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Han, X, Bibb, R and Harris, R [orcid.org/0000-0002-3425-7969](http://orcid.org/0000-0002-3425-7969) (2016) Engineering design of artificial vascular junctions for 3D printing. *Biofabrication*, 8 (2). 025018. ISSN 1758-5082

<https://doi.org/10.1088/1758-5090/8/2/025018>

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1                   **Engineering Design of Artificial Vascular Junctions for 3D Printing**

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11

12 **Abstract**

13 Vascular vessels, including arteries, veins and capillaries, are being printed using additive  
14 manufacturing technologies, also known as 3D Printing. This paper demonstrates that it is important  
15 to follow the vascular design by nature as close as possible when 3D printing artificial vascular  
16 branches. In previous work, the authors developed an algorithm of computational geometry for  
17 constructing smooth junctions for 3D printing. In this work, computational fluid dynamics (CFD) is  
18 used to compare the wall shear stress and blood velocity field for the junctions of different designs.  
19 The CFD model can reproduce the expected wall shear stress at locations remote from the junction.  
20 For large vessels such as veins, it is shown that ensuring the smoothness of the junction and using  
21 smaller joining angles as observed in nature is very important to avoid high wall shear stress and  
22 recirculation. The issue is however less significant for capillaries. Large joining angles make no  
23 difference to the hemodynamic behavior, which is also consistent with the fact that most capillary  
24 junctions have large joining angles. The combination of the CFD analysis and the junction  
25 construction method form a complete design method for artificial vascular vessels that can be 3D  
26 printed using additive manufacturing technologies.

27 **Keywords:** Vascular vessel design, Computational fluid dynamics (CFD), Additive Manufacturing,  
28 3D printing

29

30	<b>Nomenclature</b>
31	
32	$C_{\max}$ Non-dimensionalised maximum curvature of a bifurcation junction
33	$\bar{V}$ Non-dimensionalised branch Volume
34	$p_{\text{outlet}}$ Outlet pressure
35	$v_{\text{inlet}}$ Inlet velocity
36	$x, y, z$ Cartesian coordinates
37	$u_d$ Control parameters

38 **1. Introduction**

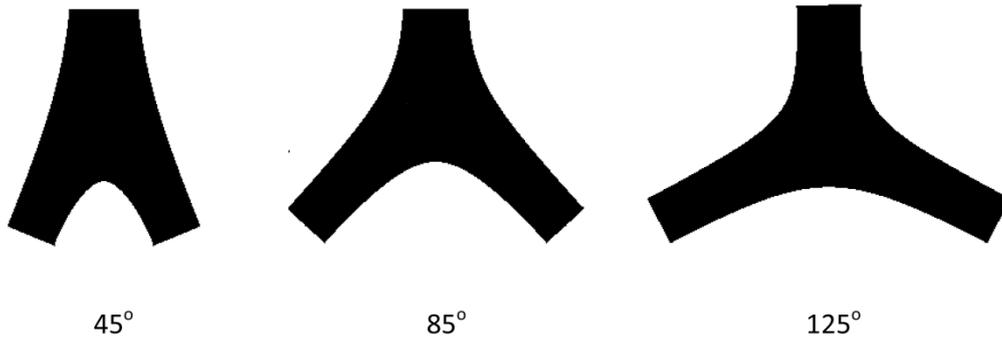
39 3D Printing has made it possible for the first time to manufacture artificial blood vessels and their  
40 networks of any sophisticated geometry and connections. The printed vascular vessels can be used as  
41 grafts to treat inadequate blood flow or alongside organ transplantation (Edelman, 1999). In this case,  
42 polymers such as poly-propylene-glycol are printed layer by layer to construct a graft which will  
43 remain permanently inside the human body. Printed vascular networks can also be used in tissue  
44 engineering for regenerative medicine. In this case, biodegradable polymers such as poly-lactide-acid  
45 are printed to fabricate a scaffold for the endothelial cells to regenerate into blood vessels. Tissue  
46 engineering using a single graft offers potential advantages over conventional autologous or synthetic  
47 grafts (Edelman, 1999; Naito and Rathore, 2011). A major issue in organ tissue engineering is that the  
48 artificial organ may not develop adequate vascularisation for long-term survival (R.A.J. and O.C.,  
49 2002). An artificial micro-vascular network can provide nutrients and soluble growth factors to cells  
50 and tissues as well as acting as a scaffold for culturing endothelial cells for capillaries (Patrick Jr,  
51 2000; Kannan et al., 2005; Kamel et al., 2013). The additive manufacturing technology offers the  
52 complete freedom to design the details of a vascular branch. Currently various research groups have  
53 successfully 3D printed and tested such vascular vessels (Wu et al., 2011; Miller et al., 2012;  
54 Kucukgul et al., 2013; Hoch et al., 2014; Kolesky et al., 2014). However, a general guidance on the  
55 design of the vascular branch is missing. There have been many published works studying the blood  
56 flow and wall shear stress in vascular networks using computational fluid dynamics (CFD). For  
57 example, the effect of branching angles on hemodynamics of bifurcations was investigated in great  
58 detail in (Friedman et al., 1983; Rabinovitz et al., 1987; Friedman and Seed, 1993; Edelman, 1999;  
59 Liu et al., 2015). It was concluded that the branching angle affects the distribution of the flow field in  
60 the branch. The flow will become fully developed downstream. The developed rate is determined by  
61 the Reynolds number. Local recirculation may occur inside a branch where the cells and nutrients  
62 aggregate. The wall shear stress (WSS) is a key hemodynamic indicator that affects the endothelial  
63 cell development (Ravensbergen et al., 1995; Ravensbergen et al., 1997; Kohler et al., 2001; Marshall  
64 et al., 2004). WSS outside a normal range between 1Pa to 7Pa (Papaioannou and Stefanadis, 2004) is  
65 considered harmful to vessel development and may lead to cardiovascular diseases (Kohler et al.,

