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25 **Abstract**

26 There is a considerable number of research publications on the characterization of porous
27 media that is carried out in accordance with the ISO 10534-2 [ISO 10534-2, 2001] and/or
28 ISO 9053 standards [ISO 9053, 1991]. According to WEB OF SCIENCE™ (last accessed on
29 22 of September 2016) there were 339 publications in the Journal of the Acoustical Society
30 of America alone which deal with the acoustics of porous media. However, the
31 reproducibility of these characterization procedures is not well understood. This paper deals
32 with the reproducibility of some standard characterization procedures for acoustic porous
33 materials. The paper is an extension of the work published in [K.V. Horoshenkov *et. al.* , J.
34 Acoust. Soc. Am. 122 (1), 2007]. One novelty of this paper is that independent laboratory
35 measurements were performed on the same material specimens so that the naturally occurring
36 inhomogeneity in materials was controlled. Another novelty of this work is that it presents
37 the reproducibility data for the characteristic impedance, complex wavenumber and for some
38 related pore structure properties. This work can be helpful to understand better the tolerances
39 of these material characterization procedures so the improvements can be developed to
40 reduce the experimental errors and improve the reproducibility between laboratories.

41

42 **I. INTRODUCTION**

43 The characterization of porous media has become a standard procedure which is carried out
44 in several laboratories worldwide to validate new models for the acoustical properties of
45 porous media, to measure the acoustical performance of new types of porous media used in
46 noise control applications, and/or to deduce the parameters of their porous micro-structure. In
47 addition, a number of industries rely heavily on their ability to model the acoustical
48 properties of porous media *in-situ*. For this purpose they need to have accurate data on the
49 acoustic impedance of porous media and propagation constant. With this in mind, it is
50 important to have a clear understanding of the dispersion of acoustical data caused by the
51 differences in the equipment and natural variation in the material formulation. However, this
52 information is scarce and the standard ISO 10534-2 procedure¹ is rather ambiguous in terms
53 of the quality and uniformity of material samples, environmental and operational conditions,
54 the quality of setup and signal processing method. It is fair to say that the reproducibility of
55 the standard acoustical method (ref. [1]) in application to the porous media characterization
56 has not been properly investigated. As a result, the uncertainties of the characterization
57 procedures are largely unknown. There are three basic questions which remain unanswered:
58 (i) 'How accurate our acoustic material data actually are?' (ii) 'Would we get the same result
59 as published by our colleagues if we test these materials in our own lab?' (iii) 'If we develop a
60 new model is it actually more accurate than existing models in terms of any potential
61 measurement errors we can incur?' A while ago the authors of this paper attempted to answer
62 some of these questions through a series of experiments designed to evaluate the
63 reproducibility in normal incidence sound absorption coefficient and surface impedance of
64 porous specimens which were cut independently from flat sheets of porous materials sent to
65 the 6 partners by 3 material manufacturers². These experiments were carried out in 2-
66 microphone impedance tubes in compliance with the ISO 10534-2¹ standard. As a general
67 summary of the results, higher variations in the measured spectra for the surface impedance
68 and acoustic absorption coefficient were observed between individual samples and individual
69 laboratories in the case of low permeability, low homogeneity, broad pore size distribution
70 and reconstituted porous rubber. The smallest variations (<20%) in the data were observed in
71 the case of high permeability porous reticulated foam, although the mounting conditions for
72 this material were difficult to reproduce in independent acoustic laboratories which resulted
73 in a shift of the frame resonance frequency affecting the absorption coefficient in a certain
74 frequency range. Finally, medium level variations in the measured acoustical absorption data
75 (> 20%) were observed in the case of fiberglass. These variations were attributed to change in
76 specimen thickness during the mounting within the measurement tube.

77 At the moment, the authors are not aware of any studies which provide experimental data
78 from independent laboratories for characteristic acoustical properties (i.e. characteristic
79 impedance and complex wavenumber) and for several physical parameters describing their

80 micro/macro structure (airflow resistivity, open porosity, tortuosity and viscous and thermal
81 characteristic lengths) measured for the same material specimens. Among physical
82 parameters, only airflow resistivity can be measured according to a standard (ISO 9053³) and
83 considerable work has been carried out by Garai and Pompoli⁴ who coordinated the European
84 Inter-Laboratory test as per that standard. The results of this work are limited to melamine
85 foam samples and show that most laboratories have good internal repeatability, particularly
86 for single sample measurements. In comparison with repeatability, the overall reproducibility
87 is not so good mainly due to systematic deviations inherent to current laboratory practice. In
88 this respect, there is a lack of reproducibility data which are obtained for the same material
89 specimen tested in independent acoustic laboratories.

90 Therefore, the aim of this paper is to determine the dispersion of surface acoustical data (i.e.
91 surface impedance, z_s , and absorption coefficient, α), characteristic properties (i.e.
92 characteristic impedance, z_c , and complex wavenumber, k_c), and related pore structure
93 parameters (airflow resistivity, σ , open porosity, ϕ , tortuosity, α_∞ , and viscous, A , and
94 thermal, A' , characteristic lengths) obtained for the same material sample, but tested in
95 different acoustic laboratories. The meaning of these parameters is detailed in ref. [5].

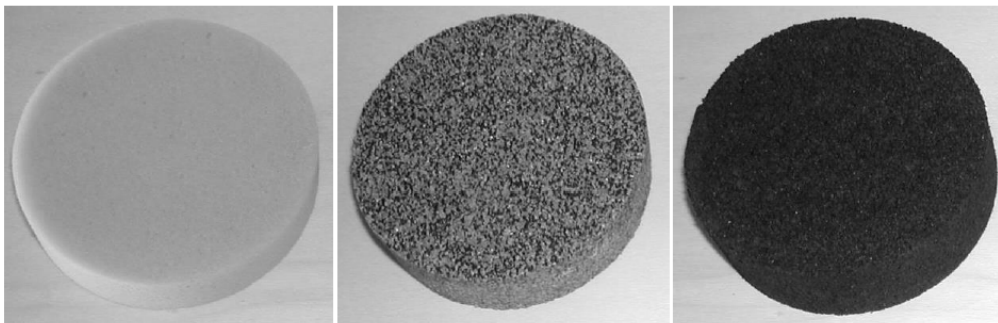
96 This paper is organised as follows: section II outlines the methodology; section III presents
97 the results of from individual laboratories and inter-laboratory data. Concluding remarks are
98 made in the last section.

99

100 II. METHODOLOGY

101 In this work, seven acoustic research centers were involved. These are: the University of
102 Ferrara (Italy), the University of Perugia (Italy), Katholieke Universiteit Leuven (Belgium),
103 Matelys/ENTPE in Lyon (France), Gesellschaft für Akustikforschung Dresden (Germany),
104 the University of Bradford (UK), and Sherbrooke University (Canada). Three different
105 porous materials were investigated: reticulated foam, consolidated flint and reconstituted
106 porous rubber, denoted materials A, B, and C, respectively (Fig.1).

107



108

109 Figure 1 – Tested materials. Tested materials. A: reticulated foam (left), B: consolidated flint (center),
110 C: reconstituted porous rubber (right).

