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eprints@whiterose.ac.uk https://eprints.whiterose.ac.uk/ Band gaps of wurtzite Sc_xGa_{1-x}N alloys

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Abstract

Optical transmittance measurements on epitaxial, phase-pure, wurtzite-structure $Sc_xGa_{1-x}N$ films with $0 \le x \le 0.26$ showed that their direct optical band gaps increased from 3.33 eV to 3.89 eV with increasing x, in agreement with theory. These films contained I₁- and I₂-type stacking faults. However, the direct optical band gaps decreased from 3.37 eV to 3.26 eV for $Sc_xGa_{1-x}N$ films which additionally contained nanoscale lamellar inclusions of the zincblende phase, as revealed by aberration-corrected scanning transmission electron microscopy. Therefore we conclude that the apparent reduction in $Sc_xGa_{1-x}N$ band gaps with increasing x is an artefact resulting from the presence of nanoscale zinc-blende inclusions.

Keywords: ScGaN, band gap, aberration-corrected STEM-HAADF, molecular beam epitaxy

PACS codes: 78.66.Fd, 61.66.Dk, 81.15.Hi

The wurtzite-structure III-nitrides AlN, GaN and InN are widely used in optoelectronic and high-power electronic applications, including high mobility electron transistors, energy harvesters, laser diodes and light emitting diodes (LEDs). Alloying nitride semiconductors enables band gap engineering, which allows the range of emission wavelengths in optoelectronic devices to be tuned, e.g. within the ultraviolet (UV) and green emission regions for $Al_xGa_{1-x}N^1$ and $In_xGa_{1-x}N^2$ respectively. However, devices made using these materials can suffer from poor efficiencies, especially in the UV region. This is attributed to difficulties in p-type doping these wide band gap materials, along with lattice parameter mismatches between different layers in the epitaxial devices, leading to the buildup of inplane stress, film cracking and high dislocation densities³. These problems are related to the fundamental electronic structure of $Al_xGa_{1-x}N$ semiconductors and to the relationship between their band gaps and lattice parameters. Consequently, there is considerable motivation to find alternative wide band gap semiconductors for use in UV optoelectronics, which have different electronic structures and different relationships between their lattice parameters and direct band gaps. Alloying GaN with ScN offers interesting possibilities in this regard⁴.

ScN is a rock-salt structure semiconductor with a direct band gap of 2.1 eV and an indirect band gap of 0.9 eV⁵⁻⁷, which can be grown easily by molecular beam epitaxy⁸ and which has already been used in as a dislocation reduction interlayer in GaN films⁹ and is of interest for thermoelectric applications^{8,10,11}. The phase stability, structural and optical properties of ScN and Sc_xGa_{1-x}N have been calculated previously. Farrer and Bellaiche predicted that ScN is unstable in the wurtzite structure but metastable in a h-BN-like non-polar structure¹². All studies have predicted that Sc_xGa_{1-x}N alloys will be metastable with respect to rock-salt

structure ScN and wurtzite-structure GaN. However, in practice Sc_xGa_{1-x}N alloys can be grown in epitaxial thin film form under a range of growth conditions^{13–15}. Recent high-quality calculations also concluded that (0001)-oriented Sc_xGa_{1-x}N films can be stabilised using inplane compressive epitaxial strain (e.g. by growing on top of GaN layers) and/or by using non-equilibrium growth conditions. In that case, $Sc_xGa_{1-x}N$ alloys are expected to retain the wurtzite structure up to an intermediate Sc content of around $x = 0.66^{16}$. However, there is controversy over the band gaps of Sc_xGa_{1-x}N alloys in the literature. Constantin et al. predicted that the band gaps of Sc_xGa_{1-x}N should decrease as the Sc content increases¹³, however, their calculations assumed that the alloys retained an undistorted wurtzite structure throughout, which does not match lattice parameter data obtained from experiment¹⁷ or from later theoretical calculations¹⁶. Farrer et al. also predicted that the band gap of Sc_xGa_{1-x}N should decrease from 3.5 eV (the band gap of GaN) to 1.55 eV (the band gap of h-BN-like ScN) as the Sc concentration increases¹². On the other hand, Zhang et al. predicted that the direct band gaps should increase from 3.5 eV (GaN) to 4.36 eV (Sc_{0.5}Ga_{0.5}N) as the Sc content increases, and will then decrease to 1.5 eV as the Sc content increases further and a phase change to the cubic structure occurs. Experimental results from Little et al.¹⁸ and Constantin et al.¹⁷ showed a decrease in the magnitude of the direct optical band gap as the Sc content increased, however, the films possessed either poor crystal quality¹⁸ and/or an extremely high density of stacking faults¹⁷. Epitaxial (0001)-oriented wurtzite-structure Sc_xGa_{1-x}N films can also contain cubic inclusions (i.e. lamellae with the zinc-blende structure oriented along the (0001) plane¹⁹), which are difficult to distinguish from stacking faults in high-resolution TEM images. All of these may lead to sub-gap absorption and misleading conclusions about the trend in band gaps as a function of Sc content. Therefore, this study aims to understand the influence of film microstructure on the direct optical band gaps of epitaxial Sc_xGa_{1-x}N films and thereby to determine which of the predicted trends in band gap versus composition is correct.

