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# A Novel Entropy-based Graph Signature from the Average Mixing Matrix

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**Abstract**—In this paper, we propose a novel entropic signature for graphs, where we probe the graphs by means of continuous-time quantum walks. More precisely, we characterise the structure of a graph through its average mixing matrix. The average mixing matrix is a doubly-stochastic matrix that encapsulates the time-averaged behaviour of a continuous-time quantum walk on the graph, i.e., the  $ij$ -th element of the average mixing matrix represents the time-averaged transition probability of a continuous-time quantum walk from the vertex  $v_i$  to the vertex  $v_j$ . With this matrix to hand, we can associate a probability distribution with each vertex of the graph. We define a novel entropic signature by concatenating the average Shannon entropy of these probability distributions with their Jensen-Shannon divergence. We show that this new entropic measure can encapsulate the rich structural information of the graphs, thus allowing to discriminate between different structures. We explore the proposed entropic measure on several graph datasets abstracted from bioinformatics databases and we compare it with alternative entropic signatures in the literature. The experimental results demonstrate the effectiveness and efficiency of our method.

## I. INTRODUCTION

Graph-based representations have proven to be powerful tools in structural pattern recognition. Compared to vector based pattern recognition, a major drawback with graph representations is the lack of a natural correspondence order for vertices. This limits the direct application of standard machine learning algorithms to problems such as classifying or clustering graphs.

One way to overcome this problem is to embed graphs onto a vectorial space by computing some entropy-based vectorial signature. Several approaches have been proposed in the literature to associate an entropic complexity measure to a graph. Dehmer et al. [1], [2] have developed a means of computing entropies of undirected graphs by using information functionals. For a given graph, the information functional is derived from its topological structure and quantify its information content. Anand et al. [3] and Passerini et al. [4] have applied the von Neumann entropy or quantum entropy to graphs, by mapping quantum states into discrete graph Laplacians [5]. In particular, Passerini et al. note that if the discrete Laplacian [6] is scaled by the inverse of the volume of the graph the resulting matrix is a density matrix, whose entropy can be computed using the spectrum of the discrete Laplacian. Moreover, they indicate that when there is degree heterogeneity then the Shannon (classical) and von Neumann (quantum) entropies are correlated. Since the von Neumann entropy relies on the computation of the normalized Laplacian spectrum, its

computational complexity is cubic in the number of vertices. To render the computation of the von Neumann entropy more efficient, Han et al. [7] have shown how the computation can be rendered quadratic in the number of vertices of the graph by using a quadratic approximation of the Shannon entropy of the Laplacian eigenvalues. An analysis of the quadratic entropy reveals that it can be computed from a number of permutation invariant matrix trace expressions. This leads to a simple expression for the approximate entropy in terms of the degrees of adjacent vertices. Furthermore, to develop this work further, Ye et al. [8] have shown how the von Neumann entropy for undirected graphs can be generalized to directed graphs. To this end, they commenced by using Chung's [9] definition of the normalized Laplacian on a directed graph. According to this definition, the normalized Laplacian matrix of a directed graph is Hermitian, and so the interpretation of Passerini et al. in [3] still holds for the domain of directed graphs. By observing that the Von Neumann entropy of the normalized Laplacian is essentially the Shannon entropy associated with the normalized Laplacian eigenvalues, the authors provide an expression for the von Neumann entropy of directed line graphs which depends on the in-degree and out-degree of pairs of connected vertices.

It should be noted that most of these entropy measures can only capture limited structural information of a graph. This is because the entropy based on information functionals of Dehmer et al. [1] and the Von Neumann entropy approximations of Han et al. [7] and Ye et al. [8] are all based on the computation of relatively simple structural characteristics.

Another way to apply standard machine learning and pattern recognition techniques to graphs is to use kernels methods [10]. The idea underpinning graph kernels is to measure the similarity between two graphs by first decomposing them into simpler substructures and then enumerating the number of substructures they share. For example, random walk graph kernels describe the structure of a graph using random walks [11]. More recently, graph kernels based on quantum walks have been shown to easily outperform their classical counterpart [12], [13]. Unlike explicit embedding methods, however, graph kernels provide an indirect embedding, where only the dot product of the graphs in the embedding space is known. With respect to entropic signature, graph kernels tend to have a higher computational cost.

