A technique for quantifying heterogeneous ice nucleation in microlitre supercooled water droplets

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Abstract. In many clouds, the formation of ice requires the presence of particles capable of nucleating ice. Ice-nucleating particles (INPs) are rare in comparison to cloud condensation nuclei. However, the fact that only a small fraction of aerosol particles can nucleate ice means that detection and quantification of INPs is challenging. This is particularly true at temperatures above about −20°C since the population of particles capable of serving as INPs decreases dramatically with increasing temperature. In this paper, we describe an experimental technique in which droplets of microlitre volume containing ice-nucleating material are cooled down at a controlled rate and their freezing temperatures recorded. The advantage of using large droplet volumes is that the surface area per droplet is vastly larger than in experiments focused on single aerosol particles or cloud-sized droplets. This increases the probability of observing the effect of less common, but important, high-temperature INPs and therefore allows the quantification of their ice nucleation efficiency. The potential artefacts which could influence data from this experiment, and other similar experiments, are mitigated and discussed. Experimentally determined heterogeneous ice nucleation efficiencies for K-feldspar (microcline), kaolinite, chlorite, NX-illite, Snomax® CE1 and silver iodide are presented.

1 Introduction

Cloud droplets can supercool to temperatures below −37°C (Rosenfeld and Woodley, 2000) but can freeze at much warmer temperatures in the presence of ice-nucleating particles (INPs). In fact, mixed-phase stratus clouds are observed to glaciate at temperatures warmer than −15°C, but only in certain locations (Kanitz et al., 2011). In addition, ice formation in the ice multiplication regime around −3 to −8°C is critically important in the formation of precipitation from convective clouds (Crawford et al., 2012; Pruppacher and Klett, 1997). However, recent assessments of laboratory ice nucleation data for a range of atmospherically relevant materials concluded that it is unclear which aerosol species trigger freezing above −15°C (Murray et al., 2012; Hoose and Möhler, 2012). Part of the problem is that many quantitative experimental techniques for determining ice nucleation efficiency are not sufficiently sensitive to quantify the efficiency of many nucleants at temperatures warmer than −15°C.

There are a number of instruments types that have been used for research into immersion mode ice nucleation, and reviews on the subject are available (Hoose and Möhler, 2012; Murray et al., 2012). These instruments include cloud chambers (e.g. Jiang et al., 2014; Niemand et al., 2012; Cotton et al., 2007), continuous-flow diffusion chambers (CFDCs; e.g. Salam et al., 2006; Rogers et al., 2001) and a wide variety of droplet-freezing experiments (e.g. Knopf and Alpert, 2013; Vali, 2008; Murray et al., 2011; Budke and Koop, 2015). Cloud chambers and CFDCs quantify the ice
nucleation ability of a dispersion of aerosol particles as a function of relative humidity and temperature. In contrast, droplet-freezing experiments tend to have multiple particles suspended in individual water droplets. A common way of quantifying the ice nucleation efficiency of a material is using the ice-active surface site density, \( n_s \), which is the cumulative number of nucleation sites per unit surface area of nucleant that become active on cooling from 0 \( ^\circ \)C to a temperature \( T \) (Connolly et al., 2009).

\[
\frac{n(T)}{N} = 1 - \exp(-n_s(T)A),
\]

(1)

where \( n(T) \) is the number of droplets frozen at temperature \( T \), \( N \) is the total number of droplets in the experiment and \( A \) is the surface area of nucleant per droplet. In using this approximation it is assumed that the time dependence of nucleation arising from the stochastic nature of ice nucleation is negligible.

In general, instrumentation employing single aerosol particles suspended in gas and droplet experiments working with cloud-sized droplets (tens of micrometres) report values of \( n_s \) down to about \( 10^5 \text{ cm}^{-2} \) (e.g. see Fig. 18 of Murray et al., 2012, for a compilation). These measurements are clearly valuable and applicable to the atmosphere, but even smaller \( n_s \) values are also relevant (Murray et al., 2012). For example, if we consider a dust-influenced atmosphere with five dust particles per cubic centimetre with a mean radius of 500 nm, then in order to generate 10 ice crystals per cubic metre an \( n_s \) of only 60 \( \text{cm}^{-2} \) would be required. Hence, it is important that we have the capacity to measure \( n_s \) smaller than \( 10^2 \text{ cm}^{-2} \) in addition to the capacity to measure larger values.

