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First measurement of nuclear recoil head-tail sense in a fiducialised WIMP dark matter detector

The DRIFT collaboration

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ABSTRACT: Recent computational results suggest that directional dark matter detectors have potential to probe for WIMP dark matter particles below the neutrino floor. The DRIFT-II detector used in this work is a leading directional WIMP search time projection chamber detector. We report the first measurements of the detection of the directional nuclear recoils in a fully fiducialised low-pressure time projection chamber. In this new operational mode, the distance between each event vertex and the readout plane is determined by the measurement of minority carriers produced by adding a small amount of oxygen to the nominal CS₂+CF₄ target gas mixture. The CS₂+CF₄+O₂ mixture has been shown to enable background-free operation at current sensitivities. Sulfur, fluorine, and carbon recoils were generated using neutrons emitted from a ²⁵²Cf source positioned at different locations around the detector. Measurement of the relative energy loss along the recoil tracks allowed the track vector sense, or the so-called head-tail asymmetry parameter, to be deduced. Results show that the previously reported observation of head-tail sensitivity in pure CS₂ is well retained after the addition of oxygen to the gas mixture.

KEYWORDS: Dark Matter detectors (WIMPs, axions, etc.); Time projection Chambers (TPC)

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1 Introduction

A central unsolved problem in physics is to unravel the nature of non-baryonic dark matter (DM), known to constitute 24% of total energy density of the Universe [1]. Many possible DM candidates have been proposed. Weakly Interacting Massive Particles (WIMPs), motivated in part by supersymmetric theories, are a popular candidate [2, 3]. There is currently a global effort towards detecting these putative WIMP dark matter particles, in particular, by direct dark matter search experiments that aim to observe the products of WIMP-nucleus elastic scattering [3]. The rate of these events, and the recoil energy spectrum produced, is expected to modulate annually. This provides one possible route to a definitive signature related to our motion through the dark matter halo. DAMA/LIBRA [4] have announced a cumulative positive result from such an analysis with over 9σ statistical significance. However, this result remains controversial [5–8], and the technique remains problematic since backgrounds may mimic a WIMP signal [6, 9–11].

A robust and potentially more powerful galactic signature is encoded in the direction of the WIMP-induced nuclear recoils. The angular distribution of WIMP-induced recoils is highly anisotropic in galactic coordinates, with the average direction oriented away from our direction of travel towards the constellation Cygnus. In contrast, the distribution of any terrestrial background would be expected to be uncorrelated with this direction. In laboratory coordinates, due to the Earth’s rotation, a sidereal oscillation in WIMP-induced recoil directions would be expected [1, 11]. By this means, directional experiments have sensitivity to DM cross-sections below the so-called neutrino floor imposed by solar and atmospheric neutrinos coherently scattering off target nuclei [12, 13]. For a recent review of the discovery reach of directional DM detectors, see ref. [14].

An idealised detector sensitive to the recoil track axis only (not the vector direction), would need a few hundred events to confirm a galactic WIMP signal [15]. If, however, the recoil track vector is measured, then only tens of events are required [15, 16]. An important component of the vector track information is the so-called head-tail information, which determines on which end of the track the collision vertex is located. Determination of head-tail information is of great

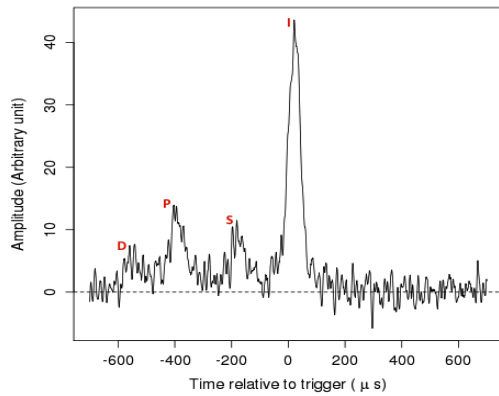
importance to the development of directional WIMP search technology. Although there has been recent progress in the use of nuclear emulsions in directional detection [17], the leading experiments all use a gaseous Time Projection Chamber (TPC). When operated at low pressure, typically 40–100 Torr, the expected recoils have a length of a few millimetres, long enough to be measured using a variety of existing charge sensitive readout technologies. Importantly, in principle this also allows measurement of any asymmetry in the ionization along the tracks and hence the head-tail of events. Directional DM search experiments including DRIFT [18], NEWAGE [19], MIMAC [20], DMTPC [21] and D³ [22] have been championing the use of this DM tracking method [14, 23]. These directional experiments are developing a joint initiative, called CYGNUS-TPC, to consider building a large modular directional experiment that would include head-tail sensitivity.

The DRIFT collaboration has already demonstrated head-tail sensitivity using pure CS₂ target gas [18, 24]. DRIFT chose to use this particular gas because it drifts signals as negative ions which results in much lower diffusion of tracks than can be achieved by conventional gases that drift electrons. Since that work, a different gas mixture has been developed comprising a 30:10:1 Torr mix of CS₂:CF₄:O₂. The relatively large ground-state nuclear spin of the fluorine nuclei provides a significant sensitivity to spin-dependent WIMP-proton interactions [25]. The addition of 1 Torr of oxygen gas to the existing CS₂+CF₄ mixture results in additional populations of faster minority charge carriers. Measurement of the relative times of arrival of these charge carriers at the x - y charge readout plane can be used to determine the absolute position of the event vertex in the drift, that is z -axis, direction [26]. This development of z -fiducialisation, adding to the existing x and y fiducialisation achieved by the MWPC (Multi-Wire Proportional Chamber) readout edge vetoes in DRIFT, has led to fully fiducialised operation, and a recently published background-free WIMP search — a first for a directional DM experiment [27]. The prime aim of the work reported here is to establish that the head-tail effect previously measured with pure CS₂ can still be detected in this new gas mode.

