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The DFT molecular modeling and experimental particle kinetics studies of CO₂-char gasification.

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Abstract

Experimental measurements and DFT atomistic modelling were conducted to elucidate the mechanisms for gasification chemistry of char with CO₂ gas. The molecular models used were based on experimental representations of coal chars derived from the vitrinite- and inertinite-rich South African coals at 1000 °C. The HRTEM and XRD techniques were used to construct parallelogram-shaped PAH stacks of highest frequency in the vitrinite-rich (7x7) and inertinite-rich (11x11) char structures. Computations were executed to get the nucleophilic Fukui functions, at DFT-DNP level, to elucidate the nature and proportions of carbon active sites and quantify their reactivity. The DFT-DNP-computed reaction pathways and transition states, to obtain the energy of reaction and activation energies for the gasification reactions of CO₂ with active carbon sites were examined. These results were compared with TGA experimental results at 900-980°C. The mean nucleophilic Fukui function of the H-terminated char models and active sites located at similar edge positions

decreased with increasing size of char molecules and followed the sequence: zigzag>armchair>tip active sites. The mean DFT-DNP values for the activation energy of 233 kJ mol⁻¹ at the reactive carbon edge was in agreement with the experimental 191±25 kJ mol⁻¹ and 210±8 kJ mol⁻¹ for the respective chars.

Nomenclature

Symbol	Definition	Units
\bar{X}	Arithmetic mean	-
Δ_1	Frequency of modal class minus that of previous class	-
Δ_2	Frequency of modal class minus that of following class	-
ΔE_b	Atomistic activation energy or energy barrier by generalized synchronous transit method ($\sim E_a$)	kJ mol^{-1}
ΔH	Enthalpy	kJ mol^{-1}
1-scf	Single point energy calculations (Simulations)	-
C	Number of aromatic carbons per cluster	-
c	Size of interval	-
C*	(Re)active carbon site	-
C _f	(Re)active carbon site	-
CP-MAS	Cross polarisation-magic angle spinning	-
d ₀₀₂	Inter-layer spacing for a group of N _{ave} parallel layers	Å
DMol ³	A modelling program that uses DFT	-
E _a	Activation energy	kJ mol^{-1}
E _{reac}	Reaction energy by generalized synchronous transit method ($\sim \Delta H$)	kJ mol^{-1}
f ⁺ (r)	Nucleophilic Fukui function	-
f(r)	Electrophilic Fukui function	-
f ⁰ (r)	Radical Fukui function	-
F ₁	Function of absolute temperature	-
F ₂	Function of CO ₂ partial pressure, order of reaction, and lumped pre-exponential factor	-
f _a	Carbon aromaticity,	-
k' _{so}	Lumped pre-exponential factor	$\text{min}^{-1} \text{bar}^{-m}$
L	Lower boundary of median interval	Å
L ₀	Total pore length per unit volume	m m^{-3}
L _a	Crystallite diameter	Å
M	Median	Å
N	number of data points used	-
n	Reaction order	-
P _{CO2}	Partial pressure of CO ₂ gas	Pa
R _i	Initial reactivity of chars	min^{-1}
r ^L	Lumped reaction rate	m s^{-1}
r _s	Reaction rate	m min^{-1}
S	Standard deviation	-
S ₀	Initial surface area	$\text{m}^2 \text{m}^{-3}$
t	Time	s and min
T	Absolute temperature	°C and °K
t _{0.9}	Time for fractional carbon conversion of 90%	min
v _{TS}	Single negative frequency from generalized synchronous transit method	cm^{-1}
X	Fractional conversion of carbon	-
SK	Skewness	-

Greek symbols

Symbol	Definition	Units
ϵ_0	Initial porosity of char samples	%
θ	Peak position / XRD angle of scan	radians
γ	Gamma band/peak	-
ψ	Dimensionless structural parameter for char pores	-

Abbreviations

Acronym	Definition
ALS	Advanced Laboratory Solutions
C-C	Carbon to carbon bond/site
CHPC	(National) Centre for high performance computing (Cape Town)
DFT	Density functional theory
DIIS	Direct inversion of the iterative subspace (for SCF convergence)
DNP	Double numerical polarization basis set
DTA	Differential thermal analyser
$f(r)$	Fukui function
$f^+(r), f(r), f^0(r)$	Nucleophilic, Electrophilic, radical Fukui functions
GGA	Generalized gradient approximation
HOMO	Highest Occupied Molecular Orbital
HPC	High performance computing (NWU)
HRTEM	High resolution transmission electron microscopy
LAMM	Laboratory of applied molecular modelling (NWU)
LUMO	Lowest Unoccupied Molecular Orbital
LST	Linear synchronous transit (method to search for TS)
MEP	Minimum energy pathway
NMR	Solid state nuclear magnetic resonance spectroscopy
NWU	North-West University
PAH	Polyaromatic hydrocarbons
PES	Potential energy surface
QOF	quality of fit
QST	Quadratic synchronous transit (method to search for TS)
RMS	Root-mean-square for convergence
RPM	Random pore model
SCF	Self-consistent forcefield
SEM	Scanning electron microscopy
TGA	Thermogravimetric analyser
TPD	Temperature-programmed desorption
TS	Transition state
XRD	X-ray diffraction

1. Introduction

Economic and environmental constraints necessitate research into higher efficiency coal utilisation, to develop advanced technologies with reduced greenhouse gas emissions. A study of coal gasification is part of this effort. Coal gasification is a complex process that involves reactions with O_2 , H_2O , H_2 and CO_2 , either as a single component reactant or a mixture of two or more of the gaseous reagents. Gasification reactions have been studied using a variety of experimental and theoretical techniques [1-14] over a lengthy period [15] because of their importance to a variety of applications.

Various aspects of coal gasification by CO_2 and its application to oxy-combustion has been reviewed by Irfan et al. [9], where the dependence on coal rank, pressure, temperature, gas composition, catalyst and the minerals present inside the coal, heating rate, particle size, and diverse reactor types was established. A large volume of literature has been devoted to mechanisms for the gasification reactions of carbon by oxygen-containing gases through experimental methods that included thermogravimetric techniques, evolved gas analysis, TPD, transient kinetics, as well as other characterisation techniques such as SEM, HRTEM and XRD [6, 16-19]. This work highlighted the effect of structural properties of chars in gasification reaction mechanism and continues to elucidate the intricacies of coal gasification. Quantum mechanics modelling studies have contributed to the evaluation of medium- to high-temperature interactions between carbonaceous material and oxygen carrying gases [20-28]. High-level molecular modelling techniques used to carry out comprehensive study of coal gasification reactions are often confronted with the complex nature of the large coal and char structures, which present challenging computational requirements. A better understanding of associated atomistic reaction mechanism may facilitate improvements in the evaluation of coal char gasification kinetics, including rates and efficiencies. Rates of gasification of porous carbons are affected by parameters such as active site concentration, accessibility of the reactant gas into the internal area of the char, and presence of active catalysts that mediate in reactions with gasification agents [24, 29-37].

While molecular modelling studies have been conducted into the reactivity of oxygen, steam and CO_2 with carbonaceous material [28, 38-43], the models used were not directly derived models from high rank coal. Other work has focused on the chemisorption of oxygen-carrying gases on carbonaceous surfaces [28, 39, 40, 42], while Montoya et al. used ab initio

techniques to model the kinetics of elementary reaction in carbonaceous material [38]. Radovic et al. compared different chars in the C-CO₂ reaction using DFT and found that char gasification may be a structure-sensitive reaction [44]. DFT was used to model the active carbon sites responsible for reaction with the oxygen-carrying gases and to highlight a dissociative-chemisorption reaction of CO₂ on active carbon sites [28, 30, 43].

The present investigation deals with reaction mechanisms between the CO₂ and coal chars, using DFT molecular modelling, and experimental techniques, to evaluate the chemistry at active sites in char, and to examine how modeling results using a simple molecular model compare with experimental data from slow-heated, or hard-cooked, and de-ashed chars derived from the South African inertinite- and vitrinite-rich coals.

2. Chars and methodology

Experimental work relevant to molecular models is based on: (i) the characterization of two de-ashed chars, and (ii) constructing a molecular model based on HRTEM data from these chars [45].

2.1 Char molecular models

The chars were prepared by heating inertinite- and vitrinite-rich South African Gondwana coals of the Permian Age to 1000 °C in a packed bed reactor, at 20 °C min⁻¹, in Nitrogen, maintaining this temperature for 60 min, and then cooling [45]. The Nitrogen flow rate was 1.5 dm³ min⁻¹ to ensure a rapid escape of volatile matter in an inert atmosphere, and averting secondary reactions [45-47]. The chars were de-ashed by HCl-HF treatments [48-52] and characterized using petrographic, HRTEM, NMR and XRD techniques, similar to Roberts et al. [45].

The char models used for modelling studies were based on the HRTEM results shown in Figure 1; the method of parallelogram catenations by Mathews et al. was used to calculate size distribution of aromatic carbon sheets [53, 54]. This technique involves intensive image processing of the photo-micrographs (Figure 1(a)) to extract skeletonised lattice fringes (Figure 1(b)) following an approach of Sharma et al. [55]. A relationship between these

lattice fringes and aromatic carbon sheets of coals (Figure 1(c)) was used to calculate the distribution of the parallelogram-shaped aromatic molecules of the chars (Figure 1(d)) [53]. The detailed table of results is presented in Roberts et al. [45]. The modal values, calculated from Equation (1), of the grouped aromatic fringe length distributions for the inertinite- and vitrinite-rich chars were 36.8 and 21.0Å, respectively. Both distributions were positively skewed at 2.4 and 3.1, respectively, according to Pearson's second coefficient of skewness [56], Equation (2), where the vitrinite-rich char was more skewed to smaller values. The modal values correspond with the 11 by 11 and 7 by 7 carbon rings of the respective chars, based on the assumption that the fringes are as deep as they are wide, of parallel catenation, and that the number of carbon atoms, and thus number of rings, can be determined as the mean from the longest and shortest possible orientations [53, 57]. It also assumes no curvature or defects [57, 58]. The results from XRD and HRTEM techniques showed a reasonable agreement. The HRTEM modal value for the 1000 °C inertinite-rich chars compares well with the crystallite diameter measured by the XRD technique, L_a (10), of 37.6Å. However, the L_a (10) of 30.7Å for the vitrinite-rich chars was higher than the corresponding modal value.

$$Mode = L + \left(\frac{\Delta_1}{\Delta_1 + \Delta_2} \right) c \quad (1)$$

$$SK = \frac{3(\bar{x} - M)}{s} \quad (2)$$

2.2 DFT computational methods

The DFT computational requirements restrict the model size for tractable calculations, and the utility of a single graphene sheet to model coal char is a major concern of the present studies. The SCF convergence for the 11x11 and 7x7 char models was difficult even if convergence were accelerated by the smearing tool from the default value of 13.1 kJ mol⁻¹, with small increases up to 525.1 kJ mol⁻¹. Thermal smearing is a setting applied to the orbital occupation to improve the SCF convergence within the specified number of iterations. Consequently, it was necessary to reduce the aromatic carbon sheets of 7x7 from the vitrinite-rich char and 11x11 from the inertinite-rich char, to smaller sizes of 3x3 and 5x5, respectively, by applying a factor of approximately two. However, the authors acknowledge that if the number of electrons in two structures are different, then one cannot compare the electronic energy directly. In addition to this, the 4x4-carbon ring structure was included to

create additional data point, since it was also present in both the vitrinite- and inertinite-rich chars [45]. While this approach was taken, it was noted that the electronic states of the graphene models are greatly dependent on the model size so that the energetic values such as adsorption energies and activation energies are affected, as pointed out by Pham and Truong [59]. Therefore, this investigation recognizes the sensitivity of the computed results to the simple molecular models, as the experimental data is derived from coal derived chars, which are expected to contain a variety of carbon active sites (and correspondingly variety of possible electronic configurations).