Our experimental approach builds on techniques employing aliquots of water much larger than the dimensions of typical cloud droplets which have been used since the very early days of ice nucleation studies (Vali, 1971, 1995) and continue to be used in the present day (Stopelli et al., 2014; Conen et al., 2011; Knopf and Forrester, 2011; Garcia et al., 2012; Budke and Koop, 2015). The advantage of this approach is that the surface area of nucleant per droplet scales with the volume of the droplet (for a constant mass fraction of nucleant in water). Hence, increasing the size of droplet from cloud droplet sizes (\( \sim 10 \mu \text{m} \); picolitre) to 1 mm (microlitre) increases the surface area per droplet by 6 orders of magnitude. This allows the quantification of \( n_s \) to much smaller values than is possible using cloud-sized droplets.

The basic concept of a droplet-freezing experiment (or droplet-freezing assay) is to take an aqueous suspension and subdivide it into multiple aliquots of ideally identical volume; however, polydisperse droplet distributions can also be used (Vali, 1971; Murray et al., 2011). The multiple droplets are then cooled identically. Experiments can be conducted using a constant cooling rate (i.e. isothermally) or a stepped temperature profile, and they have differed widely in terms of droplet volume, droplet production and cooling method (Vali, 1995). The fraction of droplets frozen at a given temperature or after a certain time interval can hence be determined. There are various methods of analysing the resultant data (Vali, 2014). These include deterministic models (such as Eq. 1) that link droplet fraction frozen directly to temperature (e.g. Vali, 1971) and time-dependent models of varying complexity (e.g. Marcolli et al., 2007; Herbert et al., 2014; Broadley et al., 2012).

In this paper, we present a method of conducting a droplet-freezing experiment using microlitre scale droplets. While the principle of the technique is not new, the specific application of the technique is, and the aim of this paper is to document the equipment, methods and analysis associated with the technique. Our instrument, the microlitre Nucleation by Immersed Particle Instrument (\( \mu \)-NIPI), is based around a Stirling cryocooler which cools a hydrophobic surface that supports microlitre volume droplets. The freezing of the droplets is monitored using a digital camera. We present new data for the nucleation efficiency of K-feldspar (muscovite), kaolinite, chlorite, NX-illite, Snomax\textsuperscript{®} and silver iodide. This data is used to illustrate several potential freezing artefacts and how to avoid them.

2 Description of the microlitre Nucleation by Immersed Particle Instrument

The \( \mu \)-NIPI forms part of a suite of instruments which are designed to make measurements of \( n_s \) over 10 orders of magnitude, thus covering the range relevant for the atmosphere. We have previously described a droplet-freezing technique using picolitre volume droplets which has been used to study mineral dusts (Atkinson et al., 2013; Broadley et al., 2012; Murray et al., 2011), soil particles (O’Sullivan et al., 2014) and homogeneous nucleation (Murray et al., 2010), as well as a technique using nanolitre droplets with soil (O’Sullivan et al., 2014) and combustion ash particles (Umo et al., 2014). The microlitre technique described in detail here has been used to study ice nucleation by mineral dusts (Atkinson et al., 2013), soil (O’Sullivan et al., 2014), nanoscale INPs (O’Sullivan et al., 2015), combustion ash (Umo et al., 2014) and time dependence of nucleation by kaolinite and K-feldspar (Herbert et al., 2014). This instrument was also included in an intercomparison between 17 instruments (Hiranuma et al., 2015). The resulting \( n_s \) values for nx-illite were within 1 order of magnitude of other immersion mode instruments, whereas the cloud chamber and FRIDGE instruments report larger \( n_s \) values in the same temperature range. The causes for this discrepancy are not clear but are discussed by Hiranuma et al. (2015).

The \( \mu \)-NIPI also offers a number of advantages over some other instruments: experiments can be performed relatively quickly; freezing events are easy to detect, and the continuous monitoring of freezing during a controlled temperature ramp allows the generation of a nucleation spectrum; and, in
addition, it has no need for cooling fluids and the equipment is portable, allowing it to be readily deployed in field settings.