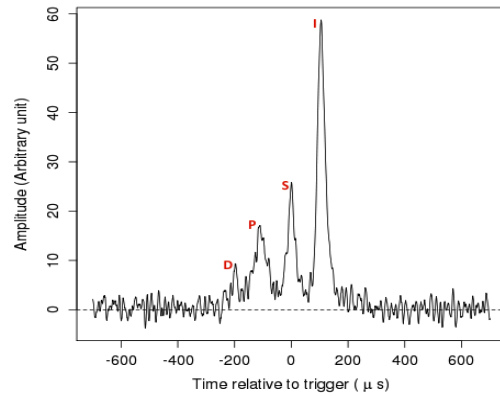
2 DRIFT-II_d detector and z -fiducialisation

The DRIFT-II_d directional dark matter detector comprises two back-to-back negative ion TPCs of 1 m³ total volume instrumented with MWPC readouts [28]. In this device, the two 50 cm long TPCs are separated by an aluminized-mylar thin-film central cathode biased at -31.9 kV, creating a drift field of 580 V cm⁻¹ for anions drifting towards the two MWPCs [27, 29]. Each MWPC comprises parallel stainless steel anode wires of 20 μ m diameter placed between and orthogonal to two grid wire planes made from parallel 100 μ m diameter stainless steel wires at 1 cm anode-grid distance. The grid potential is set at -2.884 kV, which is enough to strip off electrons from the drifting CS₂ anions and induce electron avalanches near the grounded anode wires. Each of the wire planes consist of 552 wires with 2 mm spacing. In the anode plane, 22 edge wires are grouped as guard wires, while the subsequent 82 wires are grouped as an anode veto. For the grid plane, 104 edge wires are grouped as a veto against side events. The remaining 448 wires each in the anode and grid planes are grouped down to 8 channels such that every 8th wire in each of the readout planes is read out by the same electronics channel. Each of the 8 channels is then read out with a Cremat CR-111 charge sensitive pre-amplifier, followed by a Cremat CR-200-4 μ s shaping amplifier. After signal amplification, the data are written to disk for analysis.

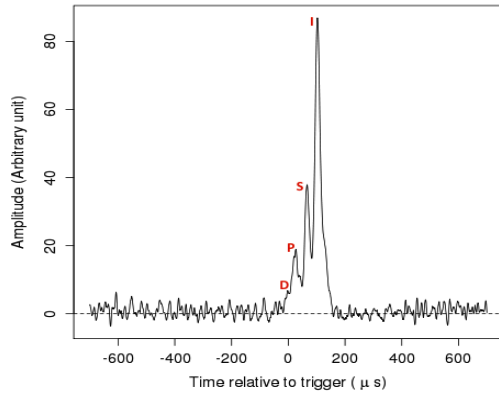
The addition of 1 Torr of oxygen to the CS_2+CF_4 mixture results in several minority carriers in addition to the main charge cloud, each drifting at a slightly different speed in the DRIFT detector. These minority carriers are believed to be due to different species of anions [26]. The three minority charge carriers and the main charge cloud are labeled as D, P, S, and I in figure 1, with I being the primary ionization, equivalent to the charge cloud observed in the absence of oxygen [26]. The minority carriers exhibit various characteristics with their behaviour depending mostly on z -position of the event as illustrated in figures 1(a) to 1(d). The four signal pulses shown in figure 1 are from four different neutron induced nuclear recoils obtained from the anode wire that recorded



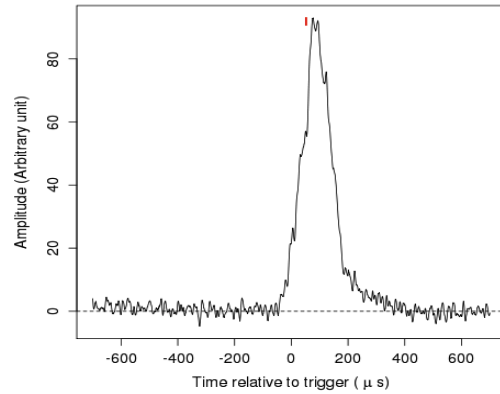
(a) High- z event at $z = 48.6$ cm, NIPs = 2300.



(b) Mid- z event at $z = 25.1$ cm, NIPs = 1504.



(c) Low- z event at $z = 9.2$ cm, NIPs = 1685.



(d) Low- z event at $z = ?$, NIPs = 1851.

Figure 1. Signal pulses from neutron-induced recoils at different z -positions within the active volume of the detector. The S specie charge carriers can change to other species of anions during drift, which reduces amplitude of the S peak (see (a)). Distinct D, P, S, and I peaks are shown for a mid- z event in (b). Typical behaviour of these peaks as they merge in the low- z region is shown in (c). The worst case in the very low- z region (closer to MWPC) is shown in (d). Events in the low- z region were rejected during analysis since their minority carriers become concealed under the main charge cloud. NIPs here stands for number of ion pairs.

the highest voltage after the respective events.