### A. Contribution

In this paper, we propose to exploit the rich expressive power of quantum walks to define a novel entropic signature

which is able to capture the structural information of graphs better than existing entropic signatures. With respect to existing quantum walk graph kernels, we propose to use the average mixing matrix [14] to summarize the time-averaged behaviour of the walk. The reason for using continuous-time quantum walks and the average mixing matrix is twofold. Firstly, quantum walks have been shown to discriminate between different structures better than their classical counterparts [12], [13]. For example, while the long-term behaviour of classical random walks is governed by the leading eigenvector of the graph Laplacian, all the eigenvectors of the Hamiltonian (i.e., the graph Laplacian or the adjacency matrix, depending on the setting) contribute to the long-term behaviour of a quantum walk. Secondly, the average mixing matrix concisely summarizes the time-averaged behaviour of continuous-time quantum walk. In this sense, the average mixing matrix represents the closest quantum analogue of the steady state of a classical random walk, which has been used to measure the Shannon entropy of a graph [15]. Therefore, the average mixing matrix represents a natural choice for our entropic signature based on quantum walks.

The proposed entropic signature is computed as follows. Given a graph, we commence by computing the average mixing matrix associated with it. We then show how this allows us to associate a probability distribution over the graph vertices, for each vertex of the graph. Then, we define our entropic signature by concatenating the average Shannon entropy of these probability distributions with their Jensen-Shannon divergence. We show that this new entropic measure can encapsulate the rich structural information of the graphs, thus allowing to discriminate between different structures. We explore the proposed entropic measure on several graph datasets abstracted from bioinformatics databases and we compare it with alternative entropic signatures in the literature. The experimental results demonstrate the effectiveness and efficiency of our method.

The remainder of this paper is organized as follows. Section II introduces the concepts of continuous-time quantum walks and average mixing matrix. We then show how this allows us to associate a probability distribution with each vertex of the graph. Section III gives the definition of the proposed entropic signature for graphs. Section IV provides experiments, while Section V concludes our work.

## II. PRELIMINARY CONCEPTS

In this section, we commence by introducing the concepts of continuous-time quantum walk and average mixing matrix. We then proceed to show how to associate a probability distribution with each vertex of a graph based on the corresponding average mixing matrix.

### A. Continuous-time Quantum Walks

The continuous-time quantum walk represents the quantum analogue of the continuous-time random walk [16]. A continuous-time random walk models a Markovian diffusion process over the vertices of a graph, where transitions are only allowed between adjacent vertices. Let  $G = (V, E)$  be an undirected graph with vertex set  $V$  and edge set  $E$ , and let  $p(t) \in \mathbb{R}^n$  denote the state of walk at time  $t$ , i.e., the  $i$ -th

component of  $p(t)$  is the probability of the random walker to be at vertex  $v_i$  at time  $t$ . In a continuous-time random walk the state vector evolves according to the equation

$$p(t) = e^{-Lt}p(0), \quad (1)$$

where the graph Laplacian  $L = D - A$  is the infinitesimal generator matrix of the underlying continuous-time Markov process,  $A$  denotes the adjacency matrix and  $D$  is the diagonal matrix with the vertex degrees as diagonal elements.

Similarly to its classical counterpart, the state space of the continuous-time quantum walks is the vertex set of the graph. The state vector of classical random walks is replaced by a vector of complex amplitudes over  $V$  whose squared norm sums to unity, and as such the state of the system is not constrained to lie in a probability space, thus allowing interference to take place. The state of the quantum walk at time  $t$  is a complex linear combination of the basis states  $|u\rangle$

$$|\psi(t)\rangle = \sum_{u \in V} \alpha_u(t) |u\rangle, \quad (2)$$

where the amplitude  $\alpha_u(t) \in \mathbb{C}$  and  $|\psi(t)\rangle \in \mathbb{C}^{|V|}$  are both complex. Moreover, we have that  $\alpha_u(t)\alpha_u^*(t)$  gives the probability that at time  $t$  the walker is at the vertex  $u$ , and thus  $\sum_{u \in V} \alpha_u(t)\alpha_u^*(t) = 1$  and  $\alpha_u(t)\alpha_u^*(t) \in [0, 1]$ , for all  $u \in V, t \in \mathbb{R}^+$ .

