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1 **Damage-tolerant architected materials inspired by crystal microstructure**

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10

11 **Summary:** Architected materials comprised of periodic arrangements of nodes and struts are
12 lightweight materials that can exhibit combinations of properties which are inaccessible to
13 conventional solids. However, with regards to their mechanical performance, they have an
14 Achilles heel in that these materials can exhibit a catastrophic post-yielding collapse, causing
15 substantial drops in strength and energy absorption during plastic deformation. This post-
16 yielding collapse is the result of the activity of single shear bands, and originates from the single
17 orientation of macro-lattices. We observe that this behaviour is analogous to deformation by slip
18 in metallic single crystals. In this study we propose that, by mimicking the microstructure
19 observed in crystalline materials, we may be able to employ hardening mechanisms found in
20 crystalline materials to help us to develop robust and damage-tolerant architected materials. This
21 study demonstrates that crystal-inspired meso-structures can play an important role in the
22 mechanical properties of architected materials as do crystallographic microstructures in the
23 case of metallic alloys. Consequently, designing meso-structures that mimic crystallographic
24 microstructure in crystalline metals enables the fusion of metallurgy and architected materials
25 to transform the way of designing a new type of materials with desired properties.

26

27 A crystal is composed of atoms that are arranged in an ordered manner in space. The ordered
28 arrangement of atoms in a crystal is called a crystal lattice which is, in turn, defined by its unit
29 cell (Figure 1a) with the unit cell being the smallest group of lattice points that can describe the
30 overall symmetry of a crystal. A single crystal consists of unit cells of the same type and same
31 orientation, however, most crystalline materials are comprised of many domains with each
32 having a lattice orientation different to that of its neighbours. Such domains are separated by
33 boundaries and termed in metallurgy a crystal grain. Crystalline materials consisting of many
34 crystal grains are called polycrystals. In addition to the crystal lattice, there are other important
35 crystallographic microstructure features (namely dislocations, grains, phases and precipitates)
36 that can be manipulated to control the plastic deformation of crystals under external load^{1,2}. In
37 particular, the slip of dislocations is the most common mechanism responsible for plastic
38 deformation¹. For single crystals, a dominant single slip mode occurs during plastic deformation
39 (Figure 1b), causing strain to become localised and the stress required for further deformation to
40 drop³. For poly-crystals, the orientation change across the grain boundary can impede or even
41 stop dislocation movement from one grain to the next one (Figure 1c), resulting in the
42 dependence of plastic deformation on the grain size³. This capacity for constraining the
43 deformation and preventing the rapid propagation of slip minimises the stress drops⁴ and
44 strengthen poly-crystals¹. The size of grains can significantly alter the mechanical strength,
45 described by the well-known Hall-Petch relationship⁵⁻⁷ that the yield strength (σ_y) of polycrystals
46 is inversely proportional to the square root of grain size (d),

$$47 \quad \sigma_y = \sigma_0 + \frac{k}{\sqrt{d}} \quad (1)$$

48 where σ_0 is the friction stress which does not depend on the grain size, and k is a material
49 constant

50 Similarly, precipitates and phases also control the slip in crystals thanks to the difference
51 in lattice parameters, leading to other strengthening sources in crystalline alloys that are termed
52 precipitation and phase hardening in metallurgy¹. The capacity to control the strength and
53 toughness of alloys via engineering the size, distribution and orientation of grains, precipitates
54 and second phases forms the foundation of physical metallurgy.

55

56 Figure 1: Lattice structures and deformation behaviour. (a) Face-centred-cubic (FCC) crystal
57 lattice, (b) Single slips in a single crystal (re-used from ⁸ with permission from Elsevier), (c)
58 Slips at grain boundary in a polycrystalline steel (re-used from ⁹ with permission from Elsevier),
59 (d) Architected FCC lattice, (e) Single slip in a single oriented lattice, (f) Unstable behaviour of
60 architected metallic lattices (reused from ¹⁰ under the Creative Commons Attribution 4.0
61 International License).

62

63 Similar to the mimicry of structures inside biological systems (or organic materials) to
64 design bio-inspired materials ¹¹⁻¹⁴, we can easily mimic the crystal structure of metals and alloys
65 on a macroscopic scale by constructing a lattice unit cell comprised of an ordered arrangement of
66 nodes (analogous to atoms) connected by struts (equivalent to atomic bonds) (e.g., Figure 1d).
67 Consequently, architected lattice materials are straightforwardly generated from such a macro-
68 lattice unit in space using a suitable computer-aided design package and can be fabricated via
69 additive manufacturing (also commonly termed 3D printing). Very similar to crystals, the
70 properties of architected materials are strongly affected by their lattice architecture¹⁵. The
71 behaviour of architected lattices can be modified through careful design of the lattice unit cells,
72 giving access to unprecedented properties, e.g. lightweight^{16,17} and negative Poisson ratio^{18,19}. By
73 further manufacturing our architected lattice designs using different materials we can develop
74 architected materials that are both lightweight and strong²⁰. To date, however, work on these
75 promising architected materials has focused only on lattice materials with a single
76 orientation^{10,15,17,21} and, these materials have been observed to suffer precisely the same drastic
77 collapse in mechanical strength when loaded beyond the yield point (i.e., post-yielding collapse)
78 because of the occurrence of shear bands¹⁰ (e.g. Figure 1f) with this deformation again becoming
79 highly localised on specific planes with defined lattice directions^{10,22}, similar to the slip activity
80 in single crystals (Figure 1e, f versus Figure 1b). Considering the similarity between single
81 crystals and singly oriented lattice structures, we hypothesised that if it should be possible to
82 develop damage-tolerant architected materials by introducing hardening mechanisms
83 analogous to those found in crystalline materials. This proposal is aligned to a recent approach
84 that aims to bring the metallurgical microstructures closer to the component-size scale to better
85 engineer the behaviour of components and structures ^{12,23}. However, if the hypothesis holds true,
86 the approach proposed in this study is distinctive on its own as it will enable a fusion of material
87 specific hardening mechanisms and architected materials. In this study, we present the

