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Dealing With Numerical Noise in CFD-Based Design Optimization

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SUMMARY

Numerical noise is an inevitable by-product of Computational Fluid Dynamics (CFD) simulations which can lead to challenges in finding optimum designs. This article draws attention to the issue, illustrating the difficulties it can cause for road vehicle aerodynamics simulations. Firstly a benchmark problem is used to assess a range of turbulence models and grid types. Large noise amplitudes up to 22% are evident for solutions computed on unstructured tetrahedral grids whereas computations on hexahedral and polyhedral grid structures exhibit substantially less noise. The Spalart Allmaras turbulence model is shown to be far less susceptible to noise levels than two other commonly-used models for this application. Secondly, multi-objective aerodynamic shape optimization is applied to a fairing for a practical road vehicle which is parameterised in terms of three design variables. Moving Least Squares (MLS) metamodels are constructed from 50 high-fidelity CFD solutions for two objective functions. Subsequent optimization is successful for the first objective, however numerical noise levels in excess of 7% give rise to difficulties for the second one. A revision to the problem leads to success and the construction of a small Pareto Front. Further analysis underlines the inherent capability of MLS metamodels in dealing with noisy CFD responses.

KEY WORDS: Numerical noise, optimization, moving least squares, metamodel, CFD, aerodynamics.

1. INTRODUCTION

In the past twenty years the speed and power of computers has increased by between 1000 and 10,000 times which is facilitating the computation of increasingly complex fluid flow systems [1]. This capability is being exploited in the rapidly-growing research field of Computational Fluid Dynamics (CFD) optimization which is being utilised across a range of areas including aerospace engineering [2,3], tribology [4], polymer moulding [5], ship design [6], vehicle aerodynamics [7-10], hospital ward ventilation [11] and jet pump design [12]. Although these examples demonstrate the versatility of CFD-based optimization, there is one aspect which can prove problematic: the presence of numerical noise in the CFD responses [10,13-19].

In fact, numerical noise has long been a hindrance for computation in general with problems first reported for finite element tidal simulations in 1974 [20] and other examples in the following decade [21-23]. There, numerically induced oscillations with small wavelengths were particularly troublesome. One of the earliest reported examples of numerical noise hindering CFD optimization was the investigation carried out by Giunta *et al.* [13] into the design of a high-speed civil airliner. Optimization on polynomial response surfaces was made difficult by spurious noise-induced local minima, which served to 'trap' the optimizer. A method of skipping over these local minima using large move limits in the initial stages of the optimization search was

beneficial, however this approach did not address the fundamental problem and no single optimum design was found. The impact of noise on response surfaces was also discussed by van Keulen *et al.* [14] for structural applications.

Later Madsen *et al.* [15] and Shyy *et al.* [16] commented that noise originating from numerical simulations is much less recognised than for physical experiments. An important point to appreciate is that numerical noise is an inherent by-product of computer simulation [15,24,25] and the observed behaviour is rather different from the noisy responses originating from experiments. In general, for a given physical experiment there will be statistical variation in the answer due to errors and uncertainties originating from both controlled and uncontrolled variables. In contrast, computer experiments produce the same output for a given set of input variables provided all aspects of the simulation are constant (i.e. identical initial/boundary conditions, solver version, hardware/architecture, grid structure etc.). The difference with numerical experiments is that the errors (and thus the noise) are repeatable due to their deterministic nature [18]. For optimization studies which require analysis of a wide range of designs (such as those for constructing response surfaces or metamodels) this characteristic can lead to problems in identifying optimum designs [10,13]. These issues have received attention; however, there are no studies dedicated exclusively to numerical noise and the negative impact it can have on CFD optimization. The purpose of this article is to draw attention to these in the context of road vehicle aerodynamics.

The article is structured as follows: Section 2 considers the benchmark Ahmed body [26] as a test case for simulating vehicle aerodynamics; a simple method for characterising numerical noise both qualitatively and quantitatively is presented. The influence of turbulence model selection and grid type on noise levels are discussed and analysed in detail. Section 3 describes the results from a practical high-fidelity CFD optimization problem, highlighting the difficulty of metamodel-based optimization in the face of noisy CFD responses. Section 4 draws on the results from sections 2 and 3 and strategies are proposed for minimising the negative impact which numerical noise has on optimization. Finally, conclusions are drawn in section 5.

2. NUMERICAL NOISE

Previous studies have shown that numerical noise resulting from CFD simulations is sensitive to the choice of turbulence model employed [15,16] and the grid used in the computations [17]. Burman and Gebart [17] showed conclusively that the component of numerical noise which is attributable to the grid (i.e. discretisation error) can be minimized with adequate grid resolution. Whilst these studies are helpful, there are no attempts to quantify numerical noise nor to determine the relative contributions from the known sources. The remainder of this section addresses this by assessing a range of turbulence model-grid combinations for a relevant test case.

2.1 Ahmed body

Moving road vehicles induce a range of aerodynamic flow features which include: separation bubbles, edge and horseshoe vortices, separating shear layers and eddying turbulent wake structures [27,28]. The presence of a

moving ground plane (relative to the vehicle) and rotating wheels introduce further complexity which makes the simulation of road vehicle aerodynamics a non-trivial exercise. Ahmed [26] conducted an experimental study using a generic road vehicle shape (the Ahmed body) inside a wind tunnel to characterise salient features of the time-averaged wake structure. Up to 85% of the total vehicle drag is attributable to pressure drag with over 90% of the latter resulting from flow separation over the rear faces. In the base region, the time-averaged wake structure consists of a pair of horseshoe vortices stacked vertically within the separation bubble. The strength of these vortices is found to depend on the slant angle at the rear of the vehicle. Furthermore, for high slant angles of 30° or greater, an additional but small separation bubble exists immediately behind the transverse edge joining the roof to the slanted rear face. These results are still being used as a benchmark for comparing against numerical simulations (see e.g. [29,30]). The Ahmed body is particularly relevant to the present work and it represents a suitable test case for analysing numerical noise in the context of steady-state vehicle aerodynamics simulations.



