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Liu, E-L, Yu, H-S, Zhou, C et al. (2 more authors) (2017) A Binary-Medium Constitutive Model for Artificially Structured Soils Based on the Disturbed State Concept and Homogenization Theory. *International Journal of Geomechanics*, 17 (7). 04016154. 04016154-04016154. ISSN 1532-3641

[https://doi.org/10.1061/\(ASCE\)GM.1943-5622.0000859](https://doi.org/10.1061/(ASCE)GM.1943-5622.0000859)

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1 A Binary-Medium Constitutive Model for Artificially Structured Soils 2 Based on the Disturbed State Concept (DSC) and Homogenization Theory

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9 **Abstract:** Triaxial compression tests were carried out on artificially structured soil samples at confining
10 pressures of 25 kPa, 37.5 kPa, 50 kPa, 100 kPa, 200 kPa and 400 kPa. A binary-medium constitutive model
11 for artificially structured soils is proposed based on the experimental results, the disturbance state concept
12 (DSC) and homogenization theory. A new constitutive model for artificially structured soils was
13 formulated by regarding the structured soils as a binary-medium consisting of bonded blocks and weakened
14 bands. The bonded blocks are idealized as bonded elements whose deformation properties are described by
15 elastic materials and the weakened bands are idealized as frictional elements whose deformation properties
16 are described by the Lade-Duncan model. By introducing the structural parameters of breakage ratio and
17 local strain coefficient, the non-uniform distribution of stress and strain within a representative volume
18 element can be given based on the homogenization theory of heterogeneous materials. The methods for
19 determination of the model parameters are given on the basis of experimental results. By making
20 comparisons of predictions with experimental data, it is demonstrated that the new model provides
21 satisfactory qualitative and quantitative modeling of many important features of artificially structured soils.

22 **Key Words:** artificially structured soils; binary-medium constitutive model; breakage ratio; local strain
23 coefficient.

24 Introduction

25 Soils in situ usually possess natural structures, referring to the combination of fabric (arrangement of
26 particles) and inter-particle bonding (Mitchell 1976), whose important influence on the mechanical features
27 of soils has been recognized for a long time, enabling soils composed of the same materials to behave
28 differently in a reconstituted state (Burland 1990; Leroueil and Vaughan 1990). The natural structure

29 conveys extra strength to natural soils, allowing them to exist at a given stress. Upon loading, the bonds
30 between soil particles may break, resulting in the so-called destructuration process. Until the present,
31 research on the geotechnical engineering properties of reconstituted soils has been relatively satisfactory,
32 and the modified Cam clay model (Schofield and Wroth 1968; Yao et al. 2009, 2015) and the Lade-Duncan
33 model (Lade and Duncan 1975; Lade 1977) have been widely used in solving geotechnical problems
34 resulting from reconstituted soils. It is widely known that the modified Cam clay model can simulate only
35 the strain hardening and volumetric contraction of remolded clays relatively well but cannot well duplicate
36 the strain softening and volumetric dilatancy of natural or structured soils at a low stress state under triaxial
37 stress conditions (Smith et al. 1992). During the process of formation, natural soils are easily deposited
38 layer by layer, which results in different mechanical properties in vertical and horizontal directions
39 (Graham and Houlsby 1983). Therefore, when formulating the constitutive model of these types of soils,
40 the influences of stress history, bonding, the fabric distribution and current stress state variables should be
41 considered concurrently to describe the stress–strain properties well.

42 There have been important developments in formulating constitutive models incorporating the influence
43 of soil structure based on comprehensive experimental studies on structured or natural soil. Many
44 researchers have investigated the mechanical properties of structured soils by laboratory experiments on
45 intact soil samples extracted from construction fields (Lo and Morin 1972; Sangrey 1972; Baracos et al.
46 1980; Schmertmann 1991; Diaz-Rodriguez et al. 1992; Callisto and Calabresi 1998; Cotecchia and
47 Chandler 2000; Dudoignon et al. 2001; Callisto et al. 2002; Rocchi et al. 2013) and on artificially
48 structured soils (Maccarini 1987; Bressani 1990; Malanraki and Toll 2001), in which the yielding, strength,
49 deformation properties, aging, anisotropy, and stress path of structured soils were investigated. When
50 formulating constitutive models for structured soils considering their mechanical properties obtained by
51 laboratory experiments, there are some widely used methods, which include revising or reformulating the
52 Cam clay model (Kavvads and Amorosi 2000; Asaoka et al. 2001; Liu and Carter 2002; Wheeler et al.
53 2003; Belokas and Kavvadas 2010; Suebsuk et al. 2011; Zhu and Yao 2013; Liu et al. 2011, 2013), damage
54 mechanical model (Shen 1997; Zhao et al. 2002; Shen 2006;), DSC model (Liu et al. 2000; Liu et al. 2003)
55 kinematic or bounding model (Rouainia and Wood 2000; Gajo and Wood 2001; Baudet and Stallebrass
56 2004; Huang et al. 2011) and micromechanical model (Yin et al. 2009; Gao and Zhao 2012). Many of these

57 existing models are formulated based on macroscopic observation on stress-strain properties of structured
58 soils and few can consider the physical and deformational mechanism of them. Furthermore, there is still
59 no widely accepted constitutive model for structured soils at the moment.

60 Compared with the remolded clay, the structured soils behave with strain softening and volumetric
61 contraction followed by dilatancy upon loading under a relatively low stress state, accompanying the
62 appearance of the shear bands under triaxial and biaxial stress states. After the peak value of stress–strain
63 curves of structured soils (Cotecchia and Chandler 2000), the yielding surface will contract gradually as a
64 result of the breaking of bonds between soil particles, which makes it difficult to describe these phenomena
65 using the conventional and widely employed elasto-plastic theory. Accompanying the bond breaking
66 between soil particles, the stress and strain distributed within a soil element will not be uniform, and the
67 higher local stress that equals the strength of the bonds will result in the breakup of these bonds between
68 soil particles. Therefore, it is necessary to formulate a constitutive model for structured soils to consider the
69 non-uniform stress and strain in the soil element and reflect the macroscopic strain softening by use of the
70 parameters considering the micro deformation mechanism. Here, a new constitutive model for structured
71 soils will be proposed to consider the damage process (or gradual bond breaking) and non-uniform
72 distribution of strain (or stress) based on test results of artificially structured soils.

73 In this paper, the triaxial tests of artificially structured soils were performed at six different confining
74 pressures ranging from 25 kPa to 400 kPa with drained conditions, and a theoretical study of the behavior
75 of artificially structured soil is presented. Based on the homogenization theory of heterogeneous materials
76 and the disturbance state concept (DSC), a new model, referred to as the binary-medium model for
77 geological materials, is formulated by regarding the structured soils as a binary-medium consisting of
78 bonded blocks and weakened bands. The determination of model parameters is provided and model
79 verification is also made by comparison with the test results of artificially structured samples.

80 **Test Conditions and Results**

81 **Sample Preparation**

82 The artificially structured soils tested here are composed of silty clay, cement, kaolin clay and salt particles,
83 in which silty clay is the main matrix material, cement can provide bonding between soil particles, kaolin
84 clay can increase the content of fine particles of the samples, and salt particles can generate large pores

85 within the samples by dissolving. The silty clay was extracted from one excavation pit located in Chengdu
86 area, approximately 5 m below the ground surface, with blocky shape and slight moisture, and its G_s is 2.72.
87 The grading curve of silty clay is shown in Fig. 1, and w_L and w_p are 29.11% and 17.06%, respectively. The
88 silty clay is dried and sieved through a 0.5 mm screen and serves as the main matrix material mixed
89 uniformly with other materials, including cement, kaolin clay and salt particles by mass (or weight), in
90 which the mass ratios of silty clay, kaolin clay, cement and salty particle are 65%, 20%, 5% and 10%,
91 respectively. The cement employed is 32.5R, which is produced in China. The uniform mixture is then
92 compacted in a mold with the three same parts by five layers with dry density of 1.49 g/cm^3 to form the
93 sample. The samples are vacuumed for approximately 3 h in a vacuum chamber before the distilled water
94 flows in slowly. After the samples are soaked for 3 h, they are removed from the vacuum chamber and
95 quickly placed in flowing water with a speed of $6.65 \text{ cm}^3/\text{s}$. After curing for seven days, the samples are
96 taken from the mold and placed in the triaxial apparatus to be tested. During the process of curing, the salt
97 content in the water is measured to ensure complete dissolution of the salt particles. Through seven days of
98 curing, the salt content in the water surrounding the samples reaches its original value, which is equal to the
99 magnitude of the flowing water; this demonstrates that salt particles are dissolved completely. The
100 Scanning Electron Microscope (SEM) photo of one prepared sample is shown in Fig. 2, which presents the
101 bonding between soil particles and the distribution of large pores within the sample. For natural soils, their
102 main properties at the mesoscale are bonding and fabric (Burland 1990). In the process of preparing the
103 samples, the hydration of cement generates some materials bonding soil particles together, and the
104 dissolution of salt particles forms the large pores within the samples; thus, the initial isotropic structured
105 samples will be prepared.

106 To investigate the influence of structure deterioration on the mechanical properties of soils, the remolded
107 samples are also prepared here, and their preparing method is described as follows. The artificially
108 structured samples tested are remolded, dried and sieved through a 0.5-mm screen. After that, the soils are
109 compacted in the mold with five layers to form remolded samples with the same dry density as the
110 structured ones. Obviously, the bonding between soil particles of the remolded samples is broken
111 completely.

112 **Test Results and Analysis**

113 Triaxial compression tests under consolidated-drained conditions are conducted on both the artificially
114 prepared samples and the remolded ones. The confining pressures applied are 25 kPa, 37.5 kPa, 50 kPa,
115 100 kPa, 200 kPa and 400 kPa, and loading rate is 0.06 mm/min. The apparatus employed is a GCTS
116 triaxial system.

117 The deviatoric stress–axial strain curves and the volumetric strain–axial strain curves of the structured
118 samples are presented in Fig. 3 (a)–(b), respectively, in which “S-CD-xx kPa” means that the structured
119 sample is tested under consolidated-drained conditions at the confining pressure of xx kPa. From Fig 3 (a)
120 and (b), we can find that (i) under lower confining pressures, the samples exhibit strain-softening behavior
121 and initially contract followed by dilatancy, and the lower the confining pressure, the more the sample
122 dilates; (ii) under higher confining pressures, the samples exhibit strain-hardening behavior and contract at
123 all times, and the larger the confining pressure, the more the sample contracts. When the confining
124 pressures are 25 kPa, 37.5 kPa, and 50 kPa, the bonds between soil particles at the end of consolidation are
125 hardly damaged, so these bonds should be destroyed gradually during the application of shear loading,
126 which causes the samples to exhibit strain softening behavior accompanied by the appearance of shear
127 bands as shown in Fig. 4 (a)–(c). Conversely, when the confining pressures are 100 kPa, 200 kPa and 400
128 kPa, the bonds between soil particles at the end of consolidation are heavily damaged, so the sliding of the
129 soil particles mainly contributes to their strength during the application of shear loading, which causes the
130 samples to exhibit strain hardening behavior and contract accompanied by a failure pattern of bulging in the
131 middle, as shown in Fig. 4 (d)–(f).

132 The deviatoric stress–axial strain–volumetric strain curves of the remolded samples are presented in Fig.
133 5 (a)–(b), in which R denotes the remolded samples. Because the bonds between soil particles of the
134 remolded samples are very weak, their mechanical properties are distinct from those of artificially
135 structured samples. From Fig. 5 (a) and (b), we can find that (i) the remolded samples exhibit strain
136 hardening behavior under the confining pressures ranging from 25 kPa to 400 kPa; and (ii) at low confining
137 pressures, they contract first and then finally tend to dilate with the overall volumetric compaction, and at
138 high confining pressures, they contract at all times. When remolding the artificially structured samples in
139 the process of preparation, the bonds between soil particles break to form larger aggregates that are
140 composed of the remolded samples, which thus behave as coarse-grained soils (Yu 2006). The failure

141 patterns of remolded samples are in the form of bulges in the middle, as shown Fig. 6 (a)–(f) under all
142 confining pressures.

143 From the test results of the structured samples and remolded ones under consolidated-drained conditions
144 with different confining pressures, we can find that (i) under the relatively lower confining pressures, the
145 deviatoric stresses of the structured samples are larger than those of remolded samples, as shown in Fig. 7
146 for the confining pressure of 50 kPa. The artificially structured soils exhibit strain-softening behavior, but
147 the remolded samples exhibit strain hardening behavior; (ii) under the relatively higher confining pressures,
148 both types of samples exhibit strain hardening behavior. In the process of strain hardening, the deviatoric
149 stresses of the structured samples are larger than those of the remolded samples, and the differences
150 between them are decreasing, as shown in Fig. 8 for the confining pressure of 200 kPa; and (iii) under the
151 confining pressures ranging from 25 kPa to 400 kPa, the volumetric compaction of the remolded samples is
152 larger than that of the structured samples.

153 **Binary-medium Constitutive Model for Artificially Structured Soils**

154 **Breakage Mechanism of Structured Soils**

155 Soil structures have a great influence on the mechanical properties of natural soils (Mitchell 1976), in
156 which the cohesive resistance and frictional resistance contribute together to the bearing capacity of the soil
157 element. It has also been long known that cohesive and frictional resistance are not mobilized
158 simultaneously at different deformation or strain levels (Lambe 1960), with the former reaching a peak
159 value within a relatively small strain and the latter making a full contribution within a relative large
160 deformation or strain. It is obvious that the cohesive component exhibits brittle behavior and the frictional
161 component exhibits nonlinear elastic behavior. The cohesion essentially comes from the cementation
162 bonding between particles, whose distribution is not uniform among geological materials. The bonded
163 blocks are formed where the cementation bonding strength is stronger, and the weakened bands are formed
164 where the cementation bonding is weaker, so the heterogeneous structured soils are developed step by step
165 via sedimentation. During the loading process, the brittle bonded blocks gradually break up, transforming
166 to elasto-plastic weakened bands, so the two components bear the loading collectively. With the
167 development of the breakage process, the bearing capacity of the bonded blocks will decrease, and that of
168 the weakened bands will increase; however, the structured soil wholly exhibits strain hardening or strain

169 softening behavior, depending on the increase of the bearing capacity of the weakened bands and the
170 decrease of the bearing capacity of the bonded blocks. In view of the understanding of the breakage
171 mechanism of structured soils mentioned previously, the structured soil can be conceptualized as a binary-
172 medium material consisting of bonded blocks and weakened bands bearing the capacity collectively (Shen
173 2006). In the following, the bonding blocks are called the bonded elements, and the weakened bands are
174 called frictional elements. There are similar concept of Disturbed State Concept (DSC) proposed by Desai
175 and coworkers (Liu et al. 2000; Desai 1974, 2001), in which the continuum element is assumed to be
176 composed of intact (RI) and adjusted (FA) states and has been used for soils (sands and clays), rocks,
177 rockfill, asphalt, concrete, silicon, polymers, and interfaces and joints. In DSC, a deforming material is a
178 mixture of (RI and FA states and similar in bonded materials) components which interact with each other to
179 lead to the observed behavior. The material mixture can undergo degradation or softening and stiffening or
180 healing. However, the basis in the damage approach is different; it starts from the assumption that a part of
181 the material is damaged or cracked. The observed behavior is then defined based essentially on behavior of
182 the undamaged part, and both do not interact because the damaged part is assumed to possess no strength.