88 methodology of designing meso-lattice features that mimic crystallographic microstructure to
89 bring metallurgical hardening principles (grain size effect, precipitation and multiphase
90 hardening) to the development of architected materials. We demonstrate that crystal-inspired
91 meso-structures can indeed strengthen the architected materials, leading to the generation of
92 highly damage-tolerant materials. The freedom in designing crystal-like meso-structures also
93 offers alternative ways of studying complex metallurgical phenomena (e.g., slip transfer) in
94 metallic alloys. We further show that the fabrication of lattices by polycrystalline materials leads
95 to the generation of a new family of materials that contain atomic lattices within meso-lattices
96 and polycrystal microstructure within polycrystal-like meso-structure; therefore, coined meta-
97 crystals. The properties of such new materials can be easily to be tailored via numerous
98 combinations of micro, meso and macro- lattices.

99 **Grain boundary hardening**

100 Tailoring the misorientation between two adjoining domains of macro-lattices creates a boundary
101 between them in the same way as the boundary between two adjoining grains of crystal
102 (Methods Section 1.1). Hereafter, each domain of the same lattice orientation is termed a meta-
103 grain. Because atoms can bond with new neighbouring atoms while physical nodes cannot, there
104 is the clear difference between polycrystal-inspired materials and polycrystals themselves. We
105 chose, as an initial step, two well-understood phenomena in metallurgy to verify whether
106 relationships observed in polycrystals are applicable to architected materials containing
107 polycrystal-like features: (i) symmetric slip in a twin-related bi-crystal (Figure 2a), and (ii) the
108 aforementioned Hall-Petch relationship^{5,6}.

109 To investigate the phenomenon (i), an architected material containing two meta-grains
110 separated by a twin boundary to mimic a twinned bi-crystal was designed and 3D printed
111 (Methods 1.1, and Extended data - Figures E1, and E2a, b). The base material used to fabricate
112 the twinned meta-grains was an elasto-plastic polymer (Methods, Materials and fabrication).
113 Later we demonstrate that a polycrystal-inspired architected material fabricated by different
114 base materials (including a stainless steel) exhibits the same behaviour in forming shear bands,
115 showing a wide applicability of this polycrystal-inspired approach. Under compression with the
116 loading parallel to the twin boundary, shear bands occurred and were parallel to the maximum
117 shear stress planes (Figure 2b). In addition, we observe that the shear bands formed are

118 symmetric about the twin boundary, confirming the shear band behaviour in meta-grain twins is
119 similar to the slip activity in crystal twins (Figures 2a, b). To assist us in understanding this
120 behaviour, finite element modelling (FEM) was used to understand the early stage of shear band
121 formation (see Methods). The simulation predicts well the location and direction of shear bands
122 in meta-grain twins as seen in the experiment (Figure 2c). By design, the (001) plane in each
123 meta-grain is parallel to the maximum shear force in that meta-grain domain and FEM shows
124 that as expected, deformation occurred via the buckling of $\langle 001 \rangle$ struts between two parallel
125 (001) planes and along a $\langle 101 \rangle$ direction, leading to the localisation through the face-centred
126 nodes parallel to the (001) plane of FCC (Extended data Video E1). The shear band system
127 family $\langle 101 \rangle \{001\}$ is different to the $\langle 101 \rangle \{111\}$ slip system one in the FCC crystal,
128 highlighting the difference between the FCC crystal lattice and architected lattice. The
129 symmetry of lattice orientation across the boundary causes the same localisation parallel to the
130 (001) plane in the other half, resulting in the symmetric shear bands in the twinned meta-grains.
131 The same twin boundary was used to further increase the number of meta-grains, in essence
132 using the twin boundary to reduce the size of meta-grains (Methods 1.1 and Extended data -
133 Figure E2) to study the phenomenon (ii) in the architected materials. Fabricated polycrystal-
134 like materials demonstrated a highly reproducible constitutive stress-strain behaviour (Extended
135 data - Figure E7). The change in orientation across boundaries effectively controls the
136 propagation of shear bands (Figure 2b, d and e). Most importantly, the yield strength of
137 architected materials substantially increases with the reduction in the meta-grain size,
138 confirming that an effect similar to the Hall-Petch relationship exists in the architected
139 materials. Fitting against experimental data using Eqn. (1) gives $\sigma_0 = 2.58 \text{ MPa}$ and $k =$
140 $14.37 \text{ MPa}\sqrt{\text{mm}}$. The value of k falls well in the range for conventional polycrystalline
141 metals²⁴. The flow stress at a given strain is also observed to increase with reducing the size of
142 meta-grains during plastic deformation (Extended data - Figure E8). In other words, the size of
143 meta-grains strongly affects both the yield strength and the work-hardening during deformation
144 of architected materials. This observation of the size effect (together with the deflecting effect,
145 Figure 2b, d and e) is significant as it confirms that the boundary strengthening is applicable to
146 architected materials. In particular, the grain size dependence in polycrystals is one of most
147 widely used mechanisms in metallurgy to achieve outstanding combination between strength and
148 ductility, e.g., development of a tough steel at low temperatures by the ultra-refinement of grains

149 ²⁵. However, it should be noted that that the grain size effect in metals is seen to be much more
150 significant with the presence of incoherently high angle boundaries²⁶. This means that the size
151 reduction of high angle boundaries should substantially strengthen architected materials. With
152 the freedom of varying the size while keeping the same type of boundary, the current approach
153 can provide clean data to obtain insights into the grain size effect of a specific type of boundary.
154 Obtaining such clean data in metallurgy are not possible because there are no effective ways to
155 reduce the grain size while maintaining the same single type of grain boundaries. This means
156 that, while the rich knowledge in polycrystals provides important underlying science to
157 engineering architected materials, this polycrystal-inspired approach offers alternative ways to
158 study complex phenomena in metallurgy.