Figure 1: Illustration of the Ahmed body with relevant dimensions for a rear slant angle of $\psi = 30^{\circ}$.

Figure 1 illustrates the Ahmed body which consists of a solid block with rounded leading edges of radius R = 0.1m, a notch at the base and support pillars to lift the body above the ground. The size of the notch is governed by the slant angle, ψ , which is kept constant at 30°. Other steady-state numerical studies cite poor results for smaller slant angles; when the slant angle is reduced to 25° for example, partial detachment of the flow behind the slant provides small scale structures which are inherently unsteady and these cannot be captured by steady-state approaches which often lead to poor results [31]. In order to simulate the flow field around the vehicle, a solid model was generated using Ansys Design Modeler (version 13.0) [32] based on the dimensions shown.

The size and shape of the air volume surrounding the vehicle is defined using the dimensions of the wind tunnel originally used by Ahmed [26]. The closed-return open working section tunnel was supplied with airflow through a square nozzle of 3m x 3m. The vehicle was mounted on a ground board of length 5m with the vehicle centre located 2.13m downstream of the inlet. To reduce the computational effort a symmetry plane was employed and the working section was assumed to have a constant cross-section matching the dimensions of the inlet nozzle. No-slip boundary conditions were used on all solid walls, whereas the side and ceiling of the

domain were assigned a zero shear stress boundary condition which is appropriate given the original opensection wind tunnel layout.

2.2 Grid structure

Recent CFD investigations of airflow past a bluff vehicle in a wind tunnel have demonstrated the importance of grid density, cell type and the choice of turbulence model for predicting aerodynamic drag [33,34]. In the present exploration three grid densities are considered for each of the following cell types: (i) hexahedral, (ii) tetrahedral and (iii) polyhedral. Each of these employ a boundary layer grid adjacent to solid walls (i.e. the vehicle and the ground) which consists of 12 layers of cells, a first cell height of 0.008m and a cell height expansion ratio of 1.2. A preliminary study showed that solutions computed using this boundary layer grid lead to average wall y^+ values of between 30 and 40 on the surfaces of the vehicle which is in the correct range for the standard wall functions used [35]. Table 1 summarises each grid and Figure 2 shows the local grid structure at the base of the vehicle for the coarse hexahedral, tetrahedral and polyhedral grids respectively. The hexahedral and tetrahedral grids were generated with AnsysMesh (version 13.0) [32] and the polyhedral grids were produced using an agglomeration procedure within Fluent (version 13.0.0-sp2) [32] which converts a standard tetrahedral grid into an equivalent polyhedral one.

		Global Cell Count			
Grid type	Local grid spacing (m)	Hexahedral	Tetrahedral	Polyhedral	
Coarse	0.015	229512	396106	173934	
Medium	0.010	479865	703887	302074	
Fine	0.007	699314	1383917	583475	

Table 1: Grid statistics



Figure 2: Local grid structure on the symmetry plane (dark cells) and the rear of the Ahmed body for (a) hexahedral, (b) tetrahedral and (c) polyhedral cells.

2.3 Turbulence models

In computing solutions to the governing incompressible Navier-Stokes equations, the choice of turbulence model is an important consideration. This is especially so for high Reynolds number turbulent flow such as the one being investigated because the dominant feature, flow separation, can be predicted somewhat differently depending on the turbulence model employed. In order to assess the impact of turbulence model on the amount of numerical noise present, three models suitable for simulating external aerodynamics, see [33], were chosen,

namely: (i) the Spalart Allmaras model (SA) [36], (ii) the realizable k- ε model (RKE) [37] and (iii) Menter's shear-stress-transport k- ω model (SSTKO) [38].

Steady-state solutions were computed using Fluent (version 13.0.0-sp2) [32] for each turbulence model on all 9 grids, giving 27 solutions in total. All solutions assumed a free-stream velocity of 60 m/s and turbulence intensity of 0.5% at the inlet [26] (Note that the relatively high inlet velocity is to compensate for the reduced scale of the vehicle which leads to a Reynolds Number of 4.3 million and is consistent with the original experiments [26]). Irrespective of the turbulence model-grid combination, every simulation employed second order discretisation of the flow equations in conjunction with the SIMPLE [39] pressure-velocity coupling algorithm. Although solution convergence was generally achieved in fewer than 1000 iterations, all simulations were run for a total of 5000 iterations to eliminate convergence errors.

2.4 Quantifying numerical noise

For each of the simulations described above, the drag coefficient of the vehicle, C_D , was monitored throughout the 5000 iteration cycle. In all cases numerical noise was evident, characterised by a combination of structured periodic cycles with seemingly random oscillations superimposed. Interpreting these characteristics is difficult from visual inspection alone and so there is a need to quantify the noise levels. Figure 3 shows a typical noise sample from one of the simulations, taken from iterations beyond the point of solution convergence.



Figure 3: Sample of numerical noise exhibited by a Spalart Allmaras solution on the fine hexahedral grid.

The noise levels can be decomposed into the frequency and the amplitude of oscillation. The former is conveniently defined by the percentage oscillation frequency, f_r , given by:

$$f_r = 100 \frac{\sum_{i=0}^{n} \Omega_i}{n},\tag{1}$$

where Ω_i is the oscillation parameter evaluated for the *i*th iteration for a sample size of *n* iterations. If C_D increases monotonically from one iteration to the next, $\Omega = 0$. For a maximum or minimum point (i.e. oscillation peak in Figure 3) the sign of the gradient dC_D/di changes and $\Omega = 1$. It follows that for $f_r = 0$ the signal is likely to be stable with no oscillations present, whereas $f_r = 100$ indicates a fully oscillatory signal where the gradient changes sign every iteration. It should be noted that f_r accounts for all local gradient changes but it does not

consider low-frequency oscillations (e.g. on the order of 100's of iterations). As CFD solutions for steady state problems are typically taken from the final iteration, low-frequency oscillations are far less influential and thus less relevant than the high frequency ones described by f_r .