2001; Caro, 2008; Coppola and Caro, 2008). However, the previous CFD studies were not aimed at providing guidance for the design of 3D printed vascular vessels. Issues related to 3D printing have not been studied. In the current practice, vascular networks are 3D printed without little understanding of their hemodynamics. “Artificial vascularised scaffolds for 3D-tissue regeneration (ArtiVasc 3D)” is a large project funded by the European Union’s Seventh Framework Programme that aims to generate fully vascularised bioartificial skin. The work described in this paper is part of ArtiVasc 3D. In previous work, the authors developed an algorithm of computational geometry for the construction of vascular branch. The algorithm uses three main geometric parameters to control the branch geometry. They are branch angle, volume and maximum curvature at the apex. Three-dimensional models of vascular vessels are generated which can be translated into STL (Stereolithography) data file for 3D Printing (Gibson et al., 2010).

The purpose of this paper is to present general guidance on the design of vascular branches. CFD simulations are performed for artificial branches targeted for veins and micro-vascular vessels respectively. For each type of vascular branch, the effect of branch angle and local curvature on the blood flow behavior and WSS are studied. A laminate CFD model in commercial software, COMSOL, was used. Confidence in the CFD model was achieved by the fact that it reproduced the measured wall shear stresses away from the branch. For the large veins, it is shown that ensuring the smoothness of the junction and keeping a relatively small branch angle, as observed in human body, is very important to avoid high wall shear stress and recirculation. The issue is however less significant for capillaries. The branch angle in capillaries has little effect on the WSS, which can also be explained by the fact that capillaries in nature often have large branch angles. The combination of the CFD analysis and the junction construction algorithm form a complete design method for the vascular vessels.

## 2. Design of vascular branches and their CFD models

We consider artificial vascular vessels targeted for veins and micro-vascular vessels, respectively. Vascular branches are constructed using the algorithm described in (Han et al., 2015). Figure 1 shows the mid-sections of the smooth branches with three different joining angles of  $45^\circ$ ,  $85^\circ$ , and  $125^\circ$ .



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Figure 1. Mid-sections of the smooth branches with joining angles of 45°, 85°, and 125°.

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The branch is made by joining three circular tubes of two different diameters. For the veins, the diameters are 5mm and 3.87mm, respectively (Ravensbergen et al., 1995; Ravensbergen et al., 1997).

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The dimensions of the model were based on a morphological study of 85 human vertebra-basilar specimens. The ratio between the total cross-sectional area of the two inlet tubes and the cross-

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sectional of the outlet tube is equal to 1.2 which is biologically realistic according to (Ravensbergen et al., 1995). For the micro-vascular vessels, the diameters are 0.05mm and 0.0387mm, respectively. The

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shape of the junction profile shown in Figure 1 is taken as a Bezier curve [1]. Bezier curve uses a parameter  $u_d$  to complete the definition of the curve, which is fully explained in [1]. Table 1 provides

102

the values of  $u_d$  for cases studied in this work. The different values of  $u_d$  lead to different maximum curvatures of the junction, which are also provided in Table 1.

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Table 1. Value of controlled variables for different cases

Parameters Cases	Branch angle	$\bar{C}_{max}$	$\bar{V}$	$u_d$
1	125°	0.44	5.4	0.89
2	85°	1.43	5.4	0.84
3	45°	1.34	5.9	0.4
4	45°	3.02	5.56	0.53
5	45°	5.37	5.4	0.6
6	45°	7.5	5.3	0.63
7	45°	14	5.2	0.674

106

107 The curvature is normalised by the diameter of the larger tube. Five different curvatures are used for  
 108 the branches with joining angle of  $45^\circ$  to study the effect of different curvatures. As an extreme  
 109 example of unnatural design, vascular junctions are also constructed using rectangular tubes of two  
 110 different sizes. Bezier curves are not used to smooth the junction. The dimensions of the rectangular  
 111 tube are provided in figure 2.

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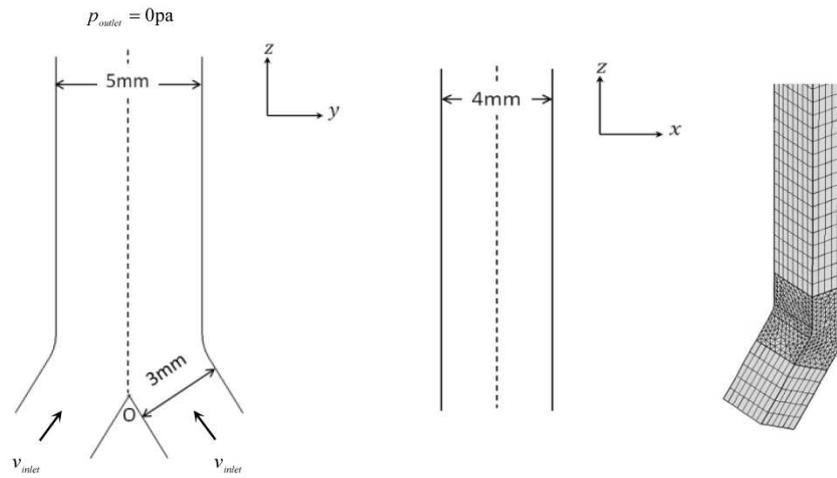
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Figure 2. Sharp branching dimensions and boundary conditions.

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These dimensions are selected to ensure the circular and rectangular tubes have roughly the same cross section area. The junction is connected to two inlet tubes of 70mm and one outlet tube of 200mm (Ravensbergen et al., 1995; Ravensbergen et al., 1997).