159

160 Figure 2: Roles of lattice orientation in the deformation behaviour of crystals and architected
161 lattices. (a) Twin bi-crystal (re-used from ²⁷ with permission from Elsevier), Shear bands in
162 meta-grain twins observed in experiment (b) and predicted by FEM (c) (Note: Cut sections
163 shows the deformation of internal struts. The sections were formed thanks to two cutting planes
164 that were parallel to {001} planes of FCC lattice), (d) – (e): Shear bands were controlled by
165 orientation of meta-grains: (d) 8 meta-grains and (e) 16 meta-grains (note the nominal strain was
166 of 30%). (f) Yield strength versus the size of meta-grains, (g) Boundaries between meta-grains
167 effectively stops cracks in brittle lattices, leading to a drastic increase in toughness of
168 architected materials.

169

170 Metallurgical studies suggest that microstructure (e.g. dislocation substructures and high
171 angle boundaries) can even be able to stop the propagation of short cracks ²⁸⁻³⁰. A series of
172 samples from a different base materials were fabricated to study the effect of Incoherent High
173 Angle Boundaries (IHABs) on crack propagation in architected materials (Methods –
174 Fabrication and Extended data Table 1). The base material used for these architected materials
175 is a brittle polymer to easily generate cracks under loading. For a single oriented lattice, cracks
176 occurred and propagated rapidly throughout the macro-lattice, causing premature and fast
177 fracture (dashed line, Figure 2g). In the case of the architected material containing 8 IHAB
178 meta-grains (Figure 2g - solid line Extended data – Figure E3), the orientation rotation of internal
179 lattice led to the creation of incomplete unit cells on the free surface, causing a lower yield
180 strength compared to the single oriented lattice. Nevertheless, cracks did stop at the IHABs,
181 preventing fast brittle fracture (solid line, Figure 2g). It is well known in metallurgy that

182 crystallographic microstructure is able to stop and deflect the propagation of cracks when the
183 size of microstructure is comparable to the crack length^{28,29}. The mimicry of microstructure on a
184 macro-scale makes polycrystal-like structures comparable to macro-cracks. The change in lattice
185 orientation across IHABs stops the propagation of cracks, and most importantly, thanks to the
186 presence of IHABs, the architected material retained its strength after yielding and was able to
187 carry load up to large deformation, substantially improving the energy absorption of the material
188 (Figure 2g): from $\sim 194 \text{ kJ/m}^3$ by the singly oriented lattice to $\sim 1309 \text{ kJ/m}^3$ absorbed by the
189 designed polycrystal-like material. The drastic increase in the energy absorption associated with
190 the presence of IHABs is in agreement with a widely accepted approach employed to strengthen
191 alloys: the grain boundary engineering²⁵. The same architected material containing 8 IHAB
192 meta-grains was also printed by two other base materials (elasto-plastic polymer and an
193 austenitic stainless steel, Methods – Materials and fabrication). To study the influence of the
194 base material on the behaviour of polycrystal-like architected materials, two other base
195 materials were used to fabricate this 8 meta-grain architected material. The selection of base
196 material does affect the behaviour of macro-lattices: Compared to the same lattice fabricated by
197 the brittle polymerisation resin, there were no brittle cracks in the lattices made by the elasto-
198 plastic polymer and the steel because of the ductile behaviour of the two latter base materials
199 whilst, in contrast, the localisation of strain in the lattice made of the two materials remains very
200 similar (Extended data - Figure E9a, b). The localised deformation of struts led to the formation
201 of shear bands in the 30°- and 60°-oriented meta-grains in both the ductile polymer and metallic
202 lattices (Extended data - Figure E9a, b). The orientation change was able to alter the localisation
203 in the architected lattice fabricated by two different base materials, confirming this crystal-
204 inspired approach is widely applicable to various materials. While the shear band activity is
205 mainly governed by the lattice architecture, the constitutive stress-strain behaviour is strongly
206 affected by the base material because different base materials have different yield strength and
207 work-hardening behaviours. Most interestingly, the metallic lattice has substantial hardening
208 during plastic deformation much more compared to that of the polymer lattice (Extended data -
209 Figure E9c and d). This is because the base material (316 steel) itself hardens during deformation
210 because of the change in its intrinsic microstructure (in particular increase in the dislocation
211 density)³¹. This means that the application of the boundary hardening to metallic base materials

212 can enable the synergistic strengthening associated with both crystallographic microstructure and
213 designed meso-structures.

214 **Precipitation hardening**

215 Precipitation hardening is widely utilised in metallurgy to design and manufacture high
216 performance alloys, e.g., Ni-based superalloys³². In precipitate-hardened alloys, precipitates act
217 as obstacles to the movement of dislocations. The influence of precipitates on the strength of
218 alloys depends on the size, shape, volume fraction and distribution of precipitates, and the
219 coherency between precipitates and matrix³³. Precipitation strengthening was incorporated into
220 lattice design by introducing embedded lattice domains (termed *meta-precipitates*) that have
221 different lattice parameters from those in the matrix. Lattice parameters (type, spacing and
222 orientation) of meta-precipitates can be tailored to assign the different degrees of coherency
223 between them and the matrix. In addition, the strut diameters of meta-precipitates can be varied
224 to account for different “atomic” bond strengths, e.g., the Ni-Al and Al-Al bondings in the
225 gamma’ L₁₂ in Ni superalloys³².

226

227 Figure 3: Precipitation and multiphase hardening in architected materials. (a) A single oriented
228 without meta-precipitates. (b) Constitutive stress-strain responses of architected materials
229 without meta-precipitates (a) and with meta-precipitates (c, d). (e) Mechanical behaviour of
230 single phase versus multiphase architected materials. (f) Pseudo-superelasticity of Kresling
231 lattice, (g) FEM simulation shows the strain localisation and local buckling of struts, and (h)
232 Energy per unit volume of the first five cycles.