183 Fig. 9 presents the sketch of Binary-Medium, where the bonded element is composed of a spring (E_b) and
184 a brittle bond (q) and the frictional element is composed of a spring (E_f) and a plastic slider (f). For the
185 brittle bond, it does not deform when the stress is less than the bond strength q and fail once the stress
186 reaches q . In a continuum of structured soil sample, there are many bonded elements and frictional
187 elements. Upon loading, some bonded elements may break up and transfer to frictional elements and bear
188 external loads collectively.

189 **Formulation of Binary-medium Constitutive Model for Artificially Structured Soils**

190 The stress–strain relation of artificially structured soils, regarded as a binary-medium material consisting of
191 bonded elements and frictional elements, can be derived by taking a representative volume element (RVE)
192 based on homogenization theory for heterogeneous materials (Wang et al. 2002) as follows.

193 For a representative volume element, or RVE, the local stress and local strain are denoted by σ_{ij}^{loc} and
194 ε_{ij}^{loc} , respectively, and thus both the average stress σ_{ij} and the average strain ε_{ij} can be written as
195 follows:

196
$$\sigma_{ij} = \frac{1}{V} \int \sigma_{ij}^{loc} dV \quad (1)$$

197
$$\varepsilon_{ij} = \frac{1}{V} \int \varepsilon_{ij}^{loc} dV \quad (2)$$

198 where V is the volume of the RVE and loc represents local stress or strain.

199 σ_{ij}^b and σ_{ij}^f are defined as the stresses of bonded elements and frictional elements, respectively, and

200 they have the following expressions:

201
$$\sigma_{ij}^b = \frac{1}{V_b} \int \sigma_{ij}^{loc} dV_b \quad (3)$$

202
$$\sigma_{ij}^f = \frac{1}{V_f} \int \sigma_{ij}^{loc} dV_f \quad (4)$$

203 where V_b and V_f are the volumes of bonded elements and frictional elements in the RVE, respectively, and

204 b and f represent the bonded and frictional elements, respectively. From equation (1), we have

205
$$\sigma_{ij} = \frac{1}{V} \int \sigma_{ij}^{loc} dV = \frac{V_b}{V} \sigma_{ij}^b + \frac{V_f}{V} \sigma_{ij}^f \quad (5)$$

206 ε_{ij}^b and ε_{ij}^f are defined as the strains of bonded elements and frictional elements, respectively, with the

207 following expressions:

208
$$\varepsilon_{ij}^b = \frac{1}{V_b} \int \varepsilon_{ij}^{loc} dV_b \quad (6)$$

209
$$\varepsilon_{ij}^f = \frac{1}{V_f} \int \varepsilon_{ij}^{loc} dV_f \quad (7)$$

210 From equation (2), we have

211
$$\varepsilon_{ij} = \frac{1}{V} \int \varepsilon_{ij}^{loc} dV = \frac{V_b}{V} \varepsilon_{ij}^b + \frac{V_f}{V} \varepsilon_{ij}^f \quad (8)$$

212 Setting λ as a breakage ratio, the ratio of volume of frictional elements to the whole volume of RVE is

213 expressed as follows:

214
$$\lambda = \frac{V_f}{V} \quad (9)$$

215 Substituting Eq. (9) into Eqs. (5) and (8), we can express the average stress and average strain as follows:

$$216 \quad \sigma_{ij} = (1 - \lambda)\sigma_{ij}^b + \lambda\sigma_{ij}^f \quad (10)$$

$$217 \quad \varepsilon_{ij} = (1 - \lambda)\varepsilon_{ij}^b + \lambda\varepsilon_{ij}^f \quad (11)$$

218 The breakage ratio is changing with strain level upon loading, which is an internal variable similar to the
 219 damage factor used in damage mechanics or hardening parameter used in plasticity. We assume here that
 220 the breakage ratio is a function of strain, namely,

$$221 \quad \lambda = f(\varepsilon_{ij}) \quad (12)$$

222 By use of Eq. (10), we can obtain the incremental expression of the stress as follows:

$$223 \quad d\sigma_{ij} = (1 - \lambda^0)d\sigma_{ij}^b + \lambda^0 d\sigma_{ij}^f + d\lambda(\sigma_{ij}^{f0} - \sigma_{ij}^{b0}) \quad (13)$$

224 where λ^0 is the current breakage ratio, and σ_{ij}^{b0} and σ_{ij}^{f0} are the current stresses of bonded elements and
 225 frictional elements, respectively. Similarly, by derivation of Eq. (11), we can obtain the incremental
 226 expression of the strain as follows:

$$227 \quad d\varepsilon_{ij} = (1 - \lambda^0)d\varepsilon_{ij}^b + \lambda^0 d\varepsilon_{ij}^f + d\lambda(\varepsilon_{ij}^{f0} - \varepsilon_{ij}^{b0}) \quad (14)$$

228 where ε_{ij}^{b0} and ε_{ij}^{f0} are the current strains of bonded elements and frictional elements, respectively.

229 The tangential stiffness matrixes of bonded elements and frictional elements are represented by D_{ijkl}^b
 230 and D_{ijkl}^f , respectively, so we have the following stress–strain relationships for bonded elements and
 231 frictional elements:

$$232 \quad d\sigma_{ij}^b = D_{ijkl}^b d\varepsilon_{kl}^b \quad (15)$$

$$233 \quad \text{and} \quad d\sigma_{ij}^f = D_{ijkl}^f d\varepsilon_{kl}^f \quad (16)$$

234 By manipulation of Eq. (14), we have

$$235 \quad d\varepsilon_{ij}^f = \frac{1}{\lambda^0} \{d\varepsilon_{ij} - (1 - \lambda^0)d\varepsilon_{ij}^b - d\lambda(\varepsilon_{ij}^{f0} - \varepsilon_{ij}^{b0})\} \quad (17)$$

236 Substituting Eq. (17) into Eq. (16), we can obtain

$$237 \quad d\sigma_{ij}^f = \frac{1}{\lambda^0} D_{ijkl}^f \{d\epsilon_{kl} - (1 - \lambda^0)d\epsilon_{kl}^b - d\lambda(\epsilon_{kl}^{f0} - \epsilon_{kl}^{b0})\} \quad (18)$$

238 Combing Eq. (13) and Eq. (18), we can have the following equation expressed as

$$239 \quad d\sigma_{ij} = (1 - \lambda^0) \{D_{ijkl}^b - D_{ijkl}^f\}d\epsilon_{kl}^b + D_{ijkl}^f d\epsilon_{kl} - d\lambda D_{ijkl}^f \{\epsilon_{kl}^{f0} - \epsilon_{kl}^{b0}\} \\ + d\lambda \{\sigma_{ij}^{f0} - \sigma_{ij}^{b0}\} \quad (19)$$

240 We introduce the local strain coefficient C_{ijkl} to establish the relationship between the strain of bonded
241 elements and the average strain of RVE as follows:

$$242 \quad \epsilon_{ij}^b = C_{ijkl} \epsilon_{kl} \quad (20)$$

243 The incremental form of Eq. (20) is expressed as

$$244 \quad d\epsilon_{ij}^b = C_{ijkl}^0 d\epsilon_{kl} + dC_{ijkl} \epsilon_{kl}^0 \quad (21)$$

245 where C_{ijkl}^0 is the current local strain coefficient matrix. Substituting Eq. (21) into Eq. (19) with some
246 manipulation, we can obtain

$$247 \quad d\sigma_{ij} = \{(1 - \lambda^0) \{D_{ijmn}^b - D_{ijmn}^f\}C_{mnkl}^0 + D_{ijkl}^f\}d\epsilon_{kl} - d\lambda D_{ijkl}^f \{\epsilon_{kl}^{f0} - \epsilon_{kl}^{b0}\} \\ + d\lambda \{\sigma_{ij}^{f0} - \sigma_{ij}^{b0}\} + (1 - \lambda^0) \{D_{ijmn}^b - D_{ijmn}^f\}dC_{mnkl} \epsilon_{kl}^0 \quad (22)$$

248 For the current stress and strain states, from Eqs. (10) and (11), we can obtain

$$249 \quad \sigma_{ij}^{f0} = \frac{\sigma_{ij}^0 - (1 - \lambda^0)\sigma_{ij}^{b0}}{\lambda^0} \quad (23)$$

$$250 \quad \text{and} \quad \epsilon_{ij}^{f0} = \frac{\epsilon_{ij}^0 - (1 - \lambda^0)\epsilon_{ij}^{b0}}{\lambda^0} \quad (24)$$

251 where σ_{ij}^0 and ϵ_{ij}^0 are the current stress and strain of RVE, respectively.

252 Substitution Eqs. (23) and (24) into Eq. (22) with some manipulations, we can obtain the general stress-
253 strain relationship as follows:

$$254 \quad d\sigma_{ij} = \{(1 - \lambda^0) \{D_{ijmn}^b - D_{ijmn}^f\}C_{mnkl}^0 + D_{ijkl}^f\}d\epsilon_{kl} + (1 - \lambda^0) \{D_{ijmn}^b - D_{ijmn}^f\}dC_{mnkl} \epsilon_{kl}^0 \\ - \frac{d\lambda}{\lambda^0} D_{ijkl}^f \{\epsilon_{kl}^{f0} - \epsilon_{kl}^{b0}\} + \frac{d\lambda}{\lambda^0} \{\sigma_{ij}^0 - \sigma_{ij}^{b0}\} \quad (25)$$

255 At the initial loading, we have $\lambda^0 = 0$, $\{\epsilon\}^0 = 0$, $\{\epsilon\}_b^0 = 0$ and $\{\epsilon\}_f^0 = 0$, which can be substituted into Eq.

256 (22) to obtain the following stress expression at initial loading:

$$d\sigma_{ij} = \{(1 - \lambda^0) \{D_{ijmn}^b - D_{ijmn}^f\} C_{mnkl}^0 + D_{ijk}^f\} d\epsilon_{kl} + (1 - \lambda^0) \{D_{ijmn}^b - D_{ijmn}^f\} dC_{mnkl} \epsilon_{kl}^0 \quad (26)$$

In Eq. (25), there are four sets of parameters that must be determined, which include the constitutive relationship of bonded elements and frictional elements, breakage parameter and local strain matrix, which will be described in the following sections.

Constitutive Relationship of Bonded Elements

The bonded elements have bonding and large pores within them, whose behavior is similar to that of artificially structured soils at the initial loading within very small strain with almost intact structures. Natural soils are formed in layers by sedimentation, whose mechanical properties are isotropic in horizontal planes and different in horizontal and vertical directions. Therefore, we assume here that the bonded elements are cross-anisotropic elastic materials. When setting the symmetry axis along the z direction and the x axis and y axis in the horizontal plane, the stress–strain relationship, Eq. (15), of the bonded elements can be rewritten in Cartesian coordinates as follows:

$$\begin{matrix} 269 \\ \left. \begin{matrix} d\sigma_x \\ d\sigma_y \\ d\sigma_z \\ d\tau_{yz} \\ d\tau_{zx} \\ d\tau_{xy} \end{matrix} \right\}_b \end{matrix} = \begin{bmatrix} D_{11} & D_{12} & D_{13} & 0 & 0 & 0 \\ D_{12} & D_{11} & D_{13} & 0 & 0 & 0 \\ D_{13} & D_{13} & D_{33} & 0 & 0 & 0 \\ 0 & 0 & 0 & D_{44} & 0 & 0 \\ 0 & 0 & 0 & 0 & D_{44} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{(D_{11} - D_{12})}{2} \end{bmatrix} \begin{matrix} \left. \begin{matrix} d\epsilon_x \\ d\epsilon_y \\ d\epsilon_z \\ d\epsilon_{yz} \\ d\epsilon_{zx} \\ d\epsilon_{xy} \end{matrix} \right\}_b \\ (27) \end{matrix}$$

where the five material constants D_{11} , D_{12} , D_{13} , D_{33} and D_{44} can be determined by the stress–strain curves at the initial loading stage of the tested samples, during which the structured samples are hardly damaged and could be regarded as bonded elements. When $D_{11} = D_{33}$, $D_{12} = D_{13}$ and $D_{44} = (D_{11} - D_{12})/2$, Eq. (27) can be reduced to the stress–strain relationship of isotropic materials with two constants (Graham and Houlsby 1983).

Constitutive Relationship of Frictional Elements

The frictional elements are transformed from bonded elements when the bonds between soil particles are broken completely, whose mechanical properties could be assumed as those of remolded soils. From the test results of the remolded soils shown in Fig. 7 (a)–(b), we know that the stress–strain relationship of

279 frictional elements can be described by the Lade-Duncan model (Lade and Duncan 1975; Lade 1977). For
 280 the Lade-Duncan model, the incremental strain of soils consists of elastic and plastic components in matrix
 281 form as follows:

$$282 \quad \{\mathcal{d}\boldsymbol{\varepsilon}\}_f = \{\mathcal{d}\boldsymbol{\varepsilon}^e\}_f + \{\mathcal{d}\boldsymbol{\varepsilon}^p\}_f \quad (28)$$

283 where $\{\mathcal{d}\boldsymbol{\varepsilon}^e\}_f$ is the incremental elastic strain and $\{\mathcal{d}\boldsymbol{\varepsilon}^p\}_f$ is the incremental plastic strain.

284 According to the Lade-Duncan model, the elastic strain can be expressed as follows:

$$285 \quad \left\{ \begin{array}{c} \mathcal{d}\boldsymbol{\varepsilon}_x^e \\ \mathcal{d}\boldsymbol{\varepsilon}_y^e \\ \mathcal{d}\boldsymbol{\varepsilon}_z^e \\ \mathcal{d}\boldsymbol{\varepsilon}_{yz}^e \\ \mathcal{d}\boldsymbol{\varepsilon}_{zx}^e \\ \mathcal{d}\boldsymbol{\varepsilon}_{xy}^e \end{array} \right\}_f = \frac{1}{E_f} \left\{ \begin{array}{c} \mathcal{d}\sigma_x - \nu_f (\mathcal{d}\sigma_y + \mathcal{d}\sigma_z) \\ \mathcal{d}\sigma_y - \nu_f (\mathcal{d}\sigma_z + \mathcal{d}\sigma_x) \\ \mathcal{d}\sigma_z - \nu_f (\mathcal{d}\sigma_x + \mathcal{d}\sigma_y) \\ 2(1 + \nu_f) \mathcal{d}\tau_{yz} \\ 2(1 + \nu_f) \mathcal{d}\tau_{zx} \\ 2(1 + \nu_f) \mathcal{d}\tau_{xy} \end{array} \right\}_f \quad (29)$$

286 where E_f and ν_f are the tangential deformational modulus and tangential Poisson ratio of the remolded
 287 samples, respectively. In the Lade-Duncan model, the failure criterion is $f_1 = I_1^3 / I_3 = K_f$, the
 288 yielding function is $f = I_1^3 / I_3 = K_0$, and the plastic potential $g = I_1^3 - K_2 I_3$, where I_1 and I_3 are
 289 the first invariant and third invariant of stress, respectively, and $K_f = K_0$ at failure. Therefore, according to
 290 the hardening elasto-plastic theory, we can obtain the incremental plastic strain as follows:

$$291 \quad \left\{ \begin{array}{c} \mathcal{d}\boldsymbol{\varepsilon}_x^p \\ \mathcal{d}\boldsymbol{\varepsilon}_y^p \\ \mathcal{d}\boldsymbol{\varepsilon}_z^p \\ \mathcal{d}\boldsymbol{\varepsilon}_{yz}^p \\ \mathcal{d}\boldsymbol{\varepsilon}_{zx}^p \\ \mathcal{d}\boldsymbol{\varepsilon}_{xy}^p \end{array} \right\}_f = d\theta \cdot K_2 \left\{ \begin{array}{c} \frac{3I_1^2}{K_2} - \sigma_y \sigma_z + \tau_{yz}^2 \\ \frac{3I_1^2}{K_2} - \sigma_z \sigma_x + \tau_{zx}^2 \\ \frac{3I_1^2}{K_2} - \sigma_x \sigma_y + \tau_{xy}^2 \\ 2\sigma_x \tau_{yz} - 2\tau_{xy} \tau_{zx} \\ 2\sigma_y \tau_{zx} - 2\tau_{xy} \tau_{yz} \\ 2\sigma_z \tau_{xy} - 2\tau_{yz} \tau_{zx} \end{array} \right\}_f \quad (30)$$

292 where $d\mathcal{G}$ is the plastic multiplier and K_2 is the model constant. A detailed description of the Lade-Duncan
 293 model can be found in the literature (Lade and Duncan 1975; Lade 1977).