Unlike the classical random walk, the evolution of the quantum walk is governed by the Schrödinger equation  $\partial/\partial t |\psi_t\rangle = -iH |\psi_t\rangle$ , where  $H$  denotes the system Hamiltonian. This is usually either the adjacency matrix or the graph Laplacian.

### B. The Average Mixing Matrix

Given a graph  $G = (V, E)$ , the behaviour of a continuous-time quantum walk at time  $t$  can be summarized using the mixing matrix [14]

$$\begin{aligned} M(t) &= U(t) \circ U(-t) \\ &= e^{iHt} \circ e^{-iHt}, \end{aligned} \quad (3)$$

where  $\circ$  denotes the Schur-Hadamard product of two matrices, i.e.,  $[A \circ B]_{uv} = A_{uv}B_{uv}$ . Since  $U$  is unitary,  $M(t)$  is a doubly stochastic matrix, i.e., a non-negative real matrix such that its rows and columns sum to 1. Intuitively, the entry  $M(t)_{uv}$  is equal to the probability that at time  $t$  a quantum walker that started from vertex  $u$  is found at vertex  $v$ .

Since  $U(t)$  is unitary and thus norm-preserving,  $M(t)$  does not converge. However, we can enforce convergence by taking a time average. More specifically, we can introduce the average mixing matrix by taking the Cesàro mean

$$\widehat{M} = \lim_{T \rightarrow \infty} \int_0^T M(t) dt. \quad (4)$$

The entry  $\widehat{M}_{v_i v_j}$  can be interpreted as the average probability that a quantum walker that started from the vertex  $v_i$  is found at  $v_j$ . Note that the average mixing matrix is still a doubly stochastic matrix. Godsil [14] also proved that its entries are rational numbers. The average mixing matrix can be easily computed from the spectrum of the Hamiltonian. Let

$H = \sum_{j=1}^M \lambda_j \mathbb{P}_j$ , where the  $\lambda_j$ s are the  $M$  distinct eigenvalues of  $H$  and  $\mathbb{P}_m$  is the matrix representation of the orthogonal projection on the eigenspace associated with  $\lambda_m$ . Then, the average mixing matrix can be written as

$$\widehat{M}(t) = \sum_{j=1}^M \mathbb{P}_j \circ \mathbb{P}_j. \quad (5)$$

### C. Shannon Entropy from the Average Mixing Matrix

In this subsection, we show how to associate a Shannon entropy with each vertex of a graph. Let  $G = (V, E)$  be a sample graph and  $\widehat{M}(t)$  be the average mixing matrix of  $G$ . For the  $i$ -th vertex  $v_i \in V$ , the corresponding row of the average mixing matrix gives time-averaged probability for the quantum walker to reach the vertices  $v_1, \dots, v_{|V|} \in V$  starting from  $v_i$ , i.e.,

$$P_i = \{p_i(v_1), \dots, p_i(v_j), \dots, p_i(v_{|V|})\}. \quad (6)$$

where  $p_i(v_j) = \widehat{M}_{v_i v_j}$  is the time-averaged probability of going from  $v_i$  to  $v_j$ .

Thus, the Shannon entropy associated with the vertex  $v_i$  can be computed as

$$H_S(v_i) = - \sum_{v_j \in V} p_i(v_j) \log p_i(v_j). \quad (7)$$

## III. AN ENTROPIC SIGNATURE FROM THE AVERAGE MIXING MATRIX

In this section, we propose to define the proposed entropic signature. We commence by introducing the concept of Jensen-Shannon divergence. With the Jensen-Shannon divergence to hand, we then proceed to show how to define a novel graph entropic signature.