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The model makes the following assumptions:

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- The flow is laminar, incompressible and Newtonian;

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- Only steady state flow is considered;

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- The vessel walls are assumed as rigid with no-slip conditions;

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When steady state and incompressible conditions are assumed, the Navier-Stoke equation can be

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written in terms of the velocity  $\mathbf{u}$  and pressure  $p$ :

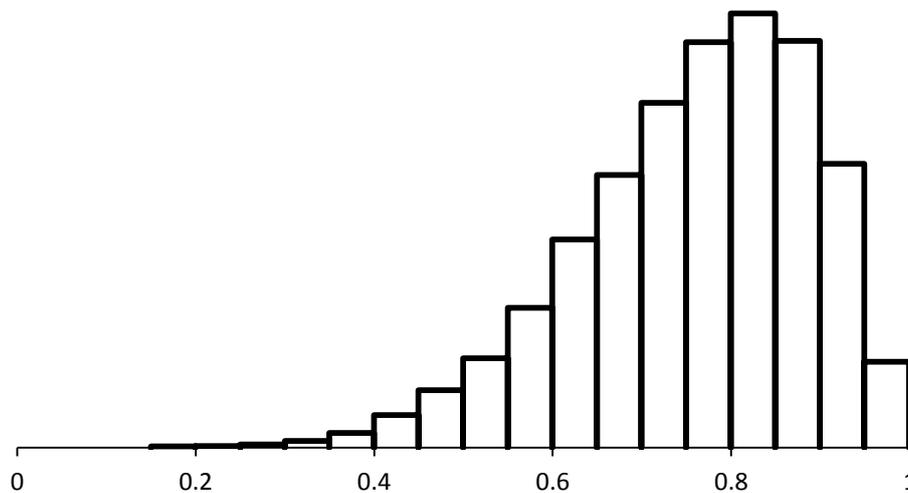
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$$-\varepsilon \Delta \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} + \nabla p = \mathbf{f} \quad \text{in } \Omega \quad (1)$$

131

$$\nabla \cdot \mathbf{u} = 0 \quad \text{in } \Omega \quad (2)$$

132 where  $\varepsilon$  is the viscosity and  $\mathbf{f}$  is the body force per unit mass. In this study,  $\varepsilon$  is set as 0.003Pa.s  
 133 following (Ravensbergen et al., 1995). As the boundary conditions, the inlet velocity at the two  
 134 smaller tubes is set as 0.465m/s (Ravensbergen et al., 1997) while the outlet pressure at the larger tube  
 135 is set as zero (Ravensbergen et al., 1997). Previous studies have shown that the flow behavior in the  
 136 junction is insensitive to the actual value of the outlet pressure because of the large length of the outlet  
 137 tube used in the model. The Reynolds number is 600 for veins and 0.06 for micro-vascular vessels  
 138 both of which are within the normal range of human blood flow (Cliff, 1976). The CFD module of  
 139 commercial software COMSOL<sup>®</sup> was used to solve the equations. A series of convergence studies  
 140 were performed to ensure the finite element meshes are fine enough to produce the correct solutions.  
 141 Two types of statistical studies were also performed here to ensure the convergence of the solutions  
 142 for bifurcation angle  $63^\circ$ .



143

Figure 3. Histogram of mesh element quality distribution

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146 Figure 3 shows a mesh element quality distribution histogram. The x axis in figure 3 is the quality  
 147 indicator for mesh elements. For a best element, the indicator equals 1 and for other meshes, the  
 148 indicator is in the range of (0,1]. Bad elements have smaller values. In statistics, the median value of  
 149 this distribution is 0.55; the mean value of this distribution is 0.76 and the mode value is 0.85. All  
 150 three values are over 0.5 which indicates a good quality of the meshes. Another statistical study is

151 convergence study of the solver we applied to solve CFD simulations. The solver applied is called a  
152 stationary solver which is used to solve steady-state problems. More precisely, Newton nonlinear  
153 method was applied in COMSOL<sup>®</sup>. In this method, the iterative solver iterates until a relative  
154 tolerance is fulfilled. The tolerance was set to be 0.0001 in the solver.



155

156 Figure 4. Relative error versus the iteration number of Newton nonlinear solver

157 Figure 4 indicates the relative error versus the iteration number. It can be seen from figure 4 that after  
158 8 iterations, the relative error was smaller than the set tolerance (0.0001). In other words, the results  
159 were converged. These two statistical studies validated that the computational results are numerically  
160 correct with the same boundary conditions and similar meshes.

161 Ravensbergen et al. (Ravensbergen et al., 1995; Ravensbergen et al., 1997) measured the velocity  
162 distributions for the junction made of rectangular tubes as described above. Figure 5 shows the  
163 comparison between the CFD simulation and their experimental measurement.

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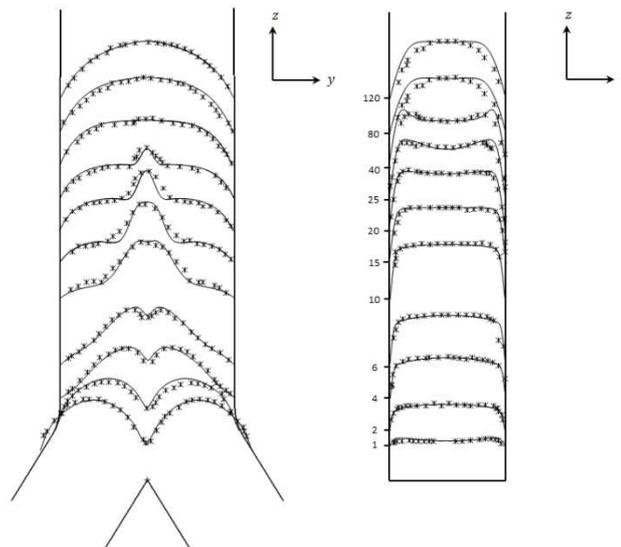
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Figure 5. Simulation results (lines) and experimental data (discrete symbols) (Ravensbergen et al., 1995) comparisons for branching angle  $63^\circ$  in (a)  $yz$  plane and (b)  $xz$  plane.