233

234 A single FCC architected material containing meta-precipitates that were face-centred
235 tetragonal and stronger than the matrix thanks to shorter lattice spacings was fabricated by 3D
236 printing (Extended data - Figure E4). Fabricated FCC lattices containing meta-precipitates
237 demonstrated a highly reproducible behaviour (Extended data - Figure E7d). It appears that the
238 presence of the harder meta-precipitates had a governing effect on the shear band propagation
239 (Figure 3a versus Figure3c-d): stabilising the overall lattice and substantially strengthening the
240 overall architected materials (Figure 3b). Shear bands were stopped at and bowed around the
241 interfaces between meta-precipitates and matrix (Figure 3d). This governing effect has the same
242 dependence as seen in the Orowan hardening effect in which the strength of crystals containing

243 precipitates is proportional to the strength of precipitates and spacing between precipitates^{33,34}.
244 This remarkable parallel between architected materials (containing meta-precipitates) and
245 alloys (containing precipitates) means that the mechanical response of architected materials
246 can be readily manipulated by varying the lattice parameters, the volume fraction and
247 distribution of meta-precipitates.

248 **Multi-phase hardening**

249 In multi-phase metals, such as martensitic stainless steels, hard phases offer high strength
250 to enhance load-bearing capacity while soft phases accommodate the plastic deformation. As a
251 crystal phase is defined by its lattice type, multiple phases can be imitated by assigning different
252 lattice types to different macro-lattice domains. An architected material comprising two phases
253 (FCC in the top and bottom layers and BCC in the middle layer, Extended data - Figure E5) was
254 designed. Because the unit cell of the FCC macro-lattice has a higher connectivity and higher
255 density than the BCC, the FCC is stronger than the BCC one (solid grey versus dashed curve,
256 Figure 3e). Similar to the behaviour of soft/hard multiphase crystals, the strength of the multi-
257 phase architected material mainly results from the strength of the hard FCC phase (black solid
258 versus grey solid lines, Figure 3e). By contrast, we observe that plastic deformation was mainly
259 accommodated by the soft BCC meta-grains, resulting in the confinement of shear bands mainly
260 to the middle layer (top left inset, Figure 3e). Therefore, the mixture of hard and soft
261 architected phases offers additional means of tailoring properties and controlling shear bands
262 along specific paths inside of an architected materials.

263 Taking this one step further, some highly engineered alloys can also exhibit a solid state
264 phase transformation whereby the crystal structure can transform by rearrangement of atoms in
265 the unit cell under external loads. In some cases, reversible phase transformation lead to a shape
266 memory effect and superelasticity³⁵. To see whether such a transformation were a feasible
267 option for crystal-inspired materials, a Kresling lattice³⁶ was designed to mimic a hexagonal
268 close packed phase (Extended data - Figure E6). Under compression, the Kresling lattice twisted
269 due to bending and local buckling of struts on prismatic and basal planes (Figure 3g and Videos
270 E2 and 3), causing a helical movement of the nodes (Extended data - Figure E10), which altered
271 the arrangement into a different one, enabling a “phase transformation”. It is worth noting that
272 for accurate mimicry of a specific type of phase transformation in crystals, one simply needs to

273 perform some additional design calculations of the helical movement of nodes to enable the
274 change in the stacking sequence (e.g., ABABAB to ABCABC for HCP to FCC transformation as
275 described Zhao *et al.*³⁷). The transformed phase was not stable during unloading, but it will be
276 straightforward to generate a lattice structure that is stable in both the original and transformed
277 configuration by incorporating the design of multi-stable cellular structures^{38,39}. Most
278 interestingly, the transformation back to the original arrangement upon unloading leads to
279 superelasticity (Figure 3f). The changes in elastic modulus observed (Figure 3f) are as a result of
280 a change from overall elastic bending to local buckling of struts on the prismatic planes (Figure
281 3g). The increasing constraint in compression caused a gradual transition from the overall elastic
282 bending to local buckling of struts near basal nodes. During unloading, the decreasing constraint
283 led to a more abrupt change: struts unbuckled at defined strain (from 14 % to 12 % during
284 unloading), leading to a change in bending moment. Such a change caused different friction on
285 the interacting faces between the Kresling lattice and compressing plates, resulting in the energy
286 dissipation during transformation (Figure 3h). This reversible alternation of node arrangements
287 demonstrates that phase transformations can be mimicked in architected phases, opening up
288 the possibility of making lattices with high energy absorption: external work can be dissipated by
289 a phase transformation, and the lattice then returns to its original shape via reverse
290 transformation upon unloading.

291 **Multiscale fractal crystal structures: Å to mm and beyond**

292 Combining all sources of tailoring properties in architected materials containing
293 crystal-like meso-structures, it will be possible to design lightweight components in which meso-
294 structures are designed in such a way that better responds to the external load, or to direct the
295 accumulated damage (shear bands and cracks) along specific structural paths as desired (Figure
296 4a-d). An exciting avenue of this approach is when a crystalline material is used to fabricate
297 crystal-like architected materials (Extended data - Figure E9a, c), it leads to the generation of
298 a new family of architected materials comprising highly scalable fractal crystal structure
299 consisting of crystallographic microstructure on μm and crystal-inspired mesostructures on mm
300 and beyond (depending on the size of build volume of manufacturing methods), i.e., micro-
301 crystals within macro-crystals. Such multi-scale fractal crystal structures are not available in
302 nature, making the new architected materials be of the meta-materials family, here coined

303 *meta-crystals*. For example, (i) a *mm*-size FCC (Figure 4e) contains an atomic FCC lattice
304 (Figure 4g), and (ii) a macro-polygrain structure on *cm* (Figure 4d) is made of a μm -size
305 polygrain microstructure of an austenitic stainless steel 316L (Figure 4f). Similarly, a fractal γ/γ'
306 super-lattice can be generated: micro- γ/γ' in Nickel superalloys (Figure 4h) can be scaled up by
307 γ/γ' -like super-lattice (note: different diameters and lengths need to be used for specific struts in
308 meta-precipitates to mimic the Ni-Al ordered atomic bonds in γ' -Ni₃Al phase). It is very easy to
309 change crystal lattices on the atomic scale and crystal microstructure (e.g., by changing the base
310 material) and architected lattice meso-structures *by design*, making it possible to engineer these
311 hierarchical lattices across multiple lengthscales. As discussed previously, the use of crystalline
312 alloys enables the synergistic strengthening associated with both crystal microstructure and
313 crystal-inspired meso/macro-structures. The interplay between microstructure and designed
314 structures on different lengthscales might lead to exciting properties and potential applications
315 for this new class of materials.

