



This is a repository copy of *Slow episodic movement driven by elevated pore-fluid pressures in shallow subaqueous slopes*.

White Rose Research Online URL for this paper:  
<http://eprints.whiterose.ac.uk/142221/>

Version: Accepted Version

---

**Article:**

Carey, J.M., Crutchley, G.J., Mountjoy, J.J. et al. (3 more authors) (2019) Slow episodic movement driven by elevated pore-fluid pressures in shallow subaqueous slopes. *Geomorphology*, 329. pp. 99-107. ISSN 0169-555X

<https://doi.org/10.1016/j.geomorph.2018.12.034>

---

Article available under the terms of the CC-BY-NC-ND licence  
(<https://creativecommons.org/licenses/by-nc-nd/4.0/>).

**Reuse**

This article is distributed under the terms of the Creative Commons Attribution-NonCommercial-NoDerivs (CC BY-NC-ND) licence. This licence only allows you to download this work and share it with others as long as you credit the authors, but you can't change the article in any way or use it commercially. More information and the full terms of the licence here: <https://creativecommons.org/licenses/>

**Takedown**

If you consider content in White Rose Research Online to be in breach of UK law, please notify us by emailing [eprints@whiterose.ac.uk](mailto:eprints@whiterose.ac.uk) including the URL of the record and the reason for the withdrawal request.



[eprints@whiterose.ac.uk](mailto:eprints@whiterose.ac.uk)  
<https://eprints.whiterose.ac.uk/>



1        1. INTRODUCTION

2        Although subaqueous mass movements can develop on very gentle slopes ( $\ll 2^\circ$ ), they  
3        can run out for extremely long distances across the deep ocean floor (Talling et al.,  
4        2007). In most cases the slope failures leave an empty scar in the source region with  
5        almost all residual strength material evacuated (Krastel et al., 2018; Mountjoy and  
6        Micallef, 2018). As source area slopes are predominantly lower than the friction angles  
7        of the landslide materials (e.g. Urlaub et al., 2015), the generation of pore pressures  
8        exceeding hydrostatic pressure are a likely loading mechanism. But, in some instances,  
9        the debris from subaqueous slope failures only moves a short distance before coming  
10       to rest despite a lack of buttressing or obvious decrease in slope gradient (Locat and  
11       Lee, 2000; Micallef et al., 2013). Such small-displacement mass failures have been  
12       found in a wide range of seafloor environments including passive margins (Baeten et  
13       al., 2014) and active margins (Micallef et al., 2016). These indicate that there are likely  
14       to be common mechanisms for the movement and arrest over a short distance of certain  
15       seafloor failures.

16  
17       The high pore pressures that often initiate movement can occur from widely recognised  
18       processes including undrained cyclic loading during earthquakes (e.g. Sassa et al.,  
19       2012), rapid sediment burial (e.g. Stigall and Dugan, 2010) and as result of focused  
20       fluid flow (e.g. Dugan and Flemings, 2000; Elger et al., 2018). In addition, high pore  
21       pressures may be generated by more complex processes involving for example gas  
22       liberation from hydrate dissociation (Riboulot et al., 2013).

23  
24       Despite the advances in understanding potential causes of submarine landslides, the  
25       fundamental processes controlling their movement remain poorly constrained

26 compared with their terrestrial counterparts. Over the last century, terrestrial landslides  
27 have been shown to display a wide array of movement behaviour, ranging from slow  
28 creep ( $\leq 1 \text{ mm a}^{-1}$ ) (e.g. Mansour et al., 2011), through episodic sliding and stick-slip  
29 ( $\approx 1 \text{ cm a}^{-1}$ ) (e.g. Allison and Brunsden, 1990) to acceleration to catastrophic failure  
30 ( $\gg 1 \text{ m s}^{-1}$ ) (e.g. Kilburn and Petley, 2003). Such complex movements are commonly  
31 associated with pore fluid pressure-induced changes, as has been hypothesised for  
32 subaqueous landslides.

33

34 The availability of high resolution landslide monitoring records onshore has allowed  
35 detailed acceleration phases in multiple landslides to be examined in a variety of  
36 materials (e.g. Schulz et al., 2009; Massey et al., 2013; Carey et al., 2015). These studies  
37 distinguish two distinct styles of movement (Petley et al., 2002). The first movement  
38 style is brittle, characterised by a distinct hyperbolic acceleration in displacement rate  
39 to failure. These movement patterns can be examined by plotting in  $1/v - t$  space (where  
40  $v$  is velocity) and yield a negative linear trend to failure which generally results in rapid  
41 accelerations and catastrophic landsliding (e.g. Voight, 1988; Fukuzono, 1990; Petley  
42 and Petley, 2006). Conversely, the second style of movement is ductile, characterised  
43 by an exponential acceleration to a constant strain rate which when analysed in  $1/v - t$   
44 space produces a distinct asymptotic trend (Petley et al., 2005b).

45

46 Specialist laboratory testing approaches have been used to simulate landslide failure  
47 conditions by increasing pore water pressures at constant normal and shear stress (e.g.  
48 Brand, 1981; Anderson and Sitar, 1995; Zhu and Anderson, 1998; Dai et al., 1999;  
49 Orense et al., 2004; Petley et al., 2005a, Ng and Petley, 2009; Carey and Petley, 2014).  
50 These approaches confirm that movement styles are controlled by mechanisms of

51 deformation occurring within shear zones (Petley et al., 2005a; Ng and Petley, 2009;  
52 Carey and Petley, 2014). In cases where brittle shear surface development occurs, the  
53 hyperbolic acceleration to final failure is observed in the experiments, whilst  
54 exponential acceleration occurs in landslides undergoing ductile deformation (Petley et  
55 al., 2005a). Numerous conceptual models have been proposed to explain these different  
56 strain responses to changes in stress state, ranging from catastrophic failure driven by  
57 micro-cracking and rapid shear-surface propagation (e.g. Petley et al., 2005a; Viesca  
58 and Rice, 2012) to slower, steady landslide motion in response to shear-zone dilation  
59 and subsequent pore-pressure feedback (e.g. Iverson, 2005).

60

61 Given increases in pore-water pressure may drive either of movement-arrest or rapid  
62 runout behaviour in terrestrial landslides, depending on the material response, similar  
63 behaviour should be expected in subaqueous landslides. The mechanisms controlling  
64 the transition from steady slow movement to rapid failure observed in specialist  
65 laboratory tests may be key to determining the behaviour of a given subaqueous  
66 landslide. However, very few small displacement or reactivated subaqueous landslides  
67 have been observed and investigated.

68

69 Two subaqueous landslides on the northern Hikurangi Margin located off the east coast  
70 of New Zealand have geomorphological characteristics that indicate movement-arrest  
71 behaviour may be occurring (Mountjoy et al., 2009; Micallef et al., 2016) making them  
72 suitable for investigating this behaviour in subaqueous slopes. In this study, we seek to  
73 constrain better the potential movement mechanisms in short displacement and  
74 reactivated subaqueous landslides by conducting novel laboratory experiments on  
75 sediment samples collected from the shallow sedimentary sequence on the northern

76 Hikurangi Margin area. We use a Dynamic Back-Pressured Shearbox (DBPSB) to  
77 accurately replicate in-situ stresses in the submarine slopes to explore the potential  
78 strain response of the landslide when subject to elevated pore water and gas pressure.

79

## 80 2. SAMPLE AREA

81 An area of the upper continental slope on the Hikurangi Subduction Margin, off the  
82 coast of Gisborne, New Zealand, hosts a number of landslides where limited  
83 displacement has occurred following their initial failure (Mountjoy et al., 2009;  
84 Micallef et al., 2016) or where repeated reactivation has been hypothesised (Mountjoy  
85 et al., 2014) (Figure 1 A and B). The landslides occur within an active subduction zone  
86 experiencing regular tectonic activity (Wallace and Bevan, 2010; Wallace et al., 2012).  
87 The upper continental slope is comprised of Miocene to Recent slope basin sequences  
88 (Mountjoy et al., 2009). A gravity core profile down the length of the slides (the  
89 Tuaheni Landslide Complex) shows that the upper few metres of sediment are  
90 dominated by mud to sand sized particles from hemipelagic drape, reworked landslide  
91 debris and airfall tephra (Kuhlmann et al., 2018).

92

93 The Tuaheni Landslide Complex (TLC) comprises an area of approximately 145 km<sup>2</sup>  
94 which is sub-divided into Tuaheni North and Tuaheni South, separated by a 2 km wide  
95 spur (Figure 1 A). While Tuaheni North is characterised by multiple evacuated  
96 landslide scarps, Tuaheni South is characterised by a large debris apron which has a  
97 distinct scarp and bench topography and features indicative of lateral, extensional and  
98 compressional deformation (Mountjoy et al., 2009; 2014). The base of gas hydrate  
99 stability has been imaged beneath the TLC's extensional domain and it has been  
100 suggested that the gas hydrate system may play a role in deformation of the landslide

101 mass (Mountjoy et al., 2014; Crutchley et al., 2016). Although no movement  
102 measurements are available, the morphology of Tuaheni South is similar to slow-  
103 moving landslide complexes observed in terrestrial environments such as earthflows  
104 and mudslides (e.g. Hungr et al., 2014) which occur in similar fine-grained sediments  
105 and are often subject to episodic remobilisation (e.g. Alison and Brunsden, 1990).  
106 Seismic-reflection surveys across the landslide immediately north of TLC (Figure 1 C)  
107 indicate that there is free gas in the proto basal failure surface, and adjacent to the  
108 landslide, but gas is not observed beneath the landslide body (Micallef et al., 2016).  
109 This suggests that free gas may have been present in the slope prior to the landslides  
110 and that the gas migrated out of the sediment sequence during and/or after failure.  
111 Micallef et al., (2016) concluded that overpressure in the slope sequence may have  
112 contributed to bringing the slope to the point of failure, and that once the landslide  
113 moved and dilated, the gas pressure reduced, and further failure was arrested. To further  
114 test this interpretation required measurement of relevant geotechnical data, and this  
115 provided an opportunity to assess different mechanisms that result in small-  
116 displacement slope failures.