The I peak contains about 50% of the total ionization products in each event while the remaining $\sim 50\%$ is shared between the D, P, and S peaks [26]. The separation in arrival times of the anions in these peaks at the $x - y$ charge readout plane allows for the absolute position of an event vertex along the z -axis to be determined, and hence the detector is fiducialised. The absolute position of an event in the z -direction is given by:

$$z = (t_i - t_p) \left(\frac{v_i v_p}{v_p - v_i} \right), \quad (2.1)$$

where z is the distance between an event vertex and the readout plane, t_i and t_p are the arrival times of the anions in the I peak and P peak respectively, after drifting with corresponding velocities v_i and v_p . The events z positions increase from 0 (MWPC) to 50 cm (central cathode) for each of the two detectors.

For the purpose of this analysis, events in the detector were grouped into three bins according to their z -position: low- z ($0 < z \leq 15$ cm), mid- z ($15 \text{ cm} < z \leq 35$ cm) and high- z ($35 \text{ cm} < z \leq 50$ cm). Relative to the I peak, the amplitude of the S peak is known to decrease with increasing z [26]. This suppression of the S peak for the high- z events can be seen in figure 1(a). In this work, many high- z events did not have a clear S peak, and so the z -position of each event was calculated from timing of the I and P peaks. For mid- z events, clear and distinct separation between the main and minority peaks was observed as depicted in figure 1(b). In this region, the amount of charge in the peaks decreases from the I to the D peak, in that order. In the low- z region, closest to the MWPC, the total drift time of the anions is short, and so the temporal spacing between the peaks is small. For many of the low- z events, the four peaks are not resolved in time and hence, were not precisely fiducialised. As a result, we omit the low- z region from this analysis. The implication is that data from 30% of the active volume of the detector were not included in this analysis. New work to allow inclusion of these events is currently underway.

As described in [27, 29], a population of nuclear recoil events near $z = 50$ cm are produced by radioactive decays on the central cathode. The effect of this contamination on the high- z sample is negligible because (1) the intrinsic decay rate is low compared to the neutron interaction rate (~ 3 events/day vs. $\sim 2.3 \times 10^4$ events/day), (2) about half of the total number of these radioactive decay events with $z > 50$ cm were not included in the head-tail analysis.

In the work presented here, the high- z and mid- z bins were selected and used as a means to investigate the possible effect of diffusion on the head-tail asymmetry for events closer to the central cathode and middle of the detector respectively.

3 Experimental set-up

The detector was exposed to neutrons from a $\sim 338 \text{ kBq}$ ^{252}Cf source from different directions for about seven days. These neutron exposures were performed without source collimation to minimise neutron back scattering. Illustration of the detector and details on the directions of the neutron exposures are depicted in figure 2. The hardware threshold was set to 15 mV throughout this experiment, resulting in a typical trigger rate of 5 Hz and minimum detected recoil energy of ~ 300 NIPs. NIPs is the number of electron-ion pairs produced by recoiling nuclei. Detector gain

the right of the MWPCs such that neutron induced recoils in the right detector were predominantly oriented towards the central cathode, while in the left detector neutron induced recoils were mainly oriented away from the central cathode. The runs where the source was centered on the central cathode from the x and y -directions are known as $-x$ and $-y$ runs. In each of these exposures, the source was placed at a distance of 32 cm (107 cm) and 132 cm (207 cm) away from the vessel (geometric center of the central cathode) for the optimal and anti-optimal directed neutron runs, respectively. These distances result in neutron interaction angular spread of about 25° (14°) for the optimal (anti-optimal) runs. In these measurements, the source neutron-vessel separation distances for the optimal and anti-optimal runs were maximised subject to available space in the underground laboratory during the exposures. The exposure from the optimal directions was done to ascertain if the sensitivity of DRIFT-II detector to signal head-tail (earlier reported in ref. [18]) was affected by addition of oxygen that aided the full fiducialisation of the detector while the anti-optimal runs were expected to provide the required reference for null results.

4 Data analysis

In this work, only the main charge clouds known as the I peaks were analysed to determine the head-tail effect. This is because the I peak retains most of an event charge compared to D, P, and S peaks. The data analysis presented here was based on the method described in [18]. In the analysis, a nuclear recoil candidate must pass a series of initial cuts described in ref. [29]. There are five main reasons to reject some events in this first stage. Events emanating from the central cathode and long alpha events may have hits (nuclear recoil signals that cross the analysis threshold) on both sides of the cathode - such events are rejected, as are events with hits on non-contiguous wires, which may be two independent events in the same time window. An analysis threshold of 9 mV allowed for identification of the D, P, and S peaks since events are typically triggered by the I peak. Apart from the I peak, at least any two of the D, P, and S peaks must pass the analysis threshold for an event to be analysed further as a nuclear recoil candidate. Events with all eight wires in a group hit are cut to reject all contained alpha events with track length in the x - y plane greater than 15 mm. This is because there are no contained alphas with x - y range < 15 mm in the DRIFT-II detector design [27]. To reject other alphas and all other events emanating from outside the active volume of the detector, events with hits on the veto wires are rejected. Also, events with rise-times less than $3 \mu\text{s}$ are rejected to remove noise (for instance sparks).

In the next stage of the analysis, the integral charge from each event in ADC units was converted to corresponding NIPs equivalent. To do the NIPs conversions, a region of interest (ROI) was defined from $+700 \mu\text{s}$ to $-700 \mu\text{s}$ relative to the trigger time for all anode channels with signal above the analysis threshold. This ROI contains the D, P, S, and I peaks. The sum of the integrated charge values was converted to the NIPs equivalent value using the most recent ^{55}Fe detector gain calibration constant and W-value of the gas mixture. The W-value is the average amount of energy required to produce an electron-ion pair in a given gas mixture.