In addition to the frequency, the amplitude of each individual oscillation accounts for the magnitude of the variations present. For a given sample size, three standard deviations, 3σ , (known as the three-sigma rule in statistics), is an adequate measure of the data spread because it accounts for 99.7% of the values recorded. This is used to define the percentage amplitude of noise, $A_{3\sigma}$, (for the sample) relative to the mean value, namely:

$$A_{3\sigma} = 100 \frac{3 \cdot \sigma}{\overline{C_D}}, \qquad (2)$$

where σ is the standard deviation and $\overline{C_D}$ is the mean drag coefficient for a sample of size *n*. As with f_r , small values of $A_{3\sigma}$ denote less noise whereas larger ones signify a noisy response. For the 150-iteration sample shown in Figure 3, $f_r = 42.0\%$ and $A_{3\sigma} = 0.4\%$; i.e. the noise is frequent but its impact is minimal because of the small amplitude.

2.5 Impact of grid type, cell type and turbulence model

The parameters defined by equations (1) and (2) are used to analyse the numerical noise present in the data obtained for C_D for the range of iterations: 2500-5000, per simulation. A sensitivity study showed that this sample size is large enough to adequately characterise both f_r and $A_{3\sigma}$ and it only considers the converged region of each solution. Table 2 summarises these parameters along with the mean drag coefficient, $\overline{C_D}$, for the range of grid-turbulence model combinations tested. In all cases the computed drag coefficients are less than the equivalent experimental value of 0.378 [26]. Overall the SA model gives the most satisfactory result with the RKE and SSTKO models generally exhibiting the smallest drag values. The computed drag coefficients generally reduce as the grid density increases regardless of the turbulence model. As mentioned previously, Reynolds-Averaged Navier-Stokes (RANS) equations such as those employed in this study do have their limitations [31], however the results obtained are in reasonable agreement with physical experiments, and the primary focus here is on numerical noise.

	-		SA			RKE			SSTKO	
Grid type	Cell type	$\overline{C_D}$	f_r (%)	$A_{3\sigma}$ (%)	$\overline{C_D}$	f_r (%)	$A_{3\sigma}$ (%)	$\overline{C_D}$	f_r (%)	$A_{3\sigma}$ (%)
Coarse	Hexahedral	0.368	55.7	0.1	0.339	79.1	0.1	0.309	14.5	1.6
Medium	Hexahedral	0.353	8.6	0.2	0.313	13.8	0.0	0.298	24.0	1.1
Fine	Hexahedral	0.341	39.2	0.4	0.342	41.4	0.5	0.300	13.2	1.6
Coarse	Tetrahedral	0.367	95.1	0.4	0.377	9.4	22.6	0.324	28.0	7.6
Medium	Tetrahedral	0.354	100.0	0.0	0.308	16.3	2.3	0.304	39.8	0.0
Fine	Tetrahedral	0.351	96.2	0.8	0.317	75.2	8.3	0.332	76.4	7.2
Coarse	Polyhedral	0.354	100.0	0.0	0.318	49.8	0.0	0.327	100.0	0.0
Medium	Polyhedral	0.346	2.6	0.1	0.302	39.8	0.0	0.309	9.2	5.1
Fine	Polyhedral	0.344	3.2	0.3	0.296	23.0	0.1	0.299	13.6	3.2

Table 2: Computed mean drag coefficients and associated numerical noise as a function of grid type, cell type and turbulence model. Note: from [26] the experimental drag coefficient, $C_D = 0.378$.

It is interesting to note that both f_r and $A_{3\sigma}$ vary considerably, depending on the grid and cell type and the turbulence model; clearly all three factors impact the noise levels which is consistent with earlier studies [15-17]. The differences in the observed values of f_r illustrate that the noise levels are not in phase from one simulation to another. In terms of the amplitude of oscillations they are below 1% for all SA solutions, however variations as high as 22.6% and 7.6% are present in the solutions for the RKE and SSTKO models respectively. Clearly, the choice of turbulence model is instrumental in determining the noise levels for this particular application.

In the majority of cases the frequency and the amplitude of oscillations are greatest for solutions computed on the tetrahedral grids compared to the hexahedral and polyhedral ones. , however there is no apparent correlation with the grid density. In some cases the noise amplitude increases as the grid becomes finer, e.g. SA solutions on the hexahedral and polyhedral grids. However, for the remaining cases the largest amplitudes (per combination of cell type and turbulence model) can occur for either the coarse, medium or fine grid densities. Whilst there is a lack of generality for these results, the fine-grid solutions are inevitably closer to being grid independent and so these are more relevant to the overall discussion.

2.6 Fine-grid solutions

Figure 4 shows a 500-iteration sample of the relative drag coefficient (with respect to the mean value, $\overline{C_D}$) as a function of turbulence model and cell type for fine-grid solutions only. In all cases the noise levels exhibited by the tetrahedral-grid solutions are substantially greater than the equivalent hexahedral and polyhedral ones with multi-modal responses clearly seen. Solutions obtained on the hexahedral and polyhedral grids show noise levels with significantly smaller amplitudes and reduced frequencies in all cases. Considering the results for each turbulence model in turn, the noise levels are smallest for the SA model with a range of values generally within $\pm 0.5\%$ of the mean with the exception of some local variations of the order $\pm 1.0\%$ for the tetrahedral-grid solutions show variations within $\pm 0.5\%$ of the mean value, however those for the tetrahedral grid are up to $\pm 8.0\%$, Figure 4(b). The same trend is seen for the SSTKO model although the hexahedral and polyhedral-grid solutions exhibit larger variations of $\pm 3.0\%$ compared to those obtained with the other models.



