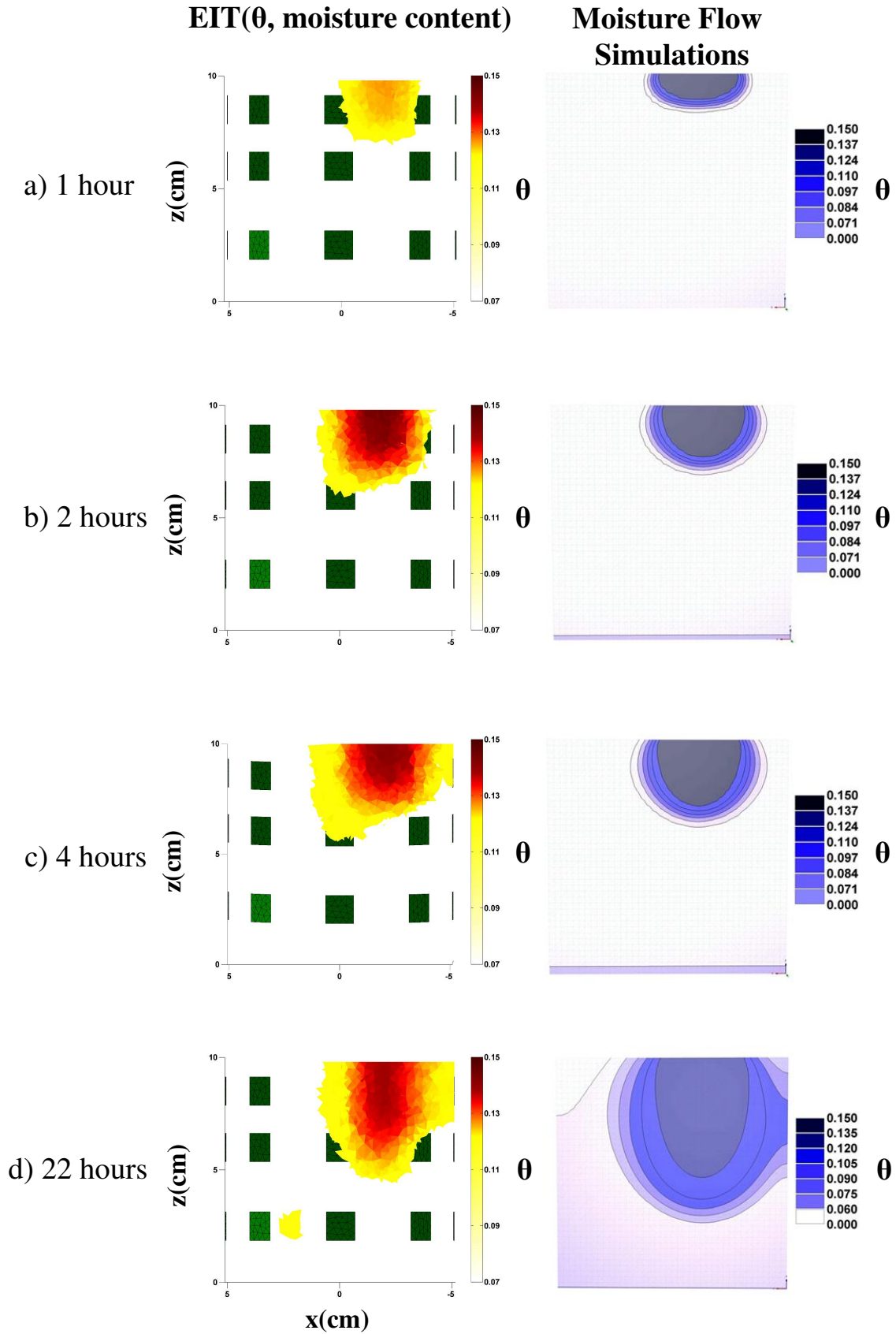


Figure 7: Results of the experiment with plain water; 2D slices of the EIT-based 3D reconstructions of the volumetric moisture content (left column) and moisture flow simulation -based images of the volumetric moisture content (right column) at: (a) 1 hour, (b) 2 hours, (c) 4 hours (d) 22 hours