The W-value for 30:10:1 Torr of $\text{CS}_2:\text{CF}_4:\text{O}_2$ gas is under study. But the addition of about 1% of O_2 to the conventional CS_2+CF_4 gas mixture is expected to change the W-value by only about 0.3 eV [30–32]. This is within the uncertainty of 25.2 ± 0.6 eV measured by ref. [33] for 30:10 Torr of $\text{CS}_2:\text{CF}_4$ mixture, adopted in this analysis. The ionization energy depends on a property

of the gas mixture known as the quenching or Lindhard factor (see ref. [34]). This gas quenching factor measures the fraction of energy that goes into ionization for a given recoil. For example, the quenching factor is 0.53 for a 50 keV fluorine recoil in a 30:10 Torr of $\text{CS}_2:\text{CF}_4$ gas mixture, and so a recoil of this energy will produce 1055 electron-ion pairs [29, 34].

Events with NIPs ≤ 6000 were considered further, no low-energy cut was applied. See figure 3 for the distribution of reconstructed neutron recoil events in the z -NIPs space before the z cuts. The absolute z position for each of the events was computed using Equation 2.1. There are no

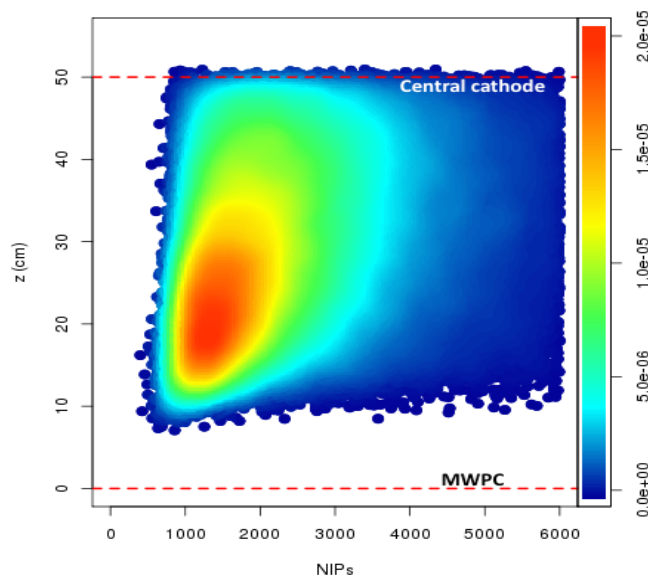


Figure 3. Event distribution in the z -NIPs space obtained from neutron exposure using a ^{252}Cf source placed in $+z$ position. This is before the z cuts described in the text. The dashed red lines at 0 and 50 cm mark the locations of the readout plane and central cathode respectively. The colour scale represents the density of the events (colour online).

events in the left part of the panel in figure 3 because charge signal pulses diffuse as they drift to the readout plane. For low-energy, high- z events, diffusion broadens the pulse width and may reduce the maximum pulse amplitude below the data acquisition trigger. The trigger has since been modified to monitor the integral charge of the waveform, rather than the amplitude. The upgraded trigger should improve the efficiency for high- z events at low-energy. In the lower part of the panel, closer to the readout plane, the reconstruction cannot identify at least two of the D, P, and S peaks because of the short drift time, and so these events do not pass nuclear recoil selection cuts. The energy distribution of these events shown in figure 3 is compared to the expected energy spectrum from ^{12}C , ^{19}F and ^{32}S recoil candidates generated with GEANT4 [35] program in figure 4. It can be seen from these two spectra that the distribution of the experimental data and theoretical expectation from ^{12}C , ^{19}F , ^{32}S candidates all reach maximum at about 1000 NIPs, each with an exponential fall-off at higher NIPs. Hence, it is likely that there are some fractions of ^{12}C , ^{19}F and ^{32}S recoil candidates in this data set. However, the energy spectrum obtained from the GEANT4 simulation (see figure 4(b)) has a slower high-energy fall-off with a steeper cut-off at low-energy when compared to the experimental data shown in figure 4(a). These discrepancies may be related