111

112 In this research, the same set of specimens for porous materials with different diameters (99
 113 mm, 44mm and 29 mm) was provided and shared amongst laboratories. Materials A and C
 114 were identical to those used in ref. [2], that are reticulated foam and reconstituted rubber,
 115 respectively. Material B was consolidated flint particles to minimize the effect of mounting
 116 thickness variations within the impedance tubes. In this way samples of each material were
 117 not exactly identical among all the partner laboratories because they were cut for a range of
 118 impedance tube diameters. Table I presents a basic description of the materials which were
 119 used in the inter-laboratory experiment. Table II lists the acoustical and pore structure
 120 parameters and partner laboratories in which these parameters were measured.

121

122 Table I. The porous materials used in the inter-laboratory experiment.

Material	Description	Thickness [mm]	Density [kg/m ³]	Diameters [mm]	Number of samples for each diameter
A	Reticulated foam	20±0.1	8,8	29/44/99	4
B	Consolidated flint	31±0.1	1500	29/44/99	6
C	Reconstituted porous rubber	28±0.1	242	29/44/99	6

123

124 Table II. The list of the acoustical and related pore structure parameters and partner laboratories in
 125 which these parameters were measured.

Partner	z_s, α	z_c, k_c	σ	ϕ	α_∞	A	A'
1	•	•	•	•	•	•	•
2	•		•	•	•	•	•
3	•	•	•	•			
4	•	•	•	•	•	•	•
5	•		•				
6	•		•	•	•	•	•
7	•	•					

126

127 **A. Measurement of acoustical properties**

128 The acoustical properties measured directly in accordance with the ISO 10534-2¹ were the
 129 normalized surface acoustic impedance z_s [-] (for plane waves at normal incidence) and the
 130 normal incidence sound absorption coefficient α [-] of the material sample backed by a rigid
 131 wall. The size and diameter of the standing wave tube, the manufacturers and the excitation
 132 stimulus used by the partners are detailed in Table III. The following methods of sample
 133 mounting conditions were adopted (see Table III): (i) the diameter of the cut samples was
 134 close to or slightly smaller than the diameter of the tube and the samples were wrapped in
 135 tape to prevent any leakage around the edge - tape constraint (TC); (ii) the diameter of the
 136 sample was exactly equal to that of the tube - perfect fit (PF).

137 All the partners applied the amplitude and phase mismatch calibration procedures before tests
 138 (with the exception of Partner 4 who used a single microphone) in accordance with ISO
 139 10534-2¹. All the microphones used in these experiments were standard 1/4 inch microphones

140 (see Table III). Partner 1 – 5 carried out tests in the frequency range consistent with that
 141 suggested in ref. [1] for a given tube diameter and microphone spacing. Partner 6 provided
 142 data in the frequency range between 200 Hz and 1600 Hz because of a low signal-to-noise
 143 ratio. It should be noted that the ISO 10534-2:2001 standard does not define the exact
 144 frequency range for a given tube diameter and microphone separation, but recommends the
 145 bounds for the lower and upper frequencies in the range (see Section 4.2 in ref. [1]).
 146 Therefore, the partners chose the frequency ranges to satisfy the standard requirements for the
 147 level of nonlinearities, frequency resolution, measurement instabilities and signal-to-noise
 148 ratio recommended in ref. [1].

149

150 Table III. The equipment and mounting conditions used to determine the acoustic absorption
 151 coefficient and surface impedance (HM: homemade equipment; TC: tape constraint; PF: perfect fit).

Partner	Tube diameter / tube manufacturer	Tube length [m] / microphone spacing [m]	Mounting conditions	Stimulus	Electronic hardware	Microphone type	Frequency range [Hz]
1	45 mm / HM	0.5 / 0.03; 0.1	TC	Sweep	NI USB 4431	PCB 377C10	100-4200
2	29 mm / HM	0.4225/0.02	PF	White noise	SR-8 Channel Analyzer (DSP Board)	BK2670	400-6900
3	29 mm / B&K 4206	0.4225/0.02	PF	White noise	Bruel and Kjaer Pulse type - 2827	BK2670	260-6400
4	29 mm / HM	0.35/0.02	PF	Pseudo random noise	NI PXI 4461	BK4187	400-6900
5	29 mm / B&K 4206	0.4225/0.02	TC	White noise	NI USB-9233	MT Gefell M 365	200-6400
6	38 mm / HM	1/0.02;0.03;0.05	PF	Sweep	GPIB-USB	GRAS40BP	200-1600
7	29 mm / B&K 4206	0.4225/0.02	TC	White noise	Brüel & Kjær PULSE Type 3560-B-030	BK4187	400-6400

152

153 Each impedance tube was driven by a single loudspeaker which was adapted to the size and
 154 the frequency range of the impedance tube and it was assumed tube vibration effect could be
 155 ignored. Regarding nonlinearity in speaker response the impedance tubes used in these
 156 experiments were designed in accordance with the ISO 10534-2:2001 [1], in which Section
 157 4.8 suggests that “*The errors in the estimated transfer function H_{12} due to nonlinearities,*
 158 *resolution, instability and temperature sensitivity of the signal processing equipment shall be*

159 *less than 0,2 dB.*” This is a very small effect and authors believe that it was insignificant in
 160 experiments given a relatively high natural inhomogeneity in the material specimens and
 161 effects of specimen mounting in the tube. The sampling frequency and the sequence length
 162 used in the Fourier analysis were chosen to cover the desired frequency range and to provide
 163 adequate frequency resolution in the transfer function spectrum as suggested in ref. [1]. The
 164 effects of temperature and variations in atmospheric pressure were compensated for as
 165 suggested in ref. [1]. The material thickness was measured to ± 0.1 mm using calibrated
 166 calipers.

167 In addition, the normalized characteristic impedance, z_c , and the complex wavenumber, k_c ,
 168 were measured using a well-established 4 microphone and transfer matrix technique as
 169 described by Song and Bolton⁶. Partner 4 used a 3 microphones technique as described in ref.
 170 [7]. The details of the equipment and measurement techniques are summarized in Table IV.
 171 The equipment used in 3- or 4-microphone tests was properly calibrated prior to the start of
 172 the experiments to compensate for microphone channel mismatch using the procedure similar
 173 to that suggested in ref. [1]. All the microphones used in these experiments were standard 1/4
 174 inch measurement microphones (See Table IV). For the frequency range for these
 175 experiments was chose to meet the recommendations for the impedance tube setup as
 176 suggested in the ISO 10534-2 [1].

177

178 Table IV. The equipment, measurement technique and sample mounting conditions used to determine
 179 the characteristic impedance and complex wavenumber (TC: tape constraint; PF: perfect fit).

Partner	Tube diameter / tube manufacturer	Measurement technique	Mounting conditions	Stimulus	Electronic hardware	Microphone type	Frequency range [Hz]
1	45 mm / HM	4 microphones techniques – refs [6-8]	TC	Sweep	NI USB 4431	PCB 377C10	100-4200
3	44 mm / HM	4 microphones techniques – refs [6]	PF	Pulse	ND	BK2670	188-3500
4	29 mm / HM	3 microphones technique – ref [7]	PF	Pseudo random noise	NI PXI 4461	BK4187	400-6800
7	29 mm / B&K 4206	4 microphones techniques – refs [6]	TC	White noise	Brüel & Kjør PULSE Type 3560-B-030	BK4187	400-6400

180

181

182

183 **B. Measurement of pore structure properties**

184 The airflow resistivity, σ , was measured by the participants using the procedure described in
 185 ISO 9053³. This standard indicates that the value of airflow resistivity has to be determined
 186 for the airflow velocity of less than 0.5 mm/s. When this is not possible the standard suggests
 187 repeating tests at different values of airflow velocity and extrapolating the value of the
 188 airflow resistivity at the nominal value of 0.5 mm/s. Table V describes the equipment, the
 189 measurement techniques and the procedures used by the partners to measure the flow
 190 resistivity.