The edge carbons of the aromatic carbon sheets from the HRTEM were terminated with hydrogen to give polyaromatic hydrocarbon (PAHs) char models, following Radovic and Bockrath assertions that structures with entirely free edge carbons were unlikely, but likely with hydrogen termination. However, hydrogen termination may be appropriate for lower temperature chars, but less so for chars at high temperatures [60]. Under gasification conditions at high temperatures, thermal decomposition and abstraction of edge hydrogen would be anticipated [41]. It is also probable that the coal-derived chars contain numerous activated C* sites, so that a simple graphene molecule containing one C* active site represents an unlikely model. This difficulty cannot be overcome with the limitations of current computational resources, and in this work, the graphene model is used with a number of active sites derived through the loss of edge hydrogen groups; we acknowledge the difficulties inherent in this approach and emphasize the limitations in our current molecular modeling capabilities. It is also understood that gasification includes other reactants (H₂, H₂O, CO₂ and O₂), but the current modelling deals solely with reactions involving CO₂ [41, 61, 62]. It is also necessary to restrict the present modeling to reactions involving CO₂, as modeling gasification involving a number of gases is outside our current capabilities. We utilize simple systems to provide insights on likely mechanisms in complicated, practical systems, with the additional input from experimental data, which enables an assessment of the computational results against experimental ones.

The modeling study discussed in this paper commences with the loss of H to form a reactive carbon edge or free active sites, followed by chemisorption of CO₂ and subsequent formation of CO. The modelling results are compared with the experimental to assert the viability of the proposed mechanisms for CO₂ gasification of char derived from bituminous coal medium C rank.

Molecular models have been proposed to simulate structures with various types of active sites, e.g. the basal plane, armchair, zigzag and tip [24, 28, 39-42, 60, 63, 64]. Solid-state ^{13}C nuclear magnetic resonance (NMR) techniques [65, 66], have shown that coal chars commonly consist of structures of aromatic stack/clusters of 12-25 aromatic carbon atoms (3-7 benzene rings) [38, 67]. NMR analysis of 1000 °C however, encounters line broadening which prevents calculation of structural parameters, and thus it cannot predict the aromatic carbon cluster size under these conditions [65]. Line broadening is attributed to aromatic free radicals that begin to change the relaxation behaviour of the aromatic carbons [68]. Most characterization results have indicated that the 1000 °C inertinite- and vitrinite-rich coal chars behaved in a similar way, except for the aromatic fringes and L_a (10) values from the HRTEM and XRD techniques, respectively. These considerations provide a rationale for using the reduced modal PAH structures of 7x7 and 11x11 PAHs as models of the respective inertinite- and vitrinite-rich chars, for DFT computations.

2.2.1 DFT reactivity modelling procedures with the Fukui function

The DFT geometry optimization of our char models were performed using with the DMol³ method in the Accelrys Material Studio[®] software (version 6.0) using the Laboratory of Applied Molecular Modelling (LAMM) high performance computing (HPC) resources of the North-West University in Potchefstroom and the National Centre for HPC in Cape Town, South Africa. The DMol³ calculations give accurate self-consistent field (SCF) densities as discussed in [69] and numerically localized basis sets yields improved enthalpies for the formation of H, N, O, F, Cl, C, Si, S atoms. The exchange-correlation potential of generalized gradient approximation (GGA) functional of Perdew and Wang (PW91) was used [70, 71]. Delley found that the DFT performance in predicting molecular enthalpies of formation of metallic solids was significantly improved for the Perdew–Burke–Ernzerhof (PBE) functional in comparison with PW91 [69]. However, several publications on investigations similar to the current (search for transition state) used PW91 functional [72-75]. This package provides electronic wave functions which are expanded in a double numerical coupled with polarization basis set (DNP) truncated at a real space cut-off of 4.4Å. The DNP is comparable to Gaussian 6-31G** basis set used for molecular and solids calculations, but the DNP is reported to be more accurate than the 6-31G** set of the same size [76, 77]. All the

calculations were spin-unrestricted. The default convergence tolerances used for the energy, force, and displacement were of 0.1 kJ mol^{-1} ($2 \times 10^{-5} \text{ Ha}$), $10.5 \text{ kJ mol}^{-1}\text{\AA}^{-1}$ ($0.004 \text{ Ha } \text{\AA}^{-1}$) and $5 \times 10^{-4} \text{ nm}$ (0.005\AA), respectively. The SCF was left on fermi, which required up to eight times the maximum number of iterations (1000) for convergence. This was coupled with the direct inversion of the iterative subspace (DIIS) technique developed by Pulay [78].

The 1-scf energy calculations were performed with the Accelrys Material Studio[®] software package on geometry optimised structures to determine the reactivity indicator using the Fukui function. The Fukui function, $f(r)$, is among the most commonly used reactivity indicators [79, 80]. The Fukui function is usually given as the change in the density function of the molecule as a consequence of changing the number of electrons in the molecule [80] and allows a chemical reaction to be understood in terms of the highest occupied molecular orbitals and lowest unoccupied molecular orbitals [81]. The function $f^+(r)$ is a measure of reactivity toward a donor reagent (nucleophilic attack), $f^-(r)$ values measure reactivity toward an acceptor reagent (electrophilic attack) and an average value of both may be taken as a measure of reactivity toward a radical $f^0(r)$ [82]. Parr and Yang showed that sites in chemical species with the largest $f(r)$ value are those with highest reactivity and are the preferred reaction site(s) [83-85]. Research has emphasised that hardness-softness $f(r)$ derived from the DFT are powerful tools to predict active sites of a molecule [84, 86, 87]. The $f(r)$ defines the reactivity of an atom in a molecule and, together with local softness, are suitable to describe reactivity of different substrates [88], including symmetrical char structures. Here, the objective is to predict reactivity of all H-terminated and selected, free edge carbon atoms from simplified coal char structures/molecules with oxygen atoms from CO_2 gas molecules, which is anticipated as the first chemical interaction step associated with the CO_2 -char gasification reaction mechanism. The O from the CO_2 atom participates as electron donor to chemisorb to the free edge carbon atom, thus making the nucleophilic Fukui function, $f^+(r)$, the preferred index to predict reactivity. The $f^+(r)$ -based reactivity analyses were first conducted on H-terminated edge carbon sites to gather the rationale for the free active sites (reactive edge) to be used in the CO_2 -char gasification reaction mechanism. The energetics of the loss of H were examined using the TS theory approaches available in the Accelrys Material Studio[®] software package. The relative accuracy of these computations were inevitably restricted by the very large computational resources required for these calculations.

2.2.2 DFT modelling of char-CO₂ gasification reaction mechanism

A general mechanism of the reaction of CO₂ gas with solid carbon surfaces is given by Equations (3)-(6) [89]. Initially, the CO₂ gas molecule is chemisorbed on the active site C_f on the char surface [28, 30, 64, 90, 91].



This is followed by dissociation of the adsorbed CO₂ to form CO according to Equation (4), and an oxygen surface complex (quinone) on the char surface, $C_f(O)$, which is [28, 91, 92].



The subsequent formation of an additional CO molecule results from the O-complex forming CO, (gasifying) with the solid carbon on the edge of the char surface, and escaping as a CO gas molecule as in Equation (5) [9, 20, 92]. Montoya et al. has shown that the six-member ring containing the O-complex ($C_f(O)$) undergoes bond weakening and breaking to form the CO gas molecule. This ring may collapse to form a five-member ring with loss of CO, or the loss of CO may accompany the formation of a defect which would act as another active carbon site, to continue the process of CO₂ gasification [38].



Equations (3)-(5) summarise the Boudouard reaction shown in Equation (6) [31, 93]; the mechanisms however include intermediate reactions associated with the bond weakening, breaking, and re-arrangements.



The Boudouard reaction [31, 93] mechanism was examined by application of TS theory, based on a generalized synchronous transit scheme, which E_{reac} and ΔE_b [94-96] from DFT modelling of the char-CO₂ gasification reaction using the 3x3 model. Transition state theory enables consideration of transition rates in diffusion reactions and also energy profiles between reactants and products, as discussed in [73, 97-99]; the TS for a reaction is the

geometry corresponding to the highest energy along the MEP, which gives the activation energy. In the context of reaction rates, the reactant(s) and product(s) have an initial and a final state, both of which are local minima on the PES scan accurately locating the transition state.

The TS were searched by LST/QST methods [94-96, 100]. Usually after the second QST optimization the CG refinement finds the transition state. However, in molecular systems such as the ones used here, refinement of TS geometries would require unrealistically lengthy computations, especially when the system has two or more imaginary frequencies. Consequently these computations were restrained to the location of likely TS within the context of the proposed mechanisms. A number of calculations on the TS showed that there was only one imaginary frequency, which allowed for the refinement of the TS energy using a TS Optimisation Task. The major thrust of this work relies on DFT results of reaction energies for each reaction site and energy barriers, because these results enable a direct comparison between modelling and measured data. The energetics were determined for the char-CO₂ interactions for the C-H bond breakdown, CO₂ chemisorption, CO formation through CO₂ dissociation and C-C decomposition reactions.

2.3 Experimentation: Char-CO₂ reactivity (Particle kinetics)

The objective of the experimentation undertaken was to validate the results (activation energy) obtained from the reaction rate modelling based on the molecular structure (atomic reaction rate) with results obtained using particle reaction kinetics. For this purpose, experimentation was carried out using a thermogravimetric analyser with the random pore model reaction modelling [6, 101, 102].

2.3.1 Thermogravimetric analyser

Isothermal char-CO₂ gasification experiment was carried out on a TA instruments' SDT Q600 TGA supplied by ALS, South Africa, using de-ashed chars derived from inertinite-rich and vitrinite-rich South African coals. The SDT Q600 TGA can work within the temperature control range of ambient +5 to 1500 °C, with a platinum-rhodium Type R thermocouple; with a temperature accuracy, precision and DTA sensitivity of ± 1 °C, ± 0.5 °C and 0.001 °C

respectively. The TGA uses an alumina horizontal beam null balance with a weight sensitivity of 0.1 μg , weight accuracy of $\pm 0.1\%$ and a repeatability of 99.9%. The controlled temperature zone of the furnace is 50 mm. Gas flow through the TGA furnace is uni-directional over the horizontal balance, thus, a purge gas counter-flow is not needed during experiments. The TGA also come with two in-built mass flow controllers, which control the gas flow rates.

2.3.2 Experimental procedure

The isothermal gasification experiments were carried out with CO_2 gas at atmospheric pressure and isothermal temperatures of 900, 920, 940, 960 and 980 $^\circ\text{C}$. Sample mass of 10 ± 1 mg of $-75 \mu\text{m}$ particle size was used for each of the experiments on the TGA. The sample was loaded into a 90 μL porous alumina sample holder carefully placed on the balance pan before closing the furnace. The set-up (furnace and sample) was flushed with nitrogen gas at $100 \text{ cm}^3 \text{ min}^{-1}$ for 25 min to evacuate any oxidizing species in the furnace, as well as to stabilize the balance. The sample was heated at $30 \text{ }^\circ\text{C min}^{-1}$ under the N_2 gas flow to the required temperatures. The N_2 gas was automatically turned off and the CO_2 gas entered at a flow rate of $45 \text{ cm}^3 \text{ min}^{-1}$ for the isothermal gasification experiments. Reaction times for the experiments varied from 36 to >48 hours, similar to Everson et al. [1, 6].

