

Adsorption of small cationic nanoparticles onto large anionic particles from aqueous solution: a model system for understanding pigment dispersion and the problem of effective particle density

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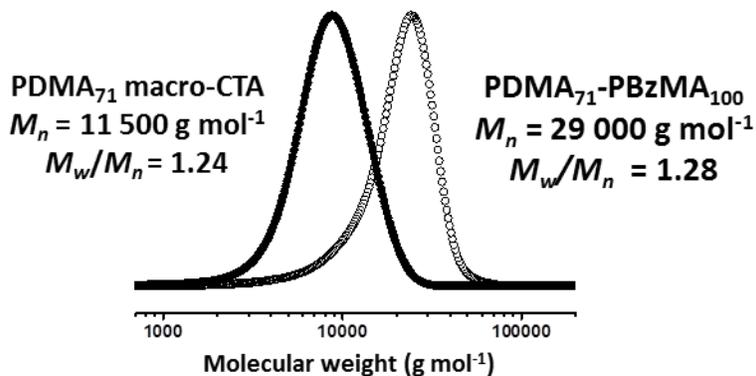


Figure S1. THF gel permeation chromatography traces obtained for the PDMA₇₁ macro-CTA and PDMA₇₁-PBzMA₁₀₀ diblock copolymer (refractive index detector; calibration using a series of ten near-monodisperse poly(methyl methacrylate) standards). A clear shift to higher molecular weight is observed following chain extension of the PDMA₇₁ macro CTA via RAFT alcoholic dispersion polymerization of BzMA.

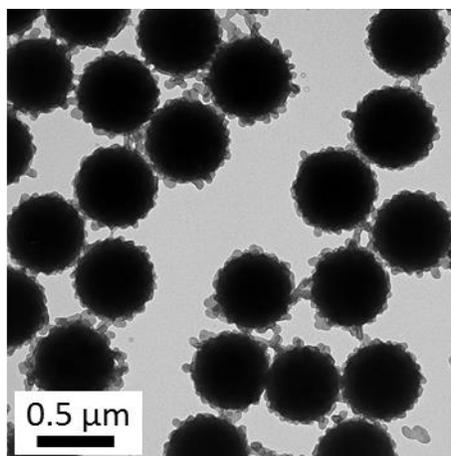


Figure S2. Representative TEM image showing PDMA₇₁-PBzMA₅₀₀ diblock copolymer nanoparticles adsorbed onto 470 nm silica particles.

Spherical micelle model used for SAXS analysis

In general, the scattering intensity of a system comprising one type of non-interacting polydisperse objects can be expressed as

$$I(q) = K \int_0^\infty \dots \int_0^\infty F(q, r_1, \dots, r_k)^2 \Psi(r_1, \dots, r_k) dr_1 \dots dr_k \quad (S1)$$

where $F(q, r_1, \dots, r_k)$ is the form factor for the scattering objects, $\Psi(r_1, \dots, r_k)$ is the distribution function, K is the number density per unit volume of the objects and r_1, \dots, r_k is a set of k parameters describing their structural morphology. The PDMA₇₁-PBzMA_x diblock copolymer chains studied in this work self-assemble in water/ethanol mixtures to form sterically stabilized spherical nanoparticles (or micelles). Thus, the form factor in equation S1 is given by an analytical expression previously reported for spherical copolymer micelles:¹

$$F(q) = [N_{agg} \beta_{mc} A_{mc}(q)]^2 + N_{agg} \beta_c^2 F_c(qR_g) + N_{agg} (N_{agg} - 1) \beta_c^2 A_c^2(q) + 2N_{agg}^2 \beta_{mc} \beta_c A_{mc}(q) A_c(q) \quad (S2)$$