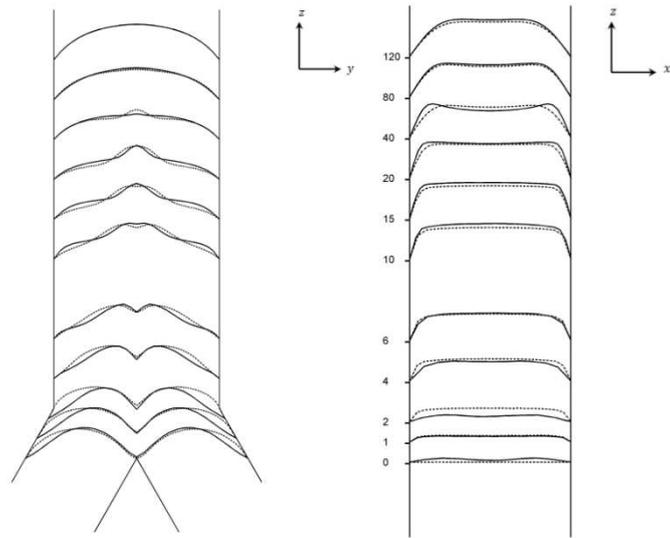
The continuous lines in the figure are simulation results while the discrete symbols are corresponding experimental measurement taken from (Ravensbergen et al., 1995). All the velocities are scaled by the mean axial velocity in the  $z$  direction. Figure 5(a) shows the velocity profiles in the  $y-z$  central plane while Figure 5(b) shows the velocity profiles in the  $x-z$  central plane. The values shown in Figure 5(b) indicate the distances (mm) in the  $z$ -direction measured from the branching point. It can be observed from the figure that the numerical and experimental results agree almost perfectly with each other. The maximum error is around 5%. The CFD model can therefore be accepted as accurate enough for the branch design.

### 3. Hemodynamic analysis for different branch designs

CFD simulations are carried out for branches of different designs outlined at the beginning of section 2. The purpose of the analysis is to compare the different designs in terms of the wall shear stress (WSS) and flow behavior. As explained in the Introduction, WSS is the most important hemodynamic factor when designing a vascular network. Local recirculation should also be minimized to avoid nutrients (for arteries) or waste (for veins) to be trapped in a junction. The CFD results are presented in this section firstly for large vein vessels and then for micro-vascular vessels.

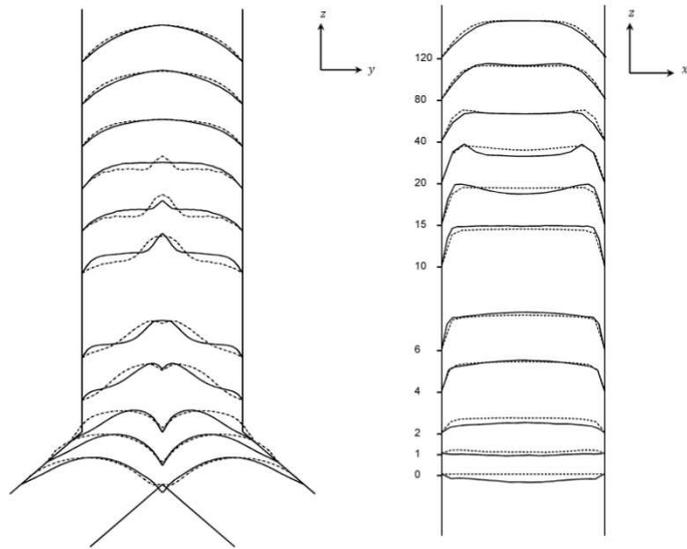
#### 3.1. CFD analysis for large vein branches

CFD simulations were performed for cases 1, 2, and 5. In all cases two flows merge from the daughter vessels into the branch leading to a volume expansion. Using the same branching volume for all the three cases, the effect of branch geometry such as the branching angles can be analyzed. Figure 6 illustrates the simulated velocity profiles for both sharp junctions (interrupted lines) and smoothed junctions (continuous lines) in the  $xz$  and  $yz$  directions, respectively.



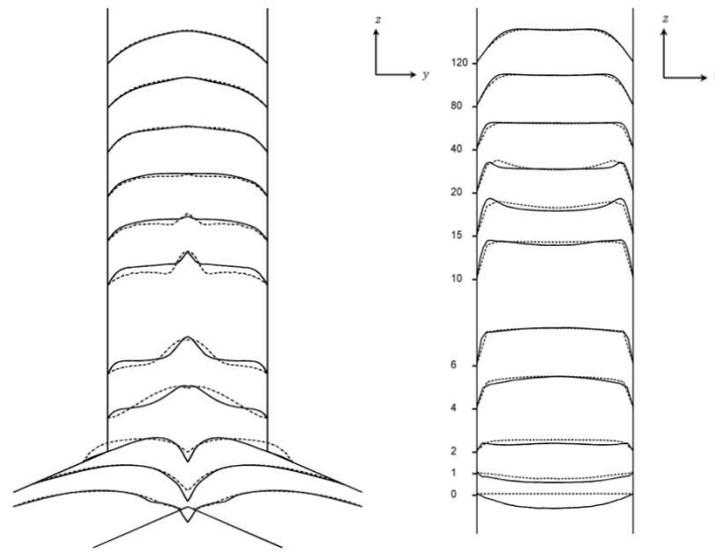
(a)  $45^\circ$ : Case 5

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(b)  $85^\circ$ : Case 2

201



(c) 125°: Case 1

202

203 Figure 6. Comparisons between simulated velocity profiles in sharp branching junction (interrupted  
 204 lines) and smoothed junctions (continuous lines) for angle (a) 45° (b) 85° (c) 125° in both yz plane  
 205 and xz planes.  
 206