294 **Structural Parameters of Breakage Ratio and Local Strain Coefficient Matrix**

295 The breakage ratio λ is a structural parameter whose evolving rules are closely related to soil type, stress
 296 and strain level, stress path and history. At the initial stage of loading, λ is very small with a value close to
 297 zero for the external loads, which are mainly borne by the bonded elements. With the process of loading,
 298 λ increases gradually, accompanied by bonded elements transferring to frictional elements, both of which
 299 bear the external loading. When the strain is very large, λ tends to be 1.0, and the external loads are
 300 mainly borne by frictional elements at the moment. In view of the determination method of the damage
 301 factor and hardening parameters (Krajcinovic and Mastilovic 1995; Yu 2006), we assume that the breakage
 302 ratio λ is a function of volumetric strain and generalized shear strain with the following expression:

$$303 \quad \lambda = 1 - \exp(-\beta(\alpha\varepsilon_z + \varepsilon_x + \varepsilon_y)^\psi - (\xi\varepsilon_s)^\theta) \quad (31)$$

304 where $\varepsilon_s = \sqrt{2e_{ij}e_{ij}/3}$, $e_{ij} = \varepsilon_{ij} - \varepsilon_{kk}\delta_{ij}/3$, δ_{ij} is the Kronecker delta, and α , β , ξ , ψ
 305 and θ are material parameters, with the symmetry axis along the z direction and the x axis and y axis in
 306 the horizontal plane.

307 The local strain coefficient bridges the strains of bonded elements and RVE, which can vary in the process
 308 of loading and be affected by loading history and strain level. We assume here that in the elements the local
 309 strain coefficient are the same and are represented by C of a function of generalized shear strain as follows:

$$310 \quad C = \exp(-\left(t_c \times \varepsilon_s\right)^{r_c}) \quad (32)$$

311 where t_c and r_c are model parameters.

312 The breakage ratio and local strain coefficient are both internal variables, which should be determined by
 313 meso-mechanics at the mesoscale. However, it is very difficult to determine the meso parameters for
 314 structured soils, so here we establish their evolving relationships using similar determination methods of
 315 hardening parameters in plasticity or damage factors in damage mechanics. Based on the analysis of the
 316 breakage mechanism of artificially structured soils from mesoscale to macroscale, we formulate their
 317 expressions in which those model parameters could be determined by test results.

318 **Determination of Model Parameters under Triaxial Stress Conditions**

319 Under conventional triaxial stress conditions in which two types of soil samples including initially isotropic
 320 structured and the remolded samples previously mentioned are tested, the vertical direction is set as the z
 321 axial direction, along which the maximal principal stress is applied, and the other two principal stresses are
 322 applied in the horizontal plane. Combining the test results provided above, we present the determination
 323 method of the model parameters under triaxial stress conditions in the following sections.

324 **(a) Parameter Determination for Bonded Elements**

325 Under conventional triaxial stress conditions, the stress–strain relationship of bonded elements, Eq. (27),
 326 can be simplified as follows:

$$327 \quad \begin{Bmatrix} d\sigma_1 \\ d\sigma_3 \end{Bmatrix}_b = \frac{E_{vb}}{(1 - \nu_{hbb})E_{vb} - 2\nu_{vbb}^2 E_{hb}} \begin{bmatrix} (1 - \nu_{hbb})E_{vb} & 2\nu_{vbb}E_{hb} \\ \nu_{vbb}E_{hb} & E_{hb} \end{bmatrix} \begin{Bmatrix} d\epsilon_1 \\ d\epsilon_3 \end{Bmatrix}_b \quad (33)$$

328 where there are four material parameters, E_{vb} , E_{hb} , ν_{vbb} and ν_{hbb} , where E_{vb} and E_{hb} represent the
 329 elastic moduli of bonded elements in the vertical and horizontal directions, respectively, and ν_{vbb} and ν_{hbb}
 330 represent the Poisson ratios of bonded elements in the vertical and horizontal directions, respectively.

331 Within a small strain range upon initial loading, there are mainly bonded elements in RVE to bear the
 332 external loads, so the stress–strain curve of the structured samples can be very similar to that of bonded
 333 elements. Here, we use the strain of 0.25% of the artificially structured samples tested to determine E_{vb} ,
 334 E_{hb} , ν_{vbb} and ν_{hbb} . Using Eq. (33), we can solve for only the values of E_{vb} and ν_{vbb} . For initially
 335 stress-induced anisotropic structured samples, when $E_{vb} = E_{hb}$ and $\nu_{vbb} = \nu_{hbb}$, they become initially
 336 isotropic structured samples. E_{vb} , E_{hb} , ν_{vbb} and ν_{hbb} are functions of confining pressure σ_3 expressed

$$337 \quad \text{as } E_{vb} \text{ (or } E_{hb}) = b_1 \ln\left(\frac{\sigma_3}{p_a}\right) + b_2 \text{ and } \nu_{hbb} \text{ (or } \nu_{vbb}) = b_3 \left(\frac{\sigma_3}{p_a}\right)^{b_4}, \text{ where } b_1, b_2, b_3, \text{ and } b_4 \text{ are}$$

338 material constants, and p_a is the atmospheric pressure of 0.1014 MPa.

339 **(b) Parameter Determination for Frictional Elements**

340 Frictional elements are transferred from bonded elements and bonding between soil particles that are fully
 341 breaking up, whose mechanical properties are similar to those of remolded soils and can be described by

342 the Lade-Duncan model (Lade and Duncan 1975; Lade 1977) as mentioned above. Based on the test results
 343 of the remolded soils, we give the parameters of the Lade-Duncan model here.

344 Under conventional triaxial stress conditions, Eq. (28) can be rewritten as follows:

$$345 \quad \begin{Bmatrix} d\sigma_1 \\ d\sigma_3 \end{Bmatrix}_f = [D]_f^{pp} \begin{Bmatrix} d\varepsilon_1 \\ d\varepsilon_3 \end{Bmatrix}_f \quad (34)$$

346 where $[D]_f^{pp}$ is the stiffness matrix of the frictional elements. According to the Lade-Duncan model, the
 347 elastic parameters of E_f and ν_f can be determined by the nonlinear elastic model of the Duncan-Chang
 348 hyperbolic model (Lade and Duncan 1975; Lade 1977) as follows:

$$349 \quad E_f = K p_a \left(\frac{\sigma_3}{p_a} \right)^n \left[1 - \frac{R_f (\sigma_1 - \sigma_3) (1 - \sin \varphi)}{2c \cos \varphi + 2\sigma_3 \sin \varphi} \right]^2 \quad (35)$$

$$350 \quad \text{and} \quad \nu_f = \frac{G - F \lg(\sigma_3 / p_a)}{\left\{ 1 - \frac{D(\sigma_1 - \sigma_3)}{K p_a \left(\frac{\sigma_3}{p_a} \right)^n \left[1 - \frac{R_f (\sigma_1 - \sigma_3) (1 - \sin \varphi)}{2c \cos \varphi + 2\sigma_3 \sin \varphi} \right]} \right\}^2} \quad (36)$$

351 where K, n, R_f, G, F and D are material constants, and c and φ are the cohesion and internal frictional
 352 angles of the remolded soils, respectively; the stress is that of the frictional elements.

353 By setting $m_1 = \frac{E_f (1 - \nu_f)}{(1 + \nu_f)(1 - 2\nu_f)}$, we can present $[D]_f^{pp}$ as follows:

$$354 \quad [D]_f^{pp} = \begin{bmatrix} m_1 - \frac{n_3}{n_9} & \frac{2m_1 \nu_f}{1 - \nu_f} - \frac{n_4}{n_9} \\ \frac{m_1 \nu_f}{1 - \nu_f} - \frac{n_5}{n_9} & \frac{m_1}{1 - \nu_f} - \frac{n_6}{n_9} \end{bmatrix} \quad (37)$$

355 where

$$356 \quad n_1 = m_1 (3I_1^3 - K_2 \sigma_3^2 + \frac{2\nu_f}{1 - \nu_f} (3I_1^2 - K_2 \sigma_1 \sigma_3)) \quad (38-1)$$

$$357 \quad n_2 = m_1 \left(\frac{\nu_f}{1 - \nu_f} (3I_1^3 - K_2 \sigma_3^2) + \frac{1}{1 - \nu_f} (3I_1^2 - K_2 \sigma_1 \sigma_3) \right) \quad (38-2)$$

$$n_3 = \frac{m_1^2 \times n_1}{I_3^2} \left[(3I_1^3 I_3 - \sigma_3^2 I_1^3) + \frac{\nu_f}{1-\nu_f} (3I_1^3 I_3 - \sigma_1 \sigma_3 I_1^3) \right] \quad (38-3)$$

$$n_4 = \frac{m_1^2 \times n_1}{I_3^2} \left[\frac{2\nu_f}{1-\nu_f} (3I_1^3 I_3 - \sigma_3^2 I_1^3) + \frac{1}{1-\nu_f} (3I_1^3 I_3 - \sigma_1 \sigma_3 I_1^3) \right] \quad (38-4)$$

$$n_5 = \frac{m_1^2 \times n_2}{I_3^2} \left[(3I_1^3 I_3 - \sigma_3^2 I_1^3) + \frac{\nu_f}{1-\nu_f} (3I_1^3 I_3 - \sigma_1 \sigma_3 I_1^3) \right] \quad (38-5)$$

$$n_6 = \frac{m_1^2 \times n_2}{I_3^2} \left[\frac{2\nu_f}{1-\nu_f} (3I_1^3 I_3 - \sigma_3^2 I_1^3) + \frac{1}{1-\nu_f} (3I_1^3 I_3 - \sigma_1 \sigma_3 I_1^3) \right] \quad (38-6)$$

$$n_7 = \frac{m_1}{I_3^2} \left[(3I_1^3 I_3 - \sigma_3^2 I_1^3) + \frac{\nu_f}{1-\nu_f} (3I_1^3 I_3 - \sigma_1 \sigma_3 I_1^3) \right] \quad (38-7)$$

$$n_8 = \frac{m_1}{I_3^2} \left[\frac{2\nu_f}{1-\nu_f} (3I_1^3 I_3 - \sigma_3^2 I_1^3) + \frac{1}{1-\nu_f} (3I_1^3 I_3 - \sigma_1 \sigma_3 I_1^3) \right] \quad (38-8)$$

$$n_9 = \frac{[1 - \beta'(f - f_t)] \sigma_3 m_1}{\alpha'} \frac{m_1}{I_3^2} \left[n_7 (3I_1^2 - K_2 \sigma_3^2) + n_8 (3I_1^2 - K_2 \sigma_1 \sigma_3) \right] \quad (38-9)$$

and the stresses in these expressions are those of the frictional elements. K_2 and stress level f have the relationship shown in Fig. 10, which can be expressed as follows:

$$K_2 = Af + 27(1-A) \quad (39)$$

where A is the materials constant and f has the relationship with the plastic work as shown in Fig. 11, which can be expressed as follows:

$$f - f_t = \frac{W_p}{\alpha' + \beta' W_p} \quad (40)$$

where $f_t=27$ for the remolded soils tested, and α' , β' are model parameters. Under conventional triaxial

stress conditions, we have $f = \frac{[(\sigma_1 - \sigma_3) + 3\sigma_3]^3}{[(\sigma_1 - \sigma_3) + \sigma_3] \sigma_3^2}$ and $W_p = \int \sigma_{ij} d\varepsilon_{ij}^p$. Substituting f , f_t and W_p

into Eq. (40), we can obtain $\beta' = 0.01$ and α' varying with the confining pressure as

$$\alpha' = r_1 \left(\frac{\sigma_3}{p_a} \right) + r_2, \text{ where } r_1 \text{ and } r_2 \text{ are material constants.}$$

(c) Parameter Determination for Structural Parameters

376 Under conventional triaxial stress conditions, the breakage ratio λ of Eq. (31) can be written as

$$377 \quad \lambda = 1 - \exp \left\{ -\beta (\alpha \varepsilon_1 + 2\varepsilon_3)^\psi - \left(\xi \frac{2}{3} (\varepsilon_1 - \varepsilon_3) \right)^\theta \right\} \quad (41)$$

378 For the artificially structured soils, the parameters β and ψ are constants; α , ξ and θ vary with the

379 confining pressure by α (or ξ, θ) = $e_1 \left(\frac{\sigma_3}{p_a} \right)^{e_2}$, where e_1 and e_2 are constants.

380 The local strain coefficient of C in Eq. (32) can be expressed under triaxial stress conditions as follows:

$$381 \quad C = \exp \left\{ - \left(t_c \frac{2}{3} (\varepsilon_1 - \varepsilon_3) \right) \right\} \quad (42)$$

382 where $t_c = s_1 \left(\frac{\sigma_3}{p_a} \right) + s_2$, and s_1 and s_2 are constants.

383 **Model Verification**

384 There are four sets of parameters, including those of bonded elements, frictional elements, and structural
385 parameters of breakage ratio and local strain coefficient, that must be provided in the proposed binary-
386 medium constitutive model for artificially structured soils. These model parameters are determined for the
387 samples tested as explained in Section “Test Conditions and Results” as follows.

388 For the bonded elements, the parameters are obtained as follows: $b_1=9.8383$ and $b_2=30.37$ for E_{vb} ,
389 $b_1=9.1511$ and $b_2=28.61$ for E_{hb} , $b_3=0.2134$ and $b_4=-0.41$ for v_{hbb} , and $b_3=0.1389$ and $b_4=-0.668$ for v_{vbb} .

390 For the frictional elements, the parameters are obtained as follows: $K=88.797$, $n=0.3425$, $R_f=0.95$,
391 $G=0.242$, $F=0.313$, $D=0.0113$, $c=0$, $\varphi=32.062^\circ$ and $A=0.3535$; $r_1=-14.0$, $r_2=-10.0$ when $\sigma_3 < 100$ kPa and
392 $r_1=-155.0$, $r_2=-66.67$ when $\sigma_3 \geq 100$ kPa.

393 For the structural parameters, ψ is 1.0 and $\beta=0.4$ at $\sigma_3 < 100$ kPa and $\beta=0.5$ at $\sigma_3 \geq 100$ kPa; when
394 determining α , $e_1=100.55$, $e_2=0.1135$; when determining ξ , $e_1=40.56$, $e_2=40.0$ at $\sigma_3 < 100$ kPa and $e_1=2.535$,
395 $e_2=100.0$ at $\sigma_3 \geq 100$ kPa; when determining θ , $e_1=0.0$, $e_2=0.15$ at $\sigma_3 < 100$ kPa and $e_1=0.0435$, $e_2=0.325$ at
396 $\sigma_3 \geq 100$ kPa; and $s_1=11.859$, $s_2=30.854$.