Figure 4: Plots of the relative drag coefficient $(C_D / \overline{C_D})$ as a function of steady-state iteration number per grid cell type for (a) SA, (b) RKE and (c) SSTKO turbulence models. Note the smaller y-axis scale for (a).

The disparity in results between respective turbulence models is particularly noteworthy, especially when considering that the same grids and initial/boundary conditions are used throughout. One possible reason for this is that solutions produced by certain turbulence models may have large discretisation errors due to inadequate grid resolution. To investigate this, the Grid Convergence Index (GCI) [40] was employed to estimate the fine-grid discretisation error, *GCI*_{FINE}, namely:

$$GCI_{FINE}(\%) = 100 \frac{F_{s}|e|}{(r^{p}-1)},$$
(3)

where F_s is the factor of safety, r is the grid refinement ratio (based on local grid spacing), p is the order of discretisation and e is given by:

$$e = \frac{\left(C_{D-MEDIUM} - C_{D-FINE}\right)}{\left(C_{D-FINE}\right)}.$$
(4)

For this study p = 2 (i.e. second order), r = 1.5 and $F_s = 1.25$ based on the recommendations of [40,41]. It is clear that the discretisation errors vary significantly depending on the turbulence model and the grid cell type, see Table 3.

	GCI _{FINE} (%)				
Cell type	SA	RKE	SSTKO		
Hexahedral	3.4	8.4	0.7		
Tetrahedral	0.9	3.1	8.4		
Polyhedral	0.6	2.0	3.2		

Table 3: Discretisation errors calculated for the fine-grid solutions using the Grid Convergence Index (GCI) [40].

Overall, the errors are small for the SA model which may explain the small noise levels shown in Figure 4a. Similarly, larger discretisation errors are evident for both the RKE and SSTKO models which correspond with the larger noise levels present in Figures 4b and 4c, respectively. Thus, on the whole, the lower the discretisation error, the smaller the level of numerical noise present in the solutions.

2.7 Impact of noise

It is clear from the results presented thus far that numerical noise can dramatically influence CFD solutions. Previous studies by Madsen *et al.* [15] and Forrester *et al.* [19] have highlighted the impact this can have when plotting functions of interest through a design region. The former study focussed on an extremely small segment of the design space where the CFD responses were oscillatory yet they should have shown a linear variation; noise levels of 1.0% were found to be responsible. This problem occurs because CFD solutions are typically taken from the final iteration (for steady state simulations) and this could correspond with any region of the noisy response. As some solutions will coincide with noise peaks and others nearer the mean value, each solution is essentially iteration dependent.

In the present work this iteration dependency is investigated by varying the slant angle on the Ahmed body through the range $\psi = 30-31^{\circ}$ in increments of 0.1°. From the fine-grid results already presented in Table 2, the largest noise levels occurred for solutions computed on the tetrahedral grid using the RKE model ($A_{3\sigma} =$ 8.34%) with the smallest occurring for the SA model in conjunction with the polyhedral grid ($A_{3\sigma} = 0.28\%$). As these two turbulence model-grid combinations produce both extremes of noise, they are suitable for investigating iteration dependency for small changes in ψ . Corresponding simulations were run for 5000 iterations with the mean value for the drag coefficient, $\overline{C_D}$, taken for the interval 2500-5000 iterations (as before) and the final value, $C_{D-final}$, taken from the last iteration. Figure 5 shows how $C_{D-final}$ (normalised with respect to $\overline{C_D}$) varies as a function of ψ for both cases. The variations are clearly seen when compared to the mean relative drag coefficient ($\overline{C_D}/\overline{C_D}$) which has error bars determined from the noise amplitude, $A_{3\sigma}$, per simulation.



Figure 5: Plots of the relative drag coefficient in terms of the mean $(\overline{C_D} / \overline{C_D})$ and the final iteration value ($C_{D-final} / \overline{C_D}$) for solutions computed by (a) RKE model on the tetrahedral-based grid and (b) SA models using the polyhedral grid. Error bars determined from $A_{3\sigma}$ (equation 2). Note the difference in the vertical scales.

Figure 5a underlines the inherent variability of solutions computed using the RKE model in combination with tetrahedral-based grids. As an example, for $\psi = 30.2^{\circ}$ the spread of solution values due to numerical noise is $\pm 13.6\%$ (based on $A_{3\sigma}$). More importantly, in a number of cases the range of final iteration solutions vary considerably from the mean: e.g. for $\psi = 30.1^{\circ}$ the final iteration value, $C_{D-final}$, is almost 5% greater than the mean. Thus these iteration dependent solutions skew the final result by a considerable margin and the root cause of the discrepancy, numerical noise, should be treated with caution in other investigations.

For the low noise case (Figure 5b), solutions computed on the polyhedral grid using the SA model lead to significantly smaller noise levels with a maximum spread of 1% observed for $\psi = 30.2^{\circ}$. In terms of the final iteration values, the largest discrepancy also occurs for $\psi = 30.2^{\circ}$ with a solution 0.3% below the mean; however, discrepancies of 0.1% are typical. Another aspect of Figure 5 which deserves attention is the *variation* in noise levels observed between geometries, especially given the small (one-degree) range of slant angles considered. This shows that noise levels are sensitive to the exact geometry and this is particularly relevant to design optimization which typically requires solutions for a range of geometries.