### A. The Jensen-Shannon Divergence

The Jensen-Shannon divergence is a dissimilarity measure defined between probability distributions [17]. Consider two (discrete) probability distributions  $\mathcal{P} = (p_1, \dots, p_m, \dots, p_M)$  and  $\mathcal{Q} = (q_1, \dots, q_m, \dots, q_M)$ , then the Jensen-Shannon divergence between  $\mathcal{P}$  and  $\mathcal{Q}$  is defined as

$$\begin{aligned} D_{JS}(\mathcal{P}, \mathcal{Q}) &= H_S\left(\frac{\mathcal{P} + \mathcal{Q}}{2}\right) - \frac{H_S(\mathcal{P}) + H_S(\mathcal{Q})}{2} \\ &= - \sum_{m=1}^M \frac{p_m + q_m}{2} \log \frac{p_m + q_m}{2} \\ &\quad + \sum_{m=1}^M p_m \log p_m + \sum_{m=1}^M q_m \log q_m. \end{aligned} \quad (8)$$

where  $H_S(\mathcal{P}) = \sum_{m=1}^M p_m \log p_m$  is the Shannon entropy of the probability distribution  $\mathcal{P}$ . For a mixture of  $N$  probability distributions  $\mathcal{P}_1, \dots, \mathcal{P}_N$  associated with mixing proportions  $\pi_1, \dots, \pi_N$ , the Jensen-Shannon divergence defined in Eq.(8) can be re-written as

$$D_{JS}(\mathcal{P}_1, \dots, \mathcal{P}_N) = H_S\left(\sum_{i=1}^N \pi_i \mathcal{P}_i\right) - \sum_{i=1}^N \pi_i H_S(\mathcal{P}_i). \quad (9)$$

Note that the Jensen-Shannon divergence  $D_{JS}$ , both in Eq.(8) or Eq.(9), is always well defined, symmetric, negative definite and bounded, i.e.,  $0 \leq D_{JS} \leq 1$ .

### B. Graph Entropic Signature

Let  $G(V, E)$  be a sample graph. We commence by computing the average mixing matrix  $\widehat{M}(t)$  for  $G$ . As observed in the previous Section, the  $i$ -th row of  $\widehat{M}(t)$  corresponds to the time-averaged probability distribution  $P_i$  induced by a continuous-time quantum walk starting from  $v_i \in V$ . We define the mixture of  $|V|$  probability distributions corresponding to the  $|V|$  vertices of  $G$  by associating a weight  $\pi_i = 1/|V|$  with each row of the average mixing matrix. The entropic signature of  $G$  is then computed as

$$E_G = [D_{JS}(P_1, \dots, P_{|V|}), \widehat{H}_S(v_1, \dots, v_{|V|})], \quad (10)$$

where  $D_{JS}(P_1, \dots, P_{|V|})$  is the entropic dissimilarity of  $G$

$$D_{JS}(P_1, \dots, P_{|V|}) = H_S\left(\sum_{i=1}^{|V|} \pi_i P_i\right) - \sum_{i=1}^{|V|} \pi_i H_S(P_i), \quad (11)$$

and  $\widehat{H}_S(v_1, \dots, v_{|V|})$  is the average entropic measure of  $G$

$$\widehat{H}_S(v_1, \dots, v_{|V|}) = \sum_{i=1}^{|V|} \pi_i H_S(P_i). \quad (12)$$

Note that, for  $E_G$  the element  $D_{JS}(P_1, \dots, P_{|V|})$  reflects the entropic dissimilarity between the probability distribution associated with the graph vertices. On the other hand, the element  $\widehat{H}_S(v_1, \dots, v_{|V|})$  encapsulates the average mixing Shannon entropy of the vertices. Both the elements of our entropic measure rely on the average mixing matrix of continuous-time quantum walks. As we have stated, quantum walks encapsulate rich interior graph information. As a result, the proposed entropic signature not only encapsulates the rich information from the Shannon entropies of the distributions induced by the quantum walk, but also captures the interior dissimilarity information between vertices.

### C. Computational Complexity

For a graph  $G(V, E)$  having  $n$  vertices, computing the entropic characteristics measure has time complexity  $O(n^3)$ . This is because the required average mixing matrix from continuous-time quantum walks relies on the eigen-decomposition of the graph Laplacian, that requires time complexity  $O(n^3)$ . This indicates that the proposed entropic characteristics measure for a graph can be computed in a polynomial time. Note that this is the same complexity of the von Neumann entropy computation. However, as our experimental results show, the proposed signature is better able to discriminate between different graph structures.