The general layout of the µl-NIPI is shown in Fig. 1. The µl-NIPI consists of a cold stage, a hydrophobic surface which supports the droplets, an enclosure in which the humidity experienced by the droplets can be controlled and a digital camera to monitor the state of the droplets. To provide cooling and temperature monitoring, a Grant-Asymptote EF600 cold stage was employed. The EF600 was developed for the purpose of cooling samples for biological cryopreservation and can control the temperature of a sample between 20 and −100 °C. For cryopreservation, a top plate capable of holding multiple cryovials is typically employed. However, a flat aluminium top plate is also available and was used for this experimental setup. To conduct a droplet-freezing experiment, a 22 mm diameter hydrophobic silanised glass slide of 0.22 mm thickness (Hampton Research HR3-231) was placed onto this flat top plate. Prior to the experiments, the slide was cleaned using water, methanol and chloroform. Around 40 droplets of 1 µL volume were pipetted onto the slide using a Picus Biohit electronic pipette while the slide was at room temperature. To ensure that individual droplets contained the same amount of material the suspensions were stirred during the pipetting process. The uncertainty in volume quoted by the manufacturer is ±0.025 µL. The droplets and slide were then covered within a Perspex chamber with ports for a camera (Microsoft Lifecam HD) and stainless-steel pipes for delivering a gas flow to the cell. A recessed rubber O-ring was used to seal the chamber to the EF600 cold stage and an O-ring is also used to seal the camera opening. Both O-rings were coated with vacuum grease. A flow of dry zero-grade nitrogen (0.2 L min\(^{-1}\)) was passed through the cell in order to prevent frost growth (see discussion in Sect. 3.1).

The EF600 was internally controlled by a Eurotherm 2416 PID controller, run via Eurotherm’s iTools control software. For all work presented here, this software was used to program and commence a 1 °C min\(^{-1}\) temperature ramp from 1 to −35 °C. Once the ramp was commenced, data logger software associated with the EF600 was started and used to produce a log of temperature against time. A LabView program was used to record an image series from the digital camera, typically at a rate of one frame per second, and to produce a timestamp for each frame. Hence, the temperature of the cold stage during each frame was known. Videos were reviewed frame by frame and the temperature of freezing of each droplet recorded. Stills from the digital camera at several stages in the freezing experiment are shown in Fig. 2. The first change in droplet structure leading to droplet freezing was taken to be the nucleation event, and this information was used to establish the fraction of droplets frozen as a function of temperature.

The EF600 has a quoted temperature uncertainty of ±0.15 °C at −7 °C. To check the reliability of temperature measurement across a range of temperatures, a variety of compounds with known melting points were frozen and then melted by heating at 0.1 °C min\(^{-1}\). The melting temperature range was determined visually. Results from this process are presented in Table 1. It is necessary to calibrate using melting points, rather than freezing points; crystallisation observed during cooling is always subject to nucleation, making them
unsuitable for calibration of temperature unless the nucleation temperatures are very well defined (Riechers et al., 2013). By propagating the temperature error of the EF600 and the melting point range seen for water, a maximum temperature error of ± 0.4 °C has been estimated.

2.1 Suspension preparation

As in previous studies suspensions of solid material in Milli-Q purified water (18.2 MΩ cm) are made up gravimetrically and are then mixed using a magnetic stirrer plate overnight (Atkinson et al., 2013; O’Sullivan et al., 2014). In this study we used four mineral dusts (K-feldspar, kaolinite, chlorite and NX-illite), Snomax® (a commercial ice nucleant derived from *Pseudomonas syringae* bacteria) and silver iodide (details in Table 2). Mineralogies and other details are given in Table 2. The specific surface area of K-feldspar sample used was measured using the Brunauer–Emmett–Teller (BET) N₂ adsorption method using a Micromeritics TriStar 3000. The data presented in Sect. 4 using these nucleants are to illustrate the utility of the µl-NIPI and also illustrate potential artefacts the user should be aware of and how to avoid them. Uncertainty in surface area per droplet has been estimated at around ± 15 % by propagating uncertainties in weighing of samples, the BET surface area and the volume error of the pipette.