3. CFD-BASED OPTIMIZATION

The results presented above for the Ahmed body provide valuable insight into numerical noise and its impact on CFD solutions. Major contributing factors include the choice of turbulence model, the grid density and cell type. In combination, these three factors can lead to significant noise levels which dramatically affect solution values from iteration to iteration. While the Ahmed body represents a valuable test case, this is a simple, constrained geometry. As described in section 1, numerical noise can be extremely problematic in the context of design optimization which often requires CFD solutions for multiple geometries. As well as highlighting the problems associated with numerical noise, the purpose of this article is to show how noise levels can impact design optimization and to establish ways of dealing with it. Considered below is the effect this can have as part of an optimization problem for a practical engineering investigation.

3.1 Small livestock trailers

In the United Kingdom the majority of animals transported between farms, markets and abattoirs are carried in small box-shaped livestock trailers such as the one depicted in Figure 6.



Figure 6: Illustration of a small livestock trailer and towing vehicle.

These trailers are towed by off-road vehicles and ventilation of the trailer is achieved primarily by virtue of vehicle movement. This leads to air exchange between the internal environment and the external free-stream through a series of vents located on either side. Although this is effective at maintaining a stable micro-environment in the upper deck, the lower deck exhibits reduced ventilation [33] with the potential for elevated temperatures and related animal welfare issues [34]. Modifying the trailer layout to improve these conditions is not viable due to practical limitations, however, implementing a retrofitted headboard fairing represents a feasible solution. One further advantage of adding such a fairing is that aerodynamic drag can be reduced leading to improved fuel economy [42,43]. It follows that the design of a headboard fairing with objectives of (i) minimising aerodynamic drag and (ii) maximizing ventilation is a suitable multi-objective optimization problem.

3.2 Problem formulation

The proposed fairing is parameterised in terms of three design variables, namely the side radius, d_1 , the lower edge extension, d_2 and the central extension of the fairing, d_3 , see Figures 7a and 7b. The purpose is to apply aerodynamic shape optimisation in satisfying the following criteria:

$$\min F_1(\boldsymbol{d}) \text{ and } \max F_2(\boldsymbol{d}), \tag{5}$$

where

$$d_i^L \le d_i \le d_i^U, \ i = 1, 2, 3; \tag{6}$$

 F_1 and F_2 are the objective functions for the aerodynamic drag coefficient (dimensionless) and the ventilation rate (m³/s) respectively and d_i is the *i*th design variable subject to relevant lower (d_i^L) and upper (d_i^U) physical constraints.



Figure 7: Parameterisation of the headboard fairing viewed from (a) the side and (b) above the trailer and (c) plot of the Design of Experiments (DoE) together with four sample fairing designs.

3.3 Design of Experiments and CFD methodology

The optimization procedure involves building metamodels for each objective function (F_1 and F_2) and then optimizing on these to obtain a Pareto front from which to determine the optimum fairing design. To achieve this, each metamodel is fitted to CFD responses from fifty fairing designs which are chosen by an Optimal Latin Hypercube (OLH) Design of Experiments (DoE) [44]. Figure 7(c) depicts the DoE and four sample fairings which exist within the design space. Though not immediately obvious from Figure 7(c), one of the defining characteristics of the DoE is the uniformity of point coverage which is governed by the Audze-Eglais potential energy criterion [45,46].

Steady-state CFD solutions to the governing incompressible Navier-Stokes equations were used to assess each fairing by obtaining values of F_1 and F_2 . A preliminary grid independence study carried out on the baseline trailer (Figure 6) showed that a hybrid hexahedral-tetrahedral grid consisting of 6.7 million cells led to small discretisation errors (equation 3) of 1.5% and 2.0% for F_1 and F_2 , respectively [10]. As shown in section 2, solutions computed on both hexahedral and polyhedral cell types lead to reduced noise levels compared to tetrahedral cells. In practice, hexahedra are easier to implement than polyhedra and they are less susceptible to numerical diffusion. For this reason hexahedral cells were placed in as many regions of the solution domain as possible which included a structured boundary layer grid adjacent to the primary livestock trailer surfaces. Inevitably tetrahedral cells were required in the remaining volume due to the geometric complexity which illustrates one of the difficulties of practical CFD application. Despite this, the results from section 2 show that the noise amplitudes for solutions computed on fine grids using the SA model are 0.4% and 0.8% for hexahedra and tetrahedral respectively. As the SA model is used in this section (the justification is explained in the next paragraph) the negative impact of tetrahedral cells is relatively small. Figure 8 shows the local grid structure in the vicinity of the baseline trailer with the symmetry plane visible in the background.



Figure 8: Illustration of the local grid structure (a) surrounding the trailer and (b) near the foremost upper vent opening.

As shown in section 2, the Spalart Allmaras turbulence model generally exhibits small noise amplitudes for road vehicle aerodynamics simulations. Initial simulations of flow around the baseline configuration (Figure 6) verified that this model performs more satisfactorily than both the SSTKO and RKE models. Furthermore, the SA model has been shown to produce accurate results when validated against wind tunnel experiments of a 1/7th scale livestock trailer [33], and this was selected together with QUICK [48] discretisation. Computations were carried out using Fluent [32] for a total of 10,000 iterations and convergence of all quantities was observed after 9000 of these, thus ensuring that no inaccuracies were present due to convergence error. The greater detail contained within this geometry required finer grids than those used for the Ahmed body (considered in section 2) which is why a larger number of iterations was required to reach convergence. The important point is that the solutions were converged.

3.4 *Optimization strategy 1: final value solutions*

Having obtained all fifty sets of CFD solutions, metamodels were built for each objective function using the Moving Least Squares (MLS) method [48,49] within HyperStudy (version 8) [50]. This technique caters for noisy responses by selecting an appropriate closeness of fit parameter, θ , which is contained within a Gaussian weight decay function, namely:

$$w_j = \exp\left(-\theta r_j^2\right),\tag{7}$$

where r_j is the Euclidean distance of the metamodel prediction location from the j^{th} DoE point [49]. High noisesmoothing is achieved if θ is small because the fit is loose due to the approximation whereas high values of θ lead to interpolation and no smoothing. Each metamodel was tuned to give the optimum value of θ to ensure the best fit to the CFD responses, see [10] for more details.