397 The curves of deviatoric stress–axial strain and volumetric strain–axial strain of artificially structured soils
398 computed and tested are shown in Fig. 12 and Fig. 13. From the deviatoric stress–axial strain curves shown

399 in Fig. 12 (a) and Fig. 13 (a), although there are some slight differences in the values computed and tested,
400 the proposed constitutive model can reflect the deformational features of artificially structured soils. At low
401 confining pressures of 25 kPa, 37.5 kPa and 50 kPa, the computed results exhibit strain-softening behavior,
402 which is in agreement with tested soils and whose peak values are very close to those of the tested results;
403 at 100 kPa of confining pressure, both the computed and tested deviatoric stresses reach the plastic flow
404 state simultaneously; at high confining pressures of 200 kPa and 400 kPa, the computed results exhibit
405 strain-hardening behavior, which is also in agreement with the tested soils. From the volumetric strain-
406 axial strain curves shown in Fig. 12 (b) and Fig. 13 (b), the computed results have similar properties to
407 those of the tested soils. At low confining pressures of 25 kPa, 37.5 kPa and 50 kPa, the computed
408 volumetric strains first contract and then dilate, with slightly larger values of contraction than those of the
409 tested soils and very close dilatancy at failure; at high confining pressures of 100 kPa, 200 kPa and 400 kPa,
410 both results computed contract continuously until failure, which agrees with the tested results with slight
411 differences in values.

412 **Discussions**

413 The performance of the model for zero breakage states and completely broken states are discussed here. For
414 zero breakage states, the bonded elements are assumed to be elastic state in the paper and bear the external
415 loading. Therefore, the structured soil sample can be represented by the bonded elements for zero breakage
416 states. When determining the parameters of the bonded elements, the artificially structured soils at the
417 initial loading within very small strain (e.g. 0.25% axial strain) are used to assure that the bonds between
418 soil particles are in elastic state and not broken. For completely broken states, the bonded elements are
419 wholly broken and transformed into frictional elements. Therefore, the structured soil sample can be
420 represented by the frictional elements for completely broken states which bear the external loading. When
421 determining the parameters of the frictional elements, the remolded soil sample prepared by remolding the
422 artificially structured sample tested with dried and sieved through a 0.5 mm screen are used to assure that
423 the bonds between soil particles are completely broken. For the micromechanical model for structured soil
424 proposed here, the structured soil sample at failure usually consists of two components or binary media of
425 bonded elements and frictional elements, and at failure the frictional elements dominate.

426 The relation between the proposed model and the bonded materials under the DSC are discussed here. In
427 the references of Desai (2001) and Liu et al. (2000), it is assumed that the RI represents “zero strain
428 state,” i.e., it is characterized as a perfectly rigid material. In the paper, however, the bonded elements are
429 assumed to be elastic materials and can be transformed to be frictional elements denoted by the evolution of
430 breakage ratio. For structured or cemented materials, Desai and coworkers (Desai 2001; Liu et al. 2000)
431 only presented the constitutive model in one-dimensional formulation. In the paper, however, we give the
432 generalized stress-strain equation for artificially structured soils and can be verified in triaxial tested results
433 of artificially structured soil samples. And thus, the model proposed here is based on the disturbed state
434 concept (DSC) and homogenization theory.

435 **Conclusions**

436 Artificially structured soil samples are tested under consolidated-drained conditions at confining pressures
437 of 25 kPa, 37.5 kPa, 50 kPa, 100 kPa, 200 kPa and 400 kPa. Based on these test results, a binary-medium
438 constitutive model for artificially structured soils is proposed in the manuscript. The conclusions can be
439 drawn as follows.

440 (i) At low confining pressures of 25 kPa, 37.5 kPa and 50 kPa, the artificially structured soil samples
441 exhibit strain-softening behavior and first contracts followed by dilatancy accompanying shear bands at
442 failure; at 100 kPa confining pressure, the deviatoric stress increases gradually and reaches a plastic flow
443 state and contracts during shear with a bulge in the middle at failure; at high confining pressures of 200 kPa
444 and 400 kPa, all samples exhibit strain-hardening behavior and contract with a bulge in the middle at
445 failure.

446 (ii) The new constitutive model, the binary-medium constitutive model proposed here for artificially
447 structured soils, idealizes the structured samples as compositions of bonded elements described by elastic
448 materials and frictional elements described by the Lade-Duncan model, whose distribution of stress and
449 strain can be considered by introducing a local strain coefficient and breakage ratio. The computed results
450 compared with the tested ones demonstrate that the new model can grasp the main mechanical properties of
451 artificially structures soils including strain-softening and contraction followed by dilatancy at low confining
452 pressures and strain-hardening and continuous contraction at high confining pressures.

453 **Acknowledgments**

454 The authors thank the reviewers and Editor very much for their comments and appreciate the financial
455 support from the National Natural Science Foundation of China (NSFC) (Grant No. 51009103 and
456 51579167).

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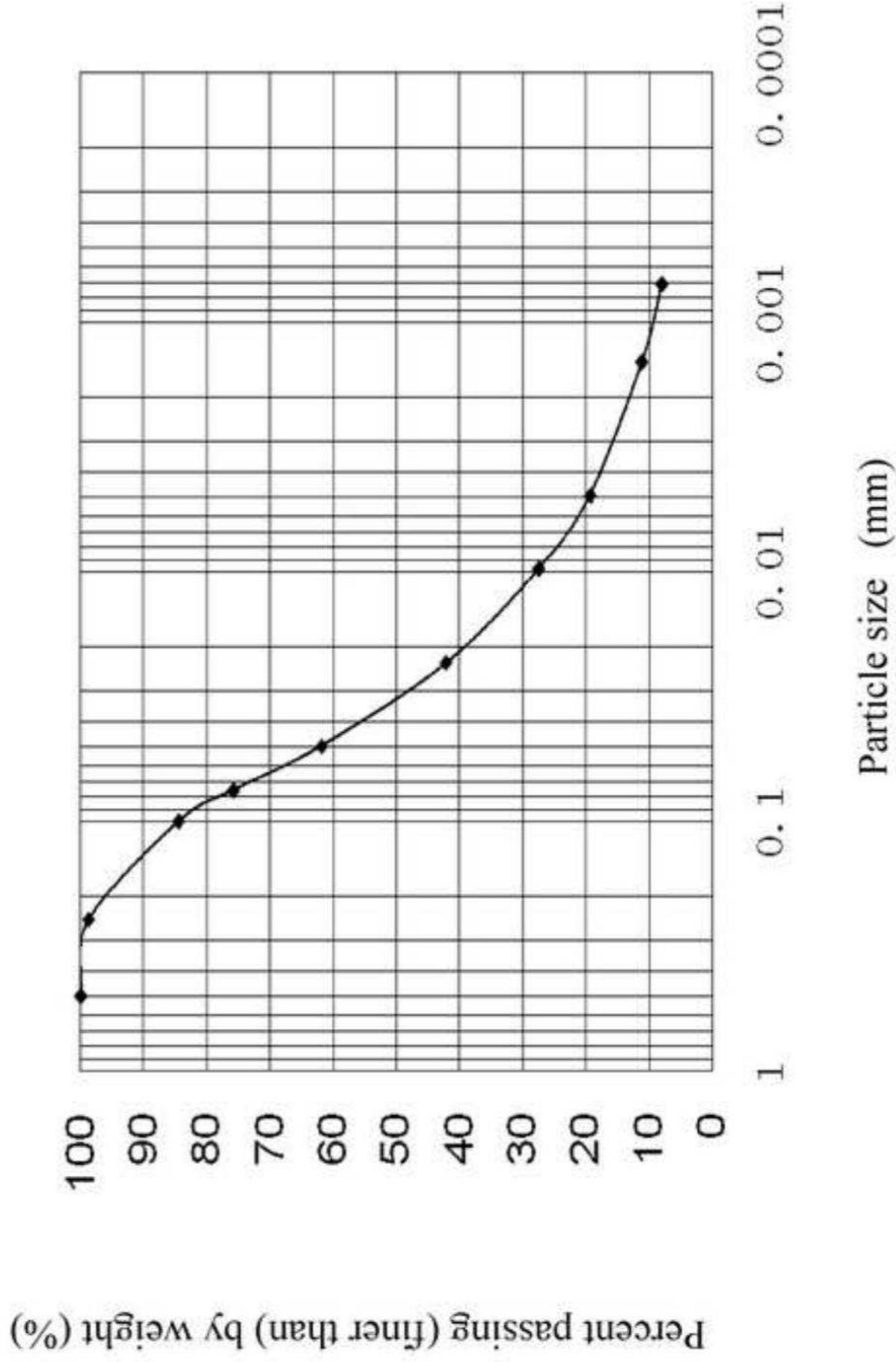