## IV. EXPERIMENTAL RESULTS

### A. Graph Datasets

We demonstrate the performance of our new kernel on six standard graph datasets from bioinformatics databases. These datasets include: MUTAG, NCI1, NCI109, ENZYMES, PPIs and PTC(MR). More details are shown in Table.I.

**MUTAG:** The MUTAG dataset consists of graphs representing 188 chemical compounds, and aims to predict whether each compound possesses mutagenicity.

TABLE I. SUMMARY STATISTICS FOR THE SELECTED GRAPH DATASETS

| Datasets         | MUTAG | NCI1  | NCI109 | ENZYMES | PPIs   | CATH1  | CATH2  |
|------------------|-------|-------|--------|---------|--------|--------|--------|
| Max # vertices   | 28    | 111   | 111    | 126     | 232    | 568    | 568    |
| Min # vertices   | 10    | 3     | 4      | 2       | 3      | 44     | 143    |
| Mean # vertices  | 17.93 | 29.87 | 29.68  | 32.63   | 109.60 | 205.70 | 308.03 |
| Number of graphs | 188   | 4110  | 4127   | 600     | 86     | 712    | 190    |

**NCI1 and NCI109:** The NCI1 and NCI109 datasets consist of graphs representing two balanced subsets of datasets of chemical compounds screened for activity against non-small cell lung cancer and ovarian cancer cell lines respectively. There are 4110 and 4127 graph based structures in NCI1 and NCI109 respectively.

**ENZYMES:** This dataset consists of graphs representing protein tertiary structures consisting of 600 enzymes from the BRENDA enzyme database. In this case the task is to correctly assign each enzyme to one of the 6 EC top-level classes.

**PPIs:** The PPIs dataset consists of protein-protein interaction networks (PPIs). The graphs describe the interaction relationships between histidine kinase in different species of bacteria. Histidine kinase is a key protein in the development of signal transduction. If two proteins have direct (physical) or indirect (functional) association, they are connected by an edge. There are 219 PPIs in this dataset and they are collected from 5 different kinds of bacteria with the following evolution order (from older to more recent) *Aquifex4* and *thermotoga4* PPIs from *Aquifex aelicus* and *Thermotoga maritima*, *Gram-Positive52* PPIs from *Staphylococcus aureus*, *Cyanobacteria73* PPIs from *Anabaena variabilis* and *Proteobacteria40* PPIs from *Acidovorax avenae*. There is an additional class (*Acidobacteria46* PPIs) which is more controversial in terms of the bacterial evolution since they were discovered. We select *Proteobacteria40* PPIs and *Acidobacteria46* PPIs as the second group test graphs.

**CATH1 and CATH2:** The CATH1 dataset consists of proteins in the same class (i.e Mixed Alpha-Beta), but the proteins have different architectures (i.e. Alpha-Beta Barrel vs. 2-layer Sandwich). CATH2 contains proteins in the same class (i.e. Mixed Alpha-Beta), architecture (i.e. Alpha-Beta Barrel), and topology (i.e. TIM Barrel), but in different homology classes (i.e. Aldolase vs. Glycosidases). The CATH2 dataset is harder to classify, since the proteins in the same topology class are structurally similar. There are 712 and 190 test graphs in the CATH1 and CATH2 datasets.

### B. Experiments on Graph Classification

**Experimental Setup:** We compare the proposed entropic signature with several state-of-the-art entropy-based graph complexity measure. These include 1) the von-Neumann graph entropy (VNGE) [7], 2) the Shannon graph entropy associated with the steady state random walk (SGE) [15], 4) the Shannon entropy associated with the information functionals,  $f^{V_2}$  (FV2) and  $f^{P_2}$  (FP2) [2], 5) the von Neumann entropy associated with the normalized Laplacian matrix [4], 6) the backtraceless random walk kernel using the Ihara zeta function based cycles (BRWK) [18], and 7) the graphlet count graph kernels with graphlet of size 3 (GCGK) [19]. For the proposed method and the entropy-based graph complexity methods, we calculate the characterization values of graphs as features. We then evaluate

the performance of the signatures in a graph classification task by performing 10-fold cross-validation using a Support Vector Machine (SVM) with Sequential Minimal Optimization (SMO) [20] and the Pearson VII Universal Kernel (PUK) [21]. For the BRWK and GCGK kernels, evaluate the classification performance by performing 10-fold cross-validation using a C-SVM associated with the kernel matrix. For any method, We use nine folds for training and one fold for testing. We evaluate the performance of our method and the alternative methods. For each method, we repeat the experiments 10 times. We report the average classification accuracies of each method in Tables II. The runtime is measured under a Matlab R2011a running on a 2.5GHz Intel 2-Core processor (i.e., i5-3210m). Table III, on the other hand, shows the runtime evaluation for each of method and dataset.