191

192 Table V. The equipment and measurement technique used to determine the airflow resistivity.

Partner	Tube diameter / tube manufacturer	Measurement technique	Pressure transducer / Pressure range	Extrapolation of d at 0.5 mm/s
1	100 mm / HM	ISO 9053-Method B	BK4186	Linear best-fit between pressure difference and velocity passing through zero
2	99 / 44 mm / HM	ISO 9053-Method A	MKS Type 698A (0.1-1000 Torr)	No extrapolation
3	100 mm / HM	ISO 9053-Method A	FCO 34 (0-10 Pa)	Linear best-fit between pressure difference and velocity passing through zero
4	29 mm / HM	ISO 9053-Method A	MKS 120AD Baratron 1 torr (0-1 Torr)	Direct measurement at 0.5 mm/s
5	99 mm / HM	ISO 9053-Method A	SET-D267MR-6 (-100-100 Pa)	Linear best-fit between pressure difference and velocity passing through zero
6	38 mm /HM	ISO 9053-Method A	Not declared	No extrapolation

193

194 Five partners measured the open porosity, ϕ , using the equipment and measurement
 195 techniques as described in Table VI. Partners 1-4 used the isothermal compression of volume
 196 (Boyle's law) experiment⁹ to measure the porosity. Partner 7 used an acoustic method based
 197 on the analysis of the wave reflected from the sample at oblique incidence¹⁰.

198 Table VII gives an overview of the measurement techniques for the measurement of high
 199 frequency limit of tortuosity α_∞ and characteristic lengths (Λ and Λ'). A majority of partners
 200 obtained the tortuosity and characteristic lengths from the curve fitting of acoustical data and
 201 theoretical modelling as described in refs. [13-15]. Partners 1 and 6 performed measurements
 202 of tortuosity by means of ultrasonic tests¹¹⁻¹². Partners 1 and 2 used samples of different
 203 diameters to measure the flow resistivity and acoustical properties. This means that two
 204 different sets of material specimens were used by this partner in the reported experiments.

205 Table VI. The equipment and measurement technique used to determine the open porosity (HM:
206 homemade equipment).

Partner	Tube diameter / tube manufacturer	Measurement technique
1	99 mm / HM	Isothermal compression of volume ⁹
2	99 mm / HM	Isothermal compression of volume ⁹
3	29 mm / HM	Isothermal compression of volume ⁹
4	29 mm / HM	Isothermal compression of volume ⁹
6	38 mm /HM	Ultrasonic reflection method ¹⁰

207

208 Table VII. The equipment, measurement techniques used to determine the tortuosity and characteristic
209 lengths.

Partner	Device	Measurement technique
1	99 / 45 mm / HM	Ultrasonic test ¹¹⁻¹² and fitting from acoustical data ¹³
2	44 mm kundt tube / HM	Fitting from acoustical data ¹⁴
4	29 mm / HM	Fitting from acoustical data ¹⁵
6	38 mm /HM	Ultrasonic test ¹¹ / fitting from acoustical data

210

211 C. Error analysis

212 Each laboratory carried out two different sets of measurements: (i) tests on different samples
213 of each material (with the exception of Partner 6), (ii) tests on the same sample for each
214 material (with the exception of Partner 4 and 6). The relative errors for a quantity (here
215 generically named as x) measured from these tests were defined as the ratio between its
216 standard deviation and mean value (and expressed in percentage):

$$217 \quad \varepsilon_x = \frac{\sigma_x}{\langle x \rangle} \times 100, \quad [\%] \quad (1)$$

218 $\langle x \rangle$ and σ_x being the mean value and the standard deviation, respectively.

219 The statistical procedures for the analysis of the sound absorption coefficient, airflow
220 resistivity and open porosity described in the ISO 5725-1 and 5725-2 standards^{15, 16} were
221 applied.

222 According to the ISO 5725-2, the repeatability standard deviation is a measurement of the
223 dispersion of the distribution of independent test results obtained with the same method on
224 identical test items in the same laboratory by the same operator using the same equipment
225 within short intervals of time. The reproducibility standard deviation is a measurement of the
226 dispersion of the distribution of test results obtained with the same method on identical test
227 items in different and independent laboratories with different operators using different
228 equipment. According to these standards it is also possible to define for each of the tested
229 materials the repeatability standard deviation in an acoustical parameter for a single sample

230 measured in laboratory i :

$$231 \quad \bar{\sigma}_{1,i} = \frac{\sum_{j=1}^{n_f} \sigma_{1,ij}}{n_f} \quad (2)$$

232 where $\sigma_{1,ij}$ is the standard deviation for laboratory i at frequency j for the measured values of
 233 the acoustical parameter for the same one sample and n_f is the number of discrete frequencies
 234 at which this parameter was measured. Such deviation depends mainly on the random error in
 235 the measurement chain, environmental factors, post-processing of data and mounting
 236 conditions for the sample in the tube. The repeatability standard deviation for all the different
 237 samples in laboratory i can be defined as

$$238 \quad \bar{\sigma}_{A,i} = \frac{\sum_{j=1}^{n_f} \sigma_{A,ij}}{n_f} \quad (3)$$

239 where $\sigma_{A,ij}$ is the standard deviation for laboratory i at frequency j for the measured values of
 240 the acoustical parameter between the all different samples. Such deviation depends on
 241 random errors, sample mounting conditions, homogeneity and sample preparation techniques.
 242 The above quantities can be used to calculate the mean material standard deviation as:

$$243 \quad \langle \sigma_M \rangle = \frac{\sum_{i=1}^{n_L} \bar{\sigma}_{M,i}}{n_L} \quad (4)$$

244 where:

$$245 \quad \bar{\sigma}_{M,i} = \sqrt{\bar{\sigma}_{A,i}^2 - \bar{\sigma}_{1,i}^2} \quad (5)$$

246 is the material standard deviation for laboratory i and n_L is the number of independent
 247 laboratories. In the above equation we assume that the total error is a combination of the
 248 natural variation in the material properties and that which results from the measurement itself.
 249 Therefore, the material standard deviation is a measure of the dispersion in the data due to
 250 natural variation in the material properties from sample to sample so that the mean material
 251 standard deviation is related mainly to homogeneity and sample preparation technique
 252 adopted in this work.

253 The inter-laboratory standard deviation for a single sample is calculated as :

$$254 \quad \langle \sigma_{I1} \rangle = \frac{1}{n_f} \sum_{j=1}^{n_f} \sqrt{\frac{\sum_{i=1}^{n_L} (m_{I1,ij} - \langle m_{I1,j} \rangle)^2}{n_L - 1}} \quad (6)$$

255 where $m_{I1,ij}$ is the mean value of the acoustic parameter measured for the same sample in the
 256 laboratory i at frequency j . Here:

$$257 \quad \langle m_{I1,j} \rangle = \frac{\sum_{i=1}^{n_L} m_{I1,ij}}{n_L} \quad (7)$$

258 is the average of the mean values among different laboratories at frequency j .

