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A Conformal Thin Boundary Model for FDTD

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Abstract—Thin layer models are widely used in the finite-difference time-domain (FDTD) technique to efficiently model boundaries in multi-scale simulations as they significantly reduce simulation run-times and memory requirements. These models often utilise surface impedance boundary conditions (SIBCs) to represent the material of the boundary. Conformal meshes are a popular method of representing curved and non-aligned surfaces in FDTD. These meshes deform cells in the FDTD grid around the boundary between bulk materials so as to more accurately represent the shape of the material. Here we present an algorithm that combines the efficiency of a thin layer model with the accuracy of a conformal mesh. The algorithm is applied to three resonant cavity models and the accuracy verified using comparisons to non-conformal meshes and analytic solutions. Improvements are shown in the accuracy of the resonant frequencies and magnitude of the shielding effectiveness (SE) of the cavities. It is also shown to reduce the prevalence of extraneous features in the frequency response of the SE that are apparent when using a stair-cased mesh.

Index Terms—Finite-Difference Time-Domain, Conformal Techniques, Surface-Impedance Boundary Condition, Thin Layer

I. INTRODUCTION

FDTD is a popular numerical method for solving Maxwell's equations [1]. The strength of the method lies in its efficiency and simplicity. Being a time domain algorithm it is possible to solve problems over a large range of frequencies in a single simulation.

Surface Impedance Boundary Conditions (SIBCs) are useful tools that can be used to model very thin materials [2], [3]. They allow a material that is much thinner than the size of the mesh to be simulated without having to simulate the field inside the material. Instead, they use a behavioural model that controls how the material interacts with the mesh. Doing this can reduce simulation run-times significantly [4].

However it is commonly known that FDTD can suffer from inaccuracies due to its reliance on a cuboid grid. Surfaces that are curved, or do not align with the grid, cannot be accurately represented and are commonly approximated using a stair-cased mesh. This can introduce a number of errors into the results of the algorithm [5]. A popular solution to this problem is the introduction of conformal algorithms. These algorithms deform individual cells in the cuboid grid to

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better conform to non-aligned and non-orthogonal geometries. Early conformal algorithms could be used to represent perfect electrical conductors [6], [7] and later methods allowed for conformal meshes of penetrable materials [8]. Current conformal algorithms rely on meshing bulk materials, where the field inside the material is simulated on the FDTD grid. This means that having a bulk conformal material boundary also means having a larger simulation to account for the cells inside the material.

In this paper we propose a conformal algorithm that utilises a SIBC for representing curved and non-aligned planar surfaces in FDTD. This combines the efficiency of using a thin boundary model and the accuracy of using a conformal mesh. The algorithm is proposed in Section II and validated using different cavity models in Sections III, IV and V. Conclusions about the approach are presented in Section VI

II. PROPOSED CONFORMAL ALGORITHM

Here the adaptation of an existing SIBC [3] to a conformal mesh is described. Consider the cell in Fig. 1, the shaded boundary has been moved from the xy -plane face at i towards the centre to conform with a surface passing through the cell, by directly modifying a stair-cased surface mesh. Where l_1, l_2, l_3 and l_4 are the distance that each corner of the boundary has moved.

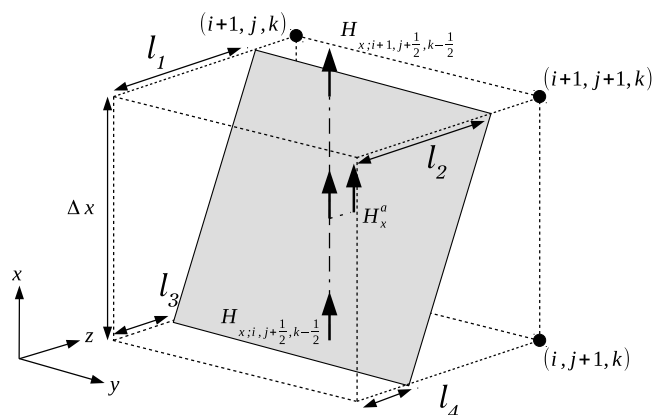


Fig. 1. Conformal algorithm applied to a single cell. The shaded surface is not parallel to any cell face.