316

317 Figure 4: Lightweight and damage-tolerant architected materials inspired by crystal
318 microstructure. (a)-(d) Lightweight lattice component. (e)-(g) FCC fractal lattices from atomic
319 up to cm scales, (h) γ/γ' scalable fractal super-lattices.

320

321

322 **Conclusions**

323 This current study presents a novel way of combining the hardening mechanisms in
324 crystalline alloys and architected materials to transform the way of designing materials with
325 desired properties. A comprehensive study was carried out to substantiate: (1) the similarities
326 between shear bands in crystal-inspired architected materials and dislocation slip in crystals
327 and (2) the applicability of strengthening mechanisms (grain boundary, precipitation and
328 multiphase) to architected materials. This study confirms that metallurgical concepts can be
329 applied to enhance the mechanical behaviour of architected materials. In return, this crystal-
330 inspired approach provides alternative ways to study complex phenomena in metallurgy. This
331 unique way of fusing physical metallurgy and architected meta-materials opens new
332 opportunities in (1) designing and engineering damage-tolerant architected materials with

333 desired strength and toughness and (2) enhancing the functionality and performance of
334 architected materials in response to external loads. Possibilities offered by this approach are
335 not limited to those presented here. The application of this approach to crystalline metallic alloys
336 opens an exciting frontier of research both experimentally and computationally to understand the
337 possibilities that varying both intrinsic microstructure and designed meso-crystal structures of
338 meta-crystals can afford us.

339

340

341 **Data availability** The datasets generated during and/or analysed during the current study are
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343

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352

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354 computer-aided designs, fabrication and mechanical tests and post analyses. J.L. performed FEM
355 simulations. I.T. discussed and contributed to the further development of the concept. All the
356 authors participated in analysing and interpreting the data. The manuscript was written and
357 approved by the authors.

358

359 **Competing interest statement** The authors declare the following competing interests: a patent
360 developed on the basis of the approach proposed in this study was filed and managed by Imperial
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444

445

Methods

446 1. Design

447 A unit cell of crystal lattice is defined by its lattice constants (which are the distances between
448 atoms along three principal axes x , y , z : a , b and c , Extended data - Figure E1a) and the angles
449 between them (Figure E1a). For cubic lattices, x , y and z axes are parallel to the three [100]
450 orientations of lattice, $a=b=c$ and $\angle(a,b)=\angle(b,c)=\angle(c,a)=90^\circ$. All macroscopic FCC unit cells in
451 this study had the same node arrangement and connection as the one shown in Extended data -
452 Figure E1a. Macroscopic cubes containing crystal-like mesostructures were designed and
453 fabricated to study the roles of such meso-structures on the behaviour of architected materials,
454 e.g., Extended data - Figure E1b. A global (X,Y,Z)-coordinate attached to each cube was defined
455 by the three orthogonal directions of the cube (Extended data - Figure E1b). The global
456 dimension of macroscopic cubes in this study was $40\text{mm} \times 40\text{mm} \times 40\text{mm}$. All lattice models
457 presented in this study were designed by Solidworks and Element softwares.

458 1.1. Mimicry of polygrains

459 The global cubic matrix was fragmented into Voronoi domains. Each domain mimics a
460 crystalline grain: It is infilled by macro-lattice whose orientation of lattice is different to those in
461 adjoining domains. To create two twin meta-grains of FCC lattice (Figure 2b), a cube of single
462 oriented FCC lattice was first created with its three $\langle 001 \rangle$ orientations being parallel to the X, Y,
463 Z axes (Extended data - Figure E1b and c). The FCC unit cell had lattice constants of $5\text{mm} \times$
464 $5\text{mm} \times 5\text{mm}$ and struts diameter of 1mm. Subsequently, a plane at the centre of the global cube
465 and perpendicular to the X axis was defined as the boundary between two twin meta-grains.
466 Lattices on each side of the boundary were rotated to the same angles but in opposite directions:
467 lattices in the right side of the boundary (Extended data - Figure E2c) were first rotated counter-
468 clockwise by 45° about X axis (Extended data - Figure E1d), and then 45° about the Y axis to
469 generate a meta-grain of FCC lattice. Such a rotation sequence for the meta-grain is shown in
470 Extended data - Figure E1c-e where x'' , y'' and z'' axes were the final coordinates of lattices
471 after the two counter-clockwise rotations. By contrast, lattices in the other side were rotated
472 clockwise by 45° about the X, following by 45° about the Y axis to create the second twinned
473 meta-grain. The two constructed meta-grains were shown in Extended data - Figure E2a with the
474 one on the right-hand side highlighted. Because of the globally constrained dimension ($40\text{mm} \times$
475 $40\text{mm} \times 40\text{mm}$), the rotation of lattices caused incomplete unit cells on free surfaces (and at the
476 boundaries between meta-grains) (Figure 2a), this weakened the architected material
477 containing multiple meta-grains with different orientations. A planar lattice frame with a 2D unit
478 cell containing a centred node was introduced to help maintain the connectivity of struts
479 (Extended data - Figure E2b). The dimension of the planar unit cell is $5\text{mm} \times 5\text{mm}$ with struts
480 diameter of 1mm. The size of meta-grains was reduced by further dividing the global cube into
481 smaller meta-grains, i.e., increasing the number of meta-grains from 2 to 4, 8, 16, 18 and 27
482 (Extended data - Figure E2b-e).