259 The inter-laboratory standard deviation for tests on all the material samples can be calculated
 260 in a similar manner as:

$$261 \quad \langle \sigma_{IA} \rangle = \frac{1}{n_f} \sum_{j=1}^{n_f} \sqrt{\frac{\sum_{i=1}^{n_L} (m_{IA,ij} - \langle m_{IA,j} \rangle)^2}{n_L - 1}} \quad (8)$$

262 where $m_{IA,ij}$ is the mean value for laboratory i and frequency j obtained for different samples.
 263 Here:

$$264 \quad \langle m_{IA,j} \rangle = \frac{\sum_{i=1}^{n_L} m_{IA,ij}}{n_L} \quad (9)$$

265 is the average of the mean values among different laboratories measured at frequency j .
 266 In this way the reproducibility standard deviations for a single sample and for all the samples
 267 can be calculated as:

$$268 \quad \sigma_{R1} = \sqrt{\langle \sigma_1 \rangle^2 + \langle \sigma_{I1} \rangle^2} \quad \text{and} \quad \sigma_{RA} = \sqrt{\langle \sigma_A \rangle^2 + \langle \sigma_{IA} \rangle^2}, \quad (10)$$

269 respectively. Here:

$$270 \quad \langle \sigma_1 \rangle = \frac{\sum_{i=1}^{n_L} \bar{\sigma}_{1,i}}{n_L} \quad \text{and} \quad \langle \sigma_A \rangle = \frac{\sum_{i=1}^{n_L} \bar{\sigma}_{A,i}}{n_L} \quad (11)$$

271 are the mean repeatability standard deviation for a single sample and for all the different
 272 samples, respectively. A similar statistical analysis was applied to other material parameters
 273 which were measured non-acoustically. In this case the value of n_f in the above equations
 274 was set to 1.

275

276 III. RESULTS

277 A. Surface impedance and sound absorption coefficient

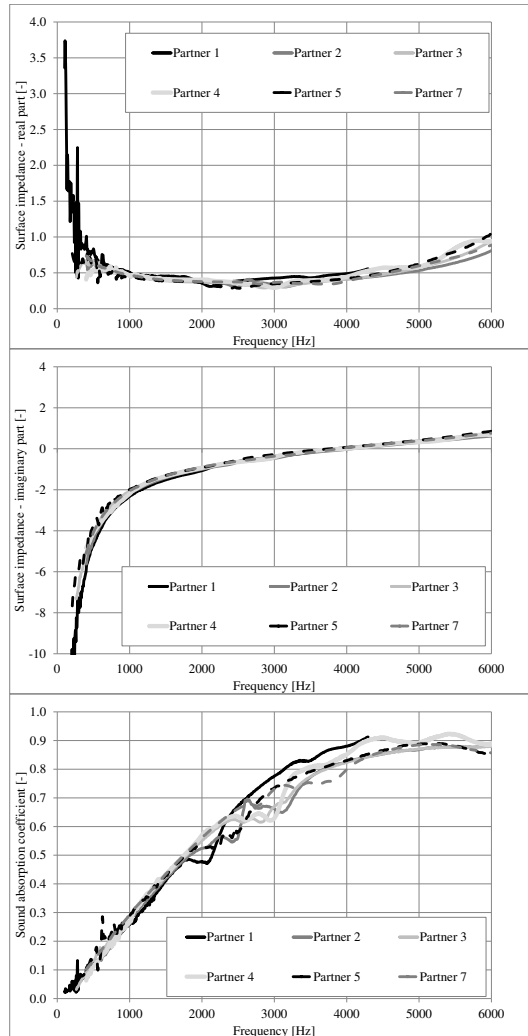
278 The error analysis was based only on the 400 - 3500 Hz range to make data from all the 6
 279 partners compatible. The following figures show the raw data in the frequency range which
 280 was actually utilized by each individual partner. The results of the inter-laboratory tests show
 281 that the relative errors (calculated using Eq. 1) in the real ($\varepsilon_{\Re(z_s)}$) and imaginary ($\varepsilon_{\Im(z_s)}$)
 282 parts of the surface impedance and that of the absorption coefficient ε_α , calculated in the
 283 frequency range between 400 Hz and 3500 Hz, were 13%, 13% and 4%, respectively. For
 284 material B these were 24%, 10% and 19%, respectively. For material C these were 29%, 9%
 285 and 7%, respectively. In the case when the same samples were measured by each laboratory,
 286 deviations were generally found lower: 11% / 9% / 7% for material A, 8% / 7% / 3% for
 287 material B and 8% / 21% / 1% for material C. Such results indicate a gain in the accuracy
 288 with respect to the previous inter-laboratory tests mainly because the same set of materials
 289 was used minimizing the effect of the variability in the pore microstructure between different

290 material slabs.

291 Figures 2 to 4 show the comparison of the measured data for the real and imaginary parts of
 292 the surface impedance and sound absorption coefficient for all the materials tested in
 293 laboratories 1 - 5 and 7. Each curve is the average of all the tests on all the different samples
 294 of the same material. The results obtained by laboratory 6 have been omitted from these
 295 figures since measurements were carried out on a single specimen for each material since
 296 accidentally destroyed some samples trying to adapt them to fit the tube.

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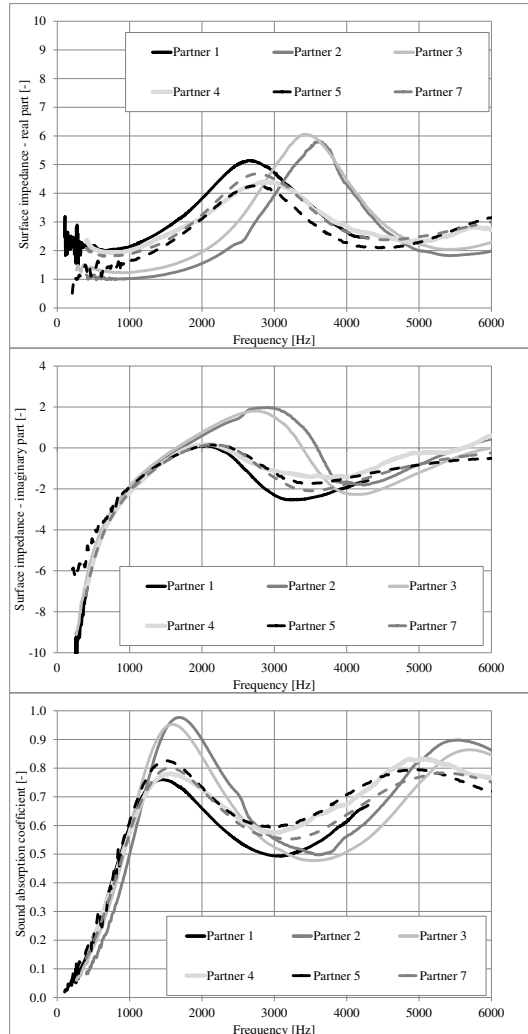
302 Figure 2 – The average of the real part of surface impedance spectra (top), imaginary part of surface
 303 impedance spectra (middle), and the sound absorption coefficient spectra (bottom) measured by the
 304 participating partners for material A.