117

### 118 3. MATERIALS AND METHODS

119 We performed a suite of conventional laboratory experiments to determine the physical  
120 and geomechanical characteristics of the shallow materials collected from the landslide  
121 complex (Tables 1 and 2). Sediment samples were collected from the crest of the  
122 landslide (Tan1404-02) and from within an extensional domain farther downslope  
123 (Tan1404-10) (Figure 1 A) using a 100 mm diameter gravity corer during RV Tangaroa  
124 voyage TAN1404 in April/May 2014.

125

126 These standard soil classification test results indicate that both materials have similar  
127 physical properties (Table 1). Natural water contents were higher in the shallow  
128 EN1285 samples (Tan1404-02) and had correspondingly higher void ratios and lower  
129 dry densities when compared with deeper EN1287 samples (Tan1404-10). Particle-size  
130 analyses (Figure 2) confirm that both materials are fine grained comprising of over 88  
131 % silt, approximately 5 % clay and with the remaining fraction consisting of mostly  
132 fine and medium grained sand. The Atterberg Limit tests performed for both samples  
133 confirmed similar plastic and liquid limits and universal soil classifications at the  
134 boundary of high plasticity silts and clays (Table 1). The results indicated that the  
135 physical properties of the sediments were similar both within the downslope extensional  
136 domain and above the current landslide crest, indicating that any sediment disturbance  
137 during sampling would have negligible impact on our study.

138

139 Conventional drained direct shear tests were undertaken on 60 mm diameter circular  
140 samples of both materials using a Wykeham Farrance direct shearbox WF2500.  
141 Shearing was conducted at three low normal stresses (Table 2) to simulate the shallow  
142 depths of burial of the samples. Shearing was initiated on completion of the  
143 consolidation phase (i.e., no further significant vertical displacement at the desired  
144 normal load for shearing). A slow shear rate ( $1.83 \times 10^{-4} \text{ mm s}^{-1}$ ), was used to avoid  
145 developing excess pore pressures within the specimens and a minimum of five shear  
146 reversals was completed for each test to ensure a representative 'residual' shear strength  
147 was reached (i.e., no further reduction in shearing stress on cyclic loading).

148

149 A suite of specialist pore-pressure reinflation experiments was performed in the  
150 Dynamic Back-Pressured Shearbox (Figure 3), constructed by GDS Instruments Ltd

151 and described in detail by Brain et al. (2015) and Carey et al. (2016). Previous studies  
152 have demonstrated that the DBPSB is able to induce a variety of styles of deformation  
153 ranging from dynamic liquefaction (Carey et al. 2017) to creep (Carey et al. 2016). Each  
154 sample was saturated prior to testing to replicate the shallow seabed conditions (Table  
155 3) using methods previously described (see Carey et al., 2016). A normal effective  
156 stress of 32 kPa was applied to each sample during consolidation by applying a back  
157 pressure of 100 kPa and a total normal stress at 132 kPa.

158

159 After consolidation, each sample was sheared to failure at a constant strain rate ( $1.83 \times$   
160  $10^{-4} \text{ mm s}^{-1}$ , chosen to avoid excess pore fluid generation) to form a shear zone  
161 representative of the base of a shallow subaqueous landslide, and to measure initial  
162 shear strength. A shear-stress of c.70 % of the undrained strength was then applied to  
163 each sample, whereupon both total normal and total shear stress were held constant  
164 whilst the normal effective stress was reduced by increasing pore pressure until the  
165 samples failed (Figure 4). The experiments were conducted using two different pore  
166 fluid conditions: water, and water plus nitrogen gas.

167

168 Pore-water pressure controlled tests (referred to as PWP) were conducted by linearly  
169 increasing the back pressure applied to the sample whilst holding both the total normal  
170 stress and shear stress constant (Figure 3B). De-aired water was used to ensure that the  
171 pore-water pressure increases and changes in fluid movement pathways anticipated in  
172 the shallow seabed were accurately simulated.

173

174 Pore-gas pressure controlled experiments (referred to as PGP) were conducted using  
175 nitrogen gas because it has similar physical properties to methane for the conditions of

176 our tests (Kossel et al. 2013), but is safer to use than methane. PGP experiments were  
177 performed by filling the volume controller (VC1) with nitrogen gas whilst a second  
178 volume controller (VC2) was used to maintain the back pressure to the sample (Figure  
179 3C). VC1 was raised to the same pressure as the vessel before being connected to the  
180 sample. Once the pressures were equal, the gas volume controller was opened to the  
181 base of the sample and monitored to ensure that no significant pressure change occurred  
182 in either volume controller, and that no sample strain occurred prior to the test run. The  
183 pressure in the gas volume controller was then increased linearly at rates of 12 kPa/hr  
184 and 30 kPa/hr to replicate increasing gas pressure within a shear zone. The porous disc,  
185 installed at the sample base was replaced with a specially designed gas test plug prior  
186 to testing to ensure gas pressures would be applied to the sample shear zone (Figure 3  
187 C). To ensure that the gas was replacing the water within the soil pores, VC1 was  
188 allowed to increase in volume (to extract water) as gas pressure was increased (Figure  
189 3 C). These PGP tests using the DBPSB were the first of their kind, representing a new  
190 methodology for testing the impact of gas pressure on sediment failure. Both the PWP  
191 and PGP testing approaches resulted in similar stress paths to failure in which  
192 increasing pore-fluid pressures resulted in a reduction in mean effective stress toward  
193 failure while shear stress and total normal stresses remained constant (Figure 4).

194

195 During each phase of pore-fluid pressure increase, shear strain was monitored by  
196 measuring the horizontal (shear) displacement of the shear box. These experiments  
197 simulated the generation of excess pore-fluid pressure, and associated shear-strain  
198 response, in near surface subaqueous landslides under a representative stress state.

199

200 4. LABORATORY RESULTS

201 The conventional shear box tests indicated no notable peak strength or shear strength  
202 reduction following repeated shear reversals (Table 2) consistent with ductile  
203 behaviour. Shear tests at three confining pressures indicated that the sediments had a  
204 linear drained failure envelope, with similar friction angles which ranged between  $34^\circ$   
205 and  $36^\circ$  with an effective cohesion of 4.0 kPa (Figure 5). The strength characteristics  
206 measured in the landslide materials were comparable, with similar materials tested from  
207 other fine-grained subaqueous landslide systems (e.g. Sassa et al., 2012).

208

209 The horizontal (shear) strain and vertical (axial) strain behaviour were found to be  
210 broadly consistent in all three PWP tests regardless of rate of pore water pressure  
211 increase (Figure 6 A and B). In each experiment horizontal deformation progressed  
212 through three movement phases as the shear zone dilated; these three phases can be  
213 observed by distinct changes in horizontal strain rate (Figure 7).

214

215 An initial low strain rate ( $< 0.05 \mu\text{S s}^{-1}$ ) was observed in each sample whilst the mean  
216 effective stress ( $p'$ ) remained high (Figure 7). The strain rate of this early phase of  
217 shear-zone deformation increased exponentially in response to the reducing mean  
218 effective stress and corresponding sample dilation. This initial phase was followed by  
219 a distinct sliding phase which was characterised by a rapid increase and then decrease  
220 in strain rate (Figure 7). This movement was not associated with a measured change in  
221 bulk sample pore-water pressure and was observed in all three samples, although more  
222 pronounced in experiments PWP 5 and PWP 3. The sliding phase developed at a similar  
223 normal effective stress (c.10 kPa) in each experiment, regardless of the rate of pore-  
224 pressure increase. The continued decrease in mean effective stress following this phase  
225 resulted in a progressive increase in horizontal strain which comprised of distinct

226 periods of increasing strain rate punctuated by periods of reducing strain rate. This  
227 cyclic strain phase produced an exponential increase in strain rate with reducing mean  
228 effective stress (Figure 7).

229

230 The development of a short, near-instantaneous, sliding phase followed by repeating  
231 cyclic strain suggested that the shear zone mobilised and sheared at a critical mean  
232 effective stress without acceleration to runaway failure despite the continued reduction  
233 in mean effective stress. Instead, the strain rate continued to increase exponentially with  
234 pore-water pressure and therefore exhibited an asymptotic trend in  $1/v$ - $p'$  space (Figure  
235 8). Despite some variability in behaviour between each experiment, the strong  
236 asymptotic trend observed in  $1/v$ - $p'$  space (Figures 8 A, B and C) demonstrated that  
237 each sample underwent ductile deformation after a critical mean effective stress was  
238 reached, regardless of the rate of pore-water pressure increase. Similar styles of  
239 behaviour have been observed in shallow terrestrial landslides (e.g. Allison and  
240 Brunsden, 1990) and have been shown to be controlled by localised pore fluid changes  
241 in laboratory experiments (e.g. Ng and Petley, 2009).