2.3.3 Calculation of activation energies

The intrinsic reaction rate equation based on the random pore model was evaluated according to the method proposed by Everson et al [101] involving essentially: (1) the determination of the lumped reaction rate consisting of the intrinsic reaction rate and the initial physical properties of the chars, and (2) the determination of the activation energies and pre-exponential factors using the Arrhenius equation using the lumped reaction rate determined from the experimental results.

The relevant random pore model is in Equations (7)-(13) with a brief description of the calculation procedure. The equation for the conversion of carbon, on ash free basis, involves the reaction rate and structural behaviour, as well as characterised by the lumped reaction rate and structural parameter.

$$X = 1 - \exp \left[-r^L t \left(1 + \frac{r^L t \psi}{4} \right) \right] \quad (7)$$

$$r^L = \frac{r_s S_o}{(1 - \varepsilon_o)} \quad (8)$$

$$r_s = k_{so} \exp \left(\frac{-E_a}{RT} \right) \cdot P_{CO_2}^n \quad (9)$$

$$\psi = \frac{4\pi L_o (1 - \varepsilon_o)}{S_o^2} \quad (10)$$

For the evaluation of the lumped reaction rate, the structural parameter needs to be evaluated, which can be accomplished with Equation (11) involving regression with the unknown structural parameter. This equation can be derived from Equation (7) and can be seen to be independent of temperature (reaction rate) and is a convenient method to determine the structural parameter.

$$\frac{t}{t_{0.9}} = \frac{\sqrt{1 - \psi \ln(1 - X)} - 1}{\sqrt{1 - \psi \ln(1 - 0.9)} - 1} \quad (11)$$

With a known structural parameter, the lumped reaction parameters can be calculated also by regression with the Equation (7) and the experimental results. From a knowledge of the lumped reaction rate as a function of temperature, the Arrhenius equation parameters can be calculated using the logarithmic form of a combination of Equations (8) and (9), to yield Equations (12) and (13) [6, 101].

$$\ln(r^L) = -\frac{E_a}{RT} + n \ln(P_{CO_2}) + \ln(k'_{so}) \quad (12)$$

$$\ln(r^L) = F_1(T) + F_2(P_{CO_2}, n, k'_{so2}) \quad (13)$$

3. Results and discussion

The utility of single graphene sheets (or ribbons) as models for coal-char gasification studies presents a number of challenges. Pham and Truong have shown the size of the graphene

ribbon increases in length and width, its orbital energy gap decreases, with the lowest energy gap obtained for ribbons of length >8 ; they also demonstrate the errors in the total energies as a function of the ribbon length at different levels of theory [59]. The present work involves the size of the graphene structure consistent with the low energy differences reported by Pham and Truong [59]. This work also uses the Fukui function as an indicator of reactivity [79, 80] for edges on the graphene structure. The standard Fukui function concept of conceptual DFT includes the contribution from HOMO and the LUMO [81], but also from other chemically relevant orbitals. While the treatment by Radovic may provide particular insights related to electronic states and spin crossover effects, some of the DFT results (e.g. suggesting exothermic chemistry for CO dissociation involving small graphene ribbons and CO₂) appear counter intuitive [28]. Of greater concern, however, is the large difference between the molecular structure of a single graphene sheet, and that of freshly formed coal-char, which is three-dimensional and disordered to varying degrees, and may form a variety of active sites during pyrolysis of coal as it undergoes heating in a gasifier. These considerations have caused the present investigation to emphasise the endothermic aspects of char gasification by CO₂, and to seek energy of reaction from the model for comparison with experimental data. Detailed studies of electronic states would need additional work with a larger char model, requiring considerably greater computational resources that are currently available and would be reported in a separate paper.

Fukui function as a reactivity indicator, allowed for the modelling sequence that commences with loss of hydrogen to form active C* sites, followed by chemisorption of CO₂, gas molecule and the energies obtained for thermal dissociation of CO₂ to form the first CO gas molecule, and thermal breakdown to yield the second CO gas molecule. Clearly, this sequence is a simplification of the actual gasification process observed in practice. The mechanism, however, is elucidated using both computed values of ΔE_b and E_{reac} and their comparison with experimental data.

3.1 DFT reactivity modelling and the Fukui function

Substantial research have found that the chemical interactions between carbon molecules and reactant gases take place at the edge active carbon atoms for coal based char models and also for non-coal char models [3, 28, 39, 43, 60, 64, 103-105], and Zhang et al. found that the

carbons in the basal plane of large, fused-ring aromatic domains of chars, were inert [106]. Figure 1 shows results from the HRTEM technique, which were used to derive the char models used here [45].

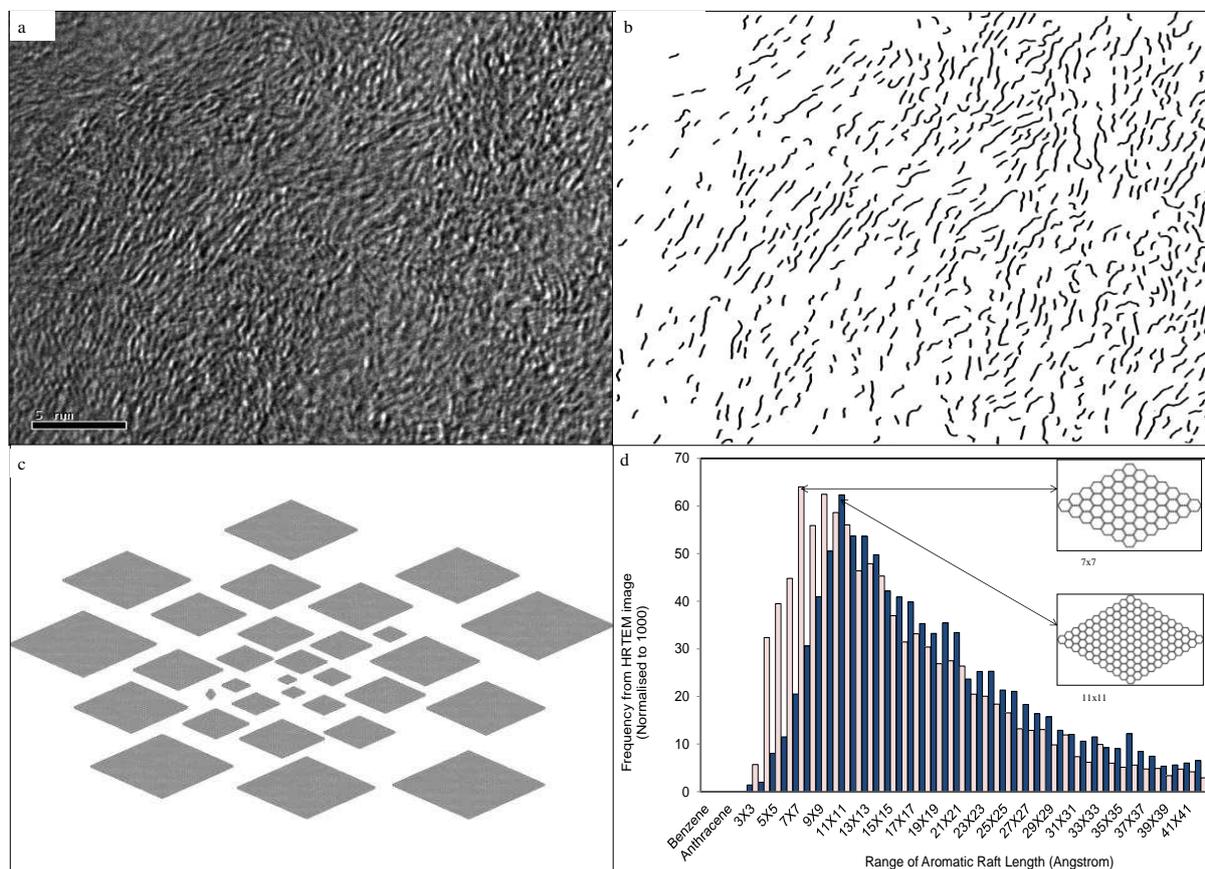
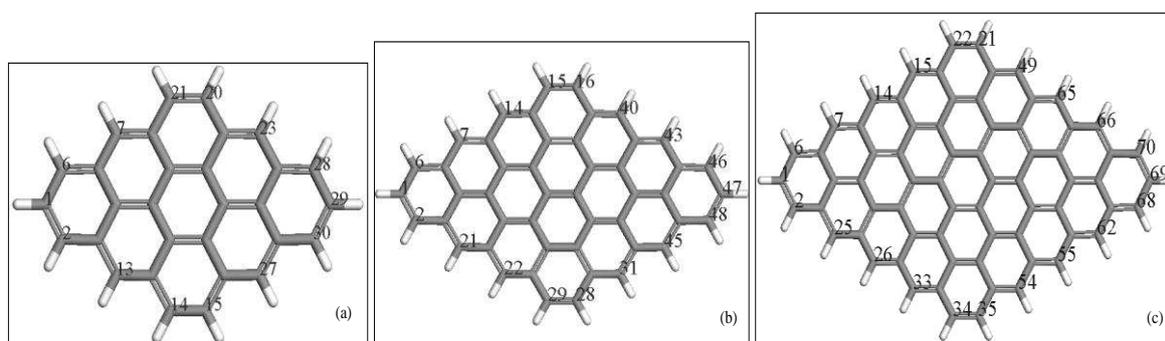


Fig 1. HRTEM results. Microphotograph for the inertinite-rich chars ((a), top left), with scale bar = 5 nm. Lattice after image processing ((b), top right). Parallelogram catenations ((c), bottom left). Distributions of aromatic raft lengths ((d), bottom right), where, dark and light bars = inertinite- and vitrinite-rich chars, respectively.

The char-CO₂ reaction mechanism commenced with reactivity analysis of the H-terminated edge carbon atoms using the $f^+(r)$ values. Results from here were used to rationalise the selection of free active edge carbon atoms, which involved the loss of H from C-H edges, using the transition state theory. The definition, position and numbering of H-terminated C active sites are presented in Table 1 accompanied by Figure 2. The $f^+(r)$ results are presented in Table 2(a)-(c) at an accuracy of ± 0.001 . The $f^+(r)$ is a dimensionless reactivity index.

Table 1. Definition of H-terminated edge carbon sites, including position and quantity.

Edge C name	Symbol	3x3	4x4	5x5
Tip	C _t	C1, C29	C1, 47	C1, C69
Zigzag, proximate to tip	C _z	C2, C6, C28, C30	C2, C6, C46, C48	C2, C6, C68, C70
Zigzag, remote from tip and armchair	C _{zi}	C7, C13, C23, C27	C7, C14, C21, C22, C31, C40, C43, C45	C7, C14, C15, C25, C26, C33, C49, C54, C55, C62, C65, C66
Armchair	C _{r1} , C _{r2}	C14, C15, C20, C21	C15, C16, C28, C29	C21, C22, C34, C35

**Fig 2.** The PAH char models used to determine the nucleophilic Fukui values of the edge carbon sites, where (a), (b) and (c) = 3x3, 4x4 and 5x5 structures, respectively.

The results indicate that reactivity of H-terminated char models and active sites located at similar edge positions decreases with increasing size. In addition to this, the reactivity analysis of the respective char models show that the $f^+(r)$ of all zigzag (C_z) carbon sites adjacent to the tip was similar, but decreased with remoteness of intermediate zigzag sites (C_{zi}) from the tip. The reactivity values follow the sequence: zigzag > armchair > tip.