Fig. 1 Particle size distribution curve

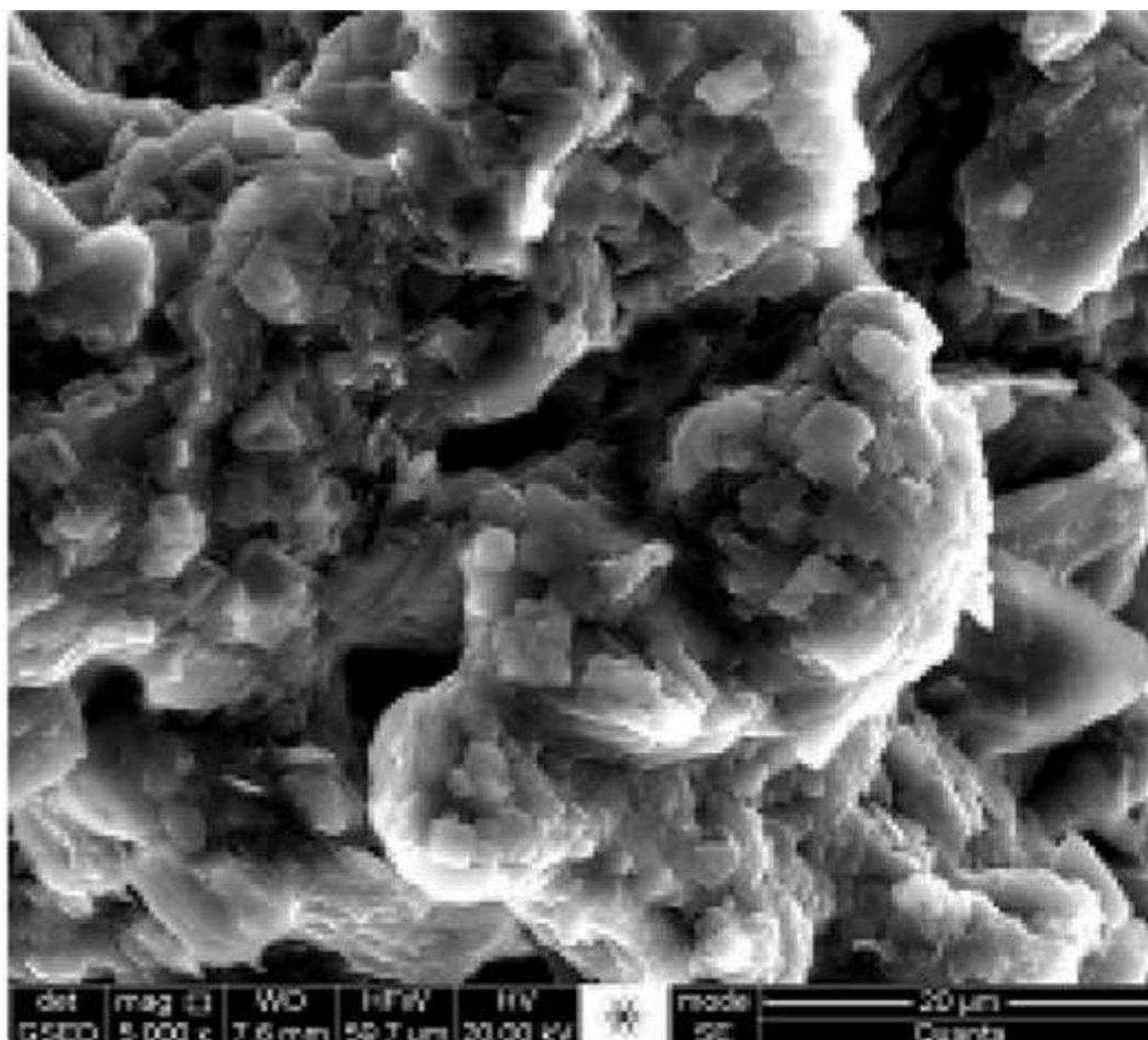


Fig. 2 SEM of artificially structured soils

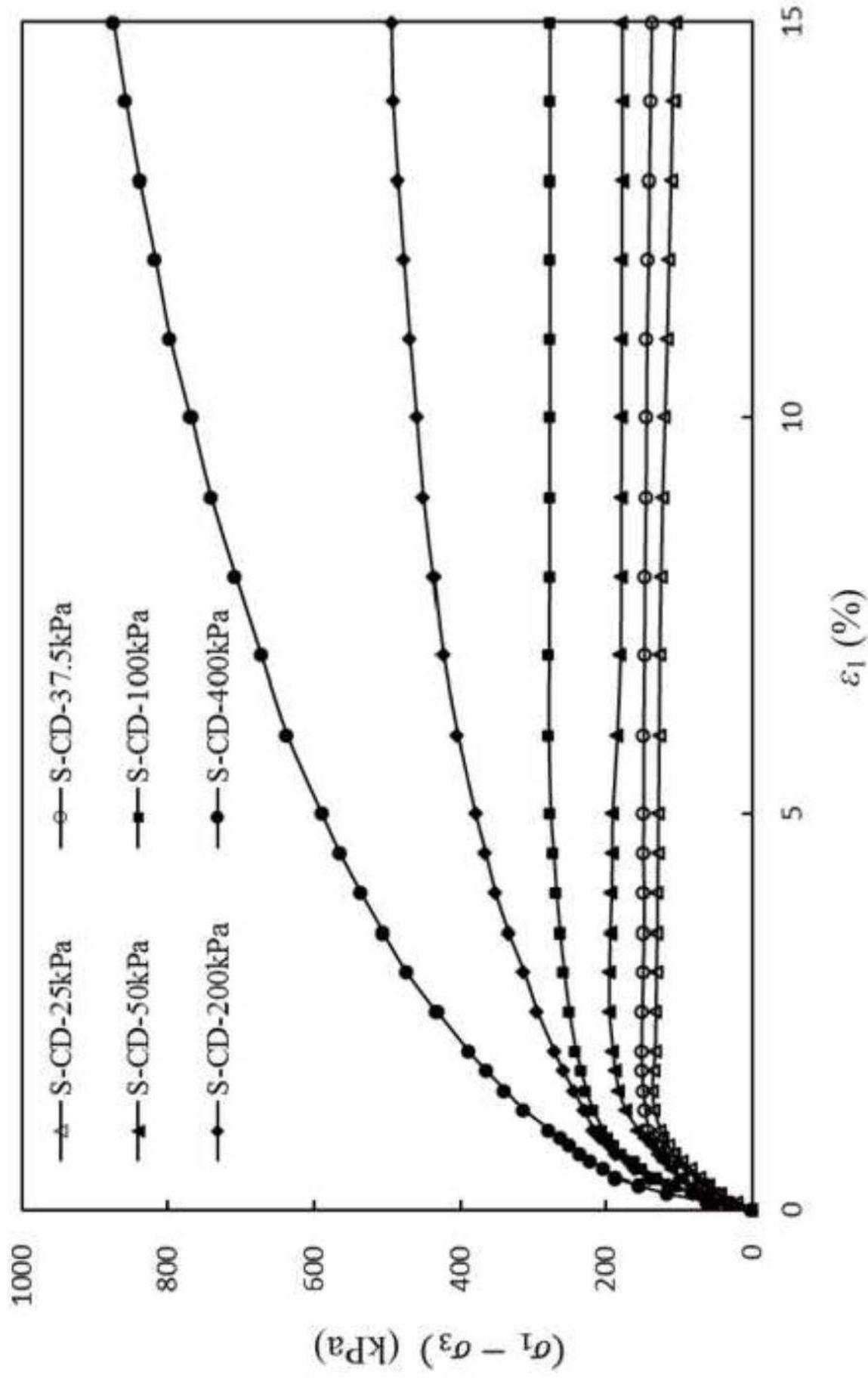


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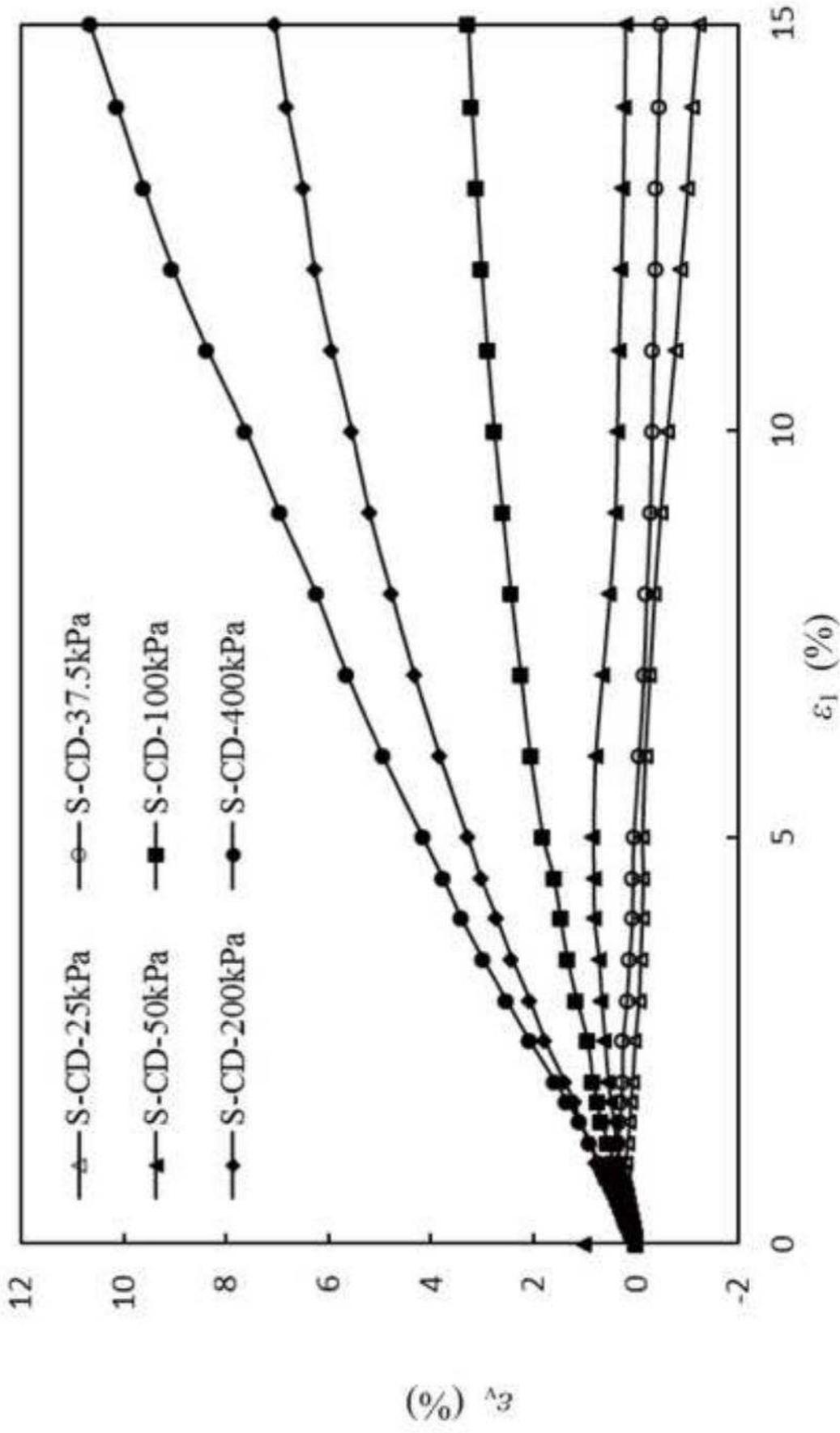


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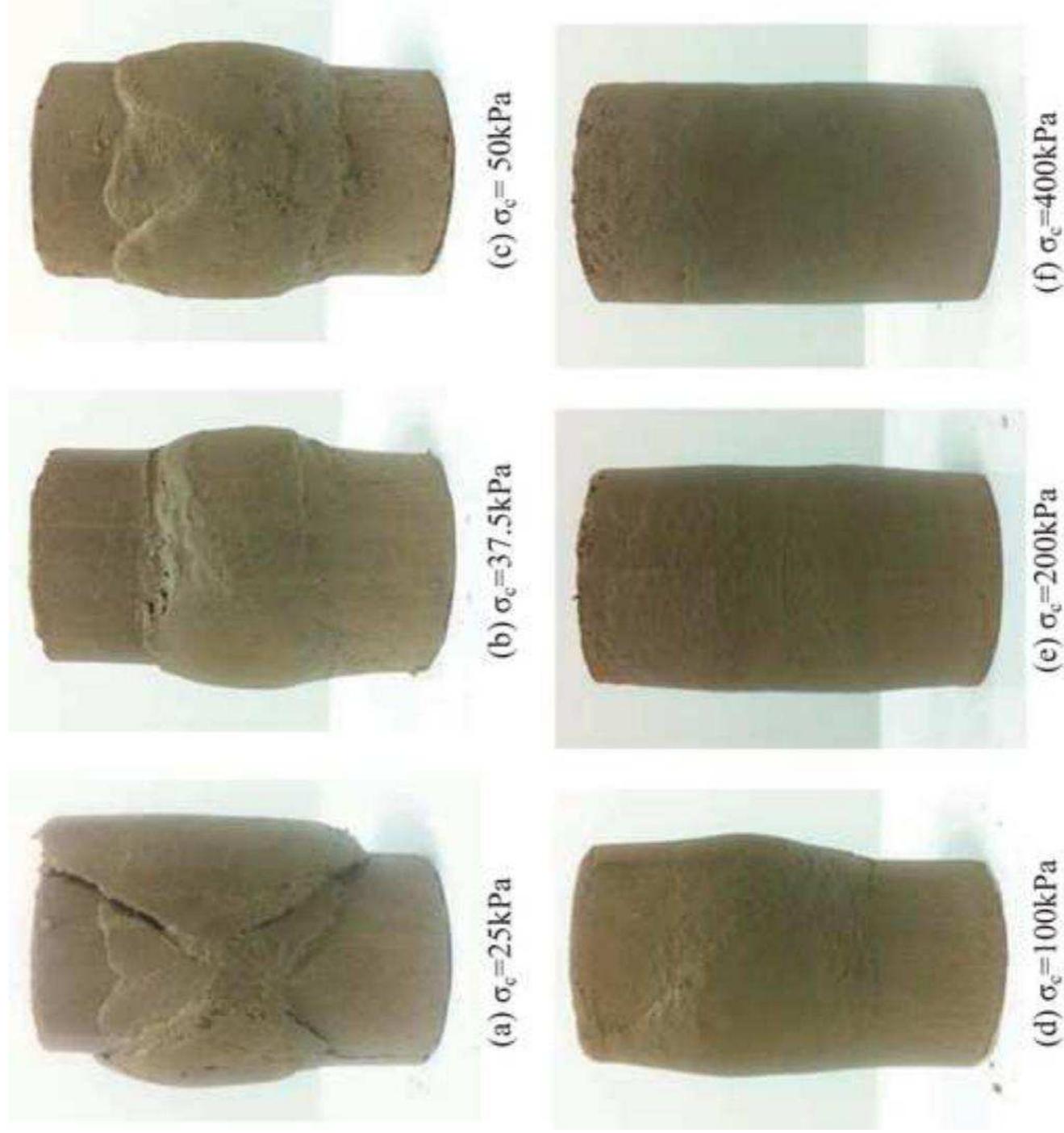


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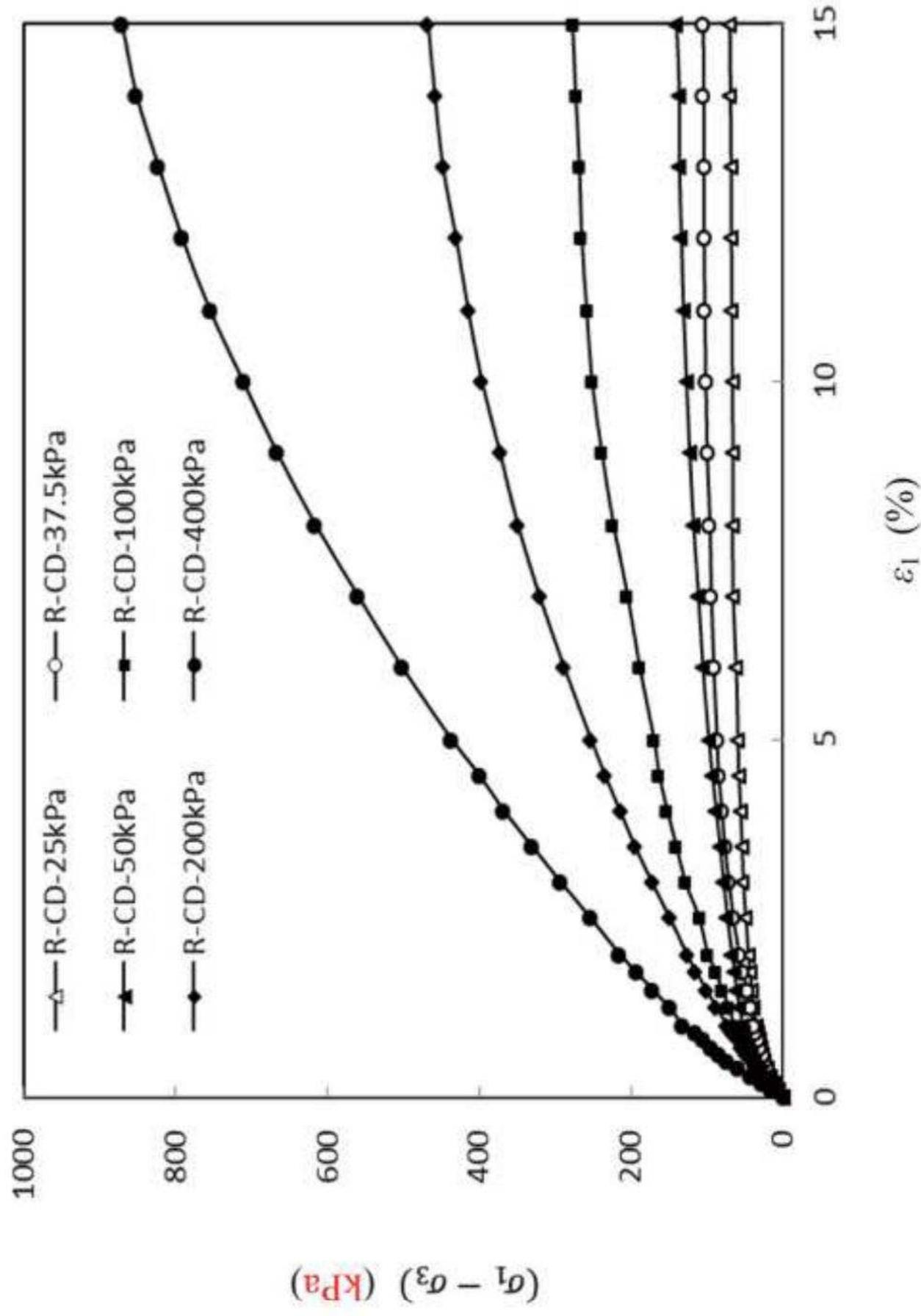


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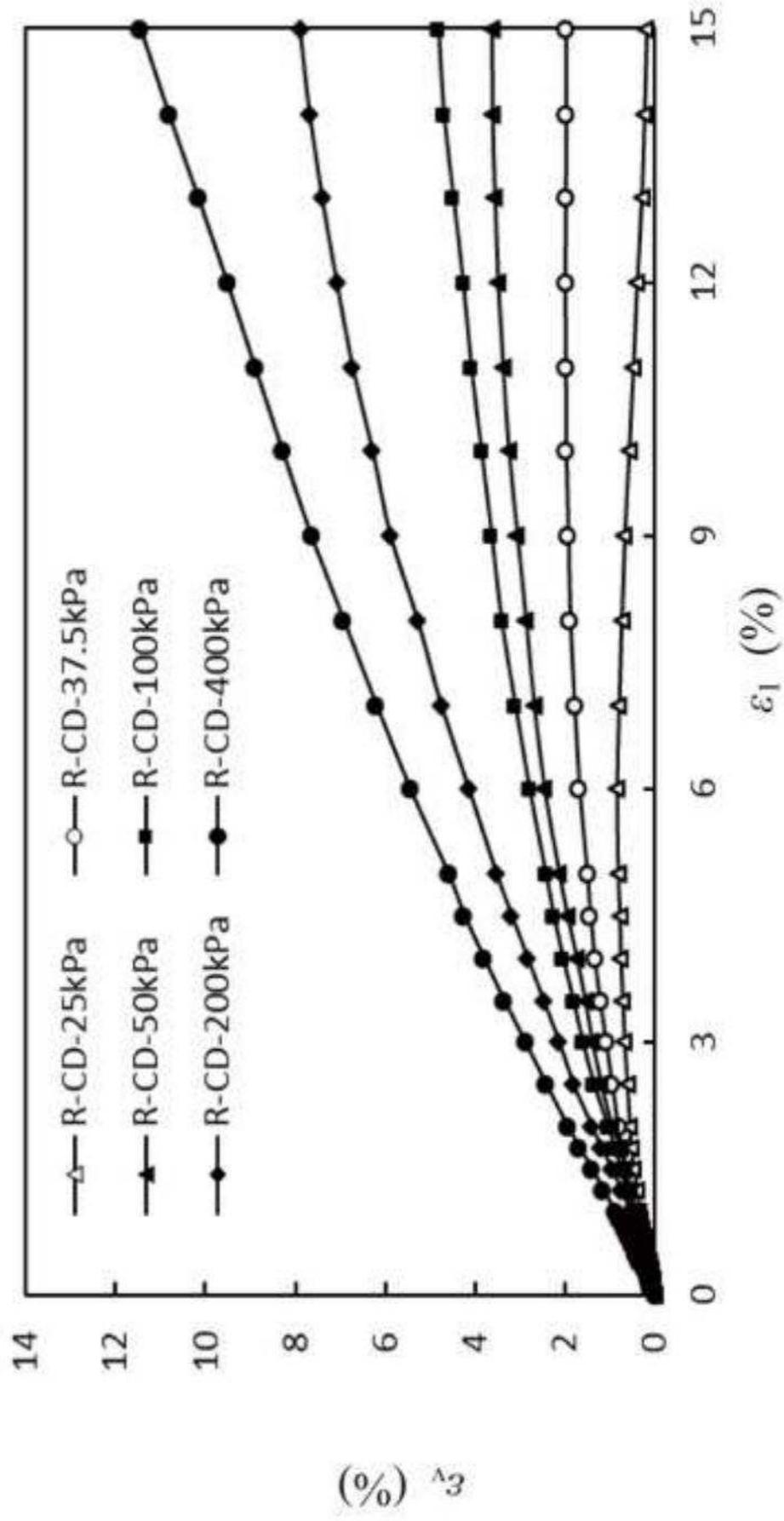