242

243 The PGP experiments showed broadly similar progressive strain development and  
244 shear-zone dilation in response to decreasing mean effective stress as observed in the  
245 PWP testing (Figure 6 A and B). However, whilst movement initiated at low horizontal  
246 strain rates ( $< 0.05 \mu\text{S s}^{-1}$ ) and high mean effective stress in both experiments (Figure  
247 9 A), only PGP 12 developed the distinct sliding phase observed during the PWP  
248 experiments. This sliding phase occurred at a similar critical effective normal stress  
249 (c.10 kPa) to that in each of the PWP experiments indicating that pore fluid phase had

250 negligible influence on the effective stress conditions required to mobilise the shear  
251 zone.

252

253 By comparison, PGP 30 experienced an exponential increase in strain rate to  
254 approximately  $0.5 \mu\text{S s}^{-1}$  at a higher effective stress (15 kPa, Figure 9A). Following  
255 this, the strain rate remained constant whilst the mean effective stress continued to  
256 reduce. The peak strain rate was then reached at a lower effective stress (c.8 kPa) and  
257 resulted from a further exponential increase in strain rate. Whilst more variability was  
258 observed in PGP 30, similar maximum strain rates were observed in both PGP 30 and  
259 PWP 30 experiments, suggesting subtle differences in rheology may have impacted the  
260 testing. Similar episodic patterns of strain-rate development were observed in both gas  
261 experiments once the shear surface had mobilised, and an asymptotic trend in  $1/v - p'$   
262 space was observed in both PGP experiments (Figure 9 B).

263

264 The results indicated that shear-zone deformation occurred through ductile deformation  
265 at a critical mean effective stress regardless of the applied rate of pore pressure increase  
266 or the phase of the pore fluid. In addition, the distinct reduction in movement rates after  
267 initial shear surface mobilisation and the exponential increase in displacement rate  
268 observed across all experiments indicated that catastrophic failure did not develop.  
269 During most of the experiments, no relationship between reduction in strain rate and  
270 decreasing sample pore pressure was observed (e.g. Figure 10 A), however, a reduction  
271 in sample pore gas pressure was measured during experiment PGP30, which coincided  
272 with the development of peak horizontal and vertical strain rates (Figure 10 B). It was  
273 inferred from this that the rapid dilation of the shear zone acted to increase its  
274 permeability, which resulted in localised dissipation of pore water pressure, temporarily

275 altering the stress state within the shear zone before pore pressure increased further.  
276 Slow episodic shear regulated by dilation provides a potential mechanism to explain  
277 the cyclic phases of increased landslide displacement rate observed in response to  
278 elevated pore-fluid pressures without the development of rapid shear failure.

279

## 280 5. DISCUSSION

281 Although a range of mechanisms have been suggested to explain various subaqueous  
282 mass-movements, including shear surface nucleation (e.g. Viesca and Rice, 2012);  
283 shear zone liquefaction and ductile extrusion (e.g. Bull et al., 2009; Sassa et al., 2012);  
284 and local lateral fluid flow (Dugan and Flemings, 2000; Fleming et al., 2002), very few  
285 mechanisms have been proposed to explain shallow subaqueous slope failures that  
286 arrest without long runout. Based on the experimental results illustrated herein, a  
287 conceptual model can be hypothesised to explain how shallow subaqueous slopes can  
288 progressively deform through episodic movement when pore-fluid pressures are  
289 elevated within a shear zone (Figure 11).

290

291 In the model, an increase in pore pressure generates localised dilation and strain within  
292 the landslide shear zone. As a consequence, very slow pre-failure deformation initiates  
293 whilst effective normal stress remains comparatively high in the slope (stage 2). Slow  
294 dilation increases the permeability of the shear zone allowing pore pressures to increase  
295 more rapidly which increases the landslide strain rate and drives further shear zone  
296 dilation (stage 3). This progressive increase in strain rate, and the associated inter-  
297 particle deformation, drives further increases in local pore-fluid pressure, within the  
298 narrow shear zone. This in turn induces a further increase in strain rate (stage 4). The  
299 feedback mechanism continues as pore water pressure increases to reach a critical mean

300 effective stress when the shear surface rapidly slides and dilates (stage 5). This process  
301 rapidly changes the shear zone properties as the rate of permeability increase through  
302 dilation exceeds the rate of pore pressure increase within the narrow shear band. This  
303 leads to pore expansion and dissipation of the excess fluid pressure (stage 5), locally  
304 increasing the effective stress and reducing the shear strain rate. These processes  
305 prevent catastrophic acceleration (stage 6).

306

307 The process of coupling local pore-pressure increase and the development of high strain  
308 rate, checked by dilation and the broadening of the shear zone, can continue while  
309 external processes drive increasing pore-fluid pressure (stage 7a). Consequently, the  
310 landslide will continue to display ductile deformation behaviour, characterised by  
311 exponential increases in strain rate punctuated by episodic phases of decreased  
312 movement rates, never leading to catastrophic failure. Should the externally derived  
313 pore-fluid pressures reduce or pore fluid dissipate fully from the shear zone (Stage 7b)  
314 the slope movement will arrest and expulsion of pore fluid from the pores will result in  
315 re-compaction of the shear zone (stage 8).

316

317 Our study has focused specifically on the movement mechanisms in subaqueous slopes  
318 using a linear increase in pore-fluid pressure within a pre-defined shear zone. Other  
319 potential mechanisms hypothesised to influence landslide motion such as state-and-  
320 rate-variable friction affects (e.g. Helmsetter et al., 2004), complex perturbations in  
321 effective stress (e.g. Hangwerger et al., 2016) and shear-surface geometry (e.g. Aryal  
322 et al., 2015) have not been analysed. The coupling of shear deformation and shear-zone  
323 dilation observed in our experiments has been shown to promote steady landslide  
324 motion, particularly within clay-rich landslide shear zones (Iverson, 2005). Similar

325 movement patterns are also observed in terrestrial landslides on shallow slopes, (e.g.  
326 Ng and Petley, 2009) and slowly deforming mudslides (e.g. Allison and Brunsten,  
327 1990). In addition, similar behaviour has been used to describe ice flow dynamics (e.g.  
328 Damsgaard et al., 2016) and the pore fluid driven cyclic fault-valve model proposed for  
329 seismic slip (Sibson, 1992), indicating that this behaviour can be expected across a  
330 broad range of geological processes.

331

332 Despite the different physical properties of nitrogen gas and liquid water, similar  
333 patterns of behaviour were observed across all the experiments, demonstrating that  
334 either fluid can generate similar movement characteristics with increasing pore  
335 pressure. The results explain how unconstrained subaqueous landslides can episodically  
336 move downslope when pore-fluid pressures at the landslide shear zone are elevated by  
337 external factors such as the injection of water or gas from below.

338

339 The experiments presented herein provide credible support for the hypothesis that over-  
340 pressuring by free gas can result in episodic/slow movements required to produce  
341 submarine spreading failures observed in different parts of the world (e.g. Micallef et  
342 al., 2007; Mountjoy et al. 2009; Micallef et al. 2016). Cyclic changes in strain rate  
343 driven by a negative feedback mechanism associated with shear-zone dilation provides  
344 a credible mechanism through which a subaqueous landslide can accumulate large  
345 strain without catastrophic failure. This cyclic process, therefore, determines the long-  
346 term behaviour of subaqueous mass movement in shallow sedimentary sequences, and  
347 as similar materials commonly form submarine slopes, it is likely to be a widespread  
348 seafloor process.

349

350 CONCLUSION

351 Geomorphological evidence suggests that some shallow subaqueous slopes can  
352 accumulate substantial amounts of downslope deformation without transitioning to  
353 catastrophic failure, even though their downslope terminations are unconstrained  
354 (Mountjoy et al., 2009; Micallef et al., 2016). Novel laboratory experiments have been  
355 used to explore this behaviour by examining the response of such slope materials to  
356 states of low effective normal stress associated with high pore-fluid pressure induced  
357 by either gas or water injected from below. The strain behaviour observed in the  
358 experiments were found to be similar to the movement patterns measured in terrestrial  
359 landslides which deform along ductile shear zones. In such circumstances movement  
360 develops when pore water pressures are sufficiently elevated in the slope and movement  
361 rates increase exponentially with increasing pore pressure. Given that this response is  
362 dominated by the effective stress conditions operating within the shear zone, such  
363 behaviour can be expected in submarine slopes regardless of the pore fluid phase (gas  
364 or liquid). The behaviour we observed provides a mechanism through which  
365 subaqueous landslides may accumulate strain without undergoing catastrophic failure.  
366 This has important implications when assessing marine geohazards since these types of  
367 slope failure will not be tsunamigenic.

368

369 ACKNOWLEDGMENT

370 We thank Peter Barker, Stuart Read and Zane Bruce for laboratory support. Samples  
371 were provided by the National Institute of Water and Atmospheric Science (NIWA).  
372 Financial support has been provided through the GNS Science Strategic Development  
373 Fund, Marsden Fund Contract NIW1603, GNS and NIWA Strategic Science

374 Investment Funding and by the NERC/ESRC Increasing Resilience to Natural Hazards

375 programme, grant NE/J01995X/1, and NERC/Newton Fund grant NE/N000315.