207

The branching angles of 45°, 85° and 125° were used in these simulations. The interrupted lines and

208 continuous lines at the bottom of in Figures 6(a-c) in xz and yz directions, compare the velocity

209 profiles at the confluence positions for the sharp junctions and the same positions for the smoothed

210 junctions, respectively. Two humps can be observed in both case 1 and the corresponding sharp

211 junction in the area of influence. In case 1, some negative velocities are observed, the values of which

212 are larger than their sharp counterparts indicating a larger backflow. The two humps begin to merge

213 downstream and start to generate peaks after  $z = 6$  mm for both models. However, the peaks in the

214 smoothed branch are narrower compared to that in the sharp branch. The flow becomes uniform

215 further downstream for both smoothed and sharp branches. The flow profile of the smoothed model

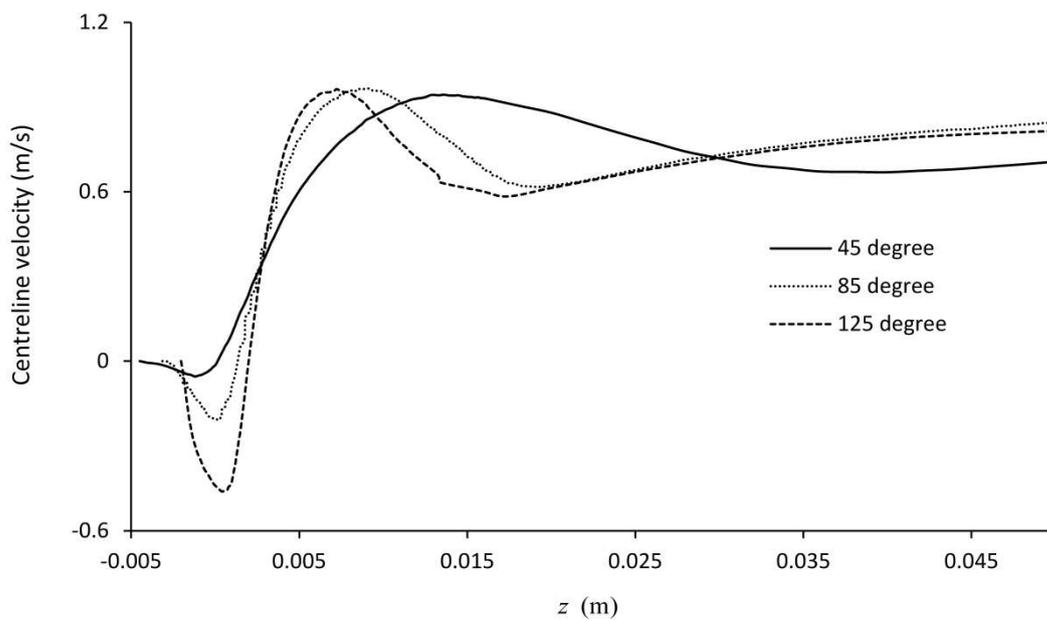
216 tends to fully develop quicker than that of the sharp one while the peak profile disappears earlier. At

217 the end of the outlet, the two flow profiles are nearly identical. Similar trends are observed in cases 2,

218 5 and their corresponding sharp counterparts. Three trends can be concluded by observing Figure 6: 1)

219 the magnitude of the negative flow at the branching point is larger for larger branching angle; 2) the  
 220 velocity profiles become smooth faster in the smoothed junctions compared with the sharp junctions;  
 221 3) the velocity profiles from rounded junctions approach those in the sharp junctions quicker for larger  
 222 branching angles. This also indicates that the branching angle has less influence in the smoothed  
 223 junctions.

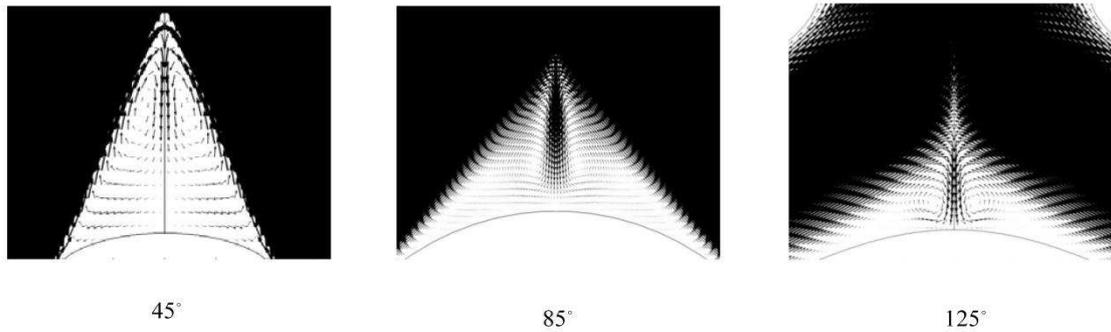
224 To further demonstrate trend one in a two dimensional manner, two dimensional velocity  
 225 distribution around the axial centerline are presented in Figure 7 for branching angles of 45°, 85° and  
 226 125°.



227  
 228 Figure 7. Centreline velocities for branching angles 45° (case 5), 85° (case 2) and 125° (case 1).