305

306 The surface impedance and absorption coefficient spectra for material A are shown in Figs.
 307 2(a)–2(c). There is better than 20% agreement in terms of relative errors between the results
 308 for the impedance obtained in the six laboratories. The maximum relative error in the real and
 309 imaginary part of the impedance spectrum of $\pm 25\%$ is observed below 3000 Hz (see Figs.
 310 2(a) and 2(b)). A noticeable increase in the dispersion in the absorption coefficient data can

311 be observed around the frequency of the frame resonance above 2000 Hz (see Fig. 2(c)). This
 312 resonance is often observed in data for low density, soft porous media¹⁸. The dispersion in the
 313 absorption coefficient due to the frame resonance can amount to values between 20% and
 314 30%.
 315

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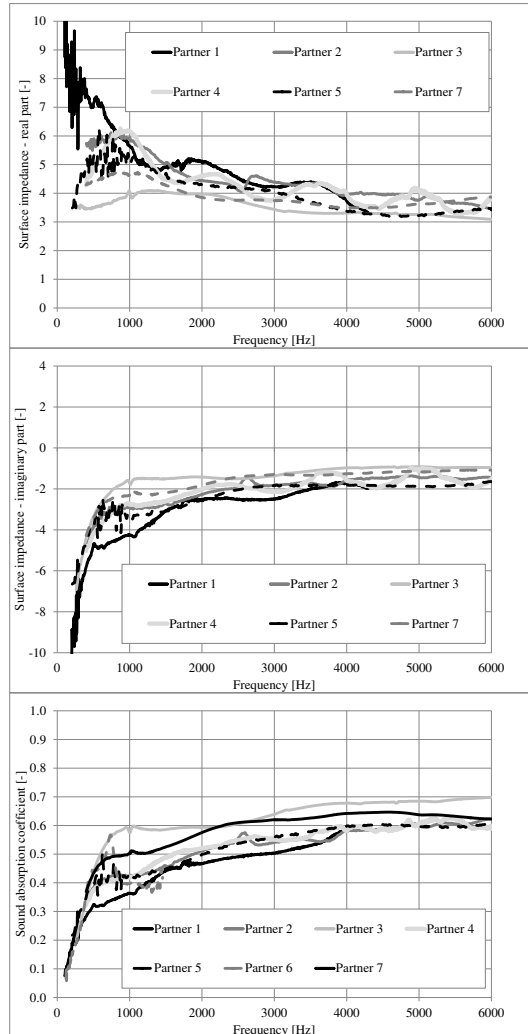
320 Figure 3 - The average of the real part of surface impedance spectra (top), imaginary part of surface
 321 impedance spectra (middle), and the sound absorption coefficient spectra (bottom) measured by each
 322 of the participating partners for material B.

323

324 In the case of material B the dispersion for all the acoustic quantities is high. The results from
 325 partners 2 and 3 are close. These partners used 29 mm diameter impedance tubes, the same
 326 type of microphones and similar excitation stimulus. Partners 5 and 7 also used the same
 327 diameter tube and similar type of acoustic stimulus. However, their results are noticeably
 328 different from those obtained in laboratories 2 and 3. The results from laboratories 1, 4, 5 and
 329 7 follow a similar trend despite some differences in the tube diameter, excitation stimulus and
 330 microphone types. The dispersion in the absorption coefficient for frequencies above 1000 Hz
 331 is between 20% and 40% (Fig. 3(c)). Given a relatively high rigidity of material B, such

332 differences are likely to be attributed to the differences in the mounting condition. Partners 1,
 333 5 and 7 wrapped the edges of their samples in tape to prevent any leakage around the edge.
 334 The other partners reported a very good fit which did not require the sample to be wrapped in
 335 tape.
 336

337



338

339

340 Figure 4 - The average of the real part of surface impedance spectra (top), imaginary part of surface
 341 impedance spectra (middle), and the sound absorption coefficient spectra (bottom) measured by each
 342 the participating partners for material C.

343

344 The results obtained for material C show that there can be a maximum of four to five fold
 345 dispersion in the value of the real part of the surface impedance in the low frequency limit
 346 below 1000 Hz (Fig. 4(a)). The agreement between the data for the imaginary part is poor
 347 across the whole frequency range (Fig. 4(b)). This dispersion is reflected in the erratic
 348 behavior of the absorption coefficient which spectra are shown in Fig. 4(c). The obtained data
 349 suggest that the absorption coefficient for this material can vary within a 10-20% range.
 350 These differences can be attributed to the variability in the mounting conditions. Partner 1
 351 wrapped the edge of their samples in tape and this could have resulted in some degree of pore
 352 deformation and increased airflow resistivity which generally leads to an underestimation of

353 the sound absorption coefficient spectrum.

354 A summary of the statistical error analysis carried out according to ISO 5725-2 can be found
 355 in Table VIII which presents the values of standard deviations for the absorption coefficient
 356 determined from this inter-laboratory experiment. These results enable us to draw the
 357 following conclusions:

- 358 • The mean repeatability standard deviation for a single sample $\langle\sigma_1\rangle$ is relatively low
 359 for all the tested materials. This can suggest that random errors and mounting
 360 conditions are not dominant (below 0.01).
- 361 • The mean repeatability standard deviation for different samples $\langle\sigma_A\rangle$ is significantly
 362 (2.8 – 7 times) higher in comparison with that for a single sample test. The lowest
 363 value is for material A and it is likely to relate to the structural resonance of the
 364 material mounted in the tube. The value of $\langle\sigma_A\rangle$ for material B is the highest,
 365 probably due to the inhomogeneity of the material itself. Material C is characterized
 366 by an intermediate value of $\langle\sigma_A\rangle$ which may relate mainly to the homogeneity of the
 367 material and variation in the mounting conditions. This material has a significantly
 368 high airflow resistivity, it is flexible and any lateral compression applied to its edge
 369 when inserted in the tube can increase the flow resistivity noticeably.
- 370 • The effect of material standard deviation, $\langle\sigma_M\rangle$, is dominant when compared with
 371 the effects due to random errors and mounting conditions for a single sample. The
 372 material standard deviation is related to the natural inhomogeneity of the material and
 373 sample preparation technique. The latter effect is on the sample mounted in the tube,
 374 that may cause a change in the sample elastic behavior (e.g. in the case of material
 375 A), a leakage between the material edge and tube walls (e.g. in the case of material
 376 B) or excessive compression of the sample effectively altering its acoustical
 377 properties (e.g. in the case of material C).
- 378 • The inter-laboratory standard deviation for a single sample $\langle\sigma_{I1}\rangle$ is approximately 2
 379 times higher than $\langle\sigma_M\rangle$, because it is calculated from the average values of $m_{IA,ij}$ for
 380 each laboratory, it is affected by the systematic errors and differences in the
 381 equipment used for the impedance tube test.
- 382 • The inter-laboratory standard deviations for a single σ_{R1} and for different samples
 383 σ_{RA} are comparable that suggests the dominant influence of different impedance
 384 tubes rather than of some systematic errors.
- 385 • The reproducibility standard deviation for single $\langle\sigma_{R1}\rangle$ and different $\langle\sigma_{RA}\rangle$ samples
 386 is lower than 0.07 for all tested materials.