376

377 REFERENCES

- 378 Allison, R., and Brunsten, D., 1990, Some mudslide movement patterns: Earth  
379 Surface Processes and Landforms, v. 15, p. 297-311.  
380
- 381 Anderson, S.A., and Sitar, N., 1995, Analysis of rainfall induced debris flows, Journal  
382 of Geotechnical Engineering. American Society of Civil Engineers. V, 121, P.544-  
383 552.  
384
- 385 Aryal, A., Brooks, B., Reid, M.E., 2015, Landslide subsurface slip geometry inferred  
386 from 3-D surface displacement fields, Geophysical Research Letters, v. 42(5), p.  
387 1411-1417.  
388
- 389 Baeten, N. J., Laberg, J. S., Vanneste, M., Forsberg, C. F., Kvalstad, T. J., Forwick,  
390 M., Vorren, T. O., Haflidason, H., 2014, Origin of shallow submarine mass  
391 movements and their glide planes—Sedimentological and geotechnical analyses from  
392 the continental slope off northern Norway, Journal of Geophysical Research: Earth  
393 Surface v.119(11), p. 2335-2360.  
394
- 395 Bull, S., Cartwright, J., Huuse, M., 2009, A subsurface evacuation model for  
396 submarine slope failure, Basin Research, v. 21, p. 433-443.  
397
- 398 Brain, M.J., Rosser, N.J., Sutton, J., Snelling, K., Tunstall, N., and Petley, D.N., 2015,  
399 The effects of normal and shear stress wave phasing on coseismic landslide  
400 displacement: Journal of Geophysical Research: Earth Surface, v 120, p. 1009-1022.  
401
- 402 Brand, E.W., 1981, Some thoughts on rain-induced slope failure, Proceedings 10th  
403 International Conference on Soil Mechanics and Foundation Engineering, San  
404 Francisco. Balkema. p. 2541-2578.  
405
- 406 Carey, J.M., McSaveney, M.J., Lyndsell, B.M., Petley, D., 2016, Laboratory  
407 simulation of a slow landslide mechanism. In: S, Aversa et al. (eds.) Landslides and  
408 engineered slopes: experience, theory and practice: proceedings of the 12th  
409 International Symposium on Landslides. Boca Raton, Fla.: CRC Press, p. 557–564.  
410
- 411 Carey, J.M., McSaveney, M.J., Petley, D.N., 2017, Dynamic liquefaction of shear  
412 zones in intact loess during simulated earthquake loading. Landslides, v. 14(3), p.  
413 789-804.  
414
- 415 Carey, J.M., Moore, R., Petley, D.N., 2015, Patterns of movement in the Ventnor  
416 landslide complex, Isle of Wight, southern England, Landslides, v. 12(6), p. 1107-  
417 1118.  
418
- 419 Carey, J.M., Petley, D.N., 2014, Progressive shear-surface development in cohesive  
420 materials; implications for landslide behaviour, Engineering Geology, v. 177: p. 54-  
421 65.  
422
- 423 Crutchley, G.J., Mountjoy, J.J., Pecher, I. A., Gorman, A., Henrys, S.A., 2016,  
424 Submarine slope instabilities coincident with shallow gas hydrate systems: Insights  
425 from New Zealand examples. In: G. Lamarche et al. (eds.), Submarine Mass

426 Movements and their Consequences, *Advances in Natural and Technological Hazards*  
427 *Research* 41, p. 401-409.

428

429 Dai, F.C., Lee, C.F., Wang, S.J. and Feng, Y.Y., 1999. Stress-strain behaviour of a  
430 loosely compacted volcanic-derived soil and its significance to rainfall-induced fill  
431 slope failures. *Engineering Geology*, v. 53, p. 359-370.

432

433 Damsgaard, A., Egholm, D. L., Beem, L. H., Tulaczyk, S., Larsen, N. K., Piotrowski,  
434 J. A., Siegfried, M. R., 2016, Ice flow dynamics forced by water pressure variations in  
435 subglacial granular beds, *Geophysical Research Letters*, v 43, p. 12-165 – 12-173.

436

437 Dugan, B., and Flemings. P.B., 2000, Overpressure and fluid flow in the New Jersey  
438 continental slope: Implications for slope failure and cold seeps: *Science*, v. 289, p.  
439 288–291.

440

441 Elger, J., Berndt, C., Rupke, L., Krastel, S., Gross, F., Geissler, W.H., 2018,  
442 submarine slope failures due to pipe structure formation, *Nature Communications*, v.  
443 9, 715.

444

445 Flemmings, P.B., Stump, B.B., Finkbeiner, T., Zoback, M., 2002, Flow focusing in  
446 overpressured sandstones: Theory, observations and applications, *American Journal*  
447 *of Science*, v. 302(10), p. 827-855.

448

449 Fukuzono, T. 1990. Recent studies on the time prediction of slope failures. *Landslide*  
450 *News*, v. 4, p. 9-12.

451

452 Hangwanger, A.I., Rempel, A.w., Skarbek, R.M., Roering, J.J., Hilley, G.E., 2016,  
453 Rate-weakening friction characterises both slow sliding and catastrophic failure of  
454 landslides, *Proceeding of the National Academy of Sciences*, v. 133(37), p. 10281-  
455 10286.

456

457 Helmsetter, A., Sornette, D., Grasso, J-R., Andersen, J., Gluzmann, S., Pisarenko, V.,  
458 2004, Slider block friction model for landslides: Application to Vaiont and La  
459 Clapiere landslides, *Journal of Geophysical Research*, 109(B2),  
460 10.1029/2002JB002160.

461

462 Hungr, O., Leroueil, S., Picarelli, L., 2014, The Varnes classification of landslide  
463 types, an update, *Landslides*, v. 11, p. 167-194.

464

465 Iverson, R.M., 2005. Regulation of landslide motion by dilatancy and pore pressure  
466 feedback: *Journal of Geophysical Research*, v. 110, F02015.

467

468 Kilburn, C.J., and Petley, D.N., 2003. Forecasting giant, catastrophic slope collapse:  
469 lessons from Vajont Northern Italy: *Geomorphology*, v. 54, p. 21-32.

470

471 Kossel, E., Bigalke, N., Piñero, E., and Haeckel, M., 2013, *The SUGAR Toolbox - A*  
472 *library of numerical algorithms and data for modelling of gas hydrate systems and*  
473 *marine environments*, Bremerhaven, PANGAEA .

474

475 Krastel, S., Li, W., Urlaub, M., Georgiopoulou, A., Wynn, R. B., Schwenk, T.,  
476 Stevenson, C., Feldens, P., 2018, Mass wasting along the NW African continental  
477 margin. In: D.G. Lintern et al., (eds). Subaqueous Mass Movements and Their  
478 Consequences: Assessing Geohazards, Environmental Implications and Economic  
479 Significance of Subaqueous Landslides. Geological Society, London, Special  
480 Publications, 477, <https://doi.org/10.1144/SP477.36>  
481

482 Kuhlman, J., Orpin, A., Mountjoy, J.J., Crutchley, G., Henrys, S., Lunenburg, R.,  
483 Huhn, K., 2018, GSL: Seismic and lithofacies characterization of a gravity core  
484 transect down the submarine Tuaheni Landslide Complex, northeastern New Zealand.  
485 In: (eds.) D. G. Lintern, D. C. Mosher, L. G. Moscardelli, P. T. Bobrowsky, C.  
486 Campbell, J. D. Chaytor, J. J. Clague, A. Georgiopoulou, P. Lajeunesse, A.  
487 Normandeau, D. J. W. Piper, M. Scherwath, C. Stacey and D. Turmel, SP477  
488 Subaqueous Mass Movements and Their Consequences: Assessing Geohazards,  
489 Environmental Implications and Economic Significance of Subaqueous Landslides  
490 Geological Society of London Special Publications, 477.  
491 <https://doi.org/10.1144/SP477.37>  
492 Locat, J., Lee, H.J., 2002, Submarine landslides: Advances and challenges, Canadian  
493 Geotechnical Journal, V. 36, p. 193-212.  
494

495 Mansour, F.M., Morgenstern, N.R., and Martin, D.C., 2011, Expected damage from  
496 displacement of slow-moving slides: Landslides, v. 8, p. 117-131.  
497

498 Massey, C.I., Petley, D.N., McSaveney, M.J., 2013, Patterns of movement in  
499 reactivated landslides. Engineering Geology, v. 159, p. 1-19.  
500

501 Micallef, A., Masson, D.G., Berndt, C., Stow, D.A.V., 2007, Morphology and  
502 mechanics of submarine spreading: A case study from the Storegga Slide, Journal of  
503 Geophysical Research, 112:F03023.  
504

505 Micallef, A., Mountjoy, J.J., Krastel, S., Crutchley, G., Koch, S., 2016, Shallow gas  
506 And the development of a weak layer in submarine spreading, Hikurangi Margin  
507 (New Zealand). In: G. Lamarche et al. (eds.), Submarine Mass Movements and their  
508 Consequences, Advances in Natural and Technological Hazards Research 41 ,p. 419-  
509 426.  
510