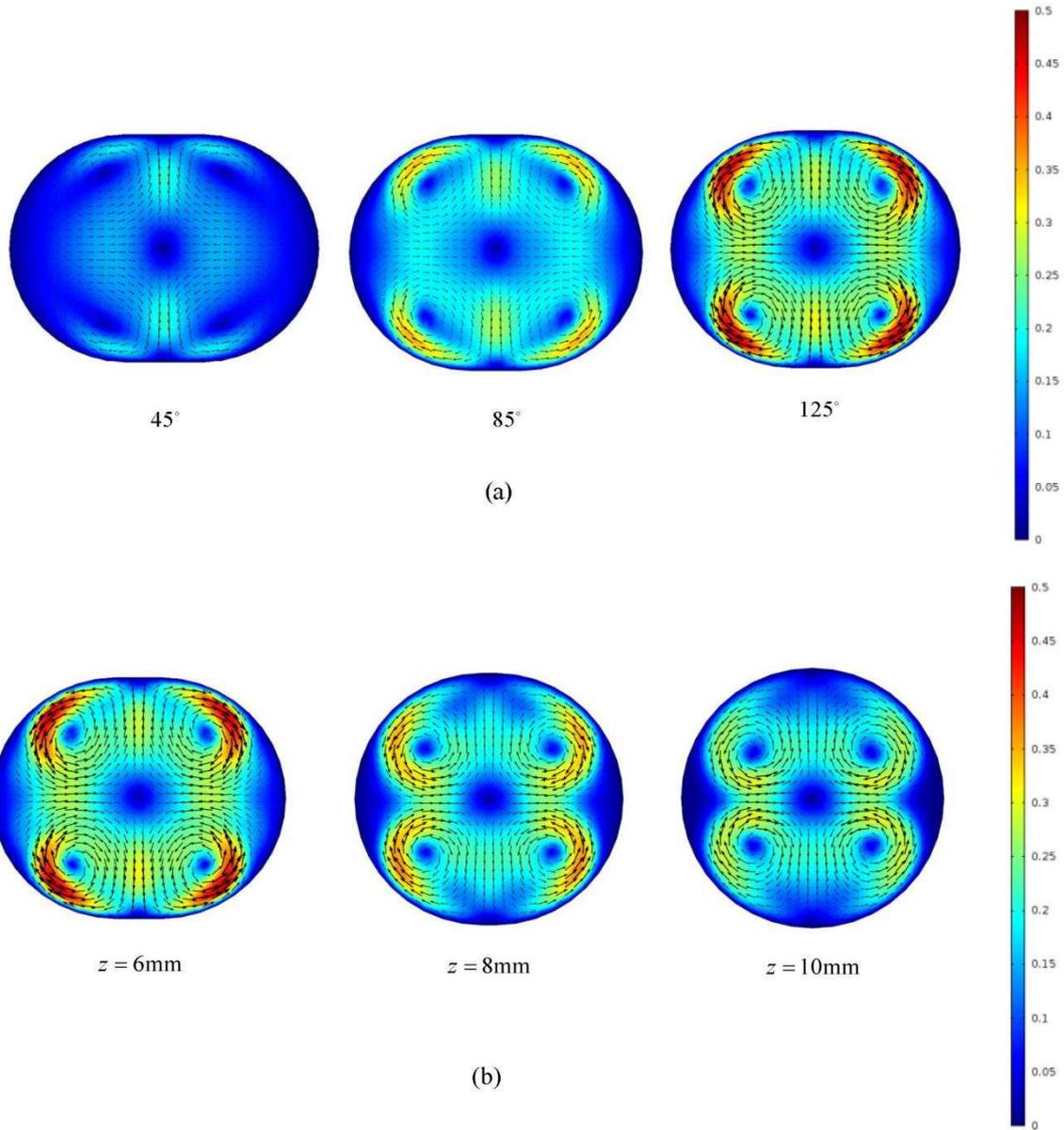
229 Negative velocities can be observed in the branching areas for all the three smoothed junctions. A  
 230 junction with a larger branching angle has a larger value of negative velocity. The larger the branching  
 231 angles, the more it increases in the z-direction. At the range of  $z \approx 0.006$  to  $0.015$  m, all the  
 232 velocities reach their peak values. The junctions with larger branching angles reach their peak values

233 first. The centerline velocities decrease afterwards and then increase gradually further downstream for  
 234 all cases. The negative velocity observed in all the cases indicates backflow. In Figure 8, it can be  
 235 observed that the backflow induces recirculation in the branching area.



237 Figure 8. Recirculation area for branching angles 45° (case 5), 85° (case 2) and 125° (case 1).  
 238

239 Two vortices can be seen in the branching area although its magnitude is small comparing with the  
 240 surrounding velocity field. A region with low velocity vortices is known as a flow recirculation area.  
 241 Nutrients for arteries or waste for veins in the blood flow can be trapped in such area. Therefore, it is  
 242 important to understand how the recirculation area of a rounded junction affects the flow velocity  
 243 profile and the wall shear stress (WSS) downstream. In figure 8, the ratios of recirculation area over  
 244 the whole branch are 1) 45° : 26.4% , 2) 85° : 24.9% and 3) 125° : 21.8%. Junction with 45° branch  
 245 angle has the largest recirculation area while junction with 125° branch angle has the smallest  
 246 recirculation area. Although the junction at 125° has the largest backflow, its recirculation area has the  
 247 least influence on the velocity profile downstream. In order to explain this, the secondary velocity  
 248 distribution for different cases at  $z = 6\text{mm}$  is shown in Figure 9 (a) while those for branching angle of  
 249 125° are shown in Figure 9 (b) at the positions of  $z = 6\text{mm}$  ,  $z = 8\text{mm}$  and  $z = 10\text{mm}$  .