Epitaxial wurtzite-structure (0001)-oriented Sc_xGa_{1-x}N films were grown using molecular beam epitaxy (MBE) with an N₂ plasma source, under metal-rich growth conditions. Three different kinds of buffer layers were used to influence the $Sc_xGa_{1-x}N$ film strain state and microstructure: (0001)-oriented GaN grown by molecular beam epitaxy on sapphire, (0001)oriented GaN grown by metal-organic vapour phase epitaxy (MOVPE) and (0001)-oriented AlN grown by metal-organic vapour phase epitaxy. The Sc_xGa_{1-x}N film compositions were controlled by varying the Sc flux while keeping the Ga flux constant and the N₂ flow rate constant values which produce a GaN growth rate of approximately 260 nm hr⁻¹. Film compositions were determined using Rutherford backscattering (RBS). RBS measurements were performed using a beam of ⁴He at 2 MeV with an incidence angle of 0°. A standard detector was placed at 140° and two pin-diode detectors located symmetrically to each other at 165°. The RBS data were analysed using the IBA DataFurnace NDF v9.6d²⁰. Optical transmittance measurements were carried out using an Agilent Technologies Cary 5000 UV-Vis-NIR spectrophotometer at room temperature, using a bare sapphire substrate as a reference. To minimise the unintentional incorporation of impurities, the MBE chamber was operated at a low base pressure of 10⁻¹⁰ mbar and the scandium metal was custom-prepared for high purity (Sc 99.999%) and no detectable fluorine.

Transmission electron microscopy (TEM) analysis was carried out using a JEOL 2100 at 200 kV. Cross-sectional TEM samples were prepared by mechanical grinding followed by ion polishing until electron transparency was reached. Aberration-corrected high angle annular dark field (HAADF) imaging in the scanning transmission electron microscopy (STEM) mode was performed on the Titan³ 80-300 at 300 kV with a probe convergence semi-angle of

15 mrad. Spherical aberrations up to the third order in the electron beam "probe" were corrected by recording Zemlin tableau diffractograms.

Figure 1 shows the direct optical band gaps of Sc_xGa_{1-x}N grown on different buffer layers. For Sc_xGa_{1-x}N grown on MBE GaN buffer layers, the band gap decreased as the Sc content increased. The same trend was found previously by Little et al.¹⁸ and Constantin et al.¹⁷ for Sc_xGa_{1-x}N films with the poor crystal quality or high densities of stacking faults. In contrast, measurements from Sc_xGa_{1-x}N films on MOVPE GaN are limited by the lower band gap of the GaN buffer layer (3.4 eV at room temperature²¹), indicating the band gaps of the Sc_xGa_1 . _xN films are higher than GaN. However, AlN has a very wide band gap of 6.2 eV at room temperature, so the band gaps of Sc_xGa_{1-x}N can be revealed in optical absorption measurements. Figure 1 shows clearly that the band gaps of Sc_xGa_{1-x}N films grown on AlN increase with increasing Sc content, in agreement with recent high-quality theoretical predictions¹⁶. However, the measured band gaps of Sc_xGa_{1-x}N are approximately 0.1 eV lower than the predicted values. This arises because in the previous theoretical study, the predicted band gap of GaN was corrected to 3.5 eV, i.e. the value it takes a temperature close to 0 K, which is appropriate for comparison to calculated data²². In contrast, if the theoretical data are corrected with respect to the room-temperature band gap of GaN of 3.4 eV^{21} , then the theoretical and experimental data match very well. Importantly, reference samples of ScN and GaN prepared in the same MBE reactor under comparable conditions had band gaps of 2.1 eV and 3.4 eV respectively, which are the literature values for pure films. This indicates minimal oxygen contamination and suggests that the films are stoichiometric, as impurities and vacancies are known to affect the band gap of $ScN^{23,24}$.