511 Mountjoy, J.J., Pecher, I., Henrys, S., Crutchley, G., Barnes, P.M., and Plaza-  
512 Faverola, A., 2014, Shallow methane hydrate system controls ongoing, downslope  
513 sediment transport in a low-velocity active submarine landslide complex, Hikurangi  
514 margin, New Zealand: Geochemistry, Geophysics, Geosystems, v. 15, p. 4137-4156.  
515

516 Mountjoy, J.J., McKean, J., Barnes, P.M., and Pettinga J.R., 2009, Terrestrial-style  
517 slow-moving earthflow kinematics in a submarine landslide complex: Marine  
518 Geology, v. 267, p.114-127.  
519

520 Mountjoy, J. and Micallef, A. 2018. Submarine Landslides. In Micallef, Krastel and  
521 Savini (eds) Submarine Geomorphology Springer International Publishing. 556p  
522

523 Ng, K-Y., and Petley, D.N., 2009, A process approach towards landslide risk  
524 management in Hong Kong, *Quarterly Journal of Engineering Geology and*  
525 *Hydrogeology*, v. 42, p. 1-13.  
526

527 Orense, R., Farooq, K. and Towhata, I., 2004, Deformation behaviour of sandy slopes  
528 during rainwater infiltration. *Soils and Foundations*, v. 44(2), p. 15-30.  
529

530 Petley, D.N., Bulmer, M.H. and Murphy, W. 2002. Patterns of movement in rotational  
531 and translational landslides. *Geology*, V.30, p.719-722  
532

533 Petley, D.N., Higuchi, T., Petley, D.J., Bulmer, M.H. and Carey. J., 2005, The  
534 development of progressive landslide failure in cohesive materials: *Geology*, v. 33, p.  
535 201-204.  
536

537 Petley, D.N., Mantovani, F., Bulmer, M.H. & Zannoni, A. 2005c. The use of surface  
538 monitoring data for the interpretation of landslide movement patterns.  
539 *Geomorphology*, 66:133-147.  
540

541 Petley, D.N. and Petley, D.J., 2006. On the initiation of large rockslides: perspectives  
542 from a new analysis of the Vaiont movement record. In: S.G Evans, G S Mugnozza, A  
543 Strom and R.L Hermanns (eds.), *Landslides from Massive Rock Slope Failure*. NATO  
544 *Science Series IV: Earth and Environmental Sciences*. Springer: Dordrecht. p. 77-84.  
545

546 Riboulot, V., Cattaneo, A., Sultan, N., Garziglia, S., Ker, S., Imbert, P. and Voisset,  
547 M. 2013. Sea-level change and free gas occurrence influencing a submarine landslide  
548 and pockmark formation and distribution in deepwater Nigeria. *Earth and Planetary*  
549 *Science Letters* v. 375, p. 78-91.  
550

551 Saito, M. (1965). Forecasting the time and occurrence of a slope failure. *Proceedings*  
552 *of the 6th International Conference on Soil Mechanics and Foundation Engineering*. 2.  
553 p. 537-541.  
554

555 Sassa, K., He, B., Miyagi, T., Strasser, M., Konagai, K., Ostric, M., Setiawan, H.,  
556 Takara, K., Nagai, O., Yamashiki, and Y., Tutumi, S., 2012, A hypothesis of the  
557 Senoumi submarine megaslide in Suruga Bay in Japan – based on the undrained  
558 dynamic loading ring shear tests and computer simulation: *Landslides*, v. 9, p. 439-  
559 455.  
560

561 Schulz, W., J. McKenna, J. Kibler and G. Biavati (2009). "Relations between  
562 hydrology and velocity of a continuously moving landslide—evidence of pore-  
563 pressure feedback regulating landslide motion?" *Landslides*, v. 6(3), p. 181-190.  
564

565 Sibson, R. H., 1992, Implications of fault-valve behaviour for rupture nucleation and  
566 recurrence, *Tectonophysics*, v 211, p. 283-293.  
567

568 Stigall, J., and Dugan, B., 2010, Overpressure and earthquake initiated slope failure in  
569 the Ursa region, northern Gulf of Mexico: *Journal of Geophysical Research*, v.115,  
570 B04101.  
571

572 Talling, P.J., Wynn, R.B., Masson, D.G., Frenz, M., Cronin, B.T., Schiebel, R.,  
573 Akhmetzhanov, A.M., Dallmeier-Tiessen, S., Benetti, S., Weaver, P.P.E.,  
574 Georgiopoulou, A., Zühlsdorff C., Amy L.A., 2007, Onset of submarine debris flow  
575 deposition far from original giant landslide, *Nature*, v. 450(7169). p. 541-544.  
576  
577 Urlaub, M., Talling, P.J., Zervos, A., Masson, D., 2015, What causes large submarine  
578 landslides on low gradient ( $<2^\circ$ ) continental slopes with slow ( $\sim 0.15$  m/Kyr) sediment  
579 accumulation? *Journal of Geophysical Research: solid Earth*, v. 120, p. 6722-6739.  
580  
581 Viesca, R. C., and Rice, J.R., 2012, Nucleation of slip-weakening rupture instability in  
582 landslides by localized increase of pore pressure: *Journal of Geophysical Research*, v.  
583 117, B03104.  
584  
585 Voight, B. 1988, A relation to describe rate dependent material failure. *Science*, v.  
586 243, p. 200-203.  
587  
588 Zhu, J.-H., and Anderson, S.A., 1998, Determination of shear strength of Hawaiian  
589 residual soil subjected to rainfall-induced landslides. *Geotechnique*, v. 48(1), p. 73-82.  
590  
591 Wallace, L. M., and Beavan, J., 2010, Diverse slow slip behavior at the Hikurangi  
592 subduction margin, New Zealand, *Journal of Geophysical Research*, v 115/  
593 B007717  
594  
595 Wallace, L. M., Beavan, J., Bannister, S., and Williams, C., 2012, Simultaneous long-  
596 term and short-term slow slip events at the Hikurangi subduction margin, New  
597 Zealand: Implications for processes that control slow slip event occurrence, duration,  
598 and migration, *Journal of Geophysical Research*, v. 112/B009489.  
599  
600

601 FIGURE CAPTIONS

602

603 Figure 1. Study area, east of New Zealand's North Island. (A) 'Hillshade' plot of  
604 Tuaheni Basin bathymetry, showing the Tuaheni Landslide Complex (TLC) and a  
605 shallow subaqueous slope failure (Inset B) which both show evidence of limited  
606 displacement despite having unconstrained toes. Yellow dots = location of Sample  
607 Tan1404-10; yellow contours = metres below sea level adapted from Mountjoy et al.  
608 (2014). (B) Bathymetric map of shallow subaqueous slope failure draped over slope  
609 gradient map and showing key morphological features including arrested debris (ad)  
610 and evacuated debris (ed) after Micallef et al. (C) Seismic section across the subaqueous  
611 slope failure headscarp after Micallef et al. (2016).

612

613 Figure 2. The particle size distribution of samples EN1285 (Tan1404-01) and EN1287  
614 (Tan1404-10) collected from the Tuaheni South section of the TLC.

615

616 Figure 3. Laboratory testing apparatus and procedures. (A) Schematic diagram of the  
617 Dynamic Back Pressure Shearbox apparatus. (B) Experimental procedure used for pore  
618 water pressure testing. (C) Experimental procedure used for pore gas pressure testing.

619

620 Figure 4. Testing parameters used during the specialist pore pressure reinflation testing.  
621 (A) Applied pore pressure (Back pressure) against mean effective stress. (B) Stress  
622 paths followed during the pore water pressure (PWP) and pore gas pressure (PGP)  
623 experiments in relation to the conventional failure envelope.

624

625 Figure 5. Residual strength envelopes constructed from conventional drained shear tests  
626 carried out on samples EN1285 and EN1287.

627

628 Figure 6. Change in horizontal and vertical strain measured during PWP and PGP  
629 experiments. (A) Horizontal strain against normal effective stress. (B) Vertical strain  
630 against normal effective stress.

631

632 Figure 7. Horizontal strain rate against normal effective stress illustrating three distinct  
633 movement phases measured during PWP experiments.

634

635 Figure 8. Analysis of movement styles observed during each PWP experiment using  $1/V$   
636 Horizontal strain rate ( $V$ ) against normal effective stress ( $p'$ ). (A) Asymptotic trend  
637 calculated during PWP 30 experiment. (B) Asymptotic trend calculated during PWP 12  
638 experiment. (C) Asymptotic trend calculated during PWP 5 experiment.

639

640 Figure 9. Analysis of movement styles observed during each PGP experiment. (A)  
641 Horizontal strain rate against normal effective stress. (B)  $1/v$  Horizontal strain rate ( $v$ )  
642 against normal effective stress ( $p'$ ). (C) Asymptotic trend calculated during PGP 30  
643 experiment. (D) Asymptotic trend calculated during PGP 12 experiment.

644

645 Figure 10. Comparison of behaviour observed during the PWP and PGP experiments.  
646 (A) Change in pore water pressure and strain rate against time measured during  
647 experiment PWP 30. (B) Change in pore gas pressure and horizontal strain rate against  
648 time measured during experiment PGP 30.

649

650

651 Figure 11. Conceptual model of the development of slow movement in the shallow  
652 subaqueous slopes driven by elevated pore-fluid pressure.