483 A cube (Extended data - Figure E3) consisted of 8 meta-grains separated by incoherently high
484 angle boundaries was designed to study the role of misorientation on crack propagation (Figure
485 2d) and the deformation of lattices in different base materials (Extended data - Figure E9). All
486 meta-grains had the same size: $20\text{mm} \times 20\text{mm} \times 20\text{mm}$ (Extended data - Figure E3a). The unit
487 cell was FCC with constants of $5\text{mm} \times 5\text{mm} \times 5\text{mm}$ and the strut diameter of 1mm. The cube
488 consisted of two layers, each comprises 4 meta-grains. In the top layer, the top right meta-grain
489 with the three [100] orientations coinciding with the global X, Y and Z axes (Extended data -
490 Figure E3b) was chosen as the reference meta-grain. The unit cells of the three neighbouring
491 meta-grains (moving anticlockwise) in the same layer were created by rotating the unit cell of the
492 reference meta-grain by 30° , 45° and 60° about the Z axis (Extended data - Figure E3b). The
493 reference point of rotation was the centre of the original unit cell. Meta-grains in the bottom
494 layer were created in the similar way, but with a different sequence order (Extended data - Figure
495 E3c) with an aim to increase the randomness of the misorientation distribution. The outer frame
496 was not introduced to reveal the stopping effect of high angle boundaries on the crack
497 propagation (Figure 2g).

498 **1.2. Mimicry of precipitates**

499 Precipitates were mimicked by embedded lattice domains (termed meta-precipitates) that have
500 different lattice parameters from those in the matrix. Lattice parameters (type, spacing,
501 orientation) of meta-precipitates can be tailored to design the different degrees of coherency
502 between precipitates and the matrix. The diameters of strut lattices in meta-precipitates can be
503 varied to account for different “atomic” bonds to mimic an ordered arrangement of atomic
504 bondings, e.g., the Ni-Al in γ' in Nickel superalloys. Meta-precipitates in Figure 3c and d were
505 designed by embedding twenty-five cubic meta-precipitates in a FCC matrix. Meta-precipitates
506 near the free surfaces of the global cube and their locations in the matrix were shown in
507 Extended data - Figures E4a and b, respectively. The [001]-orientations of the matrix unit cell
508 were parallel to the global X, Y and Z axes (Extended data - Figure E4a). The unit cell of the
509 matrix was FCC with the lattice constants of $5\text{mm} \times 5\text{mm} \times 5\text{mm}$ and the strut diameter of 1mm,
510 Extended data - Figure E4c. Each meta-precipitate had the dimension of $7.5\text{mm} \times 7.5\text{mm} \times$
511 7.5mm , and was made of a face-centred-tetragonal (FCT) unit cell with lattice constants of 3mm
512 $\times 3\text{mm} \times 4\text{mm}$ and the strut diameter of 1mm (Extended data - Figure E4d). The orientations of
513 meta-precipitates were randomly rotated. Frames were introduced as the interfaces between the
514 matrix and precipitates to increase the connectivity of struts across the interfaces.

515 **1.3. Mimicry of phases**

516 As a crystal phase is defined by its lattice type, multiple phases can be imitated by assigning
517 different lattice types to different macro-lattice domains. Single FCC (or single BCC) macro-
518 phase was generated by an FCC (or BCC) unit cell (Extended data - Figure E5). The global
519 dimension of the multiple meta-phases was $40\text{mm} \times 40\text{mm} \times 40\text{mm}$. The FCC and BCC unit
520 cells had the same lattice constants ($5\text{mm} \times 5\text{mm} \times 5\text{mm}$) and the same strut diameter (1mm).

521 The mixture of the two FCC and BCC phases consisting fifty Voronoi meta-grains was designed
522 by filling the 50 meta-grains with the FCC and BCC unit cells (Extended data - Figure E5c). The
523 global (40mm × 40mm × 40mm) cube comprising the two meta-phases was divided into three
524 layers: top, middle and bottom layers. Meta-grains in top and bottom were filled by the FCC
525 phase while the middle layer meta-grains were made of the BCC phase: 11 FCC meta-grains for
526 the top layer, 25 BCC meta-grains for the middle layer, and 14 FCC meta-grains for the bottom
527 layer. The orientations of lattices in all the meta-grains were randomly rotated to increase the
528 isotropy of the lattice cube.

529 **1.4. Kresling lattice**

530 A Kresling lattice cylinder was constructed to imitate a hexagonal close packed (HCP) phase
531 (Extended data - Figure E6a) (Methods reference 1). Struts and nodes on the bottom (*ABCDEF*)
532 and top (*A'B'C'D'E'F'*) surfaces of the Kresling unit cell formed two regular hexagons that are
533 parallel and equal to the other. Similar to the terminology in the HCP crystal phase, the
534 (*ABCDEF*) and (*A'B'C'D'E'F'*) were the two basal planes. The bottom hexagon was created
535 with the side length of 8mm. Subsequently, the top hexagon *A'B'C'D'E'F'* was formed by
536 translating the bottom hexagon along the z axis (i.e., c-axis), and then rotating the translated
537 hexagon clockwise of 30° about the c-axis (Extended data - Figure E6a). The $\angle AOH$ (with H is
538 the middle of AF) was 30°, making HA' perpendicular to AF. The c-parameter (i.e., the length of
539 OO') of the Kresling HCP unit cell was 8mm. The unit cell was created by connecting *AA'*, *A'B*,
540 *BB'*, *B'C*, *CC'*, *C'D*, *DD'*, *D'E*, *EE'*, *E'F*, *FF'* and *F'A*. The Kresling HCP-like phase was built
541 by stacking three Kresling unit cells along the c-axis. The total height of the Kresling HCP-like
542 phase was 24 mm with strut diameter of 3mm for *AA'*, *BB'*, *CC'*, *DD'*, *EE'* and *FF'*; 2mm for
543 the diagonal struts and for the basal struts (Extended data - Figure E6b).