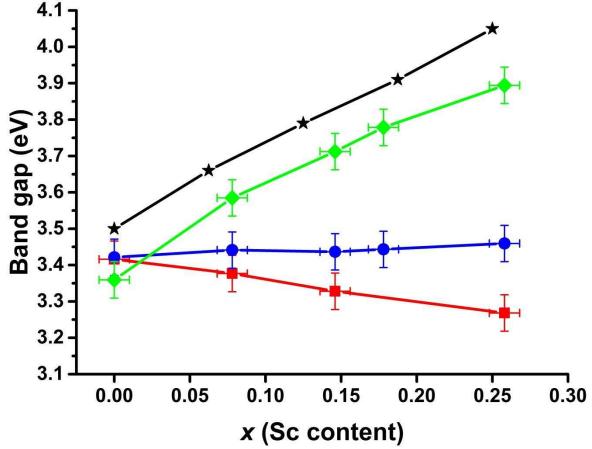


FIG. 1. Band gaps of $Sc_xGa_{1-x}N$ with $0 \le x \le 0.26$, grown on different buffer layers. Red square: $Sc_xGa_{1-x}N$ on MBE GaN; Blue circle: $Sc_xGa_{1-x}N$ on MOVPE GaN; Green diamond:

Sc_xGa_{1-x}N on MOVPE AlN; Black star: calculation results adapted from Ref. 16.

Microstructural analysis was performed to understand the effect of defects on the band gaps of Sc_xGa_{1-x}N. I₁- and I₂-type basal-plane stacking faults (BSFs) were seen in all Sc_xGa_{1-x}N films (Figure 2). This suggests that stacking faults are not responsible for the lower band gaps measured in Sc_xGa_{1-x}N on MBE GaN. On the other hand, aberration-corrected STEM-HAADF images acquired along the $< 11\overline{2}0 >$ zone axis showed cubic stacking (ABCABC) only in Sc_xGa_{1-x}N on MBE GaN (Figure 3), which has been observed previously in Sc_xGa_{1-x}N grown by NH₃-MBE and which can be distinguished from the expected rock-salt phase¹⁹. No significant differences in contrast in the STEM images can be seen between the regions with cubic stacking (these are effectively one large stacking fault) and the rest of the material, indicating minimal compositional differences. However, the zinc blende phase of GaN has a band gap approximately 0.2 eV lower than that of hexagonal GaN^{25,26}, and the band gap of zinc-blende ScN is expected to be lower than that of zinc-blende GaN²⁷, such that the band gap of the inclusions should decrease with increasing Sc content. Therefore, we conclude that inclusions of the zinc blende phase are the cause of the apparent reduction in band gap with increasing Sc content. Importantly, it is very difficult to distinguish nanoscale lamellar inclusions layered along (0001) from BSFs using high-resolution TEM imaging, as used in Ref. 13: hence zinc blende inclusions could have been present in those samples too and could account for the apparent decrease in band gap of Sc_xGa_{1-x}N with increasing Sc content, as reported in previous work.

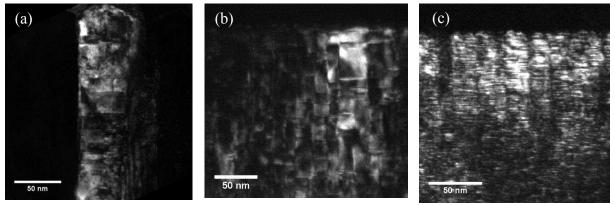


FIG. 2. Dark-field TEM images acquired along the $< 11\overline{2}0 >$ zone axis of Sc_xGa_{1-x}N on different buffer layers, (a) MBE GaN, (b) MOVPE GaN, (c) MOVPE AlN.

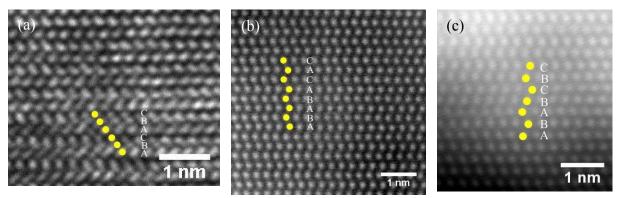


FIG. 3. (a) Cubic stacking, found only in $Sc_xGa_{1-x}N$ on MBE GaN, (b) I_1 and (c) I_2 basalplane stacking faults found in all samples using aberration corrected-STEM with specimens oriented along the $< 11\overline{2}0 > zone$ axis.

In conclusion, the band gaps of $Sc_xGa_{1-x}N$ films were found to increase with increasing Sc content for $Sc_xGa_{1-x}N$ grown on MOVPE GaN and AlN buffer layers, in agreement with recent high-quality calculations¹⁶. In contrast, band gaps of $Sc_xGa_{1-x}N$ grown on MBE GaN decrease with increasing Sc content, consistent with previous experimental reports. Although. I₁– and I₂–type basal-plane stacking faults were seen in all $Sc_xGa_{1-x}N$ films, lamellar inclusions of the cubic phase were found only in $Sc_xGa_{1-x}N$ on MBE GaN, and are believed to produce sub-gap optical absorption. Therefore, $Sc_xGa_{1-x}N$ films may prove useful in UV optoelectronic applications, as long as the buffer layers and growth conditions are selected to minimise the formation of cubic inclusions.