653

654

655

656 TABLE CAPTIONS

657 Table 1. Physical properties of the Tuaheni sediments

658 Table 2. Summary of conventional drained shear tests

659 Table 3. Summary of pore pressure reinflation tests

660

TABLE 1. PHYSICAL PROPERTIES OF THE TUAHENI SEDIMENTS

Sample Reference	EN1285 (Tan1404-02)		EN1287 (Tan1404-10)	
	Range	Average	Range	Average
Sampling depth (*mbsb)	0.32-0.49		2.23-2.40	
Moisture content (%)	77.0-80.0	78.5	70.0-72.0	70.5
Void ratio	3.11-3.17	3.14	2.00-2.18	2.06
Dry density	0.65-0.66	0.65	0.85-0.90	0.88
Atterberg limits:				
Plastic limit (%)		75		76
Liquid limit (%)		35		32
Plasticity index (%)		40		44
Particle size distribution:				
Clay (%)		5.50		4.18
Silt (%)		88.18		88.48
Sand (%)		6.32		7.33

\*mbsb = meters below sea bed

TABLE 2. SUMMARY OF CONVENTIONAL DRAINED SHEAR TESTS

Sample number	Sample depth (mbsl)	Initial water content $w_n$ (%)	Void Ratio	Dry density $\rho_d$ (t/m <sup>3</sup> )	Normal stress (kPa)	Shear stress (kPa)	
						Peak	Residual
EN1285a	0.40–0.44	78	3.14	0.65	32	27	27
EN1285c	0.36–0.40	77	3.11	0.66	76	59	59
EN1285d	0.32–0.36	79	3.17	0.65	10	11	11
EN1287a	2.31–2.35	70	2.02	0.89	32	31	31
EN1287c	2.27–2.31	72	2.18	0.85	10	12	12
EN1287e	2.23–2.27	70	2.00	0.90	76	66	66

TABLE 3. SUMMARY OF PORE PRESURE REINLFATION TESTS

Test Reference	Sample Number	Normal effective stress (kPa)	Initial shear strength (kPa)	Post failure shear stress/ Percentage of initial shear strength (kPa) / (%)	Pore fluid increase rate (kPa/ hr)	Pore fluid
PWP 5	EN1287i*	32*	22	15 (68%)	5	water
PWP 12			20	15 (75%)	12	water
PWP 30	EN1287p	32	20	13 (65%)	30	water
PGP 12	EN1287q	32	20	13 (65%)	12	nitrogen
PGP30	EN1287k	32	19	13 (68%)	30	nitrogen

\*Test PWP 5 and PWP 12 conducted on the same sample

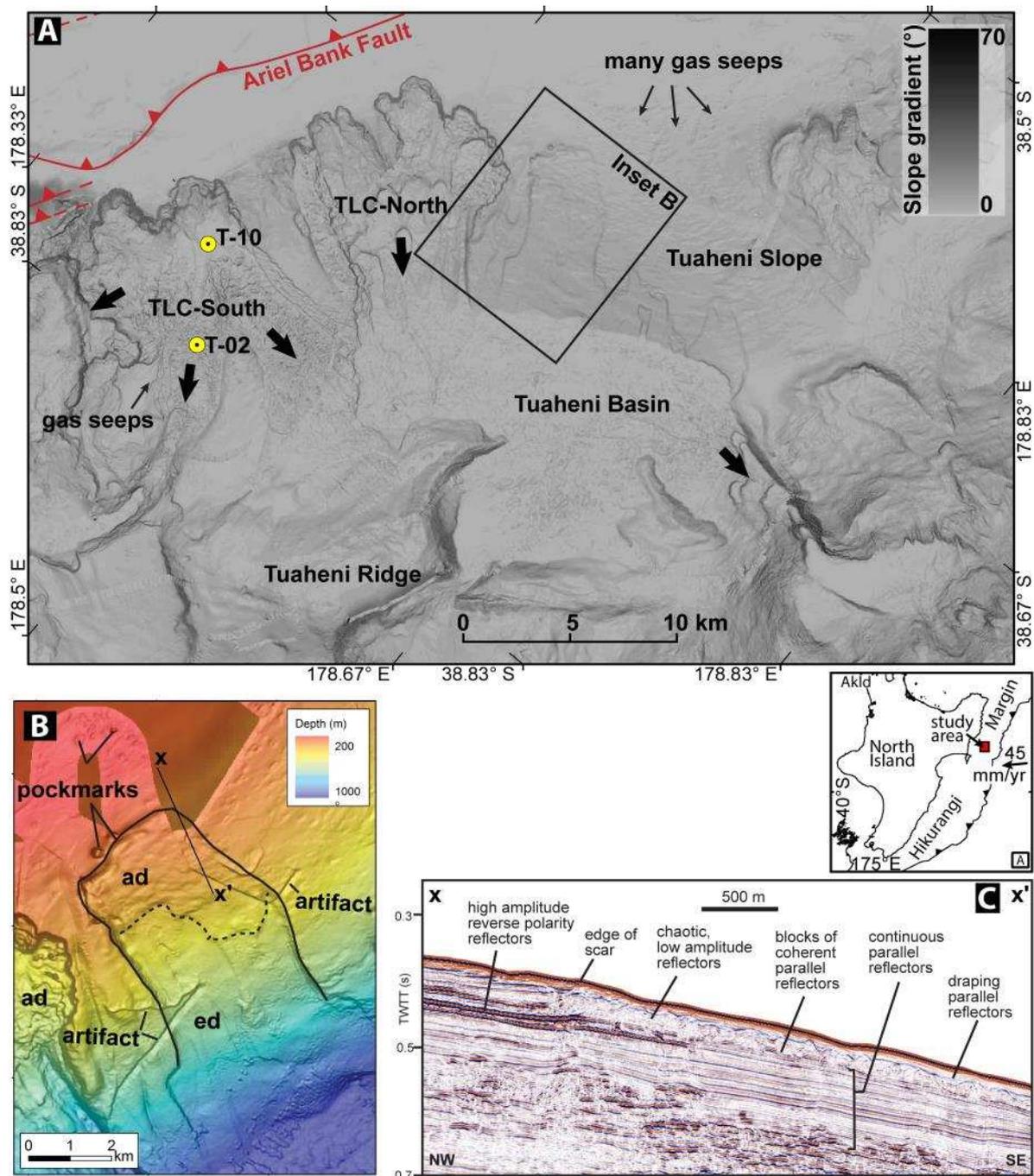


Figure 1. Study area, east of New Zealand's North Island. (A) 'Hillshade' plot of Tuaheni Basin bathymetry, showing the Tuaheni Landslide Complex (TLC) and a shallow subaqueous slope failure (Inset B) which both show evidence of limited displacement despite having unconstrained toes. Yellow dots = location of Sample Tan1404-10; yellow contours = metres below sea level adapted from Mountjoy et al. (2014). (B) Bathymetric map of shallow subaqueous slope failure draped over slope gradient map and showing key morphological features including arrested debris (ad) and evacuated debris (ed) after Micallef et al. (C) Seismic section across the subaqueous slope failure headscarp after Micallef et al. (2016).

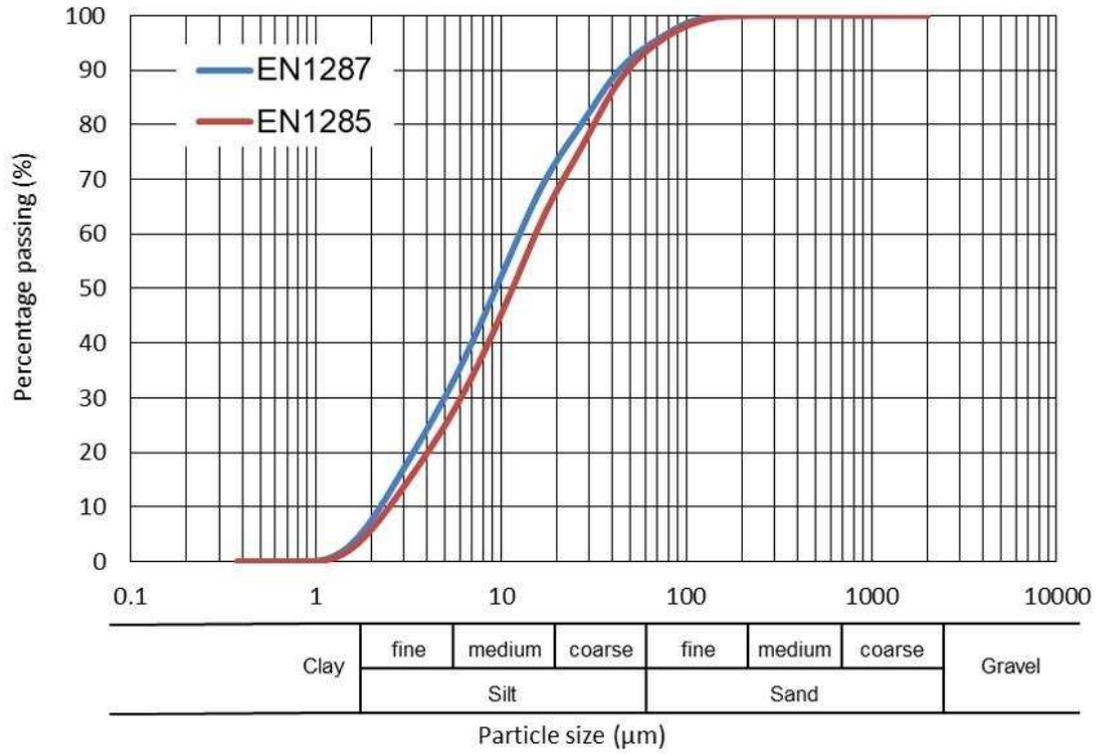


Figure 2. The particle size distribution of samples EN1285 (Tan1404-01) and EN1287 (Tan1404-10) collected from the Tuaheni South section of the TLC.