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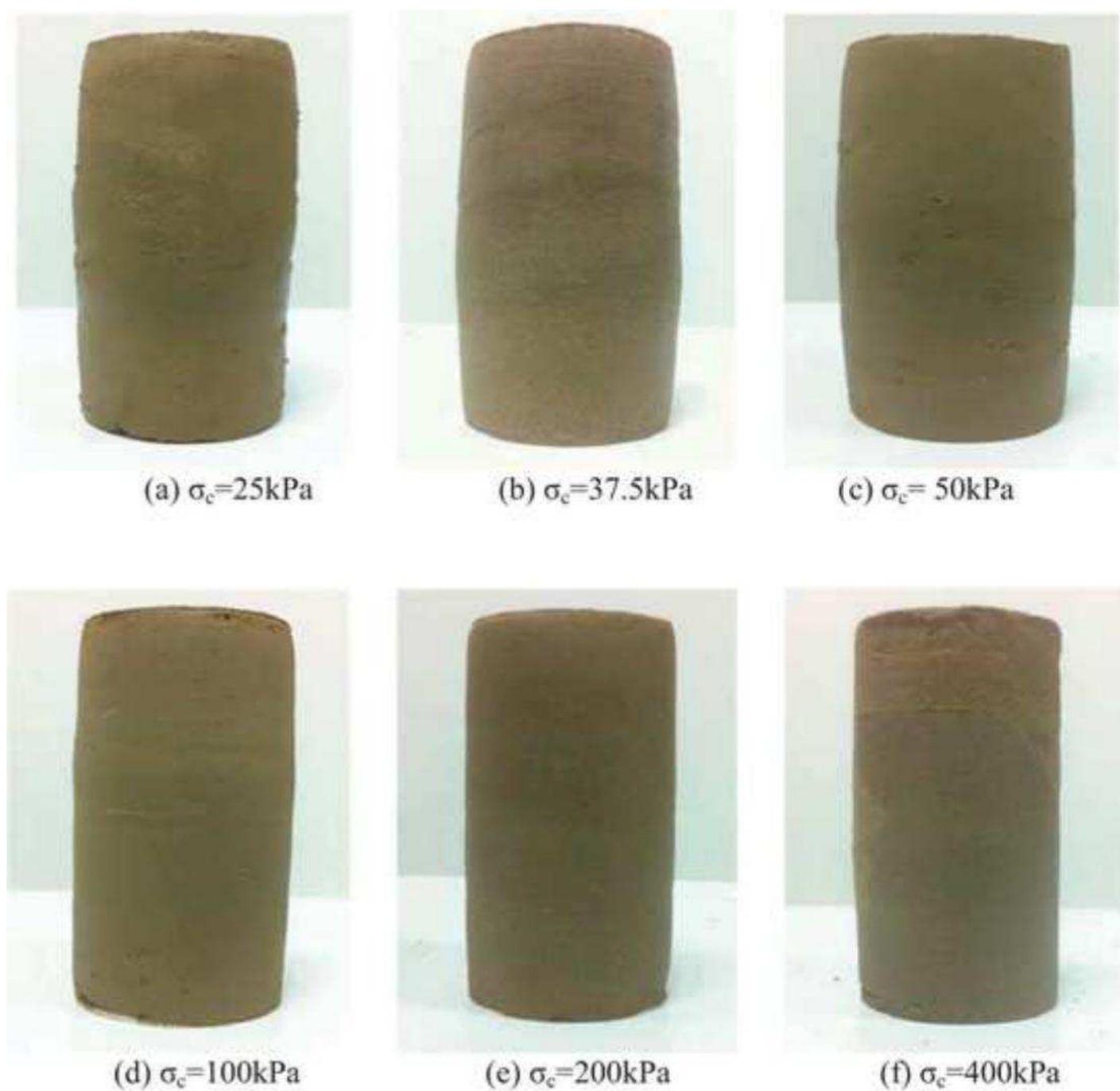


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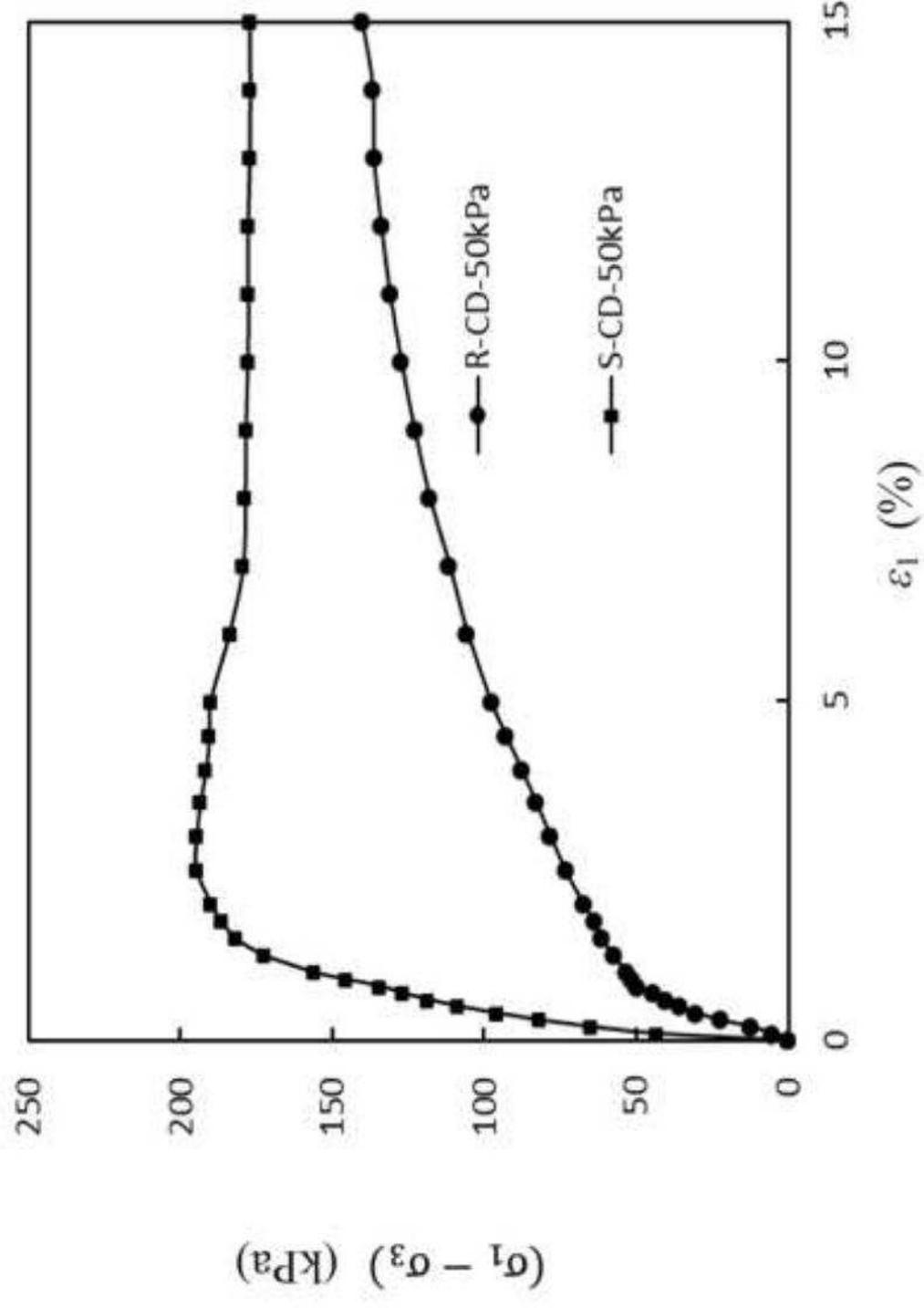


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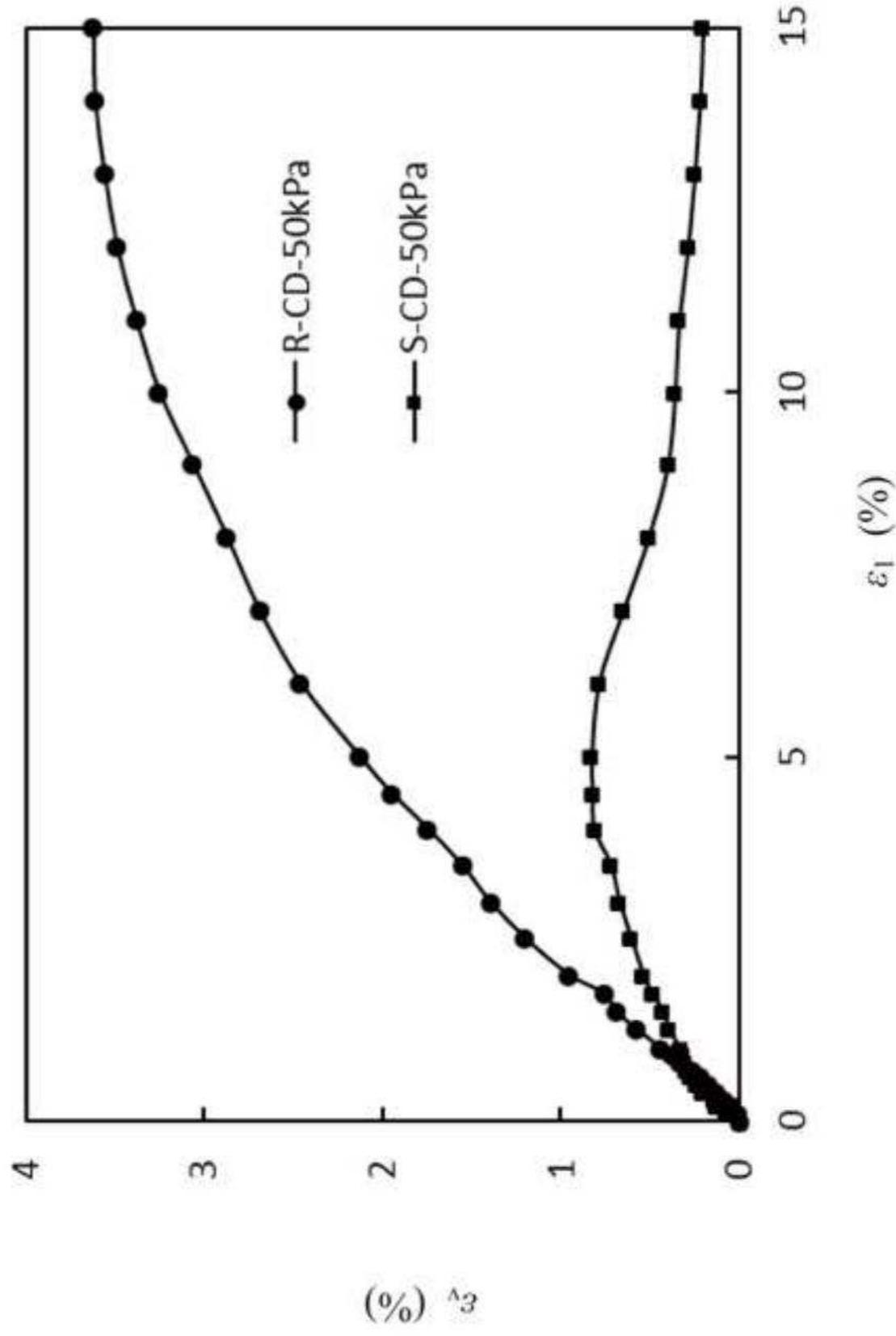


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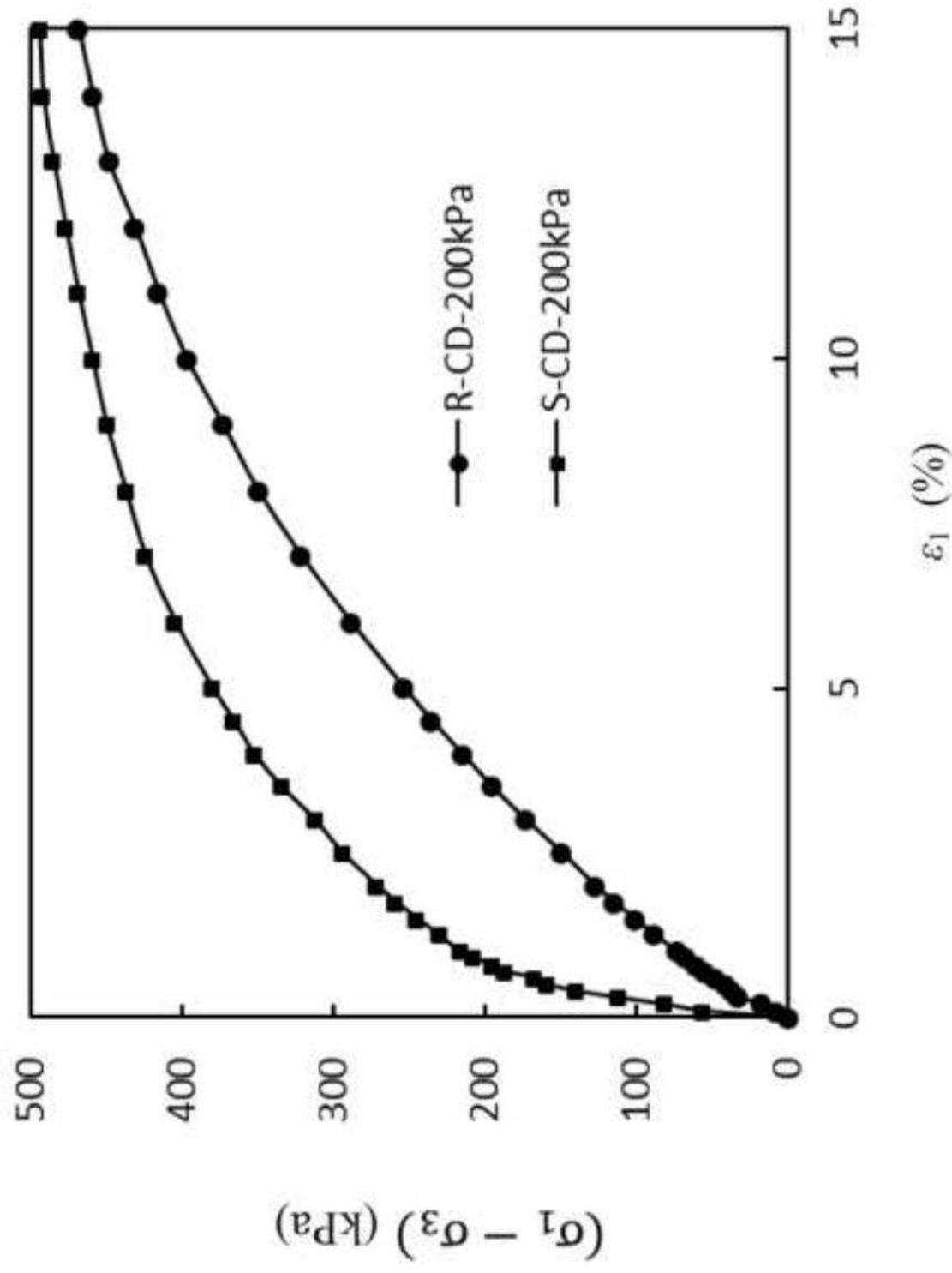