For each CFD simulation, responses for F_1 or F_2 were taken from the final iteration. Satisfying the second objective of maximizing ventilation proved difficult because all fifty fairing designs resulted in poorer ventilation with respect to the baseline case (i.e. no fairing present). Further analysis showed that the presence of any given fairing streamlined the front of the trailer which guided airflow past the vents instead of through them thereby reducing the ventilation rate. Consequently, maximizing ventilation (equation 1) is not feasible with the current problem formulation. Instead, the second objective was changed such that the precentage reduction in ventilation rate could be minimized (min $F_2(d)$) which is equivalent to minimising the negative impact that the fairing has on ventilation [10].

In order to find a fairing design for minimum drag (i.e. min $F_1(d)$), a Genetic Algorithm (GA) was employed to carry out a global search on the corresponding metamodel before using a local gradient search method, the Sequential Quadratic Programming (SQP) technique, to 'home in' on the proposed global minimum. This design (which resided in a corner point of the design space) was assessed using an additional CFD solution and the result was a net drag reduction of 6.6%. The drag metamodel was subsequently rebuilt with the additional data point but an optimization search did not yield a better design; the minimum drag design had been found in a single step. Figure 9 shows how this fairing design (labelled "Min-Drag") compares with the original fifty DoE points.



Figure 9: Objective function plot for a range of fairing designs based on final value CFD solutions.

The above optimization procedure was repeated for the ventilation metamodel in search of an optimum ventilation fairing. This proved to be problematic because the optimizer repeatedly predicted a fairing which was poorer than many of the DoE points. Nevertheless the design, which resides in another corner point of the parameter space, was tested with a CFD simulation. The result concurred with the metamodel in showing that the design exhibits poor performance, see "Max-Ventilation" label in Figure 9. Subsequent metamodel rebuilding with the additional data point made no difference to the predicted global optimum. In an effort to construct a Pareto front with the possibility of finding better ventilation designs, four additional points were suggested by the metamodel. Each of these was found by placing a constraint on F_1 to give F_2 for each design [10]. CFD solutions were obtained for each but the poor results shown in Figure 9 (see "Extra Points") underline the difficulties of optimizing the fairing for ventilation.

3.5 *Optimization strategy 2: mean solutions*

The presence of numerical noise in each CFD solution for F_2 was suspected as the root cause of the difficulties encountered above. This is explored in Figure 10a which shows a typical solution history where the drag coefficient, C_D , and the ventilation rate, Q, are plotted as a function of the iteration number. Closer inspection of C_D over the converged portion of the data (Figure 10b) shows that the noise frequency is relatively high at 36.7%, although this is accompanied by a small amplitude of 0.6% (i.e. \pm 0.3%). In contrast, Q, which is used in the calculation of F_2 , exhibits lower frequency noise of 1.3%, however, the amplitude is large at 7.0%.

This result underlines the inherent variability of solutions for F_2 and the potentially large discrepancy between mean and final value CFD solutions. In light of this fact, the difficulties in optimizing F_2 are unsurprising. A possible way of avoiding these problems is to mask the noise by building metamodels using mean values which eliminate possible data spikes. Accordingly, each CFD solution was run for a further 2000 iterations and the mean for F_1 and F_2 calculated. Both metamodels were updated using this revised data and the optimization procedure repeated (as above) in search of min $F_1(d)$ and min $F_2(d)$. As before, optimization identified the same optimum fairing for F_1 , however, the same sub-optimal design was predicted for the second objective function, F_2 , see "Max-Ventilation" point in Figure 11. Although some of the "Extra Points" show improvement in terms of drag (compared to Figure 9), the problems relating to F_2 persist and optimization of this function was not successful using our approach.



Figure 10: Plots showing typical noise levels for one CFD solution.



Figure 11: Objective function plot for a range of fairing designs based on mean value CFD solutions taken from 2000 iteration cycles.

3.6 Optimization strategy 3: problem revision

The results in Figure 11 show that simply masking pronounced noise levels using mean solutions does not remedy optimization problems. Instead, the reliability of the objective function for characterising ventilation was investigated. Whilst ventilation is a suitable measure of the air quality within the trailer, its calculation is based on 2D surface integrals of the volumetric flow rate through each of the side vents in the lower deck (Figure 6). As these regions are oblique to the free-stream, large flow gradients are present and this leads to the high noise levels cited above. Further investigation revealed that the vent openings experiencing the greatest flow gradients produced the highest variability (and thus noise) from iteration to iteration as each solution progressed. By basing the objective function on a 3D volume-averaged quantity, the impact of high flow gradients reduces substantially. Consequently, the temperature humidity index (THI) [51], which has units of °F, was chosen. This is given by the relation:

$$THI = (1.8T_{db} + 32) - [(0.55 - 0.0055rh)(1.8T_{db} - 26.8)]$$
(8)

where T_{db} is the dry bulb temperature (°C) and *rh* is the relative humidity expressed as a percentage. By taking the volume-average of this quantity throughout the lower deck of the trailer, thermal comfort and thus animal welfare is considered instead of ventilation. Accordingly the problem was revised to:

$$\min F_1(\boldsymbol{d}) \text{ and } \min F_3(\boldsymbol{d}), \tag{9}$$

where F_3 is the objective function representing THI. Using the isothermal solutions as a basis, all simulations were run for an additional 4000 iterations with extra transport equations for energy and species also solved; these account for thermal effects and humidity, respectively. Source terms for energy and moisture production were used to represent animal warmth and perspiration so that F_3 could be calculated for hot (30 °C) and humid (rh =95%) ambient conditions (see [10] for more details). With these new parameters incorporated into the CFD simulations, convergence was observed within the first 1500 iterations and so the means for F_1 and F_3 were calculated from the remaining 2500 iterations.