The original SIBC uses spatial interpolation to approximate the magnetic field at the centre of the SIBC parallel to the

face, on the mesh cell surface. This is achieved by averaging the magnetic fields on the surface of the cell:

$$H_x^a = \frac{H_{x;i,j+\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}} + H_{x;i+1,j+\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}}}{2} \quad (1)$$

To apply a conformal SIBC the magnetic fields must be weighted according to their distance to the boundary, so (1) becomes:

$$H_x^a = \frac{\frac{1}{2}(\delta_1 + \delta_2)H_{x;i,j+\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}} + \frac{1}{2}(\delta_3 + \delta_4)H_{x;i+1,j+\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}}}{2} \quad (2)$$

Where $\delta_1, \delta_2, \delta_3$ and δ_4 are the fractional edge lengths:

$$\delta_n = \frac{l_n}{\Delta z} \quad (3)$$

Similar update equations can be derived for determining the y-polarised magnetic field.

The usual FDTD equations are used to update the surrounding magnetic field components using the electric fields on the SIBC surface modified by the average fractional edge length:

$$H_{x;i,j+\frac{1}{2},k-\frac{1}{2}}^{n+\frac{1}{2}} = H_{x;i,j+\frac{1}{2},k-\frac{1}{2}}^{n-\frac{1}{2}} + C_{i,j+\frac{1}{2},k-\frac{1}{2}}^{hxe} \left[\frac{1}{\Delta z_k} \left\{ \frac{1}{2}(\delta_2 + \delta_3)E_{y;i,j+\frac{1}{2},k}^{a;n} - E_{y;i,j+\frac{1}{2},k-1}^n \right\} + \frac{1}{\Delta y_j} \left\{ E_{z;i,j,k-\frac{1}{2}}^n - E_{z;i,j+1,k-\frac{1}{2}}^n \right\} \right] \quad (4)$$

Where $C_{i,j+\frac{1}{2},k-\frac{1}{2}}^{hxe}$ is the standard magnetic field update coefficient [9]. Again, similar update equations can be derived for the y-polarised magnetic fields. Any yz-plane and zx-plane faces on the staircase mesh in the cell are dealt with in the same way and the sum of the coincident boundaries add to give the correct overall behaviour.

III. 1D RESONATOR TEST CASE

The first validation test for the proposed algorithm is a simple 1D resonator. This test consists of two thin planar boundaries separated by 3.95m of free space. The boundaries are isotropic and symmetrical with constant transmission and reflection coefficients of 0.004 and -0.99 respectively. A Gaussian pulse is used to illuminate the resonator at normal incidence from the left as shown in Fig. 2.

The electric field at a given point in the cavity can be determined analytically as:

$$E(x) = \frac{E_I \tau (e^{-\gamma x} + \rho e^{-\gamma(2d-x)})}{1 - \rho^2 e^{-2\gamma d}} \quad (5)$$

where E_I is the magnitude of the incident electric field, x is the distance from the left side of the cavity, d is the length of the cavity, γ is the propagation constant and τ and ρ are

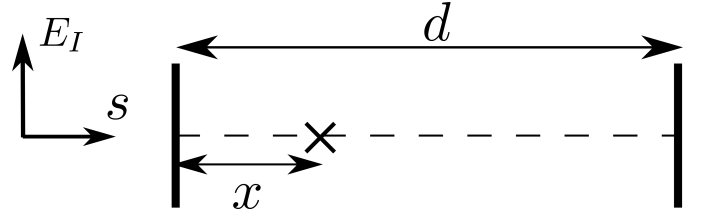


Fig. 2. 1D Resonator Diagram. E_I is the incident electric field and s is the direction of propagation.

the transmissions and reflection coefficients of the boundaries respectively.

The initial simulation for this problem uses a cavity length of 3.95m and a mesh size of 0.2m. This means that, when using a non-conformal mesh, the cavity length must be a multiple of 0.2m and therefore the cavity modelled is effectively 4m long. When using the conformal mesh one cell can be deformed allowing a 3.95m long model while using a 0.2m global mesh size. The electric field along the cavity at 37MHz is shown in Fig 3.

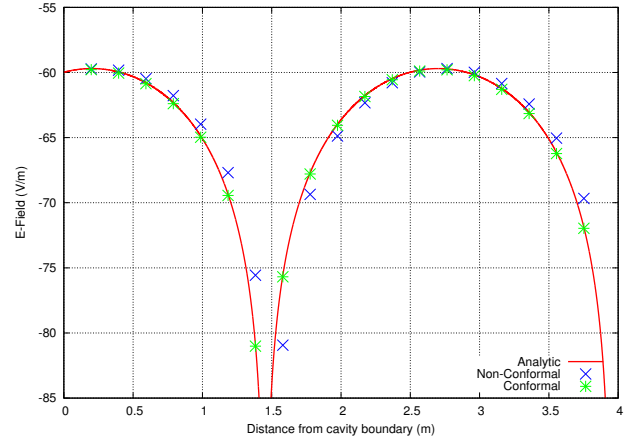


Fig. 3. E-Field along 1D Cavity at 37MHz. Comparison of analytic solution with conformal and non-conformal meshes.