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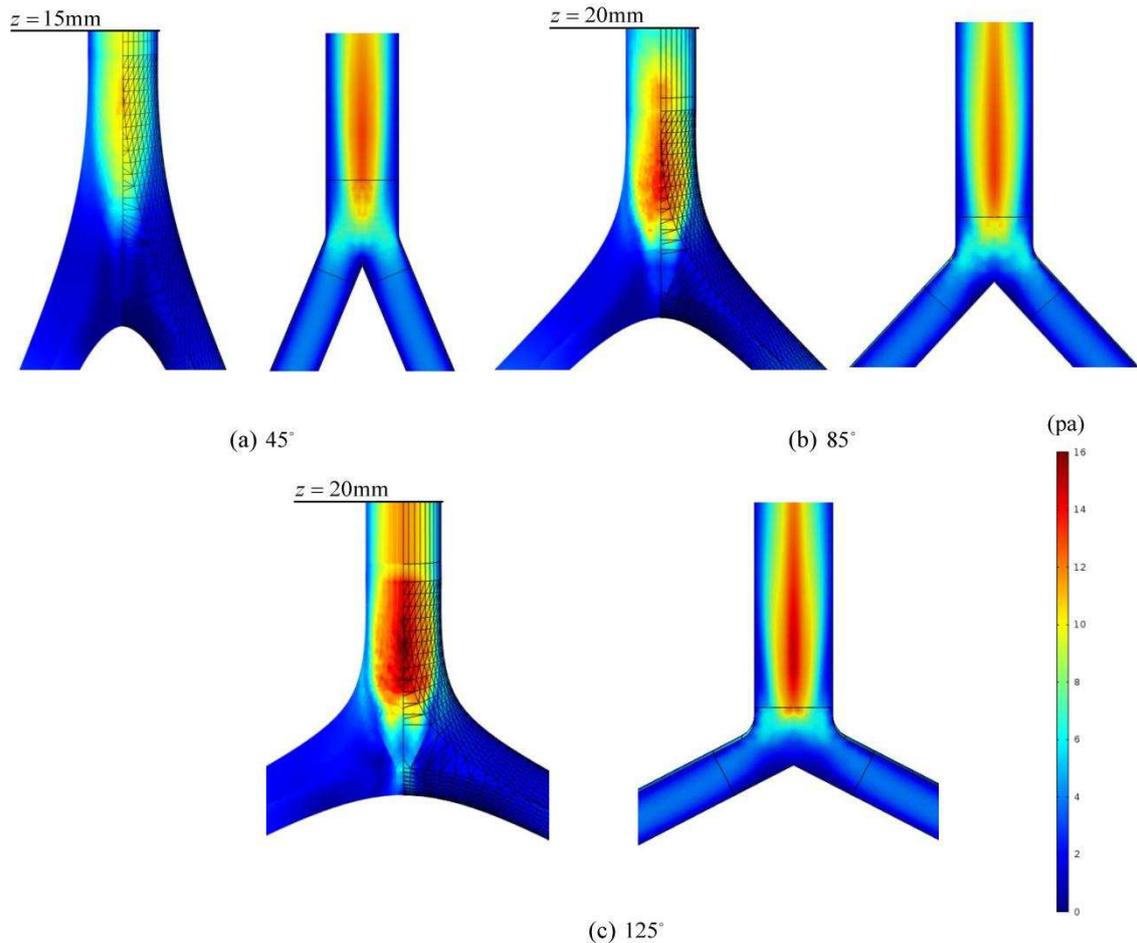
252 Figure 9. Secondary velocities for (a) branching angles  $45^\circ$  ,  $85^\circ$  and  $125^\circ$  at  $z = 6\text{mm}$  and (b)  
 253 branching angle  $125^\circ$  at  $z = 6\text{mm}$  ,  $z = 8\text{mm}$  and  $z = 10\text{mm}$   
 254

255 The vectors and magnitudes of the secondary velocity are presented in Figure 9 using colour plots for  
 256 the three models. Four vortices can be found in each cross section. Figure 9 (a) illustrates a larger  
 257 magnitude of the secondary velocity in the case of larger branching angle. A strong secondary vortex  
 258 helps the two flows from the daughter vessels to mix well, which explain why in Figure 7 the flow  
 259 profiles develop quicker in the cases of larger branching angles. A strong secondary vortex in a

260 smoothed branching junction with a large branching angle diminishes the influence of the backflow.  
261 This explains that even the backflow at the branching point of a larger angle junction is faster; its  
262 influence on the velocity profile downstream is weaker. Hence, the main factor that influences the  
263 downstream velocity profile is the size of the recirculation area rather than the strength of the  
264 backflow. Figure 9 (b) presents secondary velocities for smoothed junctions with a confluence angle  
265 of  $125^\circ$  at different downstream positions of  $z = 6\text{mm}$ ,  $z = 8\text{mm}$  and  $z = 10\text{mm}$ . It is obvious that the  
266 strength of the vortices decreases as  $z$  increases. The vortices disappear once the flow is fully  
267 developed.

268 The WSS is one of the most significant hemodynamic factors that relate to blood vessel  
269 development and cardiovascular diseases (Kohler et al., 2001; Caro, 2008; Coppola and Caro, 2008).  
270 In healthy cerebral arteries, the WSS ranges from 1 Pa to 7 Pa (Papaioannou and Stefanadis, 2004).  
271 WSS higher than 7 Pa can damage the endothelial cells during vascular remodeling while WSS lower  
272 than 1 Pa can lead to the formation of plaque due to insufficient mechanical stimulation on endothelial  
273 cells (Papaioannou and Stefanadis, 2004). WSS distributions for different smoothed cases and their  
274 sharp counterparts are shown in Figure 10.