Results from the revised simulations were found to be free of significant noise levels, justifying the problem revisions. Table 4 summarises the noise statistics for all 56 simulations with mean amplitudes of 0.3% and 0.2% evident for F_1 and F_3 , respectively. Overall, the noise frequency is greater in the results for F_3 ; however, the small amplitudes present for both objective functions underline the dramatic improvement; this is in complete contrast to the noisy responses seen earlier.

	F_{I}		1	F ₃
Value	$f_r(\%)$	$A_{3\sigma}$ (%)	$f_r(\%)$	$A_{3\sigma}$ (%)
mean	4.6	0.3	15.6	0.2
max	10.8	6.5	33.0	0.6
min	1.7	0.1	1.5	0.0

Table 4: Noise statistics for all 56 simulations obtained using the problem revision.

In addition to the reduced noise levels, the "Extra Points" generated in the previous section show a considerable improvement in terms of the second objective function, in this case F_3 . This is evident in Figure 12a which displays all 56 points based on final value solutions. In terms of mean values the objective function plot retains the same features with all 56 points residing in broadly similar locations, see Figure 12b. As mean values are more representative of the actual solutions, these were used to construct new metamodels from the revised data prior to optimizing. The drag metamodel predicted the same optimum fairing design as before, whereas the THI metamodel revealed a candidate for min $F_3(d)$. An additional CFD simulation verified that this design gave the smallest THI of all the designs tested, suggesting that the optimum for F_3 had been found, see "Min-THI" in Figure 12b. This conclusion was verified from subsequent metamodel rebuilding and optimization with the additional point; it did not lead to a better design. Note that the small Pareto front was generated using a multiobjective genetic algorithm (MOGA) [52] which was applied to both metamodels. Although this did suggest a possible improvement to the right of the "Min-THI" design, this is small and at the expense of extra drag.



Figure 12: Objective function plot for the revised problem using (a) final values and (b) mean quantities including the minimum-THI design and Pareto front.

3.7 Advantages of Moving Least Squares metamodels

Of all the data points shown in Figure 12, the locations of the "Extra Points" are particularly noteworthy. These four designs were predicted in the isothermal optimization (section 3.4) yet they exhibited sub-optimal performance in terms of ventilation and drag. With the above problem revision these very same designs now reside in the most promising region of the objective function landscape and, in terms of mean solutions, one of them is in fact Pareto optimal. Therefore, in spite of the pronounced noise levels seen in the isothermal study, the MLS metamodelling technique was in fact successful in filtering this noise and thus identifying optimal design

characteristics. It follows that the MLS metamodels in this study were in fact more accurate than the high-fidelity CFD solutions used in the initial problem formulation. This stems from the inherent noise-handling capability of the approximation-based technique employed and is consistent with the earlier findings of Papila and Haftka [53].

3.8 Optimum design

Taking the results from Figure 12 into consideration the "Min-THI" design was chosen as the overall optimum; it produces the greatest benefit in terms of THI and offers drag reduction which is close to the minimum-drag design. Compared to the baseline trailer, the optimum design offers 5.3% less drag and a small but clear 0.02% lower THI. Closer inspection of the benefits gained from the fairing show that the surface pressure distribution is the determining factor. Figure 13 shows a comparison of the surface pressure distribution (expressed in terms of the pressure coefficient, C_p) for the baseline and optimum designs. The optimum fairing reduces

the size of both the high and low pressure regions which are present around the front of the bluff, baseline design. Furthermore, the fairing effectively extends the side of the trailer upstream (circled) and the accompanying low pressure (which is absent in the baseline design) serves to extract warm, humid air through the foremost lower vent opening, thereby lowering the animal welfare indicator (THI) as desired. Finally, figure 14 illustrates how much more compact the wake behind the trailer is with the addition of the fairing. Each wake is represented by an iso-surface of constant velocity magnitude (U = 5.0 m/s) with the optimum fairing reducing the length of the wake by 22%.



Figure 13: Surface contour plots of the pressure coefficient, C_p for (a) the baseline trailer and (b) the optimum version.



Figure 14: Wake structure comparison for (a) the baseline trailer and (b) the optimum one. Wake represented using iso-surfaces of constant velocity magnitude of 5 m/s.

4. **DISCUSSION**

4.1 Numerical noise within CFD solutions

The results presented in this article demonstrate how numerical noise can affect CFD solutions for road vehicle aerodynamics simulations. Noise levels are shown to be dependent on the grid density, cell type and turbulence model which agrees with previous investigations for aerospace applications [15-17]. Flow solutions around the Ahmed body [26] which are computed on tetrahedral grid structures are by far the noisiest compared to those for hexahedral and polyhedral grids, regardless of the grid density and turbulence model used. Interestingly, noise levels do not necessarily reduce as the grid becomes finer and in fact the opposite trend is evident in some cases (e.g. solutions computed using the Spalart Allmaras turbulence model, Table 2). This is more a result of the interaction between the turbulence model and the grid type and not due to incomplete simulations; only the fully converged portion of each simulation was considered in the noise analysis. On the whole, solutions obtained using the Spalart Allmaras turbulence for the realizable k- ε and SST k- ω models.

In terms of solution accuracy, discretisation errors were computed for the fine-grid solutions and on the whole, solutions with small errors lead to reduced noise levels although this is not true for all cases. It should also be noted that for more complicated geometries, it is often necessary to employ hybrid grid structures which combine tetrahedral cells with hexahedra or polyhedra. The flexibility of tetrahedra comes with the disadvantage of greater numerical noise levels being present in the solutions obtained.