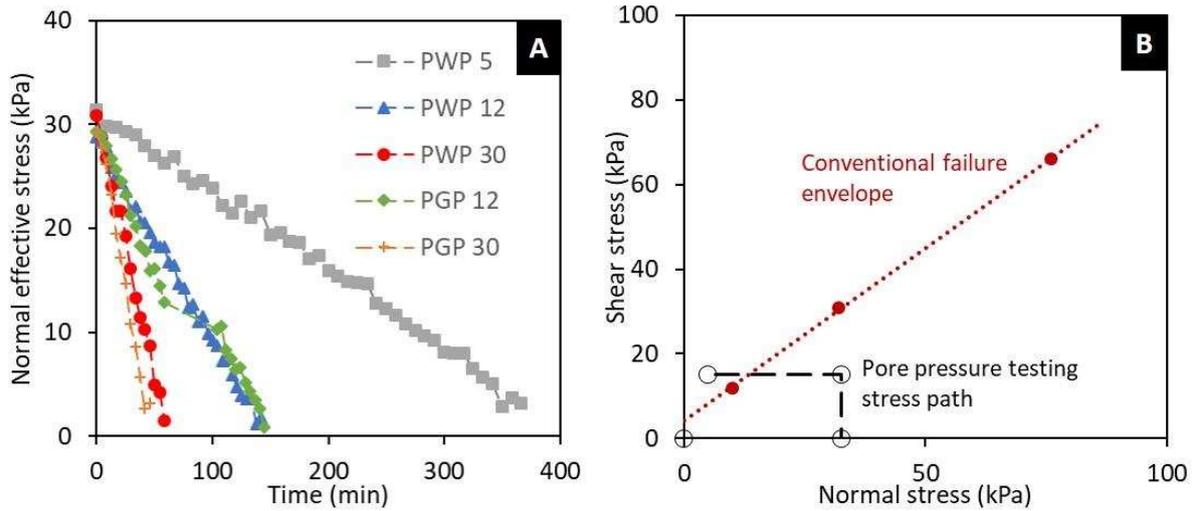


Figure 4. Testing parameters used during the specialist pore pressure reinflation testing. (A) Applied pore pressure (Back pressure) against mean effective stress. (B) Stress paths followed during the pore water pressure (PWP) and pore gas pressure (PGP) experiments in relation to the conventional failure envelope.

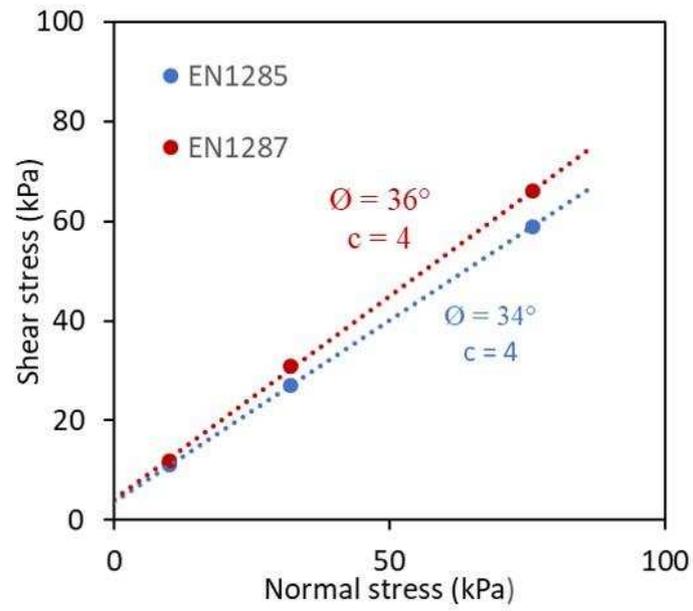


Figure 5. Residual strength envelopes constructed from conventional drained shear tests carried out on samples EN1285 and EN1287.

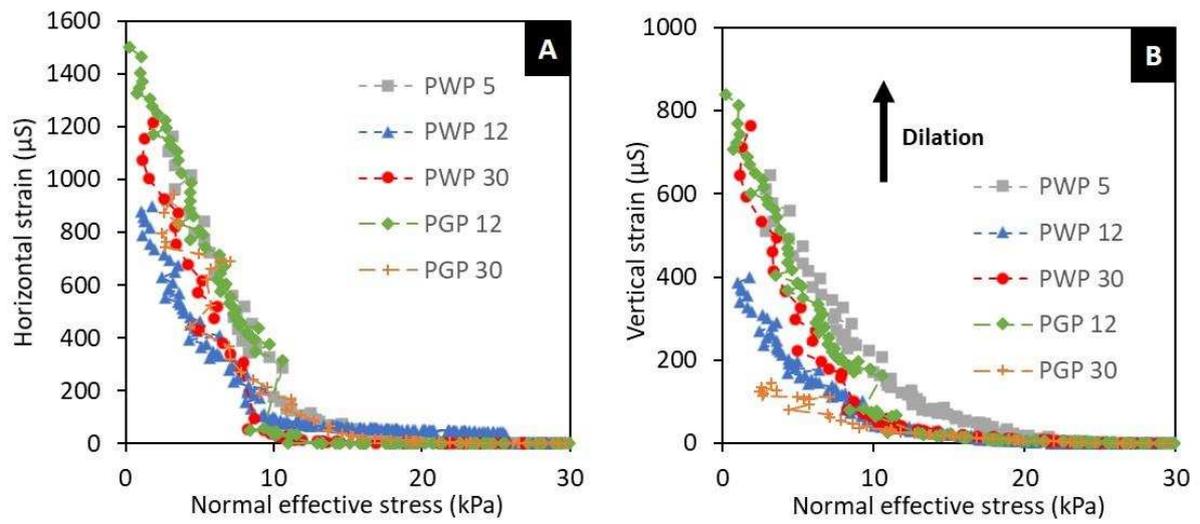


Figure 6. Change in horizontal and vertical strain measured during PWP and PGP experiments. (A) Horizontal strain against normal effective stress. (B) Vertical strain against normal effective stress.

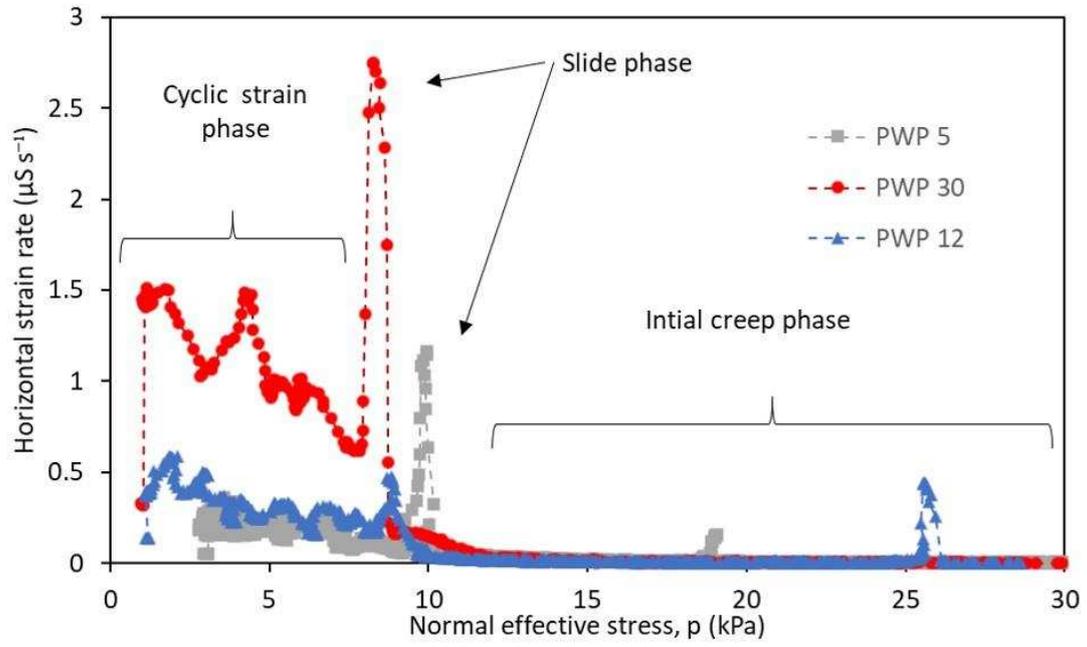


Figure 7. Horizontal strain rate against normal effective stress illustrating three distinct movement phases measured during PWP experiments.