Here the X-ray scattering length contrast for the core block and the corona block is given by $\beta_{mc} = V_{mc}(\xi_{mc} - \xi_{sol})$ and $\beta_c = V_c(\xi_c - \xi_{sol})$, where ξ_{mc} , ξ_c and ξ_{sol} are the X-ray scattering lengths of the core block ($\xi_{PBzMA} = 10.42 \times 10^{10} \text{ cm}^{-2}$), the corona block ($\xi_{PDMA} = 10.12 \times 10^{10} \text{ cm}^{-2}$) and the solvent ($\xi_{water/ethanol} = 9.35 \times 10^{10} \text{ cm}^{-2}$, where the water-to-ethanol ratio is 95/5 w/w) respectively. V_{mc} and V_c denote the volumes of the core block ($V_{PBzMA} = 25.4 \text{ nm}^3$) and the corona block ($V_{PDMA} = 17.0 \text{ nm}^3$), respectively. These volumes were obtained from $V = \frac{M_n}{N_A \rho}$ using the molecular weight (M_w) of the

corresponding block and the solid-state homopolymer densities determined by helium pycnometry ($\rho_{PBzMA} = 1.16 \text{ g cm}^{-3}$ and $\rho_{PDMA} = 1.09 \text{ g cm}^{-3}$). The spherical form factor amplitude is used for the amplitude of the micelle core self-term

$$A_{mc}(q) = \frac{3[\sin(qR_s) - qR_s \cos(qR_s)]}{(qR_s)^3} \quad (S3)$$

where R_s is the radius of the spherical micelle core. The mean aggregation number for the spherical micelles is given by $N_{agg} = (1 - x_{sol}) \frac{4\pi R_s^3}{3V_{mc}}$, where x_{sol} is the solvent fraction in the micelle core. The self-correlation term of the corona block is described by the Debye function:

$$F_c(qR_g) = \frac{2[\exp(-q^2 R_g^2) - 1 + q^2 R_g^2]}{q^4 R_g^4} \quad (S4)$$

where R_g is the radius of gyration of the corona block. For diblock copolymers with a relatively short PBzMA block DP, the corona contribution to the scattering, β_c , is comparable to the scattering from the micelle core, β_{mc} , i.e. PDMA₇₁-PBzMA₁₀₀ with $(\beta_c / \beta_{mc})^2 \approx 0.20$. Thus, in accordance with previous

work,² the amplitude of the corona self-term was calculated from a normalized Fourier transform of the radial density distribution function of the PDMA₇₁ coronal chains in the micelle:

$$A_c(q) = \frac{\int_{R_s}^{R_s+2s} \mu_c(r) \frac{\sin(qr)}{qr} r^2 dr}{\int_{R_s}^{R_s+2s} \mu_c(r) r^2 dr} \quad (S5)$$

The radial profile, $\mu_c(r)$, is expressed by a linear combination of two cubic b splines with two fitting parameters s and a corresponding to the width of the profile and the weight coefficient, respectively.³ The precise analytical expression of the integration applied in the SAXS analysis is not given in ref. 3 but it can be obtained by using a mathematical software package such as Maple or MatLab. An approximate integration can also be found elsewhere.⁴ In accordance with previous studies,³⁻⁵ the confinement $s = 2R_g$ was introduced into the model. It has to be noted that a tended to zero for this condition. For the form factor given by equation S2, a sharp, non-sigmoidal interface between the blocks with no penetration of the corona blocks into the micelle cores was assumed. It was considered for the model (equation S1) that a polydispersity of the spherical micelle core radius (R_s) is expressed as a Gaussian distribution:

$$\Psi(r_1) = \frac{1}{\sqrt{2\pi\sigma_{R_s}^2}} e^{-\frac{(r_1 - R_s)^2}{2\sigma_{R_s}^2}} \quad (S6)$$

where σ_{R_s} is the standard deviation for R_s . The number density per unit volume in equation S1 is expressed as:

$$K = \frac{c}{\int_0^\infty V(r_1) \Psi(r_1) dr_1} \quad (S7)$$

where c is the total copolymer volume fraction in the spherical micelles and $V(r_1)$ is the total copolymer volume in a spherical micelle [$V(r_1) = (V_{mc} + V_c) N_{agg}(r_1)$]. Thus, the overall number of structural parameters required for the spherical micelle model described by equations S1 and S2 is six (namely R_s , σ_{R_s} , x_{sol} , R_g , a and c).

References

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