387

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390 Table VIII. The standard deviations for the sound absorption coefficient determined in accordance with
 391 the ISO 5725-2¹⁷.

Standard deviation	Sample A	Sample B	Sample C
$\langle \sigma_i \rangle$	0.005	0.007	0.004
$\langle \sigma_A \rangle$	0.014	0.039	0.028
$\langle \sigma_M \rangle$	0.012	0.038	0.027
$\langle \sigma_{I1} \rangle$	0.03	0.054	0.044
$\langle \sigma_{IA} \rangle$	0.025	0.056	0.056
σ_{R1}	0.031	0.055	0.044
σ_{RA}	0.029	0.068	0.062

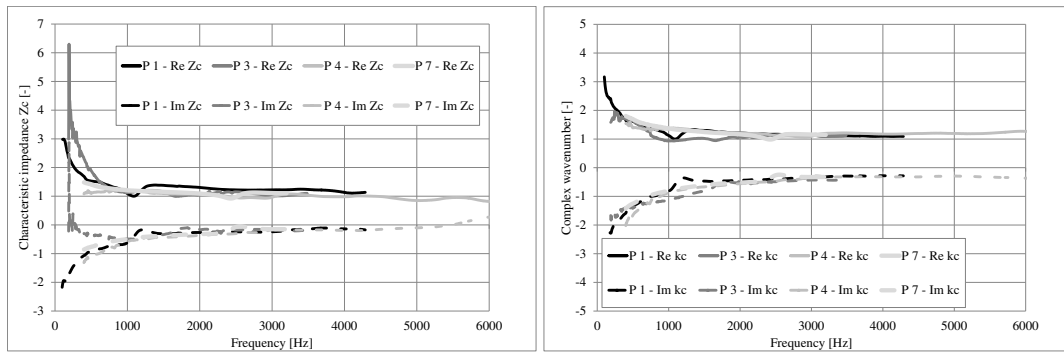
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393 B. Characteristic impedance and wavenumber

394 Partners 1, 3, 4 and 7 also measured the characteristic impedance and complex wavenumber
 395 of the same sample (with the exception of Partner 4) and of different samples of each
 396 material.

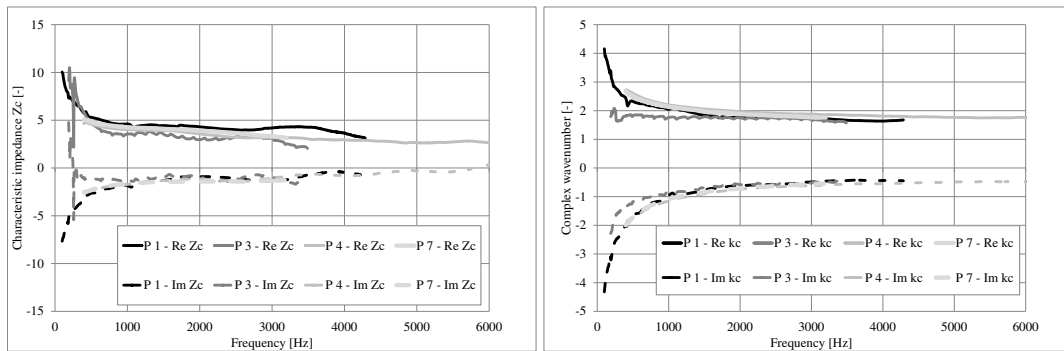
397 Figs. 5 to 7 show the comparison of the real and imaginary parts of the normalized
 398 characteristic impedance and complex wavenumber (normalized by the wavenumber for air
 399 k_0) for all the 3 tested materials. Each curve is the average of the tests on the different
 400 samples. From the data, a consistency in the results between the participating partners is
 401 observed although must be an error in the 4-microphone transfer matrix approach⁶ used by
 402 Partner 3 to invert the characteristic impedance. This approach is not regulated by a standard
 403 and it is prone to errors due to the imperfections in the quality of the anechoic termination,
 404 edge effect and microphone phase mismatch. The relative errors ($\mathcal{E}_{\Re(z_c)}$, $\mathcal{E}_{\Im(z_c)}$, $\mathcal{E}_{\Re(k_c)}$, $\mathcal{E}_{\Im(k_c)}$)
 405 calculated using Eq. 1) in the frequency range of 400-3500 Hz was found between 15% and
 406 30% for the characteristic impedance and between 10 and 30% for the complex wavenumber.
 407 The deviation in the acoustical property for material A is mainly due to the frame resonance
 408 (Figs. 5(a) and 5(b)). The leakage effect between the material edge and tube wall can be the
 409 reason for the deviation observed in the case of material B (Figs. 6(a) and 6(b)). Material C is
 410 characterized by a higher deviation in the characteristic impedance and complex wavenumber
 411 across the whole frequency range which can be attributed to the variability in the mounting
 412 conditions in the impedance tube (Figs. 7(a) and 7(b)).

413 In particular, the tests on a single sample demonstrate that the maximum relative error for all
 414 tested materials was found to be lower than 4% for real part of the characteristic impedance,
 415 14% for imaginary part of the characteristic impedance, 2% for real part of the complex
 416 wavenumber and 4% for the imaginary part of the complex wavenumber. When different
 417 samples of each material were tested, the relative error in data was found to be lower than
 418 30%.



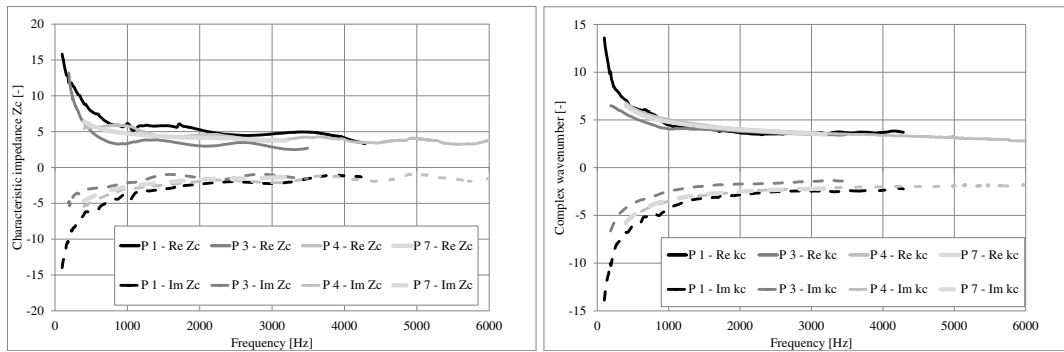
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Figure 5 - The average of the real and imaginary part of the normalized characteristic impedance spectra (left), and real and imaginary part of the normalized complex wavenumber spectra (right) measured by each of the participating partners for material A.



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Figure 6 - The average of the real and imaginary part of the normalized characteristic impedance spectra (left), and real and imaginary part of the normalized complex wavenumber spectra (right) measured by each of the participating partners for material B.



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Figure 7 - The average of the real and imaginary part of normalized characteristic impedance spectra (left), and real and imaginary part of the normalized complex wavenumber spectra (right) measured by each of the participating partners for material C.