It can be seen that for the non-conformal case the field is shifted slightly in space, however when using the conformal mesh this has been corrected and the field values lie on the analytic curve.

The simulation is repeated using the same mesh size and number of cells as before, the size of the cavity is varied from 3.81m to 4m. The error in the first resonant frequency is measured for each cavity size, the results are shown in Fig. 4. The conformal mesh is adjusted to match the cavity width.

It can be seen that the error in resonant frequency is better in all cases when using the conformal mesh except when the cavity has a width of 4m and the conformal mesh is identical to the non-conformal mesh. A width of 3.8m cannot be simulated using the conformal mesh as the boundary would lie exactly on the next surface of the cell, meaning the distance between the electric field nodes and the boundary would be zero.

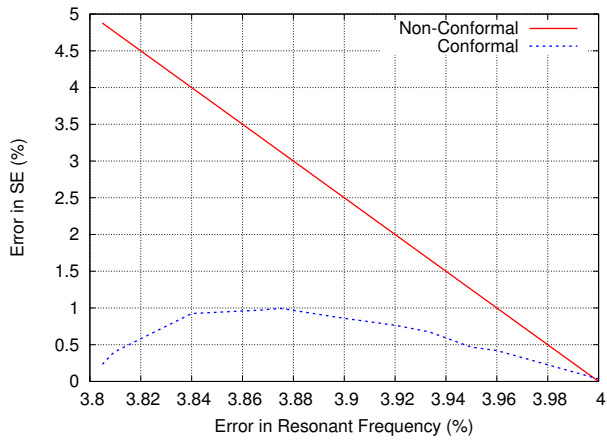


Fig. 4. Error in resonant frequency for 1D cavities of different widths. The number of cells across the cavity and the mesh size remain constant.

IV. CUBIC SHELL TEST CASE

To test the conformal algorithm on a three dimensional problem, a cubic cavity is investigated. It is possible to mesh a cubic shell so that it aligns perfectly with the grid, obviating the need for a conformal mesh. However, if the cube is not aligned with the mesh it may become necessary to use a conformal, or stair-cased, mesh to represent the faces of the cube as shown in Fig. 5.

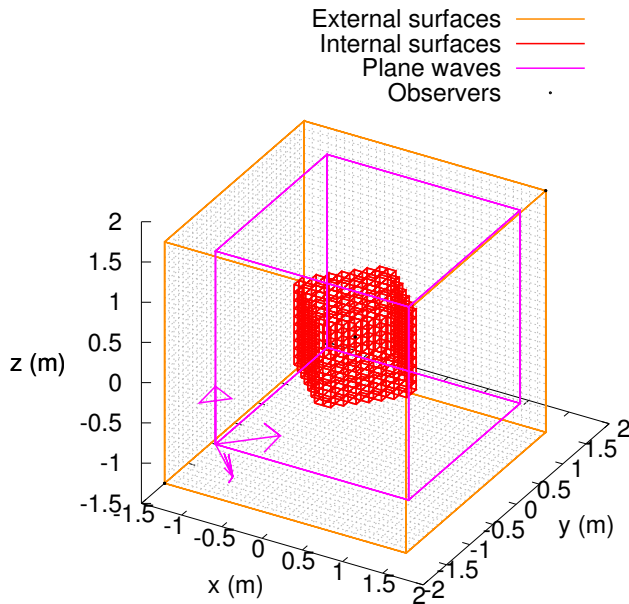


Fig. 5. Mesh of cubic shell.

In this case the cube has been rotated around a single axis causing four of the faces to become misaligned and require a stair-cased mesh. The cube is illuminated using a vertical, linearly polarised plane wave, at normal incidence to the face of the cube. The electric field at the centre of the cube is monitored. From this the shielding effectiveness (SE) at the

centre of the cube is determined and shown in Fig. 6. Here the SE is the ratio of the external and internal electric field strength. As there is no analytic solution for a cubic shell, a simulation of the cube that is aligned to the FDTD grid is used for comparison. For the aligned case no stair-casing is needed and the conformal mesh reduces to the standard algorithm.