544

545 **2. Materials and fabrication**

546 Various base materials were used to fabricate designed materials to demonstrate the applicability
547 of the polygrain-inspired approach. The choice of the base material was based on the
548 phenomenon of interest. For example, while ductile polylactic Acid (PLA) polymer and stainless
549 steel were for studying the shear band activity in architected materials, a polymerisation resin
550 that is brittle after curing and heat treatment was used to study the crack propagation in
551 architected materials.

552 Cubic (except for Figure 2g) and Kresling lattices were made of PLA filaments and a
553 thermoplastic co-polyester Natural FLEX 45 filaments, respectively. The filaments were
554 provided by the RS Limited. Lattices were fabricated by a fused deposition modelling (FDM)
555 Ultimaker 2 printer. Designed models of meta-crystals were sliced into sequential layers by a
556 pre-processing software Cura with the layer height of 0.1mm. The nozzle of the printer was
557 heated (to 210°C for PLA and to 215°C for co-polyester) to melt the filament. The nozzle

558 diameter is 0.4 mm. The molten material was extruded and deposited on a build plate layer upon
559 layer according to the sliced sequences of designed lattices. The build platform was heated (to
560 60°C for the PLA and to 100°C for co-polyester) to increase the bonding between layers. The
561 nozzle speed was 30 mm/s for PLA and 15 mm/s for co-polyester

562 To study the role of meta-grain boundaries in the propagation of macro-cracks, both single meta-
563 grain and 8 meta-grains separated by incoincidentally high angle boundaries (Extended data -
564 Figure E3) were made of a photopolymerisation resin (Visijet M2 RWT) provided by 3D
565 Systems Limited. The resin was brittle after curing and subsequently heat treatment. The meta-
566 crystals were printed by a stereolithography (SLA) printer (ProJet MJP 2500) in which resin is
567 cured by Ultraviolet (UV) light. Similar to FDM process, meta-crystal models were firstly sliced
568 into layers with thickness of 0.1mm. The UV light cured resin according to the sliced sequence
569 of designed models. A wax material was used to support built lattices, this helps improving the
570 geometrical accuracy and quality of lattices. The wax was removed from built lattices by post-
571 treatment in an oven with temperature of 60°C for a duration of 5 hours.

572 To study the influence of the base material on the behaviour of polycrystal-like lattices, the 8
573 meta-grains (Extended data - Figure E3) were fabricated by both the elasto-plastic PLA and a
574 316L austenitic stainless steel. While the PLA was printed by Ultimaker 2 with parameters given
575 above, the steel was fabricated using a Renishaw AM250 printer in argon atmosphere. The
576 dimensions of the 8 meta-grains were reduced by a factor of 2 for 316L, i.e. the global
577 dimensions: 20mm x 20mm x 20mm, the size of each meta-grain: 10mm x 10mm x 10mm,
578 lattice constants: 2.5mm x 2.5mm x 2.5mm and strut diameter: 0.5mm. A hatch pattern was used
579 for inner area, with a power intensity of 180 W, a spacing of 150 μm , an exposure time of 110
580 μs and a point distance of 65 μm . The outer skin was made by two contours with a power
581 intensity of 100 W, a spacing of 55 μm , an exposure time of 90 μs and a point distance of 40 μm .

582 The mechanical properties of the PLA and co-polyester filaments, brittle resin and 316L steel are
583 given in Table E1.

584

585 **3. Microstructure characterisation**

586 Scanning electron microscopy (SEM) imaging and electron backscatter diffraction (EBSD) were
587 carried out on an austenitic stainless steel 316L using a Zeiss Auriga SEM at a voltage of 20 kV
588 for Figure 4e, f. γ/γ' microstructure in a Ni single crystal (with a composition similar to CMSX4)
589 was revealed by chemical etching (2.5% phosphoric acid in methanol solution, applying 10 V for
590 30 s). Microscopic observation of γ/γ' microstructure (Figure 4h) was done using SEM
591 Sigma300 at a voltage of 10kV.

592 **4. Mechanical tests and analyses**

593 Mechanical properties of all macro-lattices were studied by compression tests that were carried
594 out by a 100kN Zwick machine using the displacement control at a strain rate of 10^{-3} 1/s at room
595 temperature. To minimise the effect of friction, the top and bottom compression plates were
596 lubricated with dry molybdenum disulphide. The loading direction was parallel to the build
597 direction (i.e., the Z direction) of prints. Stresses were calculated by dividing the recorded forces
598 by the nominal area, which enclosed lattices and was perpendicular to the Z direction.
599 Engineering strains were derived by dividing the change in the length along the Z direction by
600 the initial length: 40mm for the polymer lattice cubes or 24mm for the Kresling lattice. The
601 compression tests of meta-crystals were recorded by a Nikon D7100 camera with 200mm Nikkor
602 macro lenses for taking images for post-test analyses, and Canon SX210 for recording videos.
603 Captured images were characterized by digital image correlation (DIC) via a commercial
604 software DaVis. The resolution of analyzed image is $6000\text{ pixels} \times 4000\text{ pixels}$. The subset
605 dimensions used in DIC analysis were $101\text{ pixels} \times 101\text{ pixels}$ with step size of 25 pixels .
606 Matlab subroutines were written to analyse the stress-strain behaviour of tested lattices. Fitting
607 by a linear function over different ranges of the most linear part (up to 2% strain) of a stress-
608 strain curve was done to identify the elastic modulus. The yield strength was defined as the stress
609 corresponds to the 0.2% strain offset of the fitting line.