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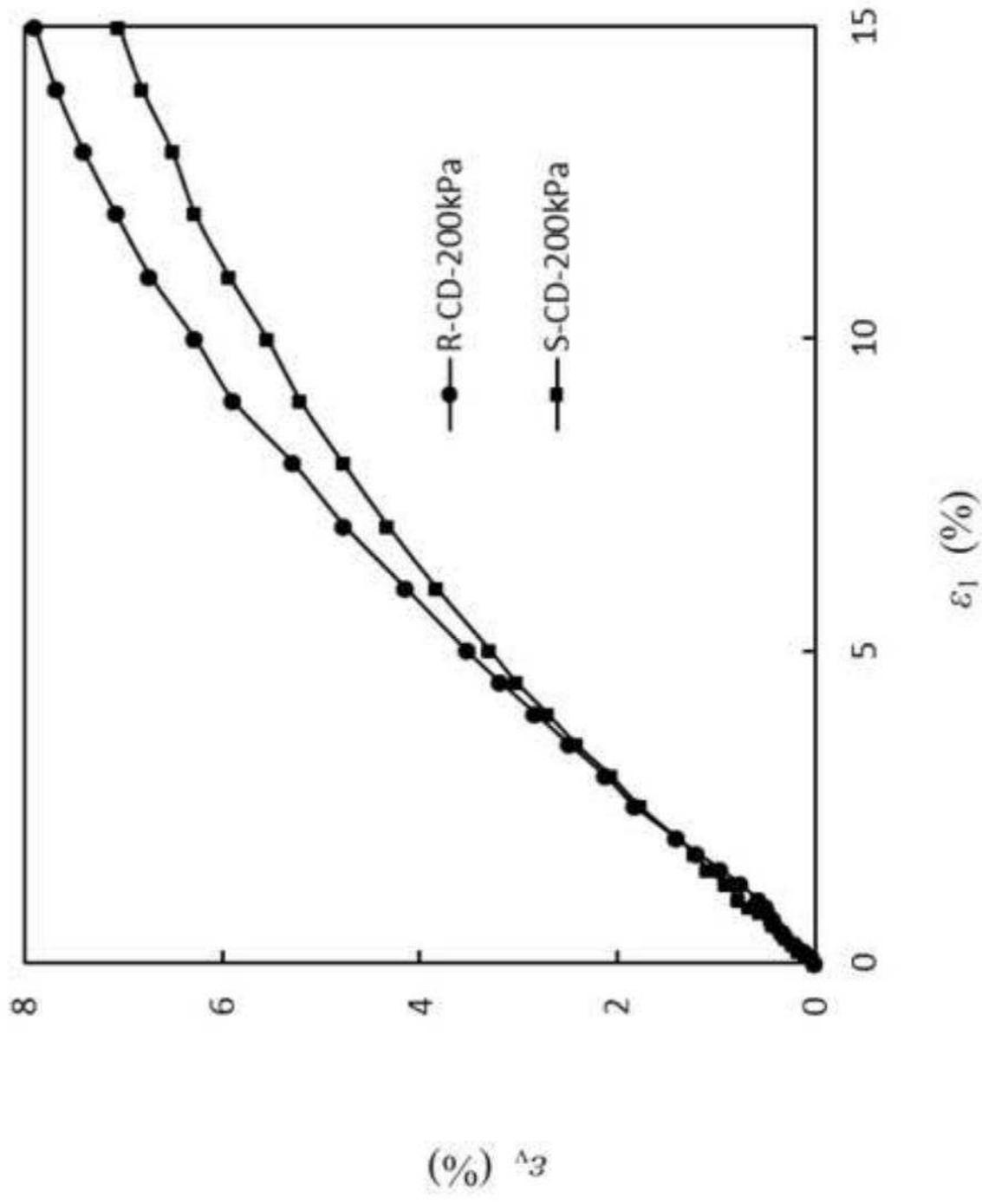


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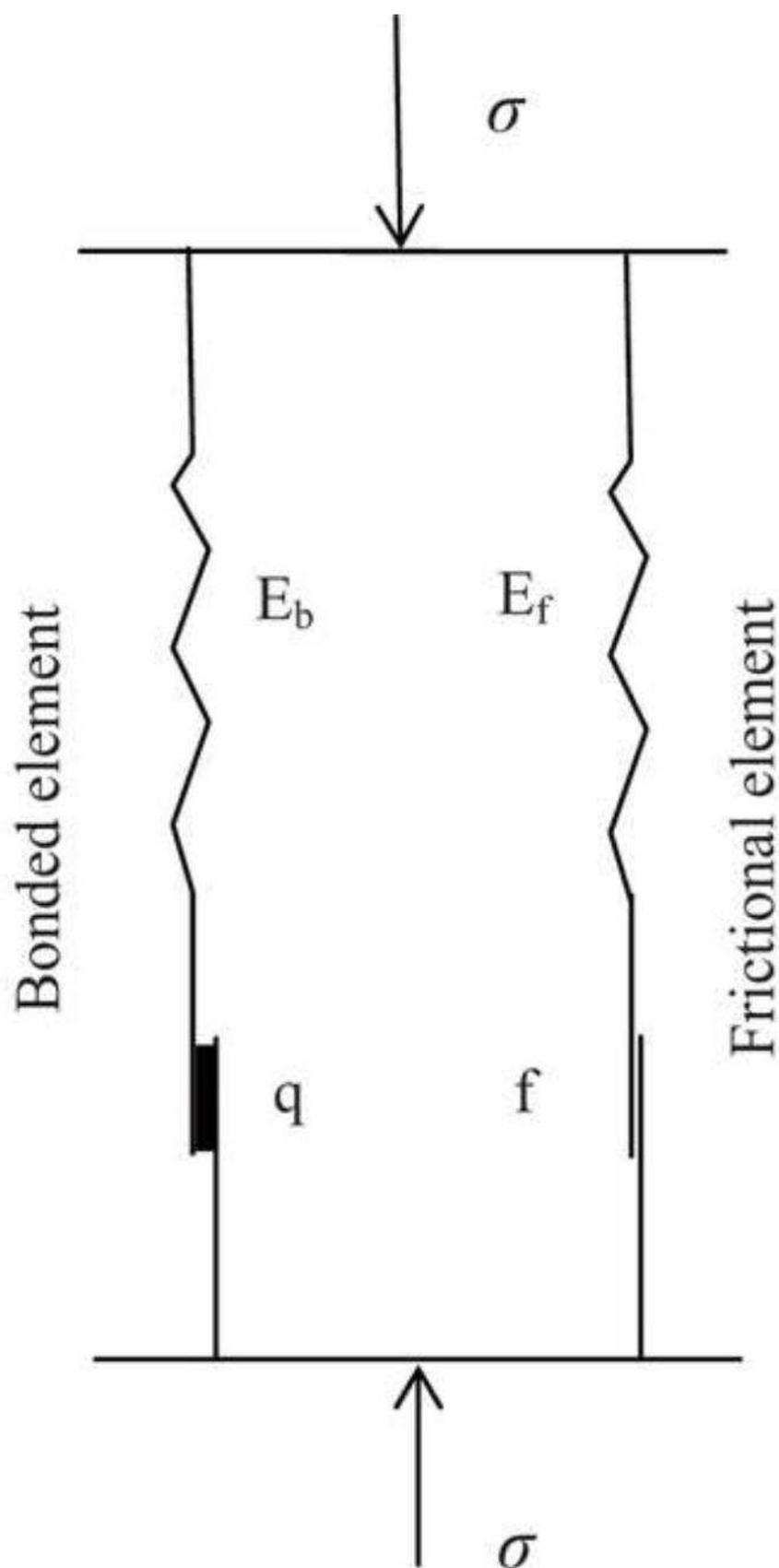


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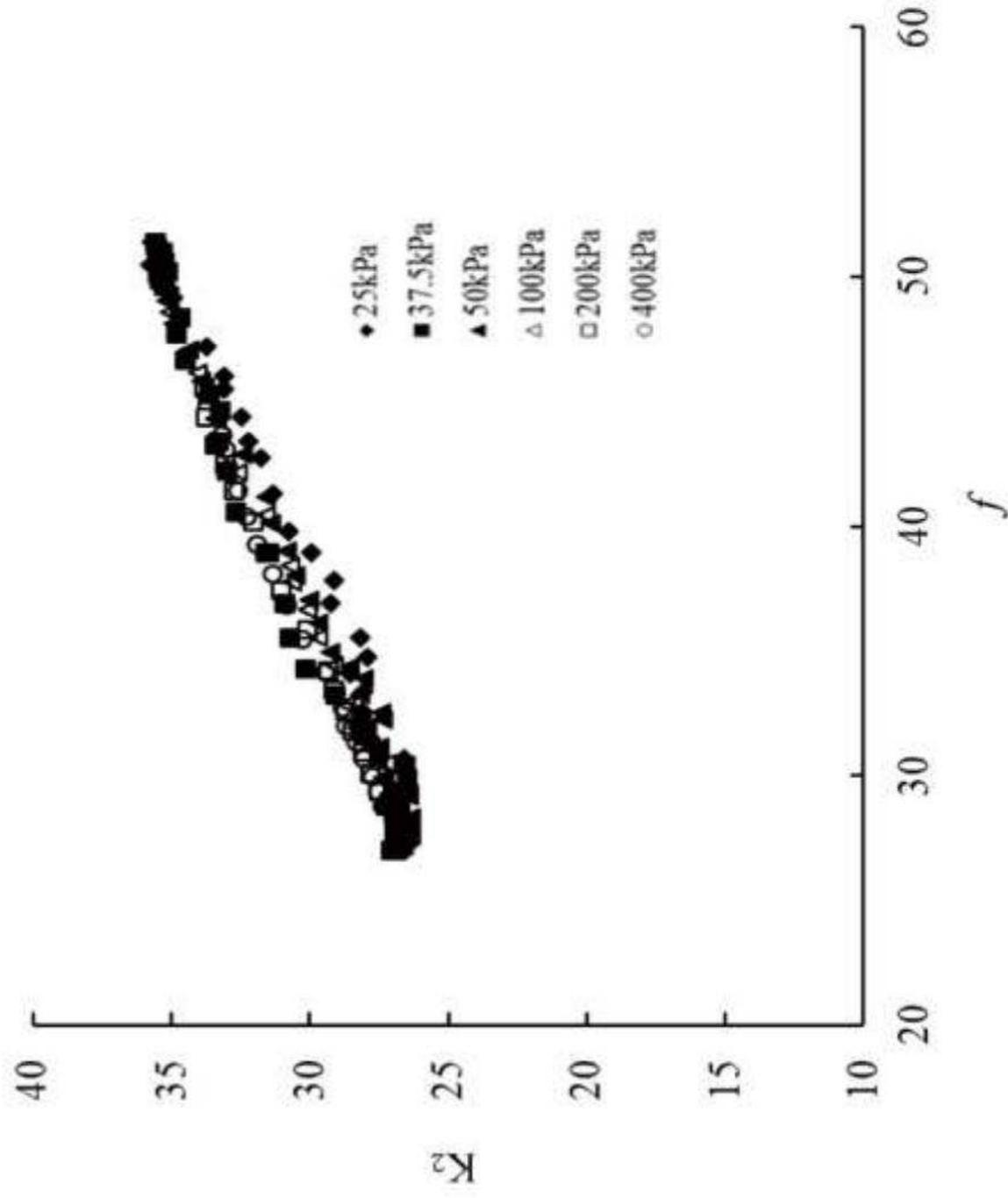