The Fukui function values on the H-terminated edge carbons are taken as indicative of the sites that would form C* active sites (reactive edge) after loss of H. This shows the reactivity of sites located at similar edge positions of the respective 3x3, 4x4 and 5x5 PAH char models was similar. This outcome was observed even if the “use symmetry” option of the program were not utilised during computations.

Table 2. Reactivity of 3x3 ((a), left-hand side), 4x4 ((b), middle) and 5x5 ((c), right-hand side) coal char models from $f^+(r)$.

(a) 3x3 PAH				(b) 4x4 PAH				(c) 5x5 PAH			
Edge site	$f^+(r)$	Milliken Reactivity level*	Edge type	Edge site	$f^+(r)$	Milliken Reactivity level*	Edge type	Edge site	$f^+(r)$	Milliken Reactivity level*	Edge type
C1	0.004	4	C _t	C1	0.003	5	C _t	C1	0.003	6	C _t
C2	0.039	1	C _z	C2	0.033	1	C _z	C2	0.026	1	C _z
C6	0.039	1	C _z	C6	0.033	1	C _z	C6	0.026	1	C _z
C7	0.033	2	C _{zi}	C7	0.028	2	C _{zi}	C7	0.024	2	C _{zi}
C13	0.033	2	C _{zi}	C14	0.016	3	C _{zi}	C14	0.017	3	C _{zi}
C14	0.014	3	C _r	C15	0.008	4	C _r	C15	0.011	4	C _{zi}
C15	0.014	3	C _r	C16	0.007	4	C _r	C21	0.005	5	C _r
C20	0.014	3	C _r	C21	0.028	2	C _{zi}	C22	0.006	5	C _r
C21	0.014	3	C _r	C22	0.016	3	C _{zi}	C25	0.023	2	C _{zi}
C23	0.033	2	C _{zi}	C28	0.007	4	C _r	C26	0.017	3	C _{zi}
C27	0.033	2	C _{zi}	C29	0.008	4	C _r	C33	0.010	4	C _{zi}
C28	0.039	1	C _z	C31	0.016	3	C _{zi}	C34	0.005	5	C _r
C29	0.004	4	C _t	C40	0.016	3	C _{zi}	C35	0.006	5	C _r
C30	0.039	1	C _z	C43	0.028	2	C _{zi}	C49	0.010	4	C _{zi}
Mean	0.025			C45	0.028	2	C _{zi}	C54	0.010	4	C _{zi}
				C46	0.032	1	C _z	C55	0.017	3	C _{zi}
				C47	0.003	5	C _t	C62	0.023	2	C _{zi}
				C48	0.032	1	C _z	C65	0.017	3	C _{zi}
				Mean	0.019			C66	0.023	2	C _{zi}
								C68	0.025	1	C _z
								C69	0.003	6	C _t
								C70	0.025	1	C _z
								Mean	0.015		

* Largest $f^+(r)$ value = highest reactivity level [83-85].

It was postulated that during CO₂ gasification, the higher reactive sites such as C_z and C_{zi}, from the 3x3, will preferentially lose H. Annealing during pyrolysis however, may give rise to active sites of lower reactivity as the 3x3 grows to 5x5 sizes, because at high temperatures, and chemically-controlled gasification process, char undergoes thermal annealing in competition with gasification [106-108]. This causes the growth of crystallite diameter, L_a, by the coalescence of crystallites along the a-axis (e.g. from 3x3 to 5x5 aromatic fringes) [109] for both graphitisable and non-graphitisable carbons [109]. Zhang et al. studied the effects of pyrolysis intensity on the reactivity of coal char in chemically-controlled CO₂ gasification experiments and found the largest difference in coal char reactivity was observed in the initial stage of gasification, but, as the process continues, the active sites of high

activity were preferentially removed by CO₂ or annealed into less active site and the active sites with lower reactivity were left in residual coal char [106]. Suuberg also found that the inherent reactivity of the chars was attributed to the changes in the active site concentrations, which were often thermally induced [110]. Hence, under high temperature regimes and chemically-controlled gasification process, the char actually undergoes a process of thermal annealing (pyrolysis) in competition with gasification [106-108]. With the increase of carbon conversion or when the higher reactive sites in the 3x3 model are removed/annealed, the reactivity, and the number of reactive carbon sites reach similar levels and the difference in char reactivities becomes smaller [106, 110].

Consequently, the carbon edge (reactive edge) composed of reactive sites C_t, C_z, C_{zi}, C_{r1} and C_{r2} shown in Figure 5(b) was used to model CO₂-char gasification reactivity and to examine relative reactivities at each site.

3.2 DFT modelling for char-CO₂ gasification reaction mechanism

The utility of single graphene sheets (or ribbons) as models for coal-char gasification studies presents a number of challenges. Pham and Truong have shown the size of the graphene ribbon increases in length and width, its orbital energy gap decreases, with the lowest energy gap obtained for ribbons of length >8; they also demonstrate the errors in the total energies as a function of the ribbon length at different levels of theory [59]. The size of the graphene structure in the current investigation is consistent with the low energy differences reported by Pham and Truong [59]. The Fukui function as an indicator of reactivity for edges on the graphene structure is also used. The model has arbitrarily been used with a number of C* (H free) sites, in an effort to illustrate different reactivities even after H loss (as may be anticipated in freshly formed coal char). Detailed studies of electronic states would require additional work with a larger char model, requiring considerably greater computational resources than are currently available.

3.2.1 The loss of H to form an active site

Char models have generally used selected active sites for DFT studies on the interactions of carbonaceous surfaces and gaseous reactants, with the remaining edge carbons capped with

hydrogen atoms [20, 24, 27, 28, 39, 111]. Radovic detailed the challenges associated with the loss of hydrogen from the C-H edge carbons to form the active sites [28], while Domazetis et al., used semi-empirical modelling techniques to describe H abstraction [33]. Radovic has proposed isolated carbene-type zigzag carbon atom and a carbyne-type armchair pair of carbon atoms [43]. While the necessity of using small molecular models may be acceptable, together with the suitability of such sites for these models, the question is how similar are these to mechanisms operating during char formation as coal particles are heated to elevated temperatures. Obviously further work is required using larger molecular models that closer reflect the properties of coal-derived chars and this requires greater computational resources than currently available. Char formation from the coals used in the present studies has been discussed elsewhere [45]. The current investigation, because of the limitations discussed, commences by examining H loss from well-known benzene molecules and progresses to larger (3x3) char molecules. The loss of H was at the following sites: C_t, C_z, C_{zi}, C_{r1} and C_{r2}. From Table 2, the initial reactions are likely to be at the C_z and C_{zi} sites.

The C-H bond breakdown for the molecule benzene molecule, as shown in Figures 3(a)-(d), was evaluated to obtain the activation energy (energy barrier), ΔE_b and reaction energies (bond dissociation energies), E_{reac} for this simple system, defined by Equation (14) and (15) respectively.

$$\Delta E_b = \text{Energy of the transition state} - \text{Energy of the reactant} \quad (14)$$

$$E_{\text{reac}} = \text{Energy of products} - \text{Energy of reactants} \quad (15)$$

Figure 3 shows the structures examined for H-loss from para-, meta-, and ortho-carbons, and the results in Table 3 show ΔE_b values of 484, 452 and 485 kJ mol⁻¹, respectively.

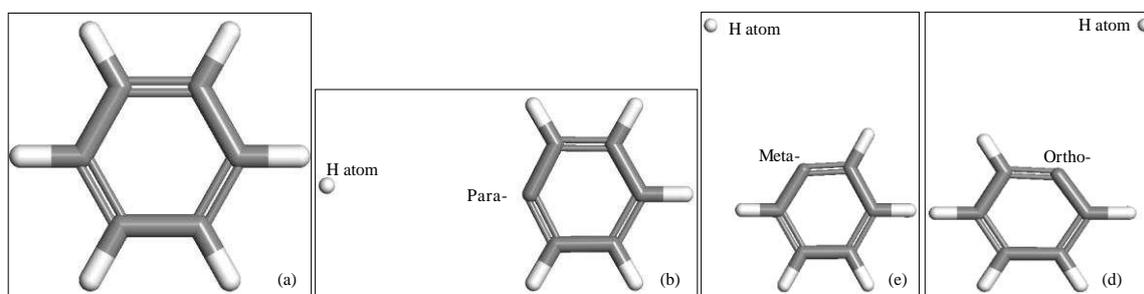


Fig 3. Geometry optimized structures for H-loss from (a) benzene molecule at (b) para-, (c) meta- and (d) ortho-carbon sites.

The mean ΔE_b value of approximately 474 kJ mol^{-1} ($113 \text{ kcal mol}^{-1}$) for a free C-edge in benzene and E_{reac} (comparable to ΔH) for the reaction in Equation (16) [112], of 440.7, 440.4 and $443.1 \text{ kJ mol}^{-1}$ respectively, was obtained.



The values compare well with the published value of $464.0 \pm 8.4 \text{ kJ mol}^{-1}$ ($110.9 \pm 2.0 \text{ kcal mol}^{-1}$) [113, 114]. These DFT-DNP procedures, used to examine loss of H in the 3x3 PAH char model shown in Figure 4(a)-(f) at C_t -, C_z -, C_{zi} -, C_{r1} -, and C_{r2} -carbons, to form free active carbon edge (shown in Figure 5b), provided ΔE_b values of 489.3, 494.2, 465.9, 455.7 and $459.1 \text{ kJ mol}^{-1}$ (Table 3). The average energy barrier ΔE_b and E_{reac} (similar to ΔH) values were respectively 472 and 446 kJ mol^{-1} , for the char structure shown in Figure 5(b). The mean value for the loss of H to form an active carbon site (reactive edge) is similar to that obtained for benzene, and similar to values for smaller and larger PAHs discussed in [113, 114]. These results illustrate that the formation of active sites by loss of H from edges incurs a high ΔE_b value.

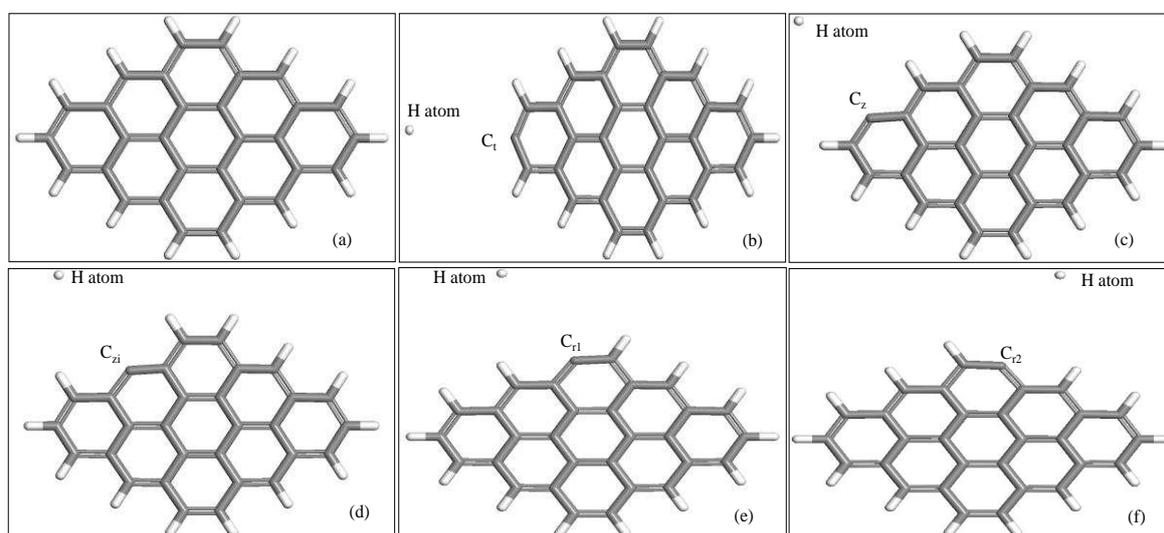


Fig 4. Geometry optimised structures for H-loss from (a) 3x3 char structure at (b)-(f) = C_t -, C_z -, C_{zi} -, C_{r1} and C_{r2} -carbons, respectively.