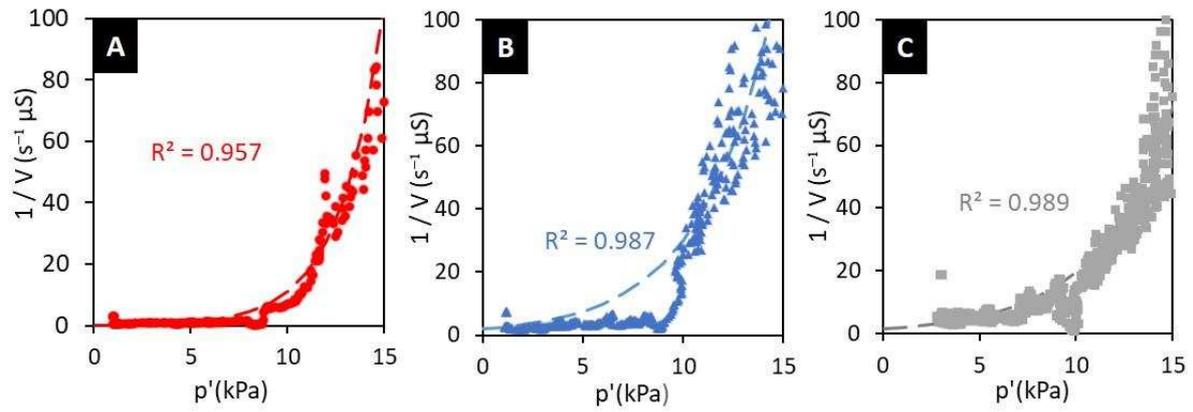


Figure 8. Analysis of movement styles observed during each PWP experiment using  $1/V$  Horizontal strain rate ( $V$ ) against normal effective stress ( $p'$ ). (A) Asymptotic trend calculated during PWP 30 experiment. (B) Asymptotic trend calculated during PWP 12 experiment. (C) Asymptotic trend calculated during PWP 5 experiment.

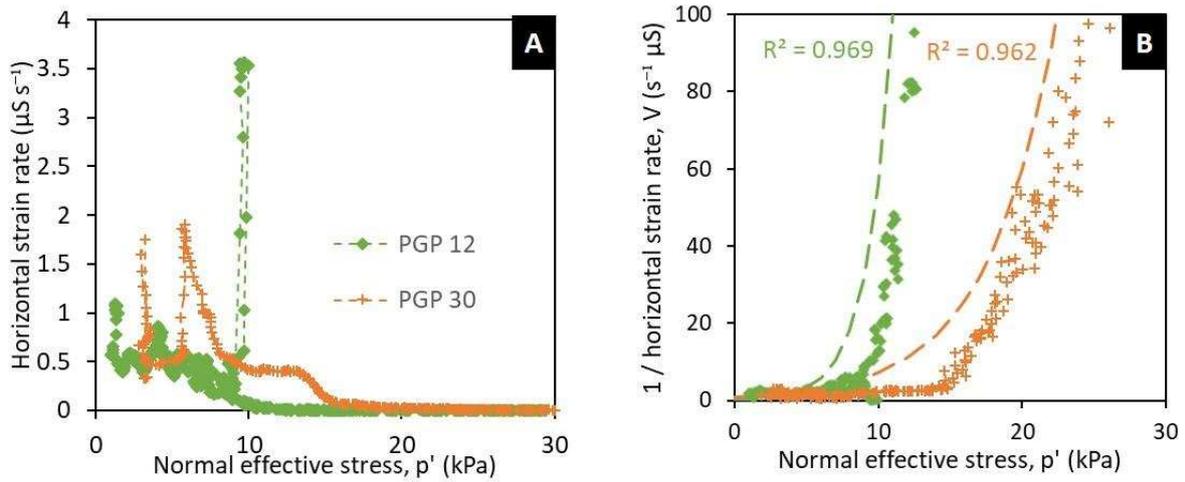


Figure 9. Analysis of movement styles observed during each PGP experiment. (A) Horizontal strain rate against normal effective stress. (B)  $1 /$  Horizontal strain rate ( $v$ ) against normal effective stress ( $p'$ ). (C) Asymptotic trend calculated during PGP 30 experiment. (D) Asymptotic trend calculated during PGP 12 experiment.

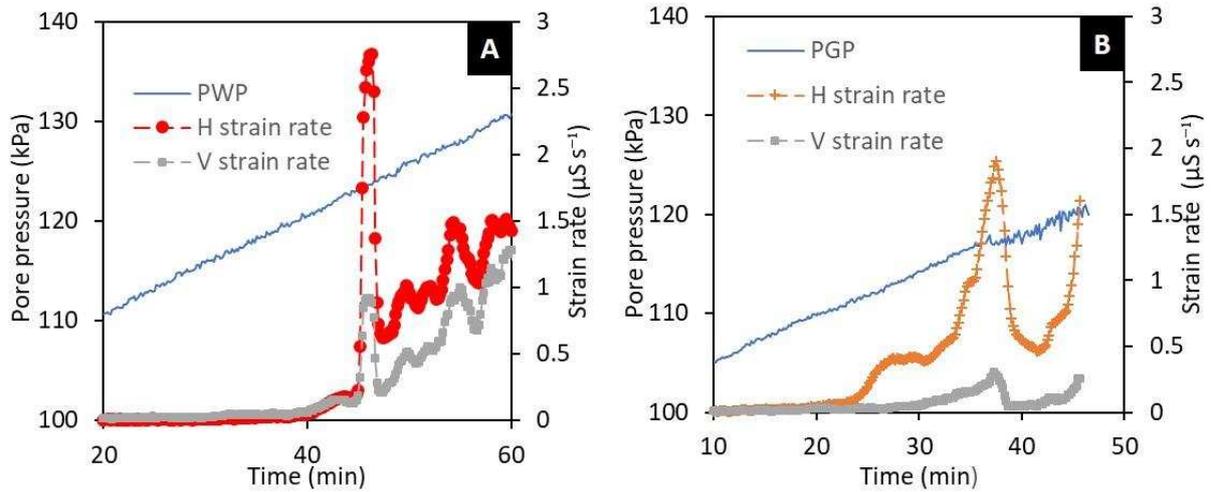


Figure 10. Comparison of behaviour observed during the PWP and PGP experiments. (A) Change in pore water pressure and strain rate against time measured during experiment PWP 30. (B) Change in pore gas pressure and horizontal strain rate against time measured during experiment PGP 30.

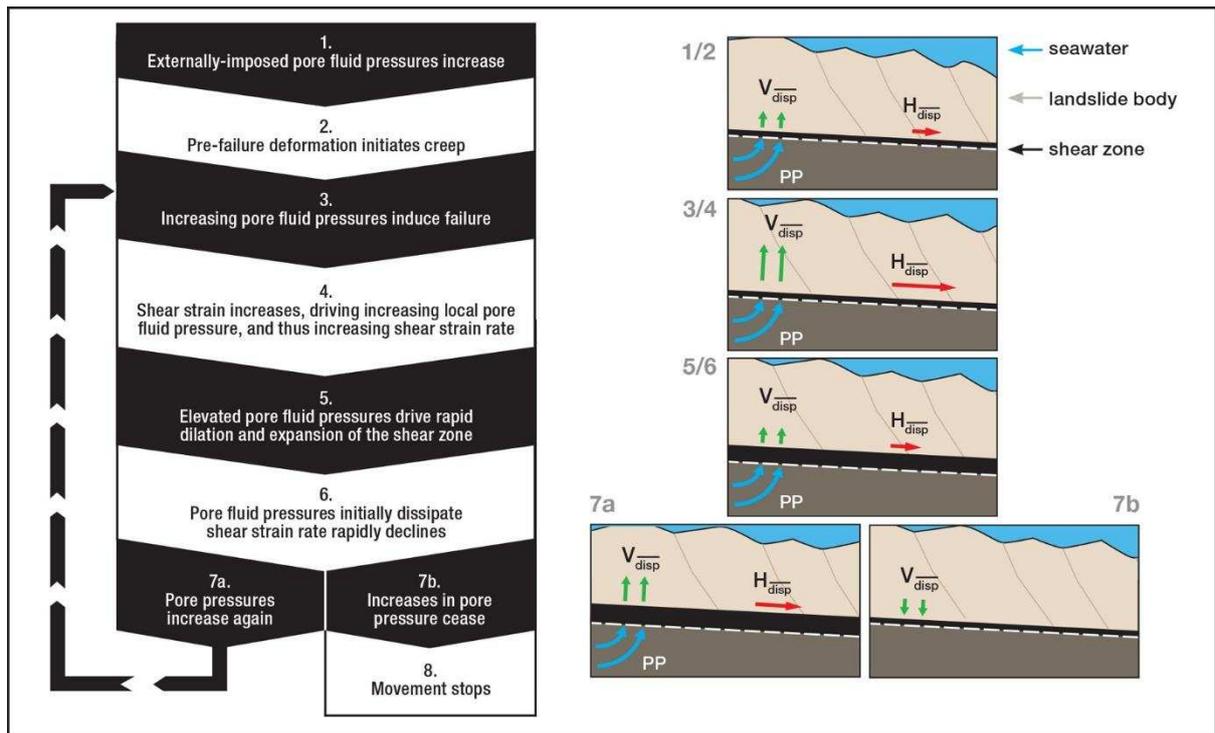


Figure 11. Conceptual model of the development of slow movement in the shallow subaqueous slopes driven by elevated pore-fluid pressure.