**Results and Discussions:** In terms of classification accuracy, the proposed entropic signature easily outperforms the alternative entropy-based complexity measures, the BRWK kernel and the GCGK kernel. Only the classification accuracy of the SGE entropy on the MUTAG and CATH1 datasets is a marginally better than that of the proposed method, but the proposed method is still competitive with respect to the SGE entropy measure on the MUTAG and CATH1 datasets. On the other hand, the classification accuracy of the proposed method is significantly greater than that of the SGE entropy measure on the remaining datasets. These observations indicate that the proposed entropic measure is more effective than the state-of-the-art entropy-based complexity measures and the graph kernels. The reason for this effectiveness is that both the two elements of the proposed entropic signature rely on the average mixing matrix of continuous-time quantum walks, which in turn can reflect rich interior graph information. By contrast, either the alternative entropy-based complexity measures or the alternative graph kernels rely on the graph vertex degrees, informational functionals, cycles and graphlet substructures, which encapsulate structurally simple features. As a result, the proposed method can reflect richer graph characteristics than the alternative methods.

In terms of runtime, it is clear that the proposed method is not the most efficient method. However, the proposed method can still finish the computation in a polynomial time on all the datasets. At the same time, the effectiveness of the proposed method is better than any alternative entropy-based complexity measure method. This indicates that the proposed method is applicable on real-world datasets.

### C. Stability Evaluation

In this subsection, we explore the relationship between the graph edit distance and the proposed entropic signature. The edit distance between two graphs  $G_p$  and  $G_q$  is the minimum edit cost taken over all sequences of edit operations that transform  $G_p$  into  $G_q$ . In our experiments, we establish

TABLE II. AVERAGE CLASSIFICATION ACCURACY (IN %  $\pm$  STANDARD ERROR).

| Datasets | MUTAG                  | NCI1                   | NCI109                 | ENZYMES                | PPIs                   | CATH1                  | CATH2                  |
|----------|------------------------|------------------------|------------------------|------------------------|------------------------|------------------------|------------------------|
| EC       | 85.63 $\pm$ .35        | <b>66.05</b> $\pm$ .46 | <b>66.32</b> $\pm$ .47 | <b>27.67</b> $\pm$ .37 | <b>77.91</b> $\pm$ .42 | 98.74 $\pm$ .11        | <b>77.89</b> $\pm$ .41 |
| VNGE     | 85.63 $\pm$ .34        | 62.15 $\pm$ .53        | 62.05 $\pm$ .51        | 25.50 $\pm$ .36        | 72.09 $\pm$ .43        | 98.59 $\pm$ .11        | 75.78 $\pm$ .42        |
| SGE      | <b>87.76</b> $\pm$ .35 | 61.84 $\pm$ .50        | 62.05 $\pm$ .51        | 23.16 $\pm$ .35        | 67.44 $\pm$ .42        | <b>98.87</b> $\pm$ .11 | 76.31 $\pm$ .42        |
| FV2      | 84.57 $\pm$ .32        | 62.04 $\pm$ .49        | 62.15 $\pm$ .49        | 24.17 $\pm$ .37        | 70.93 $\pm$ .46        | 96.91 $\pm$ .11        | 76.31 $\pm$ .40        |
| FP2      | 85.63 $\pm$ .32        | 62.37 $\pm$ .49        | 62.37 $\pm$ .49        | 23.33 $\pm$ .37        | 70.93 $\pm$ .46        | 96.91 $\pm$ .11        | 76.31 $\pm$ .40        |
| VNE      | 85.63 $\pm$ .34        | 63.18 $\pm$ .49        | 62.95 $\pm$ .51        | 25.00 $\pm$ .36        | 63.95 $\pm$ .43        | 98.45 $\pm$ .11        | 76.84 $\pm$ .42        |
| BRWK     | 77.50 $\pm$ .75        | 60.34 $\pm$ .17        | 59.89 $\pm$ .15        | 20.56 $\pm$ .45        | —                      | —                      | —                      |
| GCGK     | 82.04 $\pm$ .39        | 63.72 $\pm$ .12        | 62.33 $\pm$ .13        | 24.87 $\pm$ .22        | 46.61 $\pm$ .47        | 98.49 $\pm$ .12        | 73.68 $\pm$ 1.09       |