Table 3. Calculated ΔE_b and E_{reac} values for the loss of H from the benzene and 3x3 char structures.

<u>Benzene molecule</u>		
Edge C positions for H-loss	ΔE_b (kJ mol ⁻¹)	E_{reac} (kJ mol ⁻¹)
Para-	484.8	440.6
meta-	452.2	440.4
ortho-	484.4	443.1
Average	473.8	441.4
<u>3x3 PAH char model</u>		
C _t	489.3	440.6
C _z	494.2	445.6
C _{zi}	464.9	448.6
C _{r1}	454.7	452.6
C _{r2}	459.1	443.0
Average	472.5	446.1

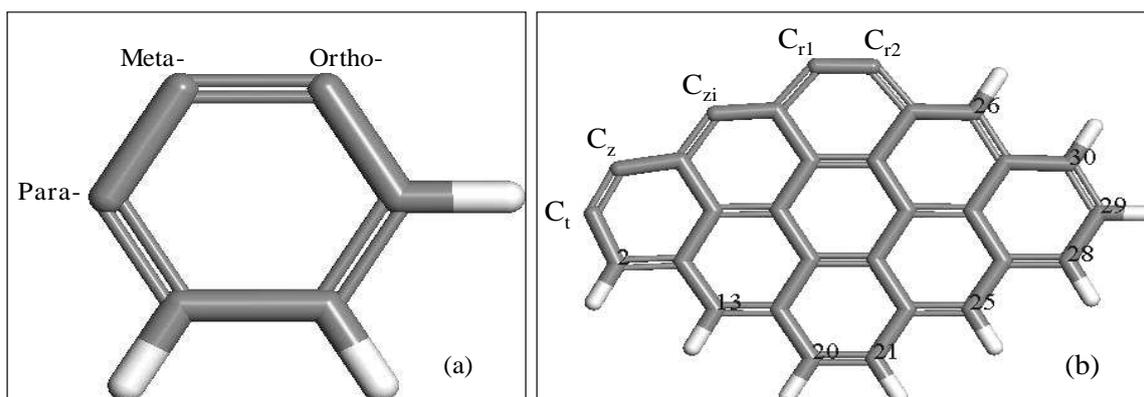


Fig 5. Structures with free edge carbon sites, where (a) = Benzene (left-hand side) and (b) = 3x3 char (right-hand side).

It is reasonable to infer from this that data on electronic states and resulting variations in computed energy values for these small molecules, but question a direct relevance to mechanisms involving coal derived char, since changes in the size of the molecular model and number of active sites, would lead to changes in electronic states. For example, it may be speculated that stacking of sheets, as found in graphite, may impact on the gammaa π delocalised electrons, and active sites located above and below edges may also impact on the electronic states and configurations of chemisorbed CO₂. Informed comment on these matters

may be enabled by computations allowing the use of large-scale char molecular models. The current work demonstrates the changes in relative reactivity for various C* active sites, and their utility and limitations of DFT computations inherent when using small graphene molecules, to elucidate details of CO₂ coal-char gasification mechanisms.

3.2.2 CO₂ chemisorption on reactive edge

Chemisorption of CO₂ on the active sites in char has been an important area of research and has been extensively discussed [28, 30, 39, 43, 64, 115-119]. The $f^+(r)$ values were computed for the respective reactive sites to evaluate the likely locations for chemisorption, and to obtain the energy of reaction (E_{reac}) and the energy barrier (ΔE_b) at each of these sites. The DFT computations for the 5x5 char model require excessive computer resources, thus effort was restricted to the 3x3 char model. However, Montoya et al. have indicated that the heats of CO₂ chemisorption reaction remained nearly constant as the model size increased, emphasizing the chemisorption energies depend significantly on nature of the active site [39]. Table 4 shows the $f^+(r)$ values for active sites in the 3x3 char model, and the values follow the sequence: $C_{zi} > C_{r1} > C_t > C_{r2} > C_z$. The $f^+(r)$ values indicate the preferred active site is C_{zi} for CO₂ chemisorption.

Table 4. Distribution of Fukui indices on the 3x3 char model with a reactive edge.

Edge C	C _t	C _z	C _{zi}	C _{r1}	C _{r2}	C2	C13	C20	C21	C25	C26	C28	C29	C30
Mulliken $f^+(r)$	0.075	0.027	0.266	0.087	0.029	0.017	0.005	0.005	0.011	0.021	0.032	0.026	0.008	0.022

The reaction pathways for gasification of coal chars with CO₂ gas are based on the gasification reaction mechanism shown in Equations (3)-(5). Figure 6 shows the DFT-DNP geometry optimised structures for (a) the 3x3 coal char with a reactive edge, (b) the CO₂ gas molecule, (c) the 3x3-coal char structure with CO₂ gas molecule approaching at 0.5 nm distance from the C_{zi} active site and (d) the CO₂ chemisorption at site C_{zi} . Structural configurations shown in Figures 6(e)-(h) depict the CO₂ chemisorption reaction at sites C_{r1} , C_t , C_z , and C_{r2} .

The MEP includes a description of the reaction and provides the energy barrier (activation energy) and the minimum energy path between a pair of stable states [74]. A single imaginary

frequency corresponds to a point somewhere on the saddle in the potential energy scan of a given structure consistent with the TS. In geometry optimisations, there is no guarantee that the optimised structures will be at a true energy minimum. Thus, finding a thermodynamically stable structure of a molecule (without soft modes) is a major topic in structural chemistry [120]. The values for E_{reac} and ΔE_{b} were computed relative to the initial ground state structure of the CO_2 adsorbed on the respective active site. The results in Table 5 indicate that these are energetically favoured, with E_{reac} value consistent with CO_2 chemisorbed on the active sites, and low ΔE_{b} energy barriers, with a mean E_{reac} value of $-11.9 \text{ kJ mol}^{-1}$ and mean ΔE_{b} of 3.4 kJ mol^{-1} respectively. The lowest ΔE_{b} for the chemisorption of CO_2 gas molecule at C_{zi} active site is indicative of a nearly barrierless reaction and shows a good correlation with the highest $f^{\dagger}(r)$ value reported in Table 4. The values are for the particular molecular model containing all of the active sites stipulated but are consistent with reports of chemisorption of CO_2 on various carbonaceous models as an exothermic process associated with low adsorption energies [28, 43, 121-123].

Table 5. Simulation energetics for CO_2 -char chemisorption at the reactive edge.

Reactive edge (active C* sites)	ΔE_{b} (kJ mol^{-1})	E_{reac} (kJ mol^{-1})
C_{zi}	1.3	-8.4
$\text{C}_{\text{r1}}^{\ddagger}$	-	-
C_{t}	4.1	-19.2
$\text{C}_{\text{r2}}^{\ddagger}$	-	-
C_{z}	4.7	-8.2
<u>Average reactive C edge</u>	<u>3.4</u>	<u>-11.9</u>

\ddagger = Unable to find potential energy maximum for transition state.

These values are in contrast to high values reported for perfect and defective basal carbon planes by Cabrera-Sanfeliix [124] and further illustrate the variations that may be obtained in these studies. The optimized configurations for CO_2 at sites $-\text{C}_{\text{zi}}$, $-\text{C}_{\text{r1}}$, $-\text{C}_{\text{t}}$, and $-\text{C}_{\text{z}}$ show approximately similar bond bending behaviour of the chemisorbed CO_2 in Figures 6(d)-(f)

and (h), highlighting that the mechanism of CO₂ chemisorption obtained from the zigzag models is not significantly different from that obtained from armchair models [111, 119].

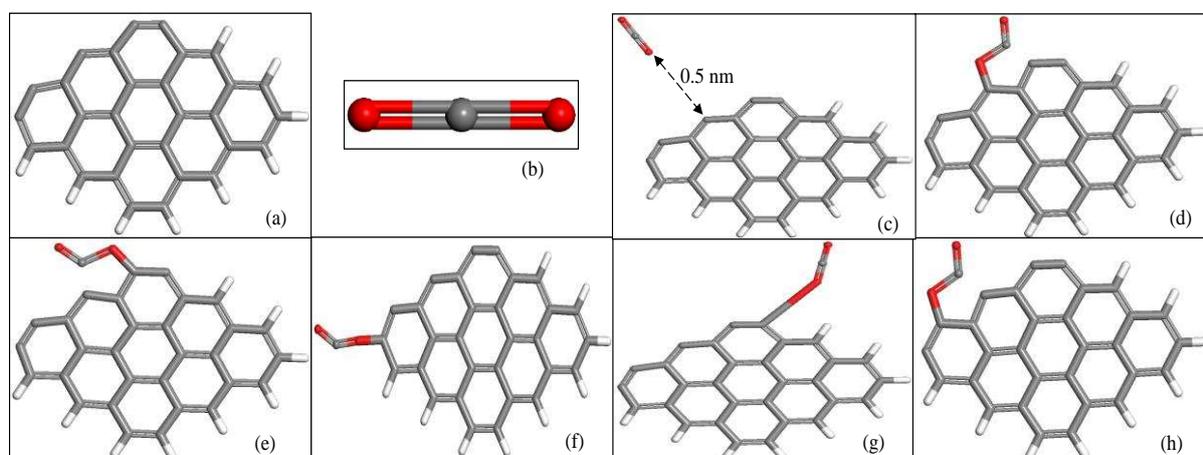


Fig 6. Various optimised structures for the CO₂ chemisorption on single active sites in the reactive C* edge. Top left is 3x3 char model with reactive edge (a). Top right is CO₂ gas model (b), (c) is the 3x3 char model with CO₂ approaching and (d)-(h) are the chemisorption structures. Red = oxygen atoms.

These E_{reac} values for the CO₂ chemisorption from the current work differ from the reported energies ranging from -355.6 to -25.1 kJ mol⁻¹ at B3LYP/6-31G(d) DFT level of theory for the different char models found by Montoya et al. [39]. The difference in these computed adsorption energies is more likely due to the physical models [63]. The results indicate CO₂ chemisorption is favoured on the C_{zi} site although the numerical ΔE_{b} and E_{reac} values may differ for different models.

Examination of bond distances on the structures of CO₂ chemisorption at C_{zi} shown in Figure 7a indicated that, upon chemisorption, the CO₂ gas molecule assumed a zigzag shape towards C_{r1} and the O-C-O bond lengths of the CO₂ gas molecule were elongated from less than 1.2Å to 1.5Å between the chemisorbed O atom and C, and more than 1.2Å between the C and the other O atom, while the CO₂-char bond length was 1.4Å. The chemisorbed CO₂ molecule also bent towards C_{r1} at a distance of 1.5Å. Analyses of these distances with respect to formation and breakage of bonds show that C_{zi} is an ideal site for the formation of a quinone char(O)-complex and CO gas molecule from CO₂ dissociation. Similar inferences could be made with CO₂ chemisorption at C_{r1} and C_z in Figure 7(b) and (d), respectively. The ground state interactions of C_t- and C_{r2}-CO₂ respectively show weaker CO₂ chemisorption, but may be suitable for physisorption processes. The E_{reac} and ΔE_{b} values were not obtained for the

CO₂-C_{r1} and CO₂-C_{r2} chemisorption, because of the occurrences of multiple imaginary frequencies (soft modes) in the geometry optimised reactants and products structures, indicative of an absence of TS. Similar observations were reported by Radovic [43].