Fig. 10 The relationship curve of $K_2 \sim f$

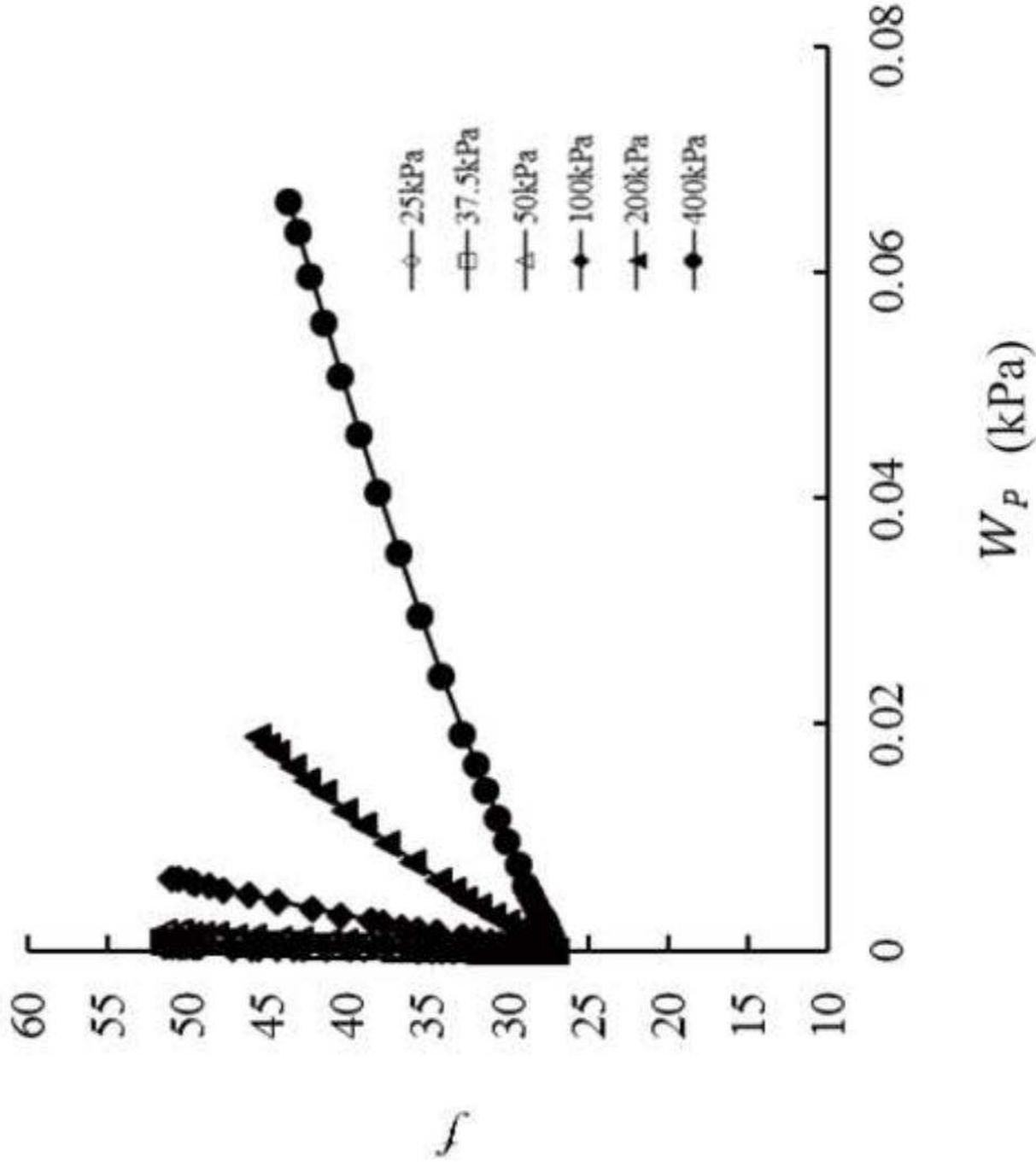


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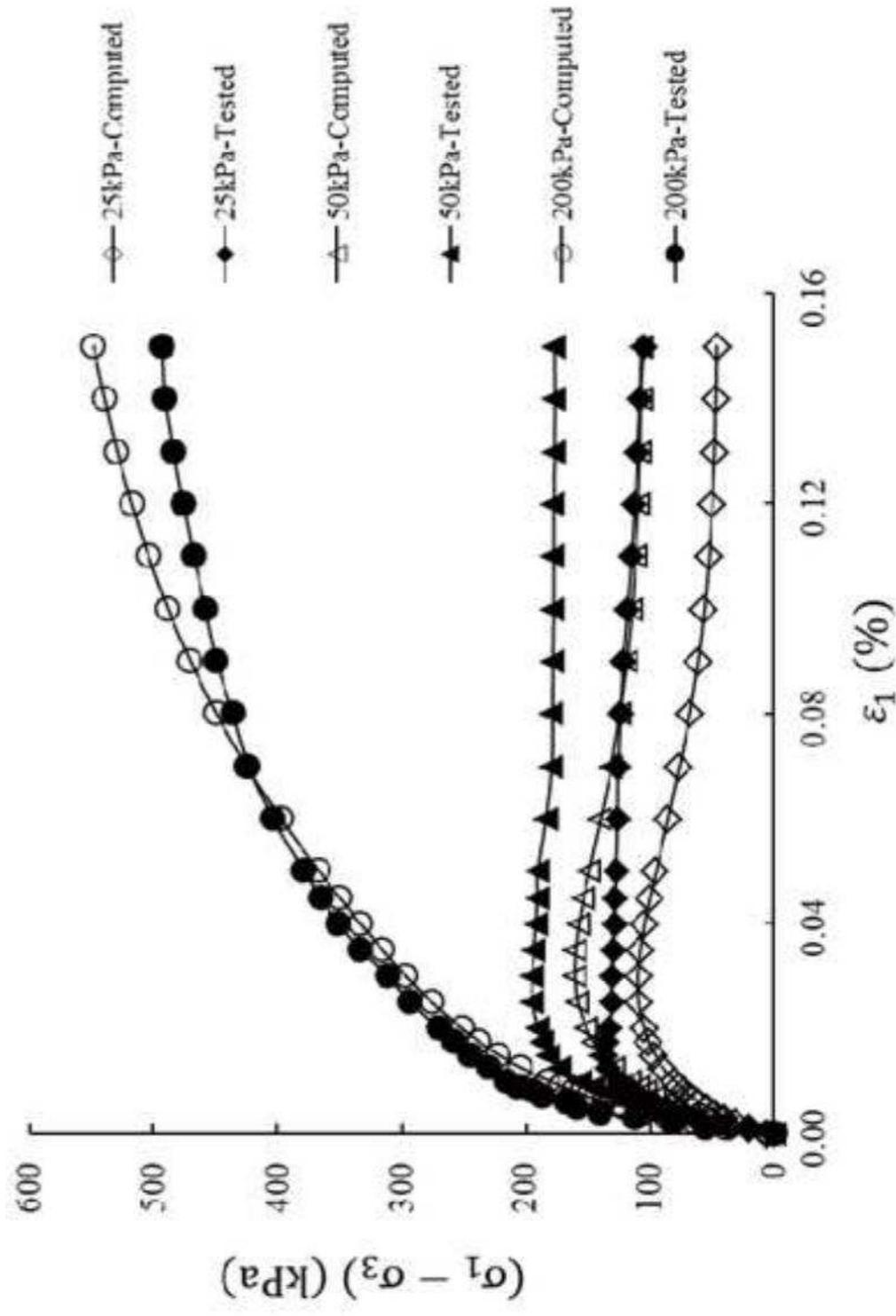


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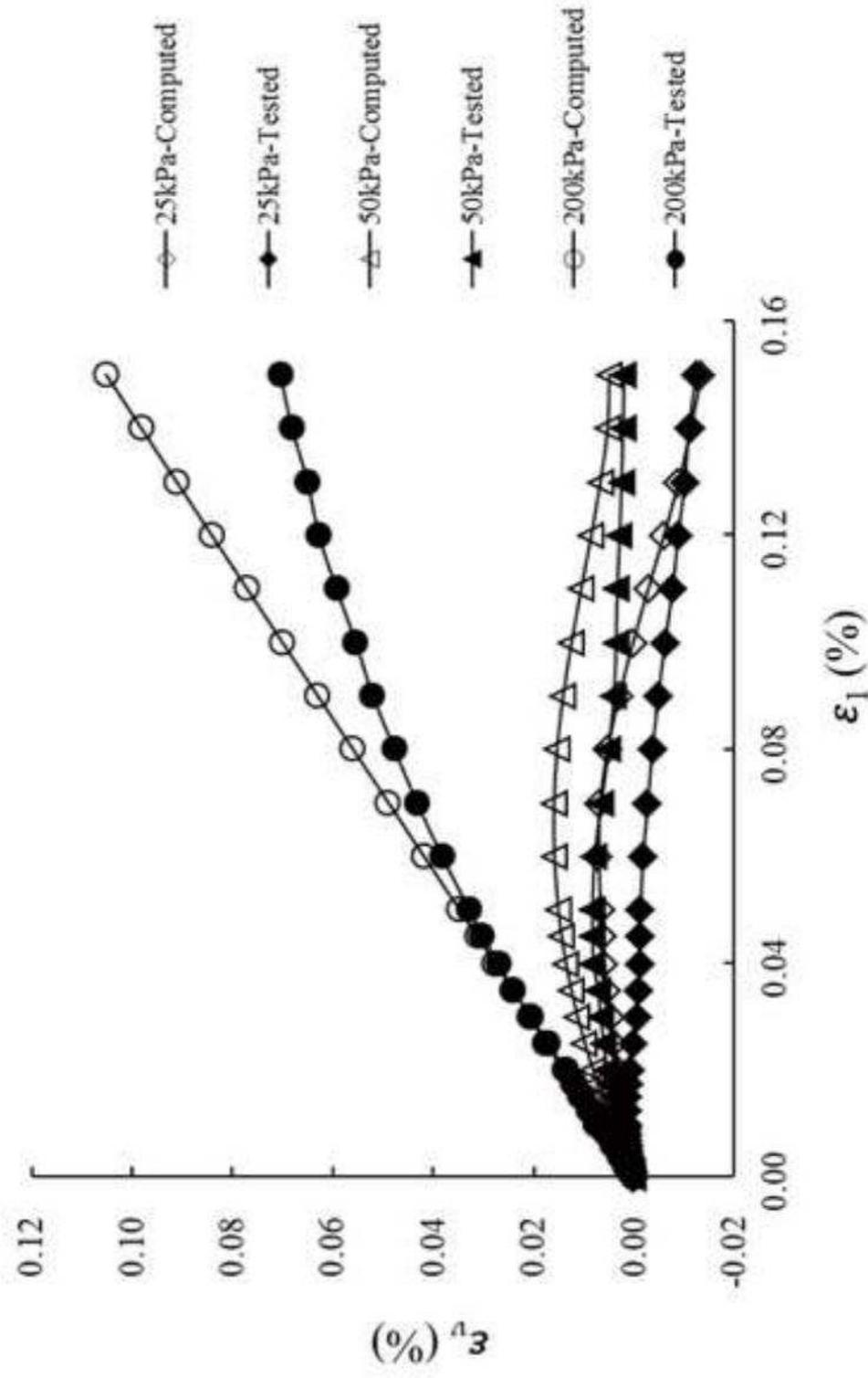


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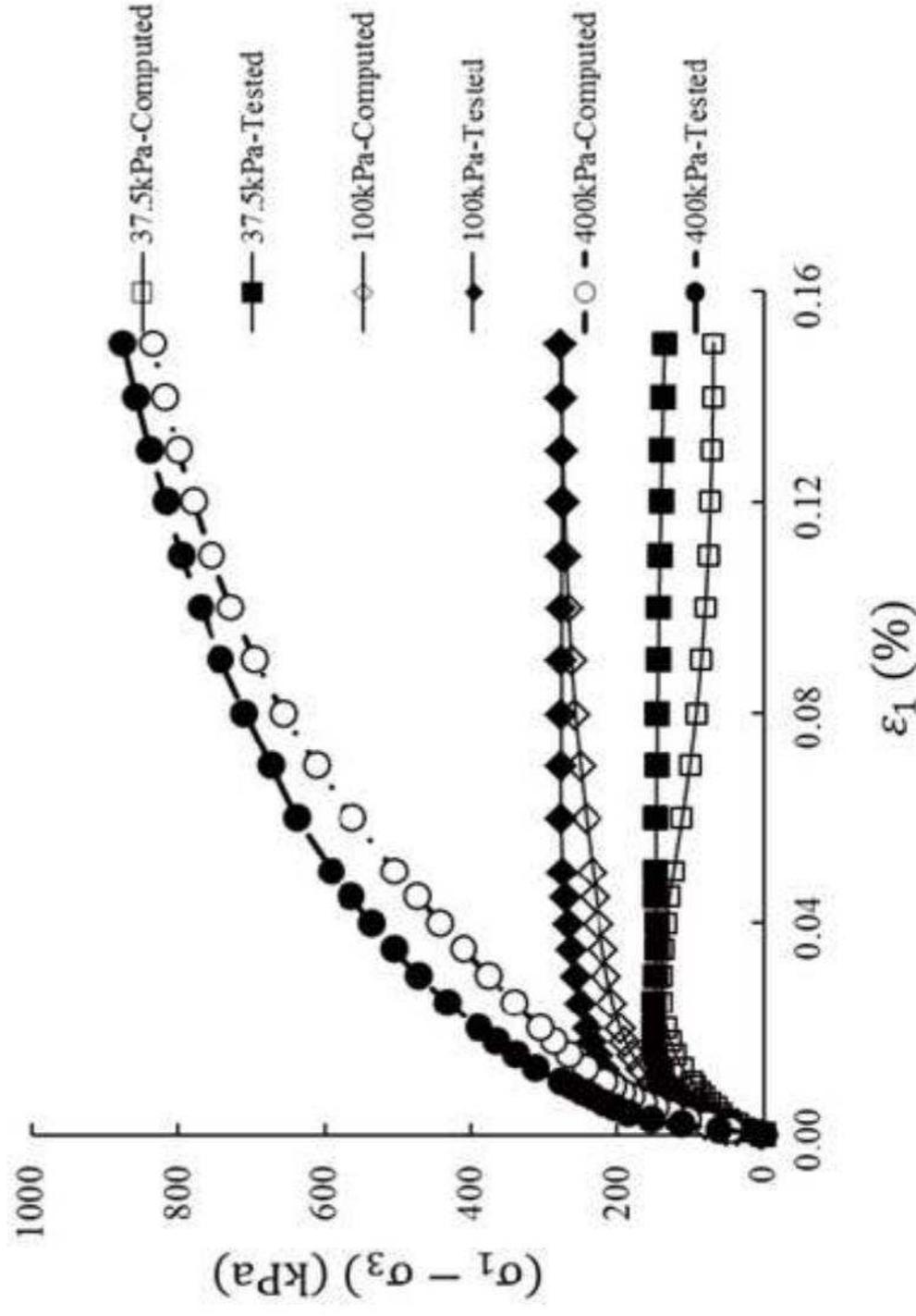


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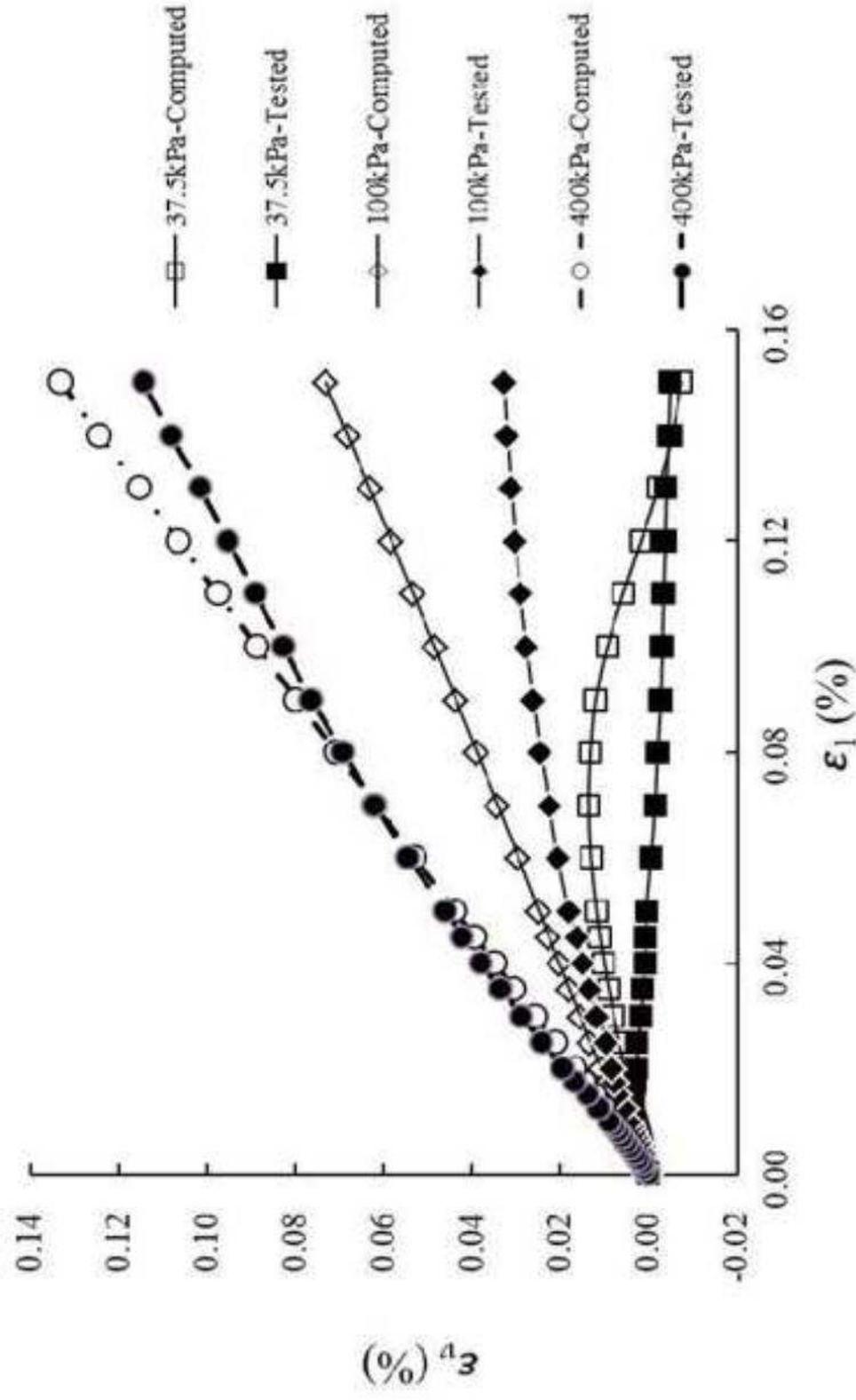


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