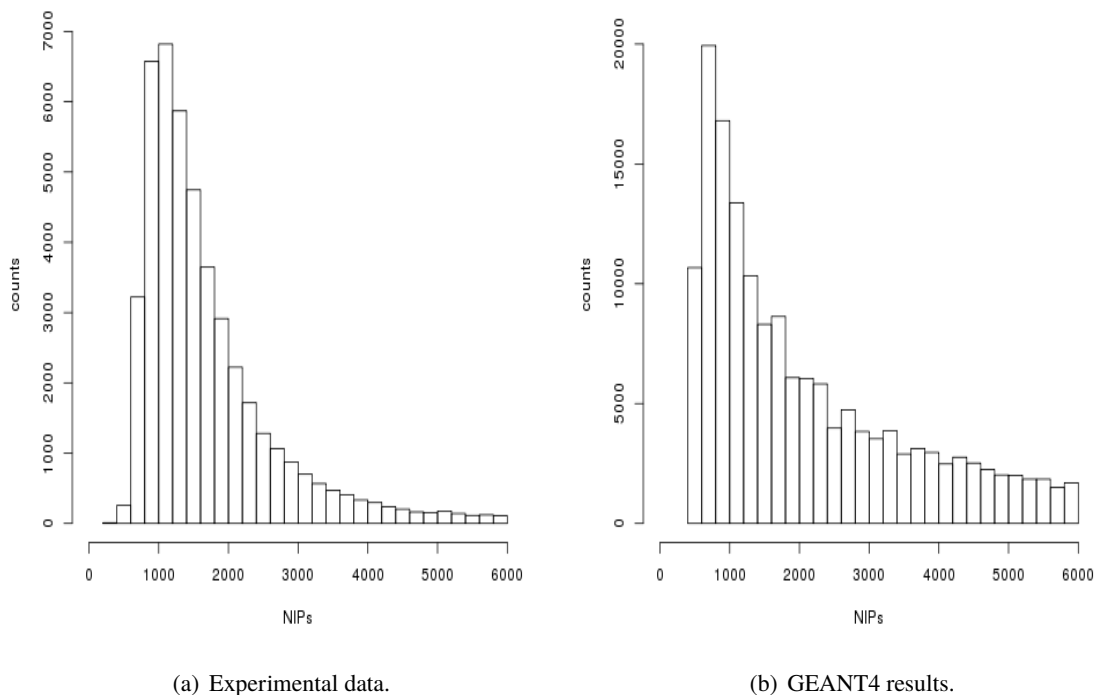


Figure 4. Energy spectra of neutron-induced nuclear recoils with the source in the optimal direction. (a) Experimental data before the z cuts (described in the text) from a $+z$ run. (b) Results of a GEANT4 Monte Carlo simulation.

to unexplained systematics discussed in section 5.

Additionally, to quantify the relative amount of ionization in the minority peaks, a parameter called the peak ratio Q , defined as:

$$Q = \frac{Q_{DPS}}{Q_T}, \quad (4.1)$$

was also computed using the signal pulse on the anode wire that recorded the highest voltage in each event, where Q_{DPS} is the integral charge in the D, P, and S peaks, and Q_T is the total integral charge in the D, P, S, and I peaks. Events whose Q values are greater than 0.30 and less than 0.65 were analysed further. This cut was found to be a powerful tool for discriminating nuclear recoils from radon progeny recoil events emanating from the MWPC charge readouts and some low- z events whose minority charge carrier peaks are completely concealed under the main charge cloud. These noise events are expected to have little or no minority charge carrier peaks, which result in small Q_{DPS} values with $Q < 0.30$. The upper cut on $Q > 0.65$ is used to remove spark events with continuous waveform pulses resulting in large Q_{DPS} values relative to expectations from nuclear recoil events.

To investigate the presence of the head-tail effect in the nuclear recoil tracks, we use the charge distribution in the main charge cloud from every nuclear recoil event. This is because of the expectation that any available head-tail signature will cause a temporal asymmetry in the charge distribution along the nuclear recoil tracks [18, 36]. The asymmetry was measured by comparing

the integral charge in the first and second halves of the I peak using the signal pulse in the anode wire with the highest voltage for that event. Specifically, the temporal extent of the I peak was defined as the region where the signal is larger than 25% of the peak amplitude (see figure 5).

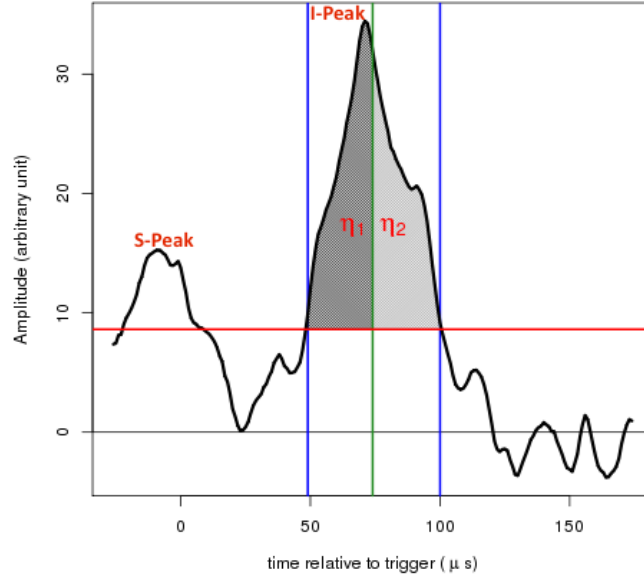


Figure 5. Event waveform from a single wire showing the analysed I peak with η_1 (integral charge in the first half of region of interest) shaded dark grey and η_2 (integral charge in the second half of the region of interest) shaded light grey. The two blue lines mark the region where the signal exceeds 25% of the peak amplitude. The green line is the mid-point time (half of signal duration above analysis threshold) between the first and second halves of the signal and the red line shows the 25% of the I peak used as the analysis threshold (colour online).

This region was then split into two equal-length regions, and the integral charges η_1 and η_2 were computed for each of the events that passed the cuts. The vector direction of the track was then determined from the ratio:

$$\alpha \equiv \frac{\eta_1}{\eta_2}. \quad (4.2)$$

Because the ionization density for nuclear recoils is larger at the start of the track than the end, a nuclear recoil whose velocity vector points towards the MWPC should have $\eta_1 < \eta_2$, and therefore $\alpha < 1$. Tracks with $\alpha > 1$ correspond to recoils that point toward the central cathode.

The head-tail asymmetry parameter α was computed for the events in each TPC. Then, the mean value of α was computed separately for events in the left and right detectors: $\langle \alpha \rangle_L$ and $\langle \alpha \rangle_R$, respectively. There is a temporal asymmetry introduced by the shaping electronics in the measured waveforms. To account for this effect, we compute $\Delta\alpha \equiv \langle \alpha \rangle_L - \langle \alpha \rangle_R$ between the mean asymmetry parameters for the left and right detectors for each neutron exposure. The expectation is that the largest head-tail effect should come from the $+z$ and $-z$ neutron source runs. In contrast, no head-tail effect is expected in the dataset from the anti-optimal $-x$ and $-y$ runs.