2.2 Control experiments

Droplets on the microlitre scale are notoriously difficult to produce free of any heterogeneous INPs. For example, Langham and Mason (1958) recorded a median freezing temperature of around −28 °C for droplets of a volume on the order of a few microlitres and only occasionally saw drops reaching what is thought to be homogeneous nucleation temperatures (−33 °C; Murray et al., 2010) in a system that suspended water droplets between two liquids of different densities. This was despite the use of a complex distillation system to obtain high-purity water. To the best of our knowledge, only Fornea et al. (2009) have reported being able to reach the freezing temperatures predicted by CNT for microlitre volume droplets repeatedly, reporting an average freezing temperature of −33.1 ± 0.6 °C for 2.0 μL droplets of high-performance liquid chromatography water. Given the difficulty in producing microlitre droplets free of suspended INPs and surfaces free of ice-active sites or contaminants, it is necessary to establish the temperature limit below which freezing cannot be assumed to have been induced by a heterogeneous nucleant.

Freezing temperatures for Milli-Q water droplets are shown in Fig. 3. These droplets froze mostly between −20 and −30 °C, substantially above the temperatures expected for homogeneous nucleation. Additionally, there is a 2–3 °C spread in the freezing temperatures of separate experiments, which is larger than would be expected if the freezing pathway was homogeneous. We conclude that there are variable quantities of heterogeneous ice-nucleating sites present that are active below about −20 °C, but it is not clear if nucleation is occurring on the slide or on impurities in the water. Milli-Q water is filtered through a 0.22 μm filter and is specified to have less than one particle per millilitre larger than this size. However, no specification is provided for smaller particles. Impurities could also be introduced at points during the experimental procedure or from airborne particles in the laboratory. In light of these control experiments, −20 °C has been taken as the lower limit for most of the heterogeneous ice nucleation experiments in this study.

3 Discussion of potential artefacts and uses of droplet-freezing experiments

Droplet-freezing experiments can suffer from a range of practical difficulties. This section addresses these problems and the solutions employed. As discussed by Stopelli et al. (2014), certain problems are encountered when using open droplet systems, where droplets are arranged on a hydrophobic surface, compared to closed droplet experiments, where each droplet is held in a separate container, or under inert oil. In an open droplet system, there is potential for droplets to evaporate over time, for contamination of droplets from airborne particles or for the freezing of droplets to influence neighbouring droplets by frost growth or splintering. As the µl-NIPI is an open droplet system, these issues must be overcome. Closed droplet systems avoid some of these prob-

![Figure 3. Temperature against fraction frozen for six different experiments using Milli-Q water. Also included is a line indicating the temperature of freezing expected for 1 µL pure water droplets from classical nucleation theory (CNT) according to the parameterisation by Murray et al. (2010).](http://www.atmos-meas-tech.net/8/1/2015/)
Table 1. Melting points of solvents used to calibrate temperature of the µl-NIPI. Recorded melting points are the average of five measurements. Literature melting points were taken from the 2007 CRC handbook.

<table>
<thead>
<tr>
<th>Chemical</th>
<th>Literature melting temperature (°C)</th>
<th>Recorded melting temperature (°C)</th>
<th>Standard deviation of melting point (°C)</th>
<th>Start (°C)</th>
<th>Finish (°C)</th>
<th>Range (°C)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Dodecane</td>
<td>−9.57</td>
<td>−9.5</td>
<td>0.09</td>
<td>−9.7</td>
<td>−9.3</td>
<td>0.4</td>
</tr>
<tr>
<td>Octanol</td>
<td>−14.8</td>
<td>−14.9</td>
<td>0.08</td>
<td>−15.1</td>
<td>−14.8</td>
<td>0.4</td>
</tr>
<tr>
<td>Undecane</td>
<td>−25.5</td>
<td>−25.4</td>
<td>0.10</td>
<td>−25.7</td>
<td>−25.3</td>
<td>0.4</td>
</tr>
<tr>
<td>Water</td>
<td>0.00</td>
<td>0.27</td>
<td>0.13</td>
<td>−0.1</td>
<td>0.6</td>
<td>0.7</td>
</tr>
</tbody>
</table>

Table 2. Characteristics of the materials used in this study. For silver iodide and Snomax® BET surface areas are not reported.