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References

¹A. Khan, K. Balakrishnan and T. Katona, Nat. Photonics **2**, 77 (2008)

²T. Mukai, M. Yamada and S. Nakamura, Jpn. J. Appl. Phys. **38**, 3976 (1999)

³M. Kneissl, T. Kolbe, C. Chua, V. Kueller, N. Lobo, J. Stellmach, A. Knauer, H. Rodriguez, S. Einfeldt, Z. Yang, N. M. Johnson and M. Weyers, Semicond. Sci. Technol. **26** 014036 (2011)

⁴M. A. Moram and S. Zhang, J. Mater. Chem. A, **2**, 6042 (2014)

⁵W. R. L. Lambrecht, Phys. Rev. B **62**, 13538 (2000)

⁶A. R. Smith, H. A. H. Al-Brithen, D. C. Ingram and D. Gall, J. Appl. Phys. **90**, 1809 (2001)

⁷D. Gall, I. Petrov, L. D. Madsen, J.-E. Sundgren and J. E. Greene, J. Vac. Sci. Technol. A **16**, 2411 (1998)

⁸S. Kerdsongpanya, N. V. Nong, N. Pryds, A. Zukauskaite, J. Jensen, J. Birch, J. Lu, L. Hultman, G. Wingqvist and P. Eklund, Appl. Phys. Lett. **99**, 232113 (2011)

⁹M. A. Moram, Y. Zhang, M. J. Kappers, Z. H. Barber and C. J. Humphreys, Appl. Phys. Lett. **91**, 152101 (2007)

¹⁰H. A. Al-Brithen, A. R. Smith and D. Gall, Phys. Rev. B **70**, 045303 (2004)

¹¹P. V. Burmistrova, J. Maassen, T. Favaloro, B. Saha, S. Salamat, Y. R. Koh, M. S. Lundstrom, A. Shakouri and T. D. Sands, J. Appl. Phys. **113**, 153704 (2013)

¹²N. Farrer and L. Bellaiche, Phys. Rev. B **66**, 201203 (2002)

¹³C. Constantin, M. B. Haider, D. Ingram, A. R. Smith, N. Sandler, K. Sun and P. Ordejon, J. Appl. Phys. **98**, 123501 (2005)

¹⁴M. A. Moram, Y. Zhang, T. B. Joyce, D. Holec, P. R. Chalker, P. H. Mayrhofer, M. J. Kappers and C. J. Humphreys, J. Appl. Phys. **106**, 113533 (2009)

¹⁵S. M. Knoll, S. Zhang, T. B. Joyce, M. J. Kappers, C. J. Humphreys and M. A. Moram, Phys. Status Solidi A **209**, 33 (2012)

¹⁶S. Zhang, D. Holec, W. Y. Fu, C. J. Humphreys and M. A. Moram, J. Appl. Phys. **114**, 133510 (2013)

¹⁷C. Constantin, H. Al-Brithen, M. B. Haider, D. Ingram and A. R. Smith, Phys. Rev. B **70**, 193309 (2004)

¹⁸M. E. Little and M. E. Kordesch, Appl. Phys. Lett. **78**, 2891 (2001)

¹⁹S. M. Knoll, S. K. Rhode, S. Zhang, T. B. Joyce and M. A. Moram, Appl. Phys. Lett. **104**, 101906 (2014)

²⁰N. P. Barradas, E. Alves, C. Jeyes and M. Tosaki, Nucl. Instrum. Methods Phys. Res. B **247** 381 (2006)

²¹S. Strite and H. Morkoc, J. Vac. Sci. Technol. B **10**, 1237 (1992)

²²R. Dingle, D. D. Sell, S. E. Stokowski and M. Ilegems, Phys. Rev. B 4, 1211 (1971)

²³S. Kerdsongpanya, B. Alling, and P. Eklund, Phys. Rev. B 86, 195140 (2012)

²⁴M.A. Moram, Z.H. Barber, C.J. Humphreys, Thin Solid Films **516**, 8569 (2008)

²⁵H. Okumura, K. Ohta, K. Ando, W. W. Ruhle, T. Nagatomo and S. Yoshida, Solid-State Electron. **41**, 201 (1997)

²⁶A. Philippe, C. Bru-Chevallier, M. Vernay, G. Guillot, J. Hubner, B. Daudin, G. Feuillet, Mater. Sci. Eng., B **59**, 168 (1999)

²⁷A. Tebboune, D. Rached, A. Benzair, N. Sekkal and A. Belbachir, Phys. Status Solidi B **243**, 2788, (2006)