To understand the head-tail effect further, the percentage difference of the head-tail asymmetry parameter was determined and termed δ . This δ parameter, given by:

$$\delta = 100 \frac{|\Delta\alpha|}{\frac{1}{2}(\langle\alpha\rangle_L + \langle\alpha\rangle_R)}, \quad (4.3)$$

is the ratio of $|\Delta\alpha|$ to the mean of $\langle\alpha\rangle_R$ and $\langle\alpha\rangle_L$, expressed as a percent. The δ parameter quantifies the measurability of the head-tail effect, including both the intrinsic head-tail signature and the sensitivity of the detector to that signature. The δ parameter was measured using events from the optimal and anti-optimal directions. The results obtained in this analysis are discussed in section 5.

5 Results and discussion

The dependence of δ on energy is shown in figure 6 for both the optimal and the anti-optimal runs, using bins of width 200 NIPs. It can be seen in that figure that the δ parameter measured in the anti-optimal runs is consistent with zero for all energies indicating that no head-tail signature was detected. This is expected because the mean directions of event tracks in these anti-optimal runs are the same for both TPCs, resulting in null $|\Delta\alpha|$ and δ parameters. Conversely, events from the optimal runs, where the mean direction of recoil tracks is oriented along the drift direction of anions, show a δ parameter that increases with energy. The δ parameter obtained from the optimal case is significant at 750 NIPs (~ 38 keV F-recoil energy equivalent). These results demonstrate the sensitivity of the detector to the head-tail directional signature and, furthermore, increased sensitivity with higher energy tracks.

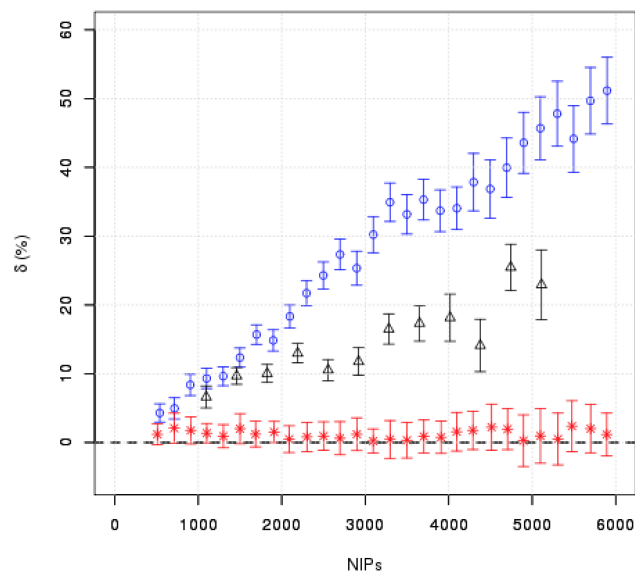


Figure 6. Percentage difference δ parameter for different recoil energies represented by NIPs equivalent. Results from this analysis (obtained with ^{12}C , ^{19}F and ^{32}S tracks) are shown with blue circle points for the optimal runs and red asterisk points for the anti-optimal runs. The black triangle points are 2009 results obtained with only ^{32}S tracks in pure CS_2 (see ref. [18]). Colour online.

In figure 6, the δ parameter obtained in this new gas mode is compared to the result obtained in a previous measurement performed in 2009 using pure CS₂ gas [18]. It was shown in that study that event tracks used in the measurement were primarily from ³²S recoils. The lower analysis threshold used in this measurement yielded the observed head-tail sensitivity below the ~ 1000 NIPs threshold obtained from the pure CS₂ data. Above 2000 NIPs, the measured values for δ in this work are approximately a factor of 2 larger than in the 2009 study.

To explain this, we hypothesize that the observed increase in the δ parameter in figure 6 is due to the inclusion of ¹²C and ¹⁹F recoils in this data set which were not included in the 2009 study. A longer recoil track is expected to yield a larger head-tail effect because it allows for a more distinct separation between the beginning and the end of the recoil track. For instance, a SRIM [37] calculation of the ranges for ¹²C, ¹⁹F, and ³²S recoils that produce 1000 NIPs in 30:10:1 Torr of CS₂:CF₄:O₂ gives 1.70 mm, 1.51 mm and 0.98 mm, respectively. However, the gas quenching for each of the ¹²C, ¹⁹F and ³²S recoil candidates is different [36]. This gas quenching effect is expected to increase with mass of a given nucleus. Hence, in the context of this experiment, ³²S nuclear recoil tracks should suffer more energy loss per given recoil distance relative to ¹²C or ¹⁹F recoil candidates. Earlier studies in ref. [19] reported the presence of head-tail effects in nuclear recoil tracks obtained with a CF₄ based TPC detector with potential ¹²C and ¹⁹F recoil candidates.