<table>
<thead>
<tr>
<th>Nucleant</th>
<th>N₂ BET surface area (m² g⁻¹)</th>
<th>Purity</th>
<th>Impurities</th>
<th>Supplier</th>
<th>Reference</th>
</tr>
</thead>
<tbody>
<tr>
<td>K-feldspar</td>
<td>1.86</td>
<td>78.1 %</td>
<td>3.9 % quartz, 16.0 % Na/Ca feldspar</td>
<td>Bureau of Analysed Samples, UK</td>
<td>Atkinson et al. (2013)</td>
</tr>
<tr>
<td>(microcline)</td>
<td></td>
<td></td>
<td>Crandallite, mica, illite, anatase</td>
<td>Clay Mineral Society</td>
<td></td>
</tr>
<tr>
<td>KGa-1b kaolinite</td>
<td>11.2</td>
<td>96 %</td>
<td></td>
<td>Chipera and Bish (2001); Murrey et al. (2011)</td>
<td></td>
</tr>
<tr>
<td>NX-illite</td>
<td>104.2</td>
<td>Mixed dust</td>
<td>–</td>
<td>Arginotec</td>
<td>Hiranuma et al. (2015); Broadley et al. (2012)</td>
</tr>
<tr>
<td>Chlorite</td>
<td>25.0</td>
<td>99.6 %</td>
<td>unknown</td>
<td>School of Earth and Environment; specimen collection</td>
<td>Atkinson et al. (2013)</td>
</tr>
<tr>
<td>Silver iodide</td>
<td>–</td>
<td>99.999 %</td>
<td>–</td>
<td>Alfa-Aesar premion</td>
<td>–</td>
</tr>
<tr>
<td>Snomax®</td>
<td>–</td>
<td>–</td>
<td>–</td>
<td>York Snow Inc.</td>
<td>–</td>
</tr>
</tbody>
</table>

3.1 Frost growth

Early experiments on the apparatus were conducted without a flow of dry N₂ over the droplets. This can cause significant issues when freezing of a droplet induces freezing in nearby droplets. Images of experiments with and without a dry flow of gas are shown in Fig. 4. In these experiments 10 droplets were seeded with silver iodide while the remainder were composed of pure water. When there was no flow of dry gas (panel b), there is extensive frost on the surface, which triggers freezing in neighbouring droplets; this artificially enhances the number of frozen droplets by 80 % at −12 °C in this case. Panel a shows an experiment conducted with a dry flow in which the seeded droplets froze at around −5 °C, as expected, while the pure water droplets remained liquid down to below −20 °C, indicating that with a dry flow there is no enhancement of freezing temperatures due to frost growth.

There are two distinct origins of the frost on the surface. Firstly, Jung et al. (2012) describe the formation of liquid condensation haloes, formed by the sudden increase in vapour pressure of water associated with the latent heat release of droplet freezing. They show that this halo can freeze if it persists. After freezing, subsequent frost growth directly from the vapour phase could come into contact with neighbouring droplets and trigger freezing in those droplets.

...
that the use of a 0.2 L min$^{-1}$ flow of dry gas eliminates the problems of frost haloes and frost growth. The condensation halo, consisting of very small droplets, evaporates rapidly in low-humidity conditions, and frost growth is reduced to the extent that it does not impact neighbouring droplets. Some small amount of frost growth is still observed in the immediate vicinity of a frozen droplet, taking the form of a very narrow ring around the droplet, but it does not extend far enough to interfere with neighbouring droplets. The flow of zero-grade nitrogen has the additional benefit that it flushes the chamber, reducing the chance of contact with aerosol particles from external sources.

3.2 Droplet evaporation

Using a dry flow as described in Sect. 3.1 has the potential to introduce additional issues. Droplets will evaporate to some extent through the course of the experiment (the extent will depend on the length of the experiment). This will increase the concentration of the contents of the droplets. Current descriptions of heterogeneous ice nucleation assume that the freezing temperature is dependent on the surface area of nucleant in a droplet, which will be constant despite evaporation. Therefore, it is assumed that changes in droplet volume will not affect experimental results. If solution droplets are being analysed, concentration changes due to droplet shrinkage will need to be accounted for as colligative effects will change the nucleation temperatures (Zobrist et al., 2008).