C. Pore structure parameters

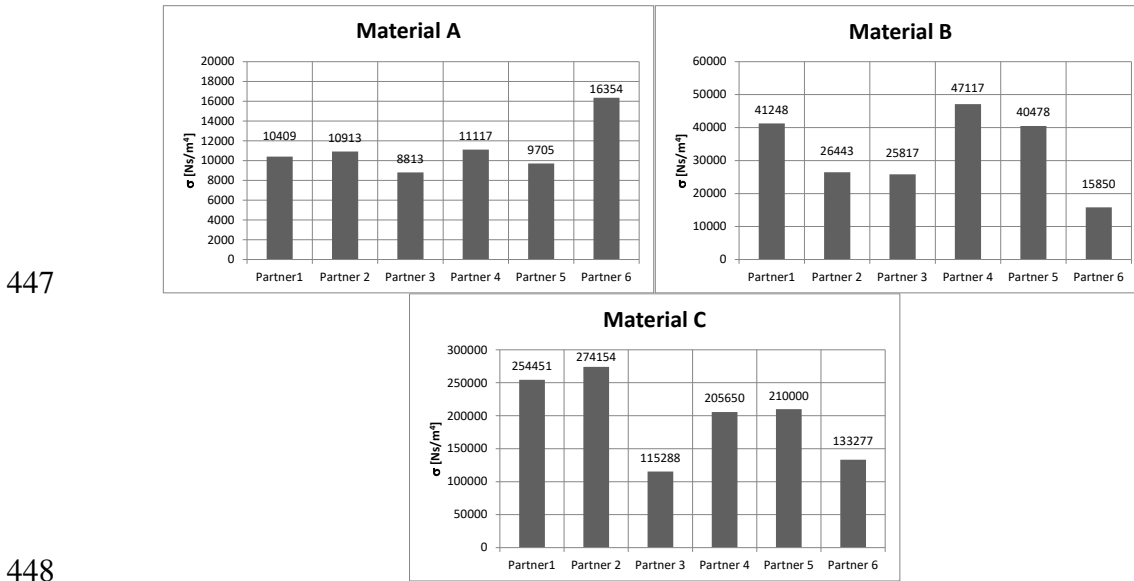
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In addition, the partners carried out tests on the same sample and on different samples for each material to determine the airflow resistivity, porosity, tortuosity and characteristic lengths. Fig. 8 shows the comparison between the average values of airflow resistivity measured for different samples by each of the participating laboratories. Table IX presents the standard deviations determined in accordance with ISO 5725-2 for airflow resistivity and

440 open porosity. Here, the standard deviations calculated according to the ISO standards have
 441 been divided by mean value of the airflow resistivity and open porosity, respectively and data
 442 are expressed in percentage. As an example the mean repeatability standard deviation for a
 443 single sample for airflow resistivity and open porosity can be written as:

$$444 \quad \varepsilon_{1,\sigma} = \frac{\langle \sigma_1 \rangle}{\bar{\sigma}} \times 100, \quad [\%] \quad \text{and} \quad \varepsilon_{1,\phi} = \frac{\langle \sigma_1 \rangle}{\bar{\phi}} \times 100, \quad [\%] \quad (12)$$

445 Similar expressions can be written for other quantities described in Eqs. (2)-(10).
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450 Figure 8 - The average of the airflow resistivity for material A (left), material B (right) and material C
 451 (center) measured by each of the participating partners.

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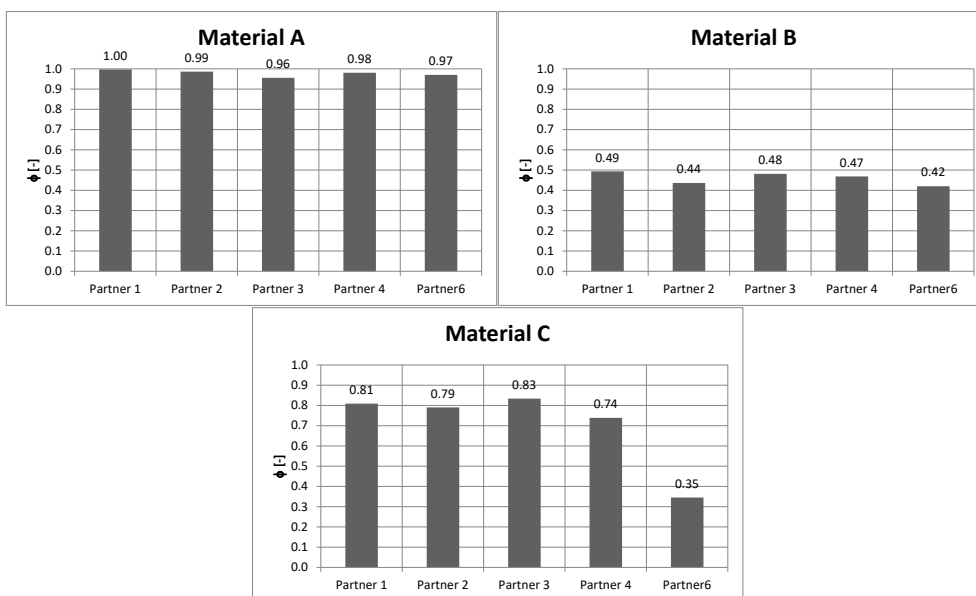
453 The in-laboratory repeatability $\varepsilon_{1,\sigma}$ for the airflow resistivity measured using the same sample
 454 is within 1%. In the case of material A the in-laboratory repeatability for different samples
 455 $\varepsilon_{A,\sigma}$ of material A are lower than 7% while they can vary between 10% and 25% for
 456 materials B and C.

457 A similar analysis is presented for open porosity tests and Fig. 9 shows the comparison
 458 between average values on different samples for each participant. Tests on the same and
 459 different samples once again revealed good internal repeatability ($\varepsilon_{1,\sigma}$ lower than 1% for the
 460 same sample and $\varepsilon_{A,\sigma}$ below 6% for different samples). Also, comparison between different
 461 laboratories is satisfactory for materials A and B (lower than 7%) while measurements on
 462 material C from partner 6 (using a method based on ultrasonic surface reflection) seems to
 463 significantly underestimate the open porosity value.

464 From the data shown in the Table IX, it is possible to come to similar conclusions as for the
 465 sound absorption coefficient. In fact, for both quantities and for all the tested materials, the
 466 mean repeatability standard deviation for a single sample is lower than the mean repeatability

467 standard deviation for several samples; in this case an important role is played by the
 468 homogeneity of materials while random errors seem to be negligible. Such results are
 469 confirmed by a relatively low value of the material standard deviation. The inter-laboratory
 470 standard deviation for a single sample is higher than material standard deviation and this
 471 suggests the occurrence of systematic errors for some of the laboratories. Reproducibility
 472 standard deviations for single and different samples range from between 10% to 45% for
 473 airflow resistivity and 1% to 10% for open porosity.
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 477

478 Figure 9 - The average of the open porosity for material A (left), material B (right) and material C
 479 (center) measured by each of the participating partners.

480

481 Table IX. The repeatability for the airflow resistivity and open porosity determined in accordance with
 482 to the ISO 5725-2¹⁷.