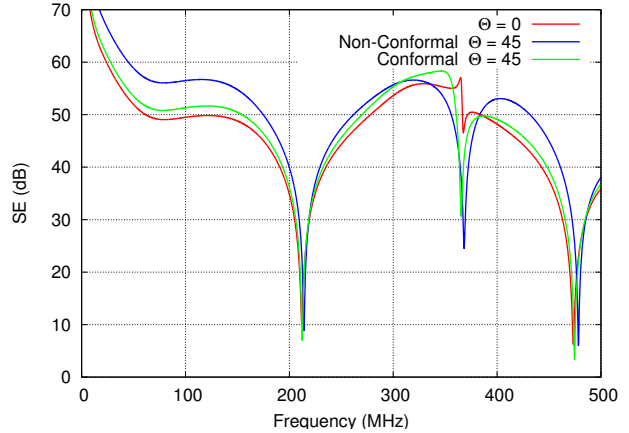


Fig. 6. Shielding effectiveness at the centre of a cubic shell rotated 45° with respect to the FDTD grid. Comparisons are made with an aligned case.

There is a significant improvement in the accuracy magnitude of the SE using the conformal mesh when compared to the stair-cased mesh, reducing the error from 7dB to approximately 1.5dB. A close up of the first resonance in Fig. 7 shows an error in the resonant frequency for the stair-cased mesh of around 1%, when using the conformal mesh this is reduced to approximately 0.1%. Although the resonant frequency is more accurate with the conformal mesh the amplitude error suggests there is still some inaccuracy in the transmission and reflection coefficients achieved. This will be further investigated in the future.

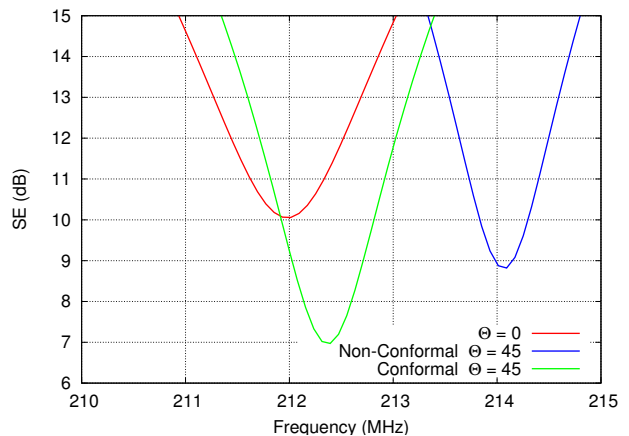


Fig. 7. Shielding effectiveness at the first resonant frequency of a cubic shell.

V. CYLINDRICAL SHELL TEST CASE

A third test case is shown to demonstrate the use of the conformal algorithm when modelling curved surfaces. An

infinite cylindrical cavity is meshed with a radius of 1m using the same boundary conditions as the previous test cases. Again the cavity is illuminated using a linearly polarised plane wave, with the incident electric field being polarised along the x-axis. Perfect Magnetic Conductor (PMC) boundaries are utilised at each end of the cylinder to simulate an infinitely long structure. The mesh size for this problem is 0.1m.

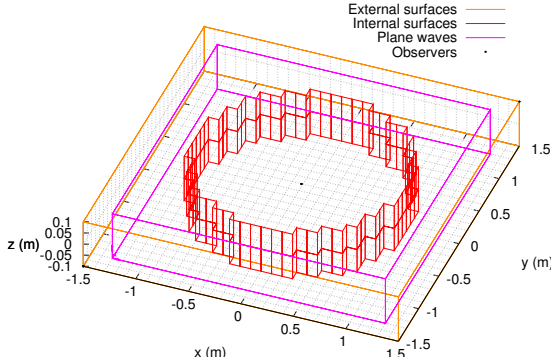


Fig. 8. Stair-cased mesh of an infinite cylinder for FDTD. Only one layer of cells is required in the Z-direction, PMC boundaries emulate an infinite length.

Again the E-field and subsequently SE is measured at the centre of the cavity, in this case there is an analytic solution [10] that can be used for comparison. The results are shown in Fig. 9

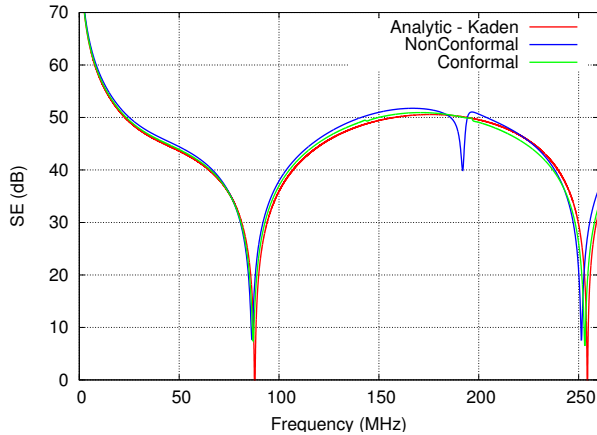


Fig. 9. Shielding effectiveness at the centre of a cylindrical cavity.