To check for the presence of nuclear recoil candidates with track ranges that are greater than the expectations from only ³²S recoils, the full width half maximum (FWHM) range $\Delta z = FWHM \times v_i$, was computed for each of the optimal (+z and -z) events. This Δz range was used because of its high resolution (< 0.1 mm) [38]. The results obtained from this analyses are shown in Δz -NIPs space in figure 7. The SRIM range (in 3D) as a function of NIPs was computed for ³²S, and the maximum

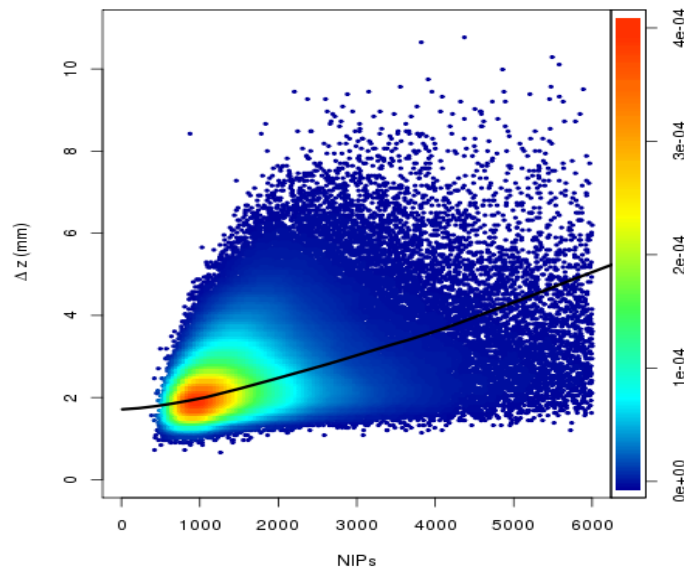


Figure 7. Experimental data overlaid with the maximal Δz predictions for ³²S recoils in 30:10:1 Torr of CS₂:CF₄:O₂ gas mix obtained using SRIM. The heat map is the experimental data while the black line is the maximum Δz predictions for ³²S recoils. The colour scale represents the density of the events (colour online).

expected Δz parameter for ^{32}S recoils was then determined by accounting for track diffusion, shaping electronics, and the differences in charge-carrier trajectories due to non-uniform electric field line distribution near the anode readout wires [39]. The measurements show a significant number of events with Δz in excess of the expectations for ^{32}S , but consistent with presence of longer recoil tracks. However, the Δz -NIPs results did not unambiguously isolate each recoil specie candidates. Isolated recoil species will be a more convincing evidence to show the number of ^{12}C and ^{19}F recoil populations in this data. Ongoing studies are aimed toward tackling the ambiguity in separating constituent recoil species in a given data set.

The head-tail asymmetry parameter $\Delta\alpha$ obtained from the high- z and mid- z bins in the $+z$ and $-z$ runs are shown in tables 1 and 2. Significant head-tail effects are apparent in both the high- z and mid- z regions as well as the combined data in table 3. Naively, it is expected that a lower value of $\Delta\alpha$ will be observed in the high- z region compared to the mid- z region of the detector due to diffusion effects. This is because diffusion can reorganise the actual charge distribution along a recoil track when it is exposed to longer drift distance which tends to conceal the head-tail effect. However, the opposite was observed because there are more high-energy (NIPs > 1500) events in the high- z region of the detector compared to the mid- z region of the detector. This is because the trigger efficiency on low-energy (NIPs ≤ 1500) events decreases for events closer to the central cathode due to accumulated diffusion effects along the drift direction.

In fact, the fraction of high-energy events in the high- z region is about an order of magnitude greater than in the mid- z region of the detector. To investigate this effect, 19800 events were selected from the optimal high- z and mid- z regions using a fixed energy spectrum generated with ^{32}S recoils in GEANT4. These selected events were analysed using the same procedure described above. The resultant $\Delta\alpha$ parameter of 0.098 ± 0.005 and 0.104 ± 0.006 were obtained from the high- z and mid- z regions, respectively. In this scenario, statistically consistent results were obtained from events in the high- z and mid- z regions. This result demonstrates that both diffusion and recoil energy affect the head-tail reconstruction.

MND	n	live-time (days)	$\langle\alpha\rangle_L$	$\langle\alpha\rangle_R$	$\Delta\alpha \pm \sigma_{\Delta\alpha}$
$+z$	19955	1.909	1.294 ± 0.003	1.146 ± 0.004	0.148 ± 0.005
$-z$	21054	3.713	1.175 ± 0.004	1.316 ± 0.003	-0.141 ± 0.005

Table 1. Head-tail parameter obtained from high- z events in the $+z$, $-z$ runs. Only statistical uncertainties are quoted. The n is the number of events, other symbols used here are defined in the text.

MND	n	live-time (days)	$\langle\alpha\rangle_L$	$\langle\alpha\rangle_R$	$\Delta\alpha \pm \sigma_{\Delta\alpha}$
$+z$	31267	1.909	1.359 ± 0.003	1.256 ± 0.004	0.103 ± 0.005
$-z$	49005	3.713	1.303 ± 0.003	1.377 ± 0.002	-0.074 ± 0.004

Table 2. Head-tail parameter obtained from mid- z events in the $+z$, $-z$ runs. Only statistical uncertainties are quoted. Symbols are defined in table 1.

The larger fraction of high-energy events in the high- z region of the detector is expected to produce longer recoil tracks with distinct track head-tail separation and higher signal to noise ratio which yield higher values of the head-tail asymmetry parameter $\Delta\alpha$, compared to the mid- z region.