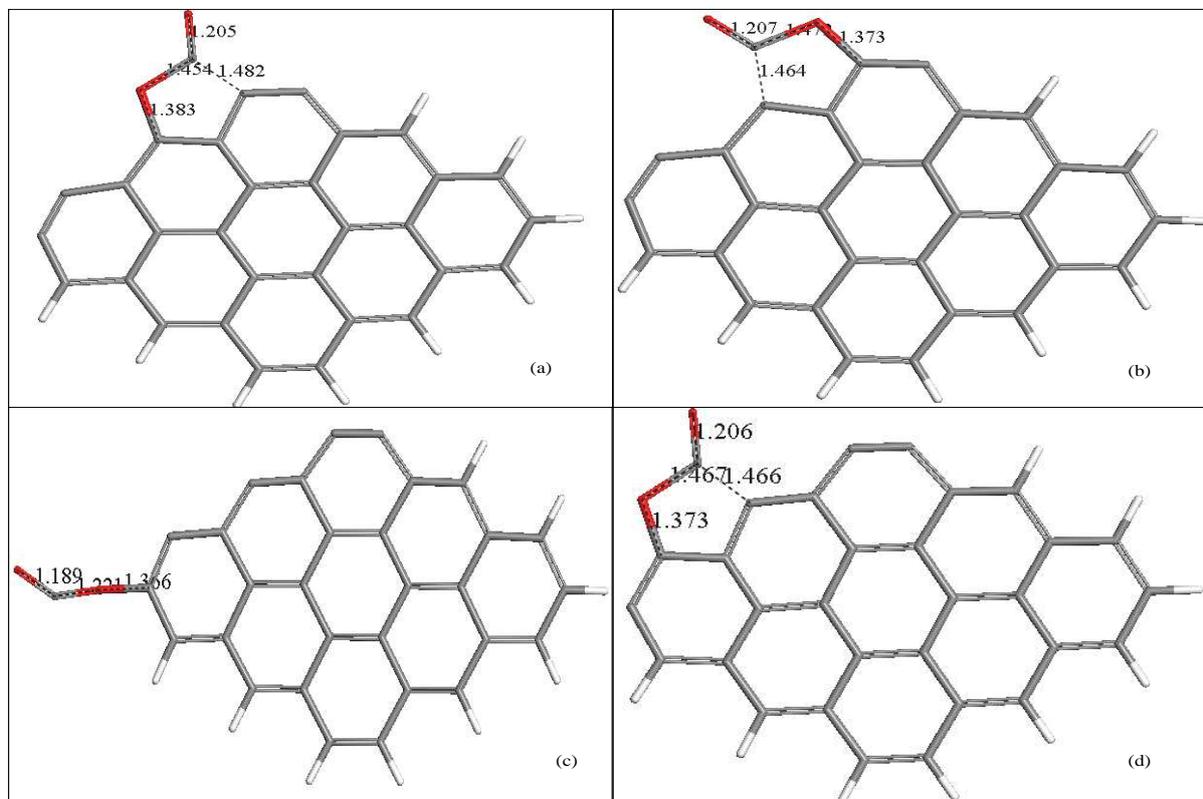


Fig 7. Bond bending and elongations of geometry optimised CO₂-C* chemisorption structures, where: red = oxygen atoms.

The change in geometry of the CO₂ chemisorbed molecule in to bent configuration, and resulting changes in the bond lengths, is indicative of additional changes in the electron configuration of the char-CO₂ molecule, and would further complicate the modeling results. Further work is required for a detailed examination of these details and to demonstrate a correlation with experimental data on char gasification.

3.2.3 The loss of first CO gas molecule through dissociation of adsorbed CO₂ molecule

In this investigation, adsorption of CO₂ on an active site is followed by the loss of CO through the dissociation of the adsorbed CO₂. This mechanism is the breakage of the char(O--CO) bond. The E_{reac} value for loss of the first CO was obtained as the difference between

the char(CO₂) structure, and the final products of char(O) plus free CO, as shown in Figure 8(a)-(b). The reaction energies for the loss of CO from the respective active sites presented in Table 6 were endothermic. The values of 576-588 kJ mol⁻¹ for C_{r1}, C_t and C_{r2} compare fairly well with CO₂ dissociation energy of 526 kJ mol⁻¹ in [125]. The respective values for E_{reac} and ΔE_{b} of 200.6 and 199.2 kJ mol⁻¹ for CO₂ chemisorbed on the active site C_{zi} are lower, and nearly the same. It can be argued that these values are consistent with an endothermic thermal breakdown of the C—(CO) bond in the chemisorbed CO₂ molecule, but the range of values for all of the active sites is unusual and may result from modelling using a number of C* active sites.

Table 6. Computed E_{reac} (energy of reaction) and ΔE_{b} (energy barrier) for loss of the first CO gas molecule.

Reactive edge (active C* sites)	ΔE_{b} (kJ mol ⁻¹)	E_{reac} (kJ mol ⁻¹)
C _{zi}	199.2	200.6
C _{r1}	587.9	142.1
C _t	575.7	155.3
C _{r2}	575.8	132.5
C _z	185.8	168.3
Average reactive C edge	424.9	159.8

The current approach has been to obtain an average value. However it is noted that the energy barrier of 193 kJ mol⁻¹ on the zigzag carbon edge, comprising C_{zi} and C_z is in good agreement with CO desorption activation energy of 130-205 kJ mol⁻¹ reported by Montoya et al. [63].

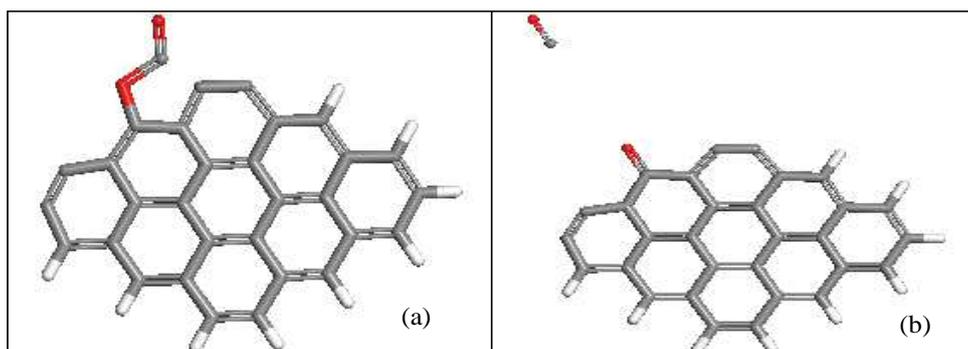


Fig 8. Ground state reactants (a), (left-hand side) and products (b), (right-hand side) structures used for the computation of CO₂ dissociation on C_{zi} active site.

When the transition state is sufficiently refined, the computational procedure will provide v_{TS} consistent with a TS complex. In view of the modelling limitations involving char molecular systems, the available computational capabilities were used to obtain values for E_{reac} and ΔE_b only. For the current purpose, the E_{reac} value is similar to ΔH values obtained by considering the reactants and products of each step in the mechanisms considered. The most likely transition state for CO₂ dissociation involves the C_{zi} site as shown in Figure 9(a). The CO molecule is situated at a distance of 2.510Å from the quinone oxygen complex, and 2.606Å from the C* atom on the adjacent aromatic ring.

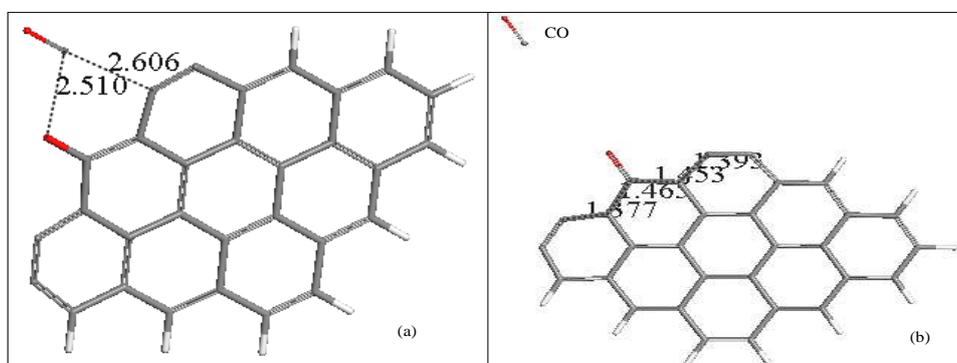


Fig 9. Transition state structure for the CO₂ dissociation at C_{zi} (a) (right-hand side). Ground state char(C=O) quinone structure and the loss of CO. Numbers = bond lengths in Å.

3.2.4 Formation of the second CO gas molecule

The chemistry considered for loss of the second CO is the thermal decomposition of the six membered carbon ring at elevated temperatures and the result is the loss of gaseous CO and formation of either a five member ring in the char substrate, or loss of CO and formation of additional active C* sites for further chemisorption of additional CO₂ molecules. The six-

member carbon ring containing the armchair and tip active sites can also form stable four- and three-membered carbon rings up on the decomposition reactions [24]. The bond lengths of the 6-member carbon ring from the geometry optimised 3x3-coal char structure are shown in Figure 9(b). Montoya et al., also investigated the loss of CO from semiquinone carbon-oxygen species in carbonaceous surfaces using ab initio/DFT (B3LYP/6-31G(d)) levels of theory [38].

Numerous computations were performed to find transition state structures for the decomposition reactions. The decomposition of the 6-member carbon ring yields a CO ($C_{zi}O$) gas molecule. However, the formation of a five-member carbon ring is considered unlikely at C_{zi} site. Instead, the formation of more active C^* sites are likely [39]) and any other ring formation is unlikely because of the long C---C distance in the resulting structure, as shown in Figure 10(a).

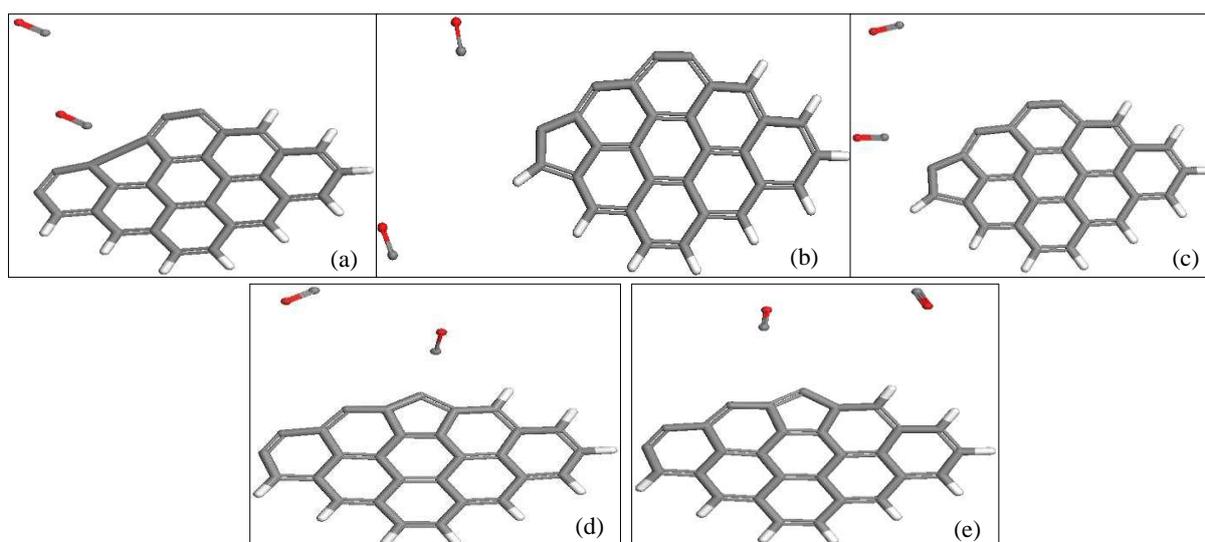


Fig 10. Transition state structures for the decomposition reactions at respective active sites (i.e. reactive edge), where, (a)-(e) = results at C_{zi} , C_t , C_z , C_{r1} and C_{r2} , respectively. Red = oxygen atoms.