A second potential problem is that evaporation of water from droplets will cause cooling and may lead to the droplets being colder than the cold-stage temperature. In order to test for this evaporative cooling effect, a series of experiments was conducted both with and without the dry flow. Droplets containing 0.1 and 1 wt% K-feldspar were tested with the dry flow switched on. They were then thawed and refrozen with the dry flow switched off. Figure 5 shows the results of this experiment. Freezing by K-feldspar is described well by a singular model (Herbert et al., 2014), suggesting that each droplet freezes with a repeatable characteristic temperature (Vali, 2008, 2014; Wright and Petters, 2013). Freezing temperatures did not shift, showing that the cooling effect of the dry flow is smaller than the temperature measurement uncertainty ($\pm 0.4 ^\circ C$). The small freezing temperature range of these experiments ($\sim 4 ^\circ C$) meant that frost growth from frozen droplets did not spread to neighbouring droplets even with the dry flow switched off.

4 Test experiments and analysis

Several example data sets are presented here to demonstrate the efficacy of the µl-NIPI. The freezing temperatures for droplets containing K-feldspar (microcline), kaolinite, chlorite, Snomax$\textsuperscript{®}$ and silver iodide are shown in Fig. 6. The concentration of K-feldspar was varied between 0.01 and 1 wt%, and as expected the droplets containing more K-feldspar froze at a higher temperature. Droplets containing 1 % chlorite and kaolinite froze at lower temperatures than droplets containing K-feldspar, which is consistent with Atkinson et al. (2013), who presented a case suggesting that K-feldspar is the most important mineral component of atmospheric mineral dusts.

Silver iodide is known to be capable of nucleating ice at high temperatures (DeMott, 1995; Vonnegut, 1947), and even with a mass concentration 100 times less than that of 1 wt% K-feldspar it still nucleated ice at higher temperature. Interestingly, when a 0.1 wt% suspension was left on a windowsill for several days in a glass vial, the freezing temperature increased by around 2 $^\circ C$. There is a precedent for light exposure improving of the ice-nucleating ability of silver iodide in the literature, albeit in deposition mode rather than immersion mode (Rowland et al., 1964). There is also evidence from older work showing that light exposure can reduce the effectiveness of silver iodide as an ice nucleus (Fletcher, 1959; Smith et al., 1955). Further work is necessary to understand and quantify these effects.

Snomax$\textsuperscript{®}$ is a commercial preparation which contains fragments of a cultivated strain of Pseudomonas syringae and is used in the production of artificial snow at ski resorts (Cochet and Widehem, 2000). Droplets containing Snomax$\textsuperscript{®}$ nucleated ice around $-4 ^\circ C$ in this set of test experiments. Again, Snomax$\textsuperscript{®}$ is likely to exhibit variability in its ice nucleation activity between batches and individual samples depending on storage history; hence quantitative comparison with previous data is difficult. Nevertheless, Möhler
et al. (2008) also report freezing by Snomax\textsuperscript{®}-containing droplets at very high temperatures, with about 1\% of particles freezing at $-5.6^\circ$C. Hartmann et al. (2013) observed freezing in Snomax\textsuperscript{®}-containing droplets at temperatures up to $-6^\circ$C. Similarly, Wood et al. (2002) report freezing up to $-6^\circ$C in free-falling micron-scaled droplets containing Snomax\textsuperscript{®}. One feature reported in these previous studies is that the probability of freezing increases very steeply with decreasing temperature, which is again qualitatively consistent with our present measurements where the fraction frozen curve is extremely steep (Fig. 6). Recently, Wex et al. (2015) have published an extensive intercomparison of instruments using Snomax\textsuperscript{®} as the test sample. The most similar instrument to µL-NIPI intercompared, BINARY (Budke and Koop, 2015), gave qualitatively very similar results to those generated here at the most similar weight fraction tested. Our experiment $(2.4 \times 10^{-4} \text{ wt\%})$ and those from from Wex et al. (2015) $(2.9 \times 10^{-4} \text{ wt\%})$ both gave very steep fraction frozen curves at $-4^\circ$C.

While droplet-freezing temperatures and fraction frozen data are useful for comparing ice nucleation abilities of different materials within the context of a single experimental setup, it is necessary to normalise the data to some measure of the amount of material per droplet, typically surface area (Murray et al., 2012). The resulting values of $n_s$ are shown in Fig. 7, and we start this discussion with the data for K-feldspar. Over the three experiments the mineral surface area per droplet was varied by 2 orders of magnitude and the $n_s$ values fall on a single line, and the data are in good agreement with our previous study by Atkinson et al. (2013).