610 **5. Finite element analysis**

611 Finite element method (FEM) was used to simulate the deformation of macroscopic lattices.
612 Model of macro-lattices generated by nTopology Element software were imported to Abaqus for
613 FEM simulations. The quasi-static compression of a macro-lattice was modelled using two
614 displacement controlled rigid-body plates. The first rigid body plate was fixed underneath the
615 macro-lattice with an encastre boundary condition. The second rigid body plate was used to
616 compress the macro-lattice with a displacement boundary condition. The General Contact
617 algorithm available in Abaqus was employed to simulate the interactions between all surfaces,
618 with a penalty friction coefficient of 0.1 and 'hard' pressure-overclosure surface behaviour
619 defined. Macro-lattices were meshed using 10-node quadratic tetrahedral elements (C3D10) with
620 isotropic elastic-plastic material behaviour that was defined according to the mechanical
621 behaviour data given in the Table E1. Because of the symmetry and periodicity of the
622 arrangement of unit cells in the meta-grain twins, a reduced FEM model of the meta-grain twins
623 was constructed. The FEM model consists the twin boundary and 3x3 unit cells in each meta-
624 grain, and was constructed by 1,184,608 elements. The Kresling lattice was built by 140,120
625 elements. The rigid body plates were meshed using 4-node bi-linear quadrilateral elements
626 (R3D4).

627

628 **Methods references**

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634

635

636 **Main figure captions**

637 Figure 1: Lattice structures and deformation behaviour. (a) Face-centred-cubic (FCC) crystal
638 lattice, (b) Single slips in a single crystal (re-used from ⁸ with permission from Elsevier), (c)
639 Slips at grain boundary in a polycrystalline steel (re-used from ⁹ with permission from Elsevier),
640 (d) Architected FCC lattice, (e) Single slip in a single oriented lattice, (f) Unstable behaviour of
641 architected metallic lattices (reused from ¹⁰ under the Creative Commons Attribution 4.0
642 International License).

643 Figure 2: Roles of lattice orientation in the deformation behaviour of crystals and architected
644 lattices. (a) Twin bi-crystal (re-used from ²⁷ with permission from Elsevier), Shear bands in
645 meta-grain twins observed in experiment (b) and predicted by FEM (c) (Note: Cut sections
646 shows the deformation of internal struts. The sections were formed thanks to two cutting planes
647 that were parallel to $\{001\}$ planes of FCC lattice), (d) – (e): Shear bands were controlled by
648 orientation of meta-grains: (d) 8 meta-grains and (e) 16 meta-grains (note the nominal strain was
649 of 30%). (f) Yield strength versus the size of meta-grains, (g) Boundaries between meta-grains
650 effectively stops cracks in brittle lattices, leading to a drastic increase in toughness of
651 architected materials.

652

653 Figure 3: Precipitation and multiphase hardening in architected materials. (a) A single oriented
654 without meta-precipitates. (b) Constitutive stress-strain responses of architected materials
655 without meta-precipitates (a) and with meta-precipitates (c, d). (e) Mechanical behaviour of
656 single phase versus multiphase architected materials. (f) Pseudo-superelasticity of Kresling
657 lattice, (g) FEM simulation shows the strain localisation and local buckling of struts, and (h)
658 Energy per unit volume of the first five cycles.

659 Figure 4: Lightweight and damage-tolerant architected materials inspired by crystal
660 microstructure. (a)-(d) Lightweight lattice component. (e)-(g) FCC fractal lattices from atomic
661 up to cm scales, (h) γ/γ' scalable fractal super-lattices.

662

663 **Extended data table captions**

664 Table E1: Mechanical properties of printed polymers. * data provided by RS Limited and 3D
665 Systems Limited. ** data obtained from 3 tensile tests of solid cylindrical samples (fabricated by
666 powder-bed selective laser fusion) at room temperature and a strain rate of 10^{-3} 1/s).

667

668 **Extended data figure captions**

669 Figure E1: Mimicry of crystal lattice. (a) Unit cell of lattice, (b) A macro-lattice cube consisting
670 of 8 x 8 x 8 macro-unit cells. (c) – (e) The rotation sequence to form a twin meta-grain of lattice.

671

672 Figure E2: A different number of meta-grains within the same global volume (40mm x 40mm x
673 40mm). (a) 1 meta-grain, (b)-(c) 2 twinned meta-grains: (b) with outer frame and (c) without the
674 outer frame, (d) 4 meta-grains, (e) 8 meta-grains, (f) 16 meta-grains, (g) 18 meta-grains and (e)
675 27 meta-grains. The locations of boundaries were highlighted.

676

677 Figure E3: Mimicry of crystalline grains separated by incoherent high angle boundaries. (a) Model
678 of 8 meta-grains. The orientations of lattices in the four meta-grains in (b) the top layer and (c)
679 the bottom layer

680

681 Figure E4: Mimicry of precipitates. (a) Meta-precipitate lattice. (b) Cubic morphology and
682 locations of meta-precipitates inside the FCC meta-phase. (c) FCC unit cell of the matrix and (d)
683 FCT unit cell of meta-precipitate.

684

685 Figure E5: Mimicry of multi-phases. (a) Single meta-grain of FCC meta-phase, (b) Single meta-
686 grain of BCC meta-phase, (c) A cube of meta-polygrains consisting of two meta-phases: FCC
687 (top and bottom layers) and BCC (middle layer).

688

689 Figure E6: Kresling lattice. (a) unit cell, (b) HCP meta-phase

690

691 Figure E7: The repeatability of mechanical behaviour of architected materials. (a), (b) and (c)
692 materials consist of 2, 8 and 16 meta-grains, respectively; and (d) materials contain 25 meta-
693 precipitates.

694

695 Figure E8: Effect of the size of meta-grains. (a) Stress-strain curves of architected materials
696 consisting a different number of Voronoi domains of lattices. (b) Flow stress of architected
697 materials containing meta-grains at a given nominal strain of 40% increases with reducing the
698 size of meta-grains.

699

700 Figure E9: Deformation behaviours of an architected material containing 8 meta-grains
701 separated by incoherent high angle boundaries. (a) and (b) the macro-lattice was fabricated by
702 316L stainless steel and elasto-plastic polymer, respectively. (c) and (d) stress-strain constitutive
703 behaviour of the macro-lattices fabricated by the steel and polymer.

704

705 Figure E10: Helical movement enables the change in the stack sequence of nodes. Red lines
706 represent helical movements of basal nodes. Note: only the movement trajectory of basal nodes
707 on the top plane were shown by the red curves.