	Airflow resistivity			%	Open porosity		
	A	B	C		A	B	C
$\mathcal{E}_{1,\sigma}$	1	1	1	$\mathcal{E}_{1,\phi}$	0,5	1,1	0,4
$\mathcal{E}_{A,\sigma}$	5	14	22	$\mathcal{E}_{A,\phi}$	1	6	1
$\mathcal{E}_{M,\sigma}$	5	14	22	$\mathcal{E}_{M,\phi}$	0,4	6	1
$\mathcal{E}_{I1,\sigma}$	10	31	29	$\mathcal{E}_{I1,\phi}$	2	10	1
$\mathcal{E}_{IA,\sigma}$	9	25	30	$\mathcal{E}_{IA,\phi}$	2	6	3
$\mathcal{E}_{R1,\sigma}$	15	30	45	$\mathcal{E}_{R1,\phi}$	2	10	1
$\mathcal{E}_{RA,\sigma}$	10	29	37	$\mathcal{E}_{RA,\phi}$	2	9	3

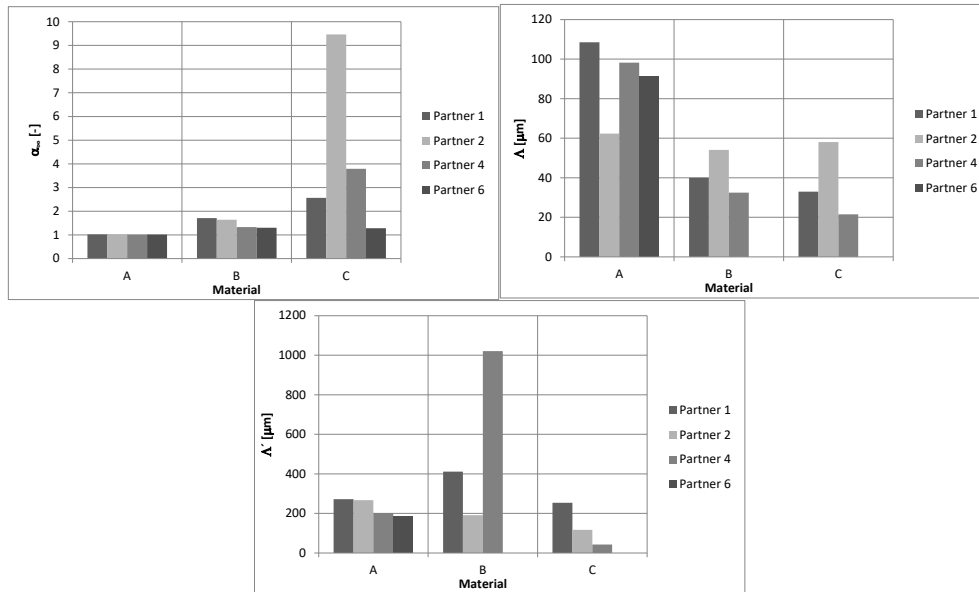
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485 Finally, Fig. 10 shows the comparison for average values of tortuosity and characteristic
 486 lengths obtained by participants. Here it is worth remembering that the direct tortuosity
 487 measurements were executed by partners 1 (on materials A and B) and by partner 6 (only one
 488 sample). The remaining data were obtained from the inverse estimation from acoustic data. In
 489 any case, the dispersions between different institutions for tortuosity are not negligible for
 490 material C (around 85%) while for materials A and B, the dispersion is lower than 15%. The
 491 dispersion for characteristic lengths varies between 20% and 80%.

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496 Figure 10 - The average of tortuosity (left), viscous characteristic length (centre) and thermal
 497 characteristic length (right) for all the materials measured by each of the participating partners.

498

499 IV. CONCLUSIONS

500 The inter-laboratory tests on the acoustical and pore structure properties suggest a poor
 501 reproducibility between laboratories especially for the acoustical properties of highly resistive
 502 materials and granular materials with a rigid frame. The maximum relative errors in the
 503 absorption coefficient, real and imaginary parts of the surface impedance were found to be
 504 $\varepsilon_{\alpha} = 19\%$, $\varepsilon_{\Re(z_s)} = 29\%$ and $\varepsilon_{\Im(z_s)} = 13\%$, respectively. A major cause is likely to be the
 505 natural inhomogeneity in the material slab from which the samples were cut. Other causes
 506 can be the way the sample was actually cut and mounted in the impedance tube. These can
 507 lead to systematic errors between laboratories.

508 There is an obvious need for revision of the current standard¹ where no discussion of
 509 potential measurement problems, and no guidance on the installation of the samples is
 510 provided, no instrument calibration procedures or procedures for periodic verification of the
 511 instruments are detailed, no indications of the number of samples to be measured for the

512 characterization of a material are given and the acceptability of a certain standard deviation
513 on the tests conducted is not discussed.

514 No ISO standard exists to measure characteristic impedance and complex wavenumber. The
515 inter-laboratory errors reach 30% and the causes are likely to be similar to those discussed
516 earlier in these conclusions. It would be appropriate to extend the standards in refs. [1] to
517 include the methodology detailed in ref. [6] for a more complete characterization of the
518 materials in an impedance tube with 3 or more microphones.

519 There is a lack of standard to measure those pore structure parameters which are used
520 routinely to predict the characteristic impedance and complex wavenumber of porous media.

521 The only ISO standard in existence is to measure the air flow resistivity³. For this parameter,
522 the in-laboratory repeatability is high ($\varepsilon_{1,\sigma}=1\%$). However, the reproducibility is reduced

523 considerably to $\varepsilon_{RA,\sigma}=10\%$ for a common poro-elastic material (material A) and to

524 $\varepsilon_{RA,\sigma}=37\%$ for a material with high airflow resistivity (e.g. material C).

525 The values of the inter-laboratory standard deviation determined in our experiments highlight
526 the presence of systematic errors between laboratories, which may be due to the absence of

527 periodic calibration of the static pressure transducers. This procedure is not included in the
528 ISO 9053 standard³. This omission suggests that a revision of the ISO 9053 standard is

529 desirable to reduce errors in the airflow resistivity measurements. One recommendation is to
530 introduce a standardized porous sample with known and well predicted flow resistivity.

531 Modern methods of 3D printing enable manufacturing of samples with highly reproducible
532 porous structure and dimensions which enable the sample to fit in the flow resistivity tube

533 perfectly.

534 The measurement of open porosity of poro-elastic materials is not described by any standard.

535 In this paper, the isothermal compression of volume (Boyle's law) method (ref. [9]) was used
536 by participating partners 1-4 to measure the porosity. The results show an excellent internal

537 repeatability $\varepsilon_{1,\phi} < 1.1\%$. The reproducibility error is $\varepsilon_{RA,\phi} < 9\%$. Partner 6 used the ultrasonic

538 reflection method (ref. [10]), which seems to underestimate the porosity systematically by up
539 to 45% in the case of material C (see Figure 9).

540 Similarly, the measurement of tortuosity and characteristic lengths of porous media is not
541 described by any standard. In this work some of the partners used acoustical inversion

542 methods to determine these parameters¹³⁻¹⁵. The reproducibility was relatively poor because
543 of large dispersion in the tortuosity was observed in the case of material C. A considerable

544 dispersion in the results was observed. As a general conclusion for such parameters, when a
545 direct measurement method was applied errors were lower than 15%. On the contrary, the use

546 of inverse method could lead to errors which could reach up to 80 %. These findings suggest
547 that new standards are needed to define procedures for measurement of the related pore

548 structure parameters of porous media.

549

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