The decomposition reaction energetics presented in Table 7, show both the high-energy barrier, ΔE_b , of $587.4 \text{ kJ mol}^{-1}$ to the formation of the second CO from C_{zi} site and endothermic reaction energy E_{reac} of $187.9 \text{ kJ mol}^{-1}$. Loss of the second CO at C_{zi} site would require relatively higher temperatures. The high activation energy suggests that this type of semiquinone complex has an important effect on the rate-limiting step in the char gasification

mechanism and, in essence, defines gasification [119, 126]. The energy barriers ΔE_b obtained for the C_t , C_z , C_{r1} and C_{r2} sites of 68.4, 136.0, 189.5 and 183.5 kJ mol^{-1} , respectively, are substantially lower than ΔE_b for the C_{zi} site. Activation energy of approximately 360 kJ mol^{-1} was reported by of Chen and Yang [126] and approximately 340 kJ mol^{-1} was found by Sendt and Haynes [42], and values for similar models were 83-369 kJ mol^{-1} and 142-334 kJ mol^{-1} [24, 119]. The mean value for ΔE_b of 233 kJ mol^{-1} is in good agreement with the experimental value of $191 \pm 25 \text{ kJ mol}^{-1}$ and $210 \pm 8 \text{ kJ mol}^{-1}$ for the de-ashed vitrinite- and inertinite-rich chars, respectively (Section 4).

Table 7. Computed ΔE_b (energy barrier) and E_{reac} (energy of reaction) for loss of second CO gas molecule.

Reactive edge (active C* sites)	ΔE_b (kJ mol^{-1})	E_{reac} (kJ mol^{-1})
C_{zi}	587.4	187.9
C_{r1}	189.5	169.5
C_t	68.4	64.5
C_{r2}	183.5	173.5
C_z	136.0	139.9
Average reactive C edge	233.0	147.0

The TS structures for the C_t , C_z , C_{r1} and C_{r2} sites, which underwent decomposition of six-membered rings provided five-membered rings with conventional C-C bond lengths. The respective distances between the newly formed CO and the char varied (4.3Å, 2.5Å, 2.8Å and 3.8Å). These results are consistent with the ΔE_b values and indicative of the ease with which CO is lost from the various sites.

The surface transformations for the formation of both the first CO gas molecule from the CO₂ dissociation reactions and the second CO from the decomposition reactions were accompanied by large endothermic energies. These endothermic reactions generally validate the bond-breaking occurrences for the yield of CO gas molecules. On the other hand, the

exothermic energies experienced by the CO₂ chemisorption reactions are generally associated with bond-forming rather than bond-breaking events. However, a mean activation energy ΔE_b for the CO₂ dissociation reactions is significantly larger than that for both the chemisorption and decomposition reactions. Inferences can be made from here that the CO₂ chemisorption was not always accompanied by dissociation and may nearly be impossible to experience dissociation at certain active sites during gasification. Zhu et al. also made similar observations [119]. There is, therefore, a likelihood that in the event that the dissociation of CO₂ did not take place, the gas remained chemisorbed on active site and simply bounced back (desorped).

The major focus of the current investigation has been to compare computer molecular modeling results obtained using a very simple molecular model with experimental results from char gasification. In this way the likely mechanisms for char gasification was identified.

4. Experimental reactivity and parameters

The conversion versus-time results for the char-CO₂ gasification reactions are shown in Figure 11, together with the predictions from the random pore model, Equation (7). The results are consistent with increasing conversion with increasing temperature and there is good agreement between the experimental results and the RPM prediction with the regressed parameters (structural parameter and lumped reaction parameters), in Equation 7. The structural parameters were found to be 2.08 and 1.08 for the inertinite- and vitrinite-rich chars, respectively, which depend on the initial properties of the chars, Equation (8). Since the limiting value for pore growth according to the random pore model is greater than a value of 2 [6, 101, 102, 106, 127] both chars can be considered to have experienced mainly pore growth and pore coalescence during the initial stage of the char-CO₂ gasification reaction. The determined values for the lumped reaction rates as a function of temperature, Figure 12, show that the vitrinite-rich char is more reactive than the inertinite-rich char. This difference is attributed to the effect of the char physical properties, Equation (8), in addition to the intrinsic reaction rate, which needs to be examined.

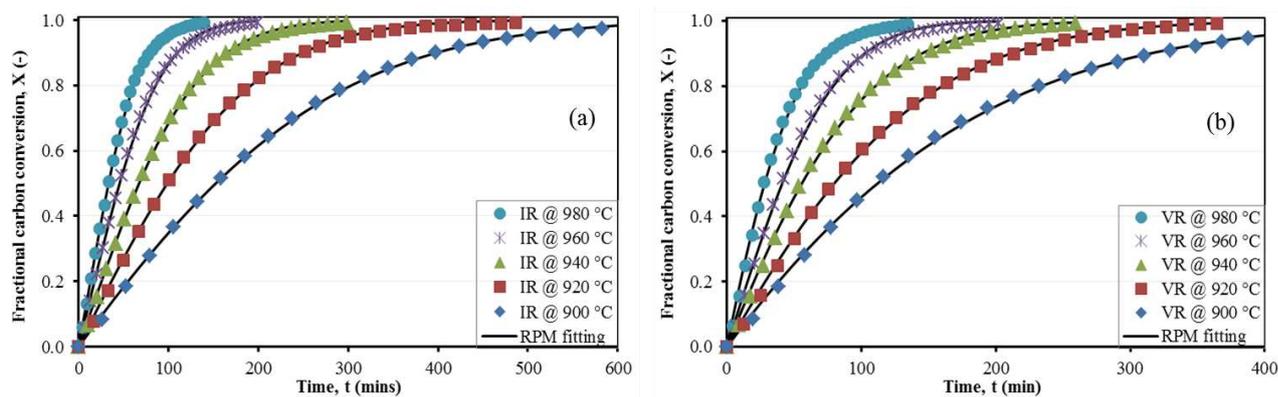


Fig 11. Conversion-time plot for inertinite- and vitrinite-rich chars, where, (a), left-hand side = inertinite-rich char and (b), right-hand side = vitrinite-rich chars.

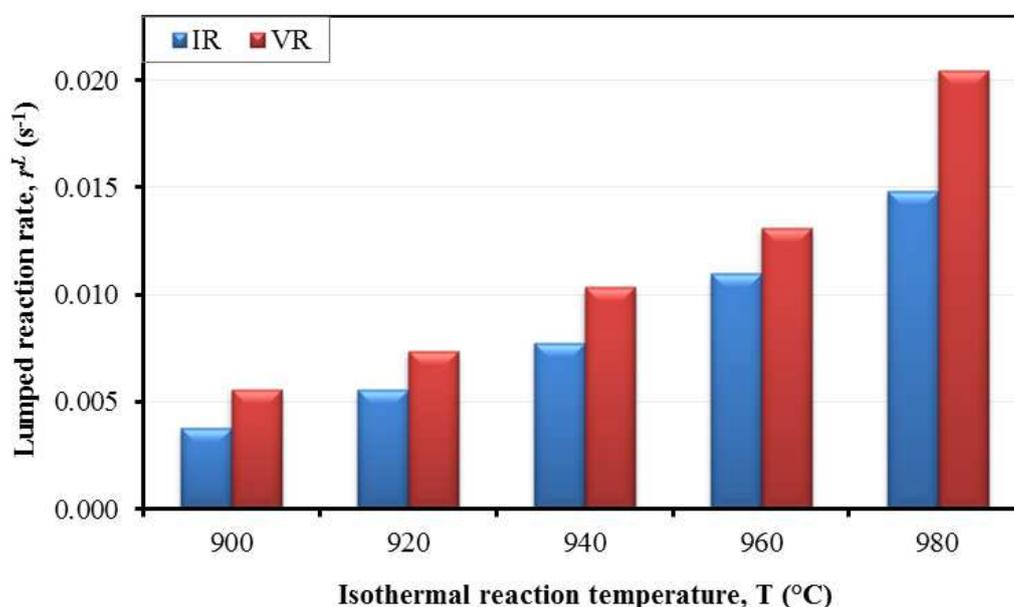


Fig 12. Influence of isothermal reaction temperatures on the lumped reaction rates of the chars.

4.1 Validation of activation energy

The experimental results consisting of the lumped reaction rates derived with the application of the random pore model, given in Figure 12, was used to validate the activation energy obtained by the DFT-DNP molecular modelling technique. The result for the two chars are shown in Figure 13. The calculation consisted of plotting the results on a graph with semi logarithmic coordinates ($\ln r^L$ versus $1/T$) and comparing the results obtained from Equation (10), using the activation energy of

233 kJ mol⁻¹ reported in Section 3.2.4 together with suitable values for r^L in Equation (12), to ensure the best fit according to linear regression procedure. The agreement was found to be very good within the limits of experimental uncertainty at 95% confidence interval, as shown with the estimated activation energies in Figure 15.

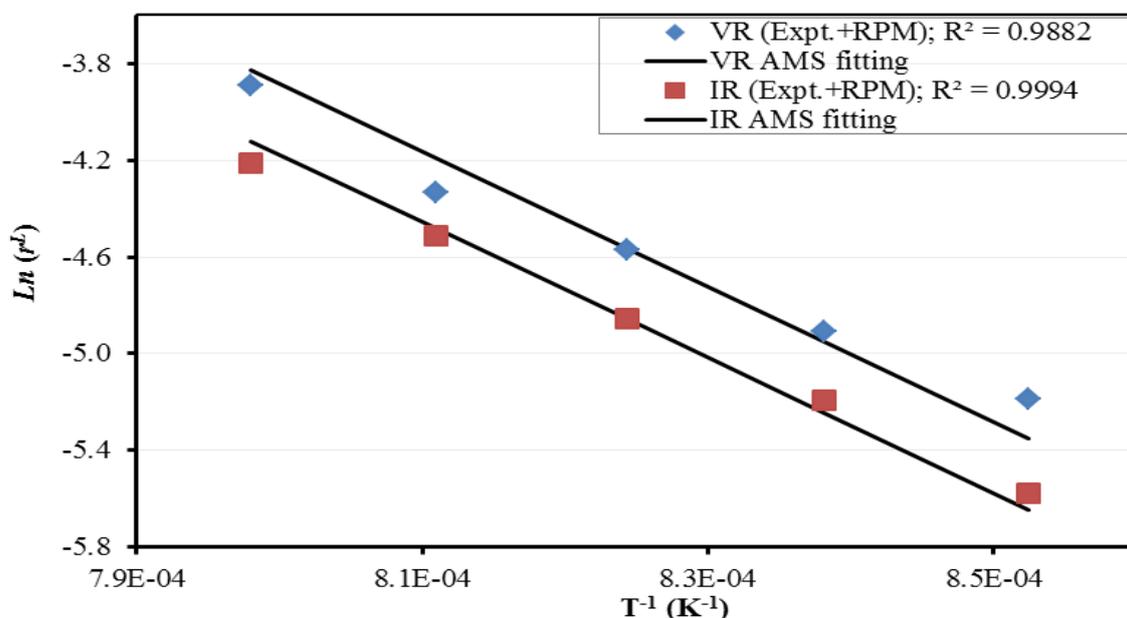


Fig 13. Validation of activation energy calculated according to atomistic reaction kinetics with experimental results obtained from thermogravimetric experimentation.