TABLE III. RUNTIME COMPARISONS

| Datasets | MUTAG | NCI1  | NCI109 | ENZYMES | PPIs   | CATH1  | CATH2  |
|----------|-------|-------|--------|---------|--------|--------|--------|
| EC       | 1"    | 10"   | 10"    | 2"      | 3"     | 1'46"  | 54"    |
| VNGE     | 1"    | 1"    | 1"     | 1"      | 1"     | 1"     | 1"     |
| SGE      | 1"    | 1"    | 1"     | 1"      | 1"     | 1"     | 1"     |
| FV2      | 1"    | 4"    | 4"     | 1"      | 1"     | 8"     | 3"     |
| FP2      | 1"    | 4"    | 4"     | 1"      | 1"     | 8"     | 3"     |
| VNE      | 1"    | 6"    | 6"     | 1"      | 1"     | 16"    | 8"     |
| BRWK     | 1"    | 6'49" | 6'49"  | 13"     | > 1day | > 1day | > 1day |
| GCGK     | 1"    | 5"    | 5"     | 2"      | 4"     | 2"     | 8"     |

a new graph by deleting a fraction of vertices from a seed graph. The evaluation employs two seed graphs each having 400 vertices. For each of the seed graph, we randomly delete a predetermined fraction of vertices to simulate the effects of noise. We apply the edit operation on the seed graph 35 times, and generate 35 edit operated graphs as the noise corrupted counterparts. Given the seed graph, we delete 4 vertices (i.e., 1% vertices) each time, respectively. Recall that the proposed entropic signature of a graph consists of two elements. One is the Jensen-Shannon divergence between the graph vertices, the other one is the average Shannon entropy of the graph vertices. To evaluate the effect of structural perturbations on the proposed entropic signature, we investigate how these two components vary under the edit operations. Fig.1 and Fig.2 shows the fluctuations of the signature elements for increasing amounts of structural noise. In each plot, the x-axis denotes the fraction of vertices deleted, and the y-axis denotes the value of the value of the entropic signature components.

We observe that there is an approximately linear relationship between both the entropic characteristics elements and the fraction of the deleted vertices. Moreover, we observe that when less than 10% of vertices are deleted the fluctuations are relatively small and the value of the signature is stable. However, when around 20% of vertices are deleted the fluctuations become moderate. These indicate that the proposed entropic signature measure is robust even when the seed graph structure undergoes relatively large perturbations. Furthermore, the proposed method also possesses a good ability to distinguish graphs under controlled amounts of structural errors.

## V. CONCLUSION AND FUTURE WORK

In this paper, we have proposed a novel entropic signature for graphs, based on the average mixing matrix from continuous-time quantum walks. More specifically, we have associated a probability distribution with each vertex of a graph using the average mixing matrix. The resulting entropic signature is computed by concatenating the average Shannon entropy of these probability distributions with their Jensen-Shannon divergence. The new entropic measure encapsulates

the rich structural information of graphs, and thus has a good ability of discriminating different graph structures. We have empirically demonstrated the effectiveness of the proposed entropic measure on several standard graph-based datasets abstracted from bioinformatics databases.

In the future work, we will extend our analysis in a number of ways. First, in our previous work [15], [12] we have developed a family of graph kernels using both the classical and quantum Jensen-Shannon divergence associated with classical and quantum walks. It would be interesting to use the average mixing matrix to develop a new Jensen-Shannon graph kernel. Second, in our previous work [22], we have developed a novel framework for computing depth-based complexity traces of graphs associated with entropy measures. It would be interesting to develop a new complexity trace method for graphs using the average mixing matrix.