MND	n	live-time (days)	$\langle\alpha\rangle_L$	$\langle\alpha\rangle_R$	$\Delta\alpha \pm \sigma_{\Delta\alpha}$
+z	51222	1.909	1.336 ± 0.002	1.203 ± 0.003	0.133 ± 0.003
-z	70059	3.713	1.258 ± 0.003	1.360 ± 0.002	-0.102 ± 0.003
-x	5094	0.915	1.408 ± 0.009	1.411 ± 0.009	-0.003 ± 0.013
-y	2873	0.561	1.345 ± 0.011	1.350 ± 0.012	-0.005 ± 0.016

Table 3. Head-tail parameter obtained from combined high- z and mid- z events in the + z , - z , - x , and - y runs. Only statistical uncertainties are quoted. Symbols are defined in table 1.

Results obtained from all the events in the optimal (+ z and - z data) and anti-optimal (- x and - y data) runs are presented in table 3. It can be seen from this table that a higher value of the head-tail parameter $\Delta\alpha$ was obtained from + z and - z directions when compared to - x and - y runs. This confirms that + z and - z are the optimal directions for head-tail detection in the DRIFT-II detector design as reported in ref. [18]. Also as expected, - x and - y exposures returned null head-tail effect. It can be seen in tables 1 to 3 that $\langle\alpha\rangle_R$ and $\langle\alpha\rangle_L$ for + z and - z neutrons, respectively, are not less than 1 as expected. As described in section 4, this is due to the effect of the shaping electronics, which is accounted for in the computation of $\Delta\alpha$ parameter. Even with the effect of the shaping electronics, it can be seen that $\langle\alpha\rangle_L > \langle\alpha\rangle_R$ for the + z run and $\langle\alpha\rangle_L < \langle\alpha\rangle_R$ for the - z run, as expected.

Some cross-checks can be done on the results presented in tables 1, 2 and 3. For each of the tables, the difference between the magnitude of the $\Delta\alpha$ parameter for the + z and - z runs should be zero within uncertainties. The differences are 0.007 ± 0.007 (table 1, high- z region), 0.029 ± 0.006 (table 2, mid- z region) and 0.031 ± 0.004 (table 3, combined mid- z and high- z regions). These checks indicate agreement of the two measurements of $\Delta\alpha$ for the high- z events, but also indicate the presence of a systematic effect for the two measurements of $\Delta\alpha$ for the mid- z events of table 2 and the combined data in table 3. Although, the observed head-tail asymmetry parameter remains statistically significant after allowing for this systematic effect.

Additionally, two other alternative measures of the $\Delta\alpha$ parameter can be formulated by using only one side of the detector (left or right) with both + z and - z runs. This can be done by subtracting $\langle\alpha\rangle_L$ (for a - z run) from $\langle\alpha\rangle_L$ (for a + z run) and also by subtracting $\langle\alpha\rangle_R$ (for a - z run) from $\langle\alpha\rangle_R$ (for a + z run). These two alternate calculations yielded 0.119 ± 0.005 and -0.170 ± 0.005 for the high- z region of the detector. For the mid- z events, the alternate measures were found to be 0.056 ± 0.004 and -0.121 ± 0.004 , while for the combined runs the alternate measures are 0.078 ± 0.004 and -0.157 ± 0.004 . The differences between the magnitude of each pair of these numbers obtained from the alternate measures are 0.051 ± 0.007 (high- z), 0.065 ± 0.006 (mid- z) and 0.079 ± 0.006 (combined mid- z and high- z). The presence of these differences again indicate systematic effects, since the magnitude of the numbers obtained from the alternative measures are not equal within uncertainty. The differences reported above are estimates of the size of the systematic effects, which are larger than the statistical uncertainties reported in tables 1–3.

After allowing for the systematic effects described above, the head-tail effect remains clearly present for the optimal runs. It should be noted that these systematic effects were not observed in [18], however, the experimental conditions for this work were different in comparison to [18]. For mechanical reasons, the neutron source had to be placed closer to the detector for this work. In addition, for this work there existed an asymmetric distribution of materials surrounding the

detector. GEANT4 simulations suggest that neutrons can be scattered from these materials, which likely contributed to the observed systematic effects. A detailed understanding of the systematics is beyond the scope of this paper, as the aim is to verify that the head-tail effect, observed in [18], is still present while operating in the new fiducialised mode.

6 Conclusion

The data presented here indicate firstly that the DRIFT-II detector is still sensitive to signal head-tail detection after addition of oxygen gas for fiducialisation. This is after the head-tail sensitivity was previously reported considering only ^{32}S recoils in pure CS_2 target gas. The longer recoil tracks due to ^{12}C and ^{19}F components in the new data were found to yield a more distinct separation between signal track head and tail. Studies done at different z regions of the detector yielded higher values of the head-tail asymmetry parameter in the high- z region compared to the mid- z region. This was found to be because the high- z bin is dominated by high-energy events mostly because low-energy events in this region of the detector do not pass the analysis threshold due to larger diffusion effects.

The $+z$ and $-z$ directions were confirmed as the optimal directional axes for the DRIFT-II detector while the $-x$ and $-y$ also produced null head-tail results, as expected. This observed head-tail effect from the optimal direction and null effects from the anti-optimal direction can be combined to track the modulation in signal directionality from WIMP events reaching the detector from the direction of the CYGNUS constellation. This method will be better utilised if the optimal axes of the detector are aligned along the earth north-south axis as proposed in previous studies.

These new results are encouraging for the possibility of developing a large volume TPC for directional dark matter detection where an appropriate readout technology, sensitive to head-tail effect can be combined with full detector fiducialisation. Further work is underway to develop a more sophisticated analysis for extracting the concealed head-tail information due to peak pile-up in the low- z region which is expected to improve the head-tail sensitivity further. Studies of alternative head-tail analyses are also underway to improve the overall directional sensitivity.

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