The results obtained from Equation (12) with respect to the lumped reaction rate (experimentally calculated and RPM evaluated) and the activation energy are also in good agreement with the results obtained from the atomistic modelling as shown in Figure 14. The pre-exponential factor was calculated from Equation (12) for CO₂ gas pressure of 875 kPa and using an order of reaction, n of 0.6 [6, 101]. The reaction kinetic parameters for the calculations are shown in Table 9. A comparison between the activation energy obtained from the respective procedures is presented in Figure 15.

Table 9. Estimated reaction parameters.

	Regressed Results		Atomistic modelling results (Activation energy validation)	
	IR	VR	IR	VR
E_a (kJ/mol)	210 ± 8	191 ± 25	233	233
k'_{S_0}	1.3E+05	9.7E+04	8.4E+07	1.1E+08

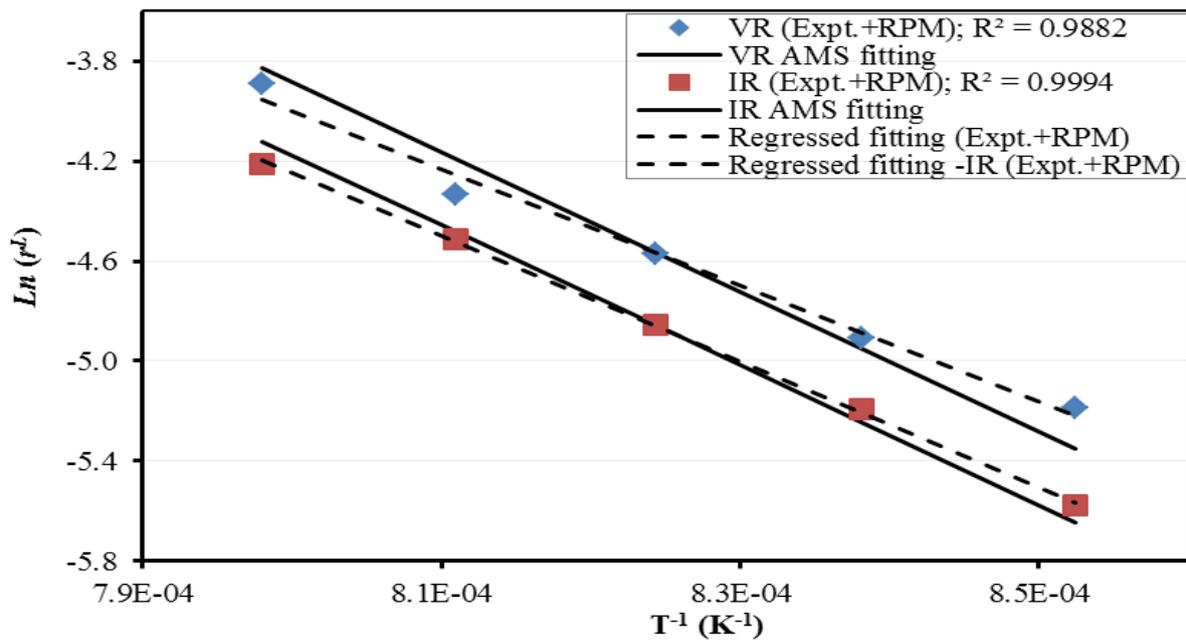


Fig 14. Comparison of experimental and model predictions.

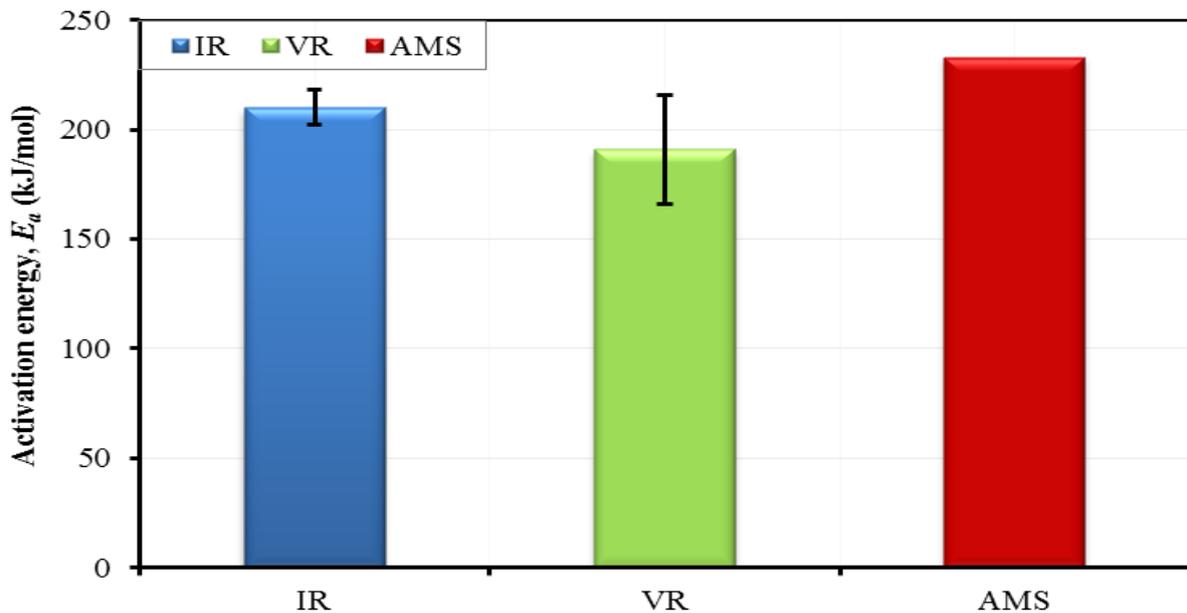


Fig 15. Comparison of activation energies.

4.2 Mechanisms of char gasification by CO_2

The CO_2 gasification occurs in regime I, which is chemically controlled. The salient features are the formation and accessibility of active sites, which in turn depends mostly on the

surface area formed during pyrolysis and formation of each char. The ΔE_b value for loss of H from C-H edges to form C* active sites was calculated at a mean value of 472.5 kJ mol⁻¹. This value is higher than the computed mean barrier ΔE_b for dissociation of CO₂ to yield the first CO molecule (424.9 kJ mol⁻¹) for all active sites comprising the reactive edge. It is also higher than the mean value for loss of the second CO, ΔE_b of 233 kJ mol⁻¹. The experimental values E_a , was 191 ± 25 kJ mol⁻¹ and 210 ± 8 kJ mol⁻¹ for the de-ashed vitrinite- and inertinite-rich chars, respectively. The experimental results are in agreement with the computed values, and support mechanism that commences with the presence of active sites on the char, formed during pyrolysis and char formation. If we assume that active sites may form purely through thermal decomposition of edge C-H groups to yield C* active sites, the major step in the mechanism of CO₂ gasification would be the formation of active sites in char (and a larger value of experimentally determined activation energy). The experimental value also may indicate the likely sites for the thermal decomposition of CO₂ to yield the first CO at sites Czi and Cz (Table 6). The difficulties associated with a small graphene molecule to model coal char however, are amply demonstrated in the present work. The data supports a mechanism for CO₂ gasification of coal-char begins when active sites are liberated as the coal is heated to high temperatures. Therefore, the following descriptive mechanism is proposed and presented in Figure 16:

- A. Chemisorption of CO₂ to active sites in char: The various C* sites located at a reactive edge may be classified as weak and strong. (This should not be confused with the strong and weak adsorption sites for CO₂ reported by Xu et al [109], related to the presence of inorganics in char). Strong adsorption is related to the formation of five member rings between two adjacent C* sites and adsorbed CO₂, whereas weakly adsorbed CO₂ occurs at a single C* site.
- B. Yield of the first CO molecule: Two structures may form; one is a stable chemisorbed CO₂ molecule which encounters a lower energy barrier to dissociation, and the other a linear CO₂ molecule adsorbed on a single C* sites, which encounters a higher energy barrier to dissociation.
- C. Yield of the second CO: A number of quinone-type char(CO) structures form after loss of the first CO. Those structures may decompose to form CO and relatively stable

five member carbon rings in char; however where stable five-membered rings are not energetically favoured, these do not undergo ring formation but instead form additional C* active sites for further gasification with CO₂. This decomposition reaction is defined as gasification because of its distinction as rate limiting [119, 126].

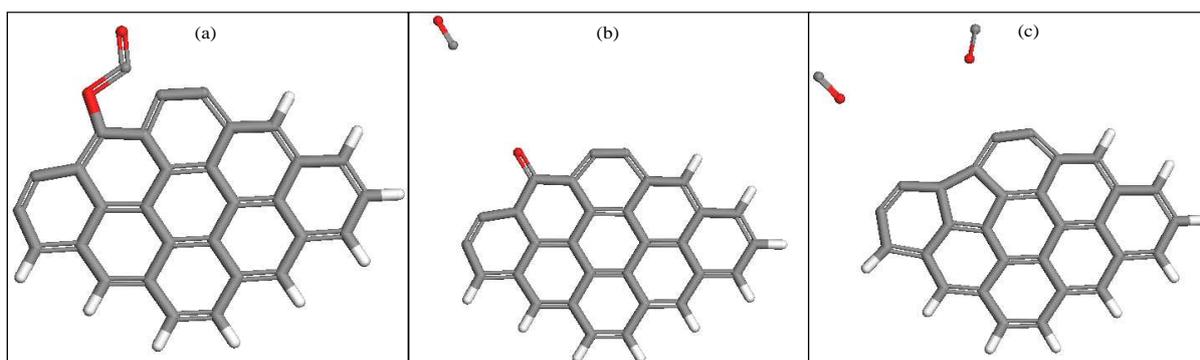


Fig 16. Mechanisms of char gasification by CO₂, where (a) = chemisorption of CO₂ to active sites in char, (b) = yield of first CO and (c) = yield of second CO.

5. Conclusions

The DFT modelling of CO₂ char gasification for the de-ashed vitrinite- and inertinite-rich chars provides details of the nature of reactive sites in chars structures, and also demonstrates the limitations inherent in using small molecular models in these studies. A combination of computer modelling and experimental techniques however, elucidated the diverse nature of active sites in these structures, the various mechanisms for chemisorption of CO₂ on active sites, dissociation of CO₂ into CO and O-complex; the loss of O-complex as CO and resulting char structures. The Fukui function $f^+(r)$ was a useful indicator of the preferred site for chemisorption of CO₂. The mean reactivity of H-saturated edge carbons decreased with increasing char molecules and was distributed according to the sequence: zigzag > armchair > tip active sites. In particular, the reactivity of the zigzag active sites decreased with increasing distances from the tip carbons. The average computed energy barrier of 233 kJ mol⁻¹ compares well with the experimentally measured activation energy of 191 ± 25 kJ mol⁻¹ and 210 ± 8 kJ mol⁻¹ for the de-ashed vitrinite- and inertinite-rich chars, respectively. Consequently, the results show a reasonable agreement between DFT modelling outcomes and experimental results. Further work is required using a larger molecular model (provided

sufficient computer resources are available) to provide additional insights on the mechanisms of coal-char gasification.

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