Section 2 highlighted how noise levels are conveniently described in terms of the frequency and amplitude (equations 1 and 2) with variations in the latter essentially placing an error band on each solution. As steady-state solutions develop, large amplitudes lead to iteration dependence and this can be particularly problematic when multiple designs are tested. Incremental changes in the slant angle of the Ahmed body

highlighted the great variability of solutions and this sensitivity was also evident to a lesser degree in a separate parametric study of flow past an aerofoil [19]. In fact the true solution per geometry should be taken from the average of a suitably large sample using solution monitors (post convergence); this serves to remove the fluctuations which make up numerical noise.

4.2 *Optimizing with noisy solutions*

Although mean values are an effective way of masking noise levels, they cannot address fundamental problems with the choice of objective function. The livestock trailer optimization study detailed in the previous section illustrated how an inappropriate objective function can in fact prevent optimization. Here, high flow gradients led to pronounced noise levels in each CFD solution for the ventilation rate within the trailer. Optimization of the ventilation rate was attempted using both final value and mean solutions but neither approach succeeded; a series of designs proposed by an MLS metamodel were shown to be ineffective. In spite of this, a subsequent change to the choice of objective function landscape (Figure 12). The difference with the revised problem was that the objective function was based on a volume-averaged quantity, the temperature humidity index, which yielded solutions exhibiting low noise levels. Whilst it took the problem revision to identify an optimum design, the earlier MLS metamodels had in fact found optimal design characteristics despite noise levels in excess of 7%.

4.3 Noise-smoothing metamodels

The conclusion that MLS approximations can deal with significant noise levels and still manage to identify optimum designs is one of the key points of this article. Allied to this, it has long been known that approximation-based metamodels can handle numerical noise; the study by Giunta *et al.* in 1994 demonstrated this [13] and other authors have discussed the advantages of approximations when applied to CFD optimization [7,10,17,53]. More traditional interpolation-based techniques such as Radial Basis Functions [54] and Kriging [54] can force metamodels into unnatural behavior, particularly when fitting to noisy data [19]. Clearly, the MLS metamodelling technique used in this study was well-suited to the noisy CFD responses and interpolation would not have been inappropriate.

Although approximations have their advantages, the criterion used to fit to data points is extremely important. For the MLS methodology adopted in this study, the closeness of fit parameter, θ , was optimized for the given data set using a build-validate-rebuild technique [10]. This is only made possible by decomposing the DoE into a primary model building DoE and a smaller validation DoE [44]. The purpose is to build an initial metamodel from the build DoE before minimising the RMS error between the metamodel prediction and the responses at the validation locations. This yields the optimal θ value to be used in the merged metamodel which consists of both the build and validation points. Such an approach ensures the accuracy of the final metamodel which cannot be guaranteed if θ is specified by the user *a priori*. Thus the choice of DoE is another important consideration is seeking optimum designs.

4.4 Control of simulation errors and validation

In addition to the optimization strategy, the quality of the CFD methodology is equally important. Maximizing solution quality can be achieved by following verification and validation (V&V) procedures such as the widely adopted guidelines from the AIAA [55], ERCOFTAC [56] and ASME [41]. They advocate great care in preparing CFD simulations to minimize the errors present. The simulations described in this study utilised double precision real-number representation to reduce round-off error, they were run to full convergence so that convergence errors could be eliminated and discretization errors were calculated using the Grid Convergence Index [40] as part of a rigorous grid independence study. Grids were also produced using as many structured hexahedral cells as possible (in section 3) to limit the negative impact of numerical diffusion. Considering the fact that discretisation errors are generally the most dominant in CFD solutions [59] it is essential that great care is taken in producing high-quality grid structures.

As well as reducing errors, the quality of CFD solutions can be improved using experimental data which is useful for minimizing uncertainties when prescribing boundary conditions for example. Although physical experiments are also subjected to errors, data obtained from them can be extremely valuable for validating the performance of individual numerical models including those designed for simulating turbulence, multiphase and combustion [59]. In the present work the Spalart Allmaras turbulence model was selected based on favourable comparison with relevant physical wind tunnel tests [33].

5. CONCLUSIONS AND RECOMMENDATIONS

A number of steps can be taken to improve the chances of success in metamodel-based CFD optimization. As already described, it is essential to base any optimization on the highest quality CFD responses which require minimizing the errors where possible. Double-precision real number representation helps reduce round-off error, convergence errors can be avoided altogether if simulations are run for a sufficient number of iterations and grid independence studies can be used to select the most appropriate grid density and to provide an estimate of the discretization error (e.g. using the Grid Convergence Index [40]). Furthermore, validation data from relevant experiments is extremely valuable in ensuring that the fundamental flow physics is being adequately represented by the computations.

Whilst these steps are beneficial, the aforementioned errors contribute to fluctuations in a given solution which can be defined as numerical noise. As such, noise levels should be monitored for quantities of interest (i.e. objective functions) and the degree of variation observed. In this study it was necessary to run simulations for more than 2000 iterations (post convergence) to visualise both high and low frequency oscillations. Although it is difficult to quantify how much noise will cause optimization difficulties in other investigations, the 7% variations seen in the present work had a negative impact. In the event of optimization problems it is advised, following [15], that a small region of the design space is explored to determine the sensitivity of the CFD responses to slight changes in the design variables. For a small enough region this procedure should result in

almost linear variations; strong non-linearity (as was the case in Figure 5) may be a sign of potentially destructive levels of numerical noise.

Emphasis should also be placed on the problem formulation and particularly on the choice of objective functions to be used in optimization studies. Basing these quantities on flow parameters which are measured in regions exhibiting high flow gradients can dramatically skew solutions. Where possible, such quantities should be based on solutions from a number of cells so that the average of these is representative of the objective function whilst retaining an element of stability (i.e. less noise). Finally, the benefits of approximation-based metamodels such as MLS are ideally suited to dealing with numerical noise and they can be readily incorporated into optimization studies.

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