For the non-conformal mesh there is a clear resonance around 180MHz, this is not a spurious peak, rather it corresponds to a resonant mode that has a node at the centre of the cavity. For points within the cavity that are not exactly at the centre this resonance would be visible in the SE, the imprecise stair-cased approximation makes the centre of the mesh a slightly ambiguous position causing this resonance to become apparent at the centre as well. Using the conformal mesh has reduced this resonance significantly, making it almost non-existent at the centre of the cylinder. However, it should be noted that there is still a slight variation in the SE around this

frequency as the conformal mesh is still an approximation, albeit a more accurate one than the stair-cased mesh.

The error in the first resonant frequency is around 1.9% when using the stair-cased mesh, using the conformal mesh gives an error of approximately 0.4%. The error in magnitude at 50MHz is approximately 0.9dB using the stair-cased mesh, this error is improved to approximately 0.3dB using the conformal mesh.

It is worth mentioning that for the stair-cased mesh the extra resonance could be reduced, and the overall accuracy in frequency increased, using a higher resolution mesh. However, to achieve a similar result to the conformal mesh the runtime would be increased by a factor of 125. There would also be no improvement in the magnitude of the SE as errors in the boundary transmission and reflection due to stair-casing are independent of mesh size [5].

VI. CONCLUSION

A new algorithm for FDTD has been presented combining a surface impedance boundary with a conformal mesh. Multiple test cases have been introduced that demonstrate the new algorithm leads to improvements in accuracy, both in the frequency of resonant modes and the magnitude of shielding effectiveness of the cavities. We have demonstrated that the conformal thin layer algorithm can also be used to reduce and potentially eliminate the presence of extraneous resonances that occur using stair-cased meshing. The method is also easy to mesh as an existing staircase mesh is modified using information on the face position in each cell.

REFERENCES

- [1] K. Yee, "Numerical solution of initial boundary value problems involving Maxwell's equations in isotropic media," *IEEE Transactions on Antennas and Propagation*, vol. 14, pp. 302–307, May 1966.
- [2] V. Nayyeri, M. Soleimani, and O. Ramahi, "A method to model thin conductive layers in the finite-difference time-domain method," *Electromagnetic Compatibility, IEEE Transactions on*, vol. 56, pp. 385–392, April 2014.
- [3] I. D. Flintoft, S. A. Bourke, J. F. Dawson, J. Alvarez, M. R. Cabello, M. P. Robinson, and S. G. Garcia, "Face-centered anisotropic surface impedance boundary conditions in FDTD," *IEEE Transactions on Microwave Theory and Techniques*, vol. PP, no. 99, pp. 1–8, 2017.
- [4] J. Maloney and G. Smith, "The use of surface impedance concepts in the finite-difference time-domain method," *Antennas and Propagation, IEEE Transactions on*, vol. 40, pp. 38–48, Jan 1992.
- [5] S. A. Bourke, J. F. Dawson, I. D. Flintoft, and M. P. Robinson, "Errors in the shielding effectiveness of cavities due to stair-cased meshing in FDTD: Application of empirical correction factors," in *2017 International Symposium on Electromagnetic Compatibility - EMC EUROPE*, pp. 1–6, Sept 2017.
- [6] S. Dey and R. Mittra, "A modified locally conformal finite-difference time-domain algorithm for modeling three-dimensional perfectly conducting objects," *Microwave and Optical Technology Letters*, vol. 17, no. 6, pp. 349–352, 1998.
- [7] S. Benkler, N. Chavannes, and N. Kuster, "A new 3-D conformal PECFDTD scheme with user-defined geometric precision and derived stability criterion," *IEEE Transactions on Antennas and Propagation*, vol. 54, pp. 1843–1849, June 2006.
- [8] W. Yu and R. Mittra, "A conformal finite difference time domain technique for modeling curved dielectric surfaces," *IEEE Microwave and Wireless Components Letters*, vol. 11, pp. 25–27, Jan 2001.
- [9] S. D. Gedney, *Introduction to the finite-difference time-domain (FDTD) method for electromagnetics*. Morgan & Claypool, 2011.
- [10] H. Kaden, *Wirbelströme und Schirmung in der Nachrichtentechnik*, pp. 118–123. Technische Physik in Einzeldarstellungen, Springer, 1959.