Atkinson et al. (2013) produced this parameterisation for K-feldspar through a combination of experiments with the µL-NIPI and pL-NIPI. The K-feldspar used in the present study is from the same stock sample used by Atkinson et al. (2013) but was ground separately and has a different specific surface area: $1.86 \text{ m}^2\text{ g}^{-1}$ rather than $3.15 \text{ m}^2\text{ g}^{-1}$ for the powder used by Atkinson et al. (2013). The chlorite and kaolinite $n_s$ values are much lower than the K-feldspar values, as expected, and are consistent with the conclusions of Atkinson et al. (2013). The data in Fig. 7 also illustrate the reproducibility of the derived $n_s$ values with this technique, with repeat runs for kaolinite and chlorite being identical within uncertainty.

By using relatively large droplets containing large surface areas of heterogeneous nucleants, higher freezing temperatures can be accessed using the µL-NIPI than are possible in smaller-scale droplet-freezing experiments or techniques employing particles dispersed in gas. A consequence of this is very small amounts of a relatively efficient nucleant can dominate the ice-nucleating ability of a given sample. For instance, if an impurity in a material is 1000 times more active than the bulk material, i.e. its $n_s$ is 1000 times higher than that of the bulk material, then the probability of freezing due to a 1 wt\% contamination would be 10 times higher than of the bulk material and the impurity would dominate observed ice nucleation. This effect is illustrated in Fig. 8, where a sample containing a 60-times-larger surface area of kaolinite than K-feldspar is shown to nucleate ice at essentially the same temperatures that the K-feldspar component of the mixture would be expected to do. In this example it
would easily attribute the ice nucleation activity observed to kaolinite, when in fact it all comes from the feldspar “contaminant”. If the aim of a study is to further fundamental understanding of ice nucleation by a particular material, it is important to recognize the potential impact of small impurities when working with larger droplets. While it is true that such an issue might also be observed with systems investigating smaller droplets, the smaller the amount of material in a droplet the less chance there is of a contaminating particle being present. Hence, there is a trade-off between measuring smaller \( n_s \) values and increasing the risk of contamination by minor components.

Some materials such as mineral dusts are known to aggregate when in aqueous suspension, and it has been hypothesized that this could lead to a reduced surface area and therefore underestimate in \( n_s \) (Hiranuma, 2015; Emersic, 2015 [15]). It is known that aggregation of mineral dusts is concentration dependent (Emersic et al. 2015); hence reducing the dust concentration should reduce the impact on \( n_s \). In order to test if aggregation is affecting our measured \( n_s \), we have performed a sequence of experiments with NX-illite suspensions of concentrations ranging from 0.05 to 2 wt% (see Fig. 9). It can be seen that the resulting \( n_s \) values are self-consistent, indicating that occlusion of surface by aggregation is not significant. Aggregates that do form are sufficiently porous that water can still reach all relevant surfaces of the dust. In addition, \( n_s \) for droplets containing 0.01–1 wt% K-feldspar also falls on a single line (Fig. 7), suggesting that aggregation does not lead to the occlusion of surface area in K-feldspar either. Although aggregation does not affect \( n_s \) in these mineral dusts, it is worth bearing in mind that aggregation could conceivably reduce surface area in some nucleants, and it is therefore well worth testing for this by varying nucleant concentration and checking for consistency.

5 Summary

A new design of droplet-freezing experiment for ice nucleation experiments has been constructed and tested. The µl-NIPI uses a Stirling engine-based cryocooler to cool microlitre volume droplets at a controlled rate in order to determine the efficiency with which immersed particles nucleate ice. Many modern atmospheric ice nucleation measurement techniques measure the nucleation by aerosolised particles or particles suspended in cloud-sized droplets. The advantage of working with droplets much larger than cloud droplets is that they contain a far larger surface area of nucleant. This allows the determination of nucleation efficiencies over a wider range of temperatures than is possible using only smaller droplets and complements the flowing aerosol and cloud-sized droplet techniques in widespread use.
In most ice nucleation experiments, efficient minor components in samples can dominate observed results, meaning that great care is needed in interpretation of results. While this is particularly true for experiments using such relatively large amounts of sample as the technique presented here, the ability to detect the activity of relatively rare, high-temperature ice nucleation events is valuable. The µl-NPI may be particularly useful in a field setting, where its low detection limit and simplicity of operation offer advantages over, and complementarity to, more complex instruments.

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