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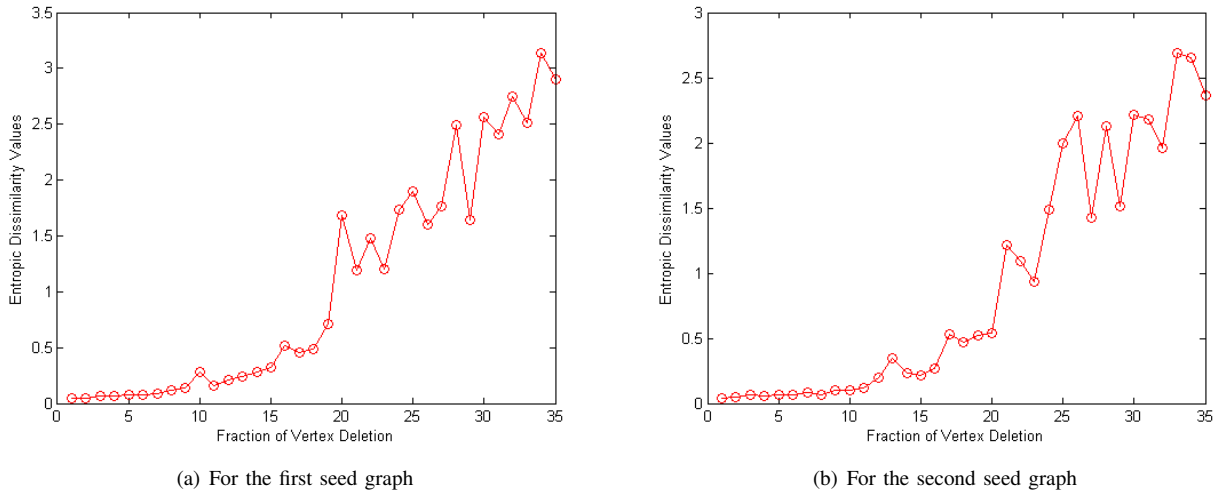


Fig. 1. Stability evaluation of Jensen-Shannon divergence under vertex edit operations

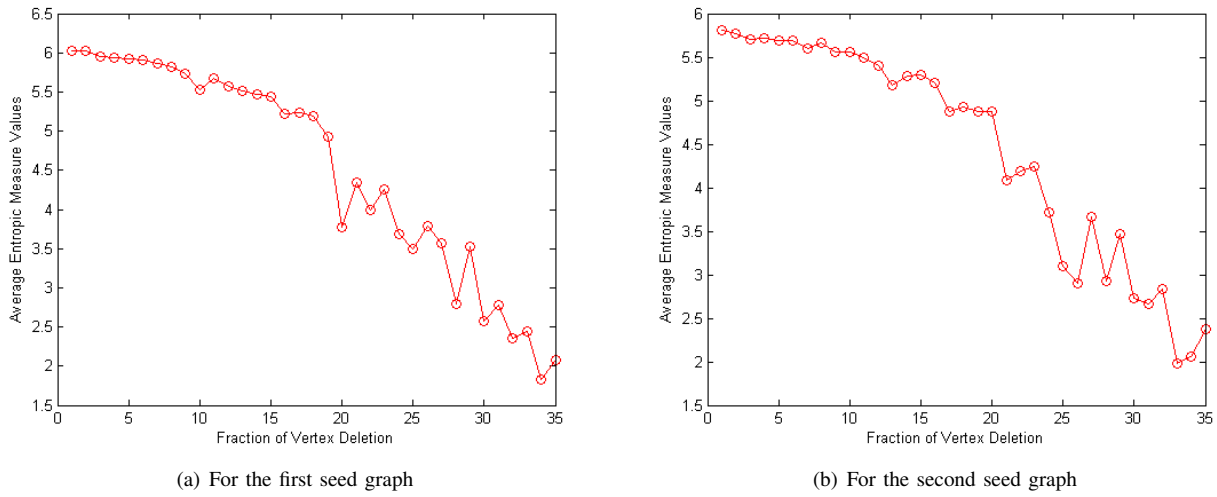


Fig. 2. Stability evaluation of the average Shannon entropy under vertex edit operations

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