275



276

277 Figure 10. Wall shear stress distribution for branching angles (a) 45°, (b) 85° and (c) 125° for the  
 278 smoothed model (left) and sharp model (right).  
 279

280 In the junction, the WSS can be many times higher than that in the straight vessel. WSS in the

281 smoothed junction has a different distribution compared with those in the sharp junctions. High WSS

282 (12 Pa) is found on the sharp junction compared to the rounded one (10 Pa) as shown in Figure 10 (a).

283 The area of low WSS in the smoothed junction is larger than that in the sharp one due to recirculation.

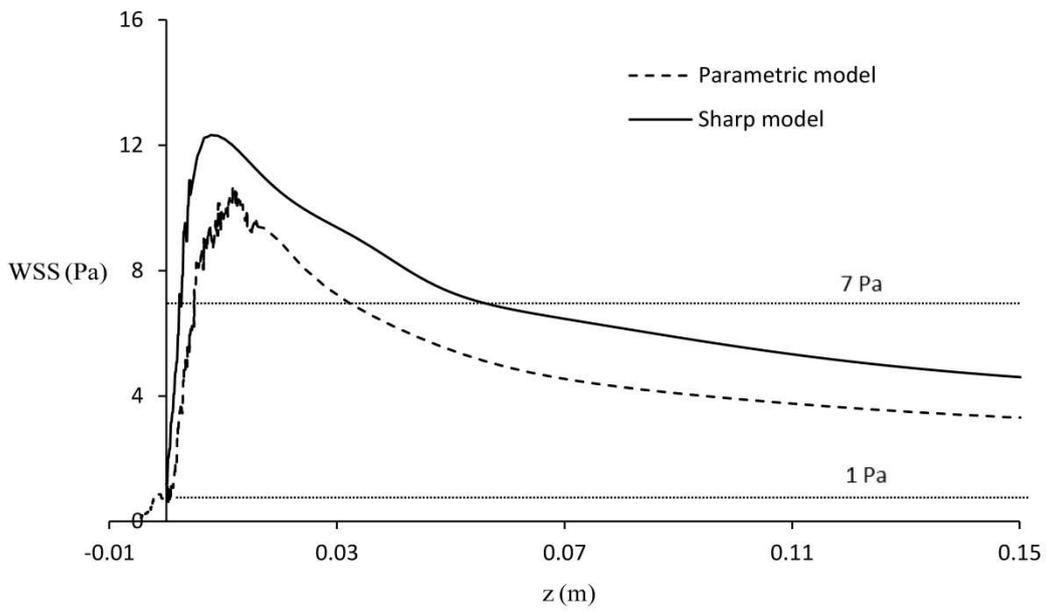
284 Figure 10 (a) shows that at  $z = 15\text{mm}$ , the WSS distribution is more uniform with a low average value

285 of 4 Pa in comparison with 5 Pa in a sharp junction. In Figure 10 (b), similar values of maximum WSS

286 can be observed for both models (~14 Pa). In the smoothed model, the distribution of WSS is more

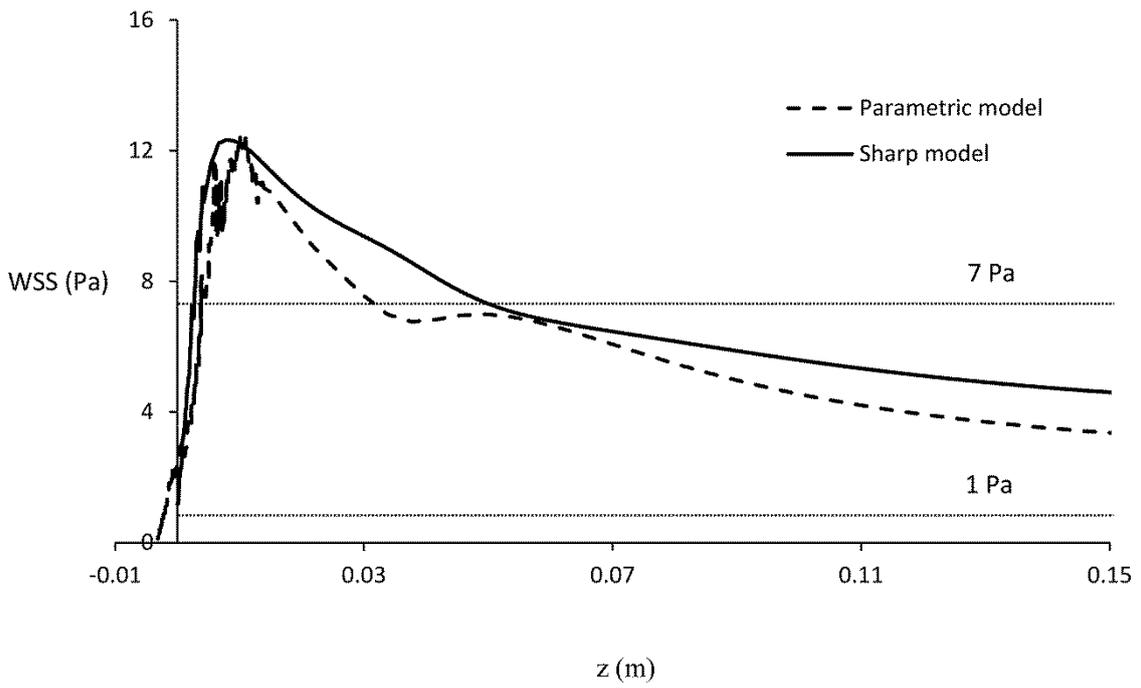
287 intense at the beginning of the downstream flow, but a more uniform distribution of low values is

288 found at  $z = 20\text{mm}$  in comparison with the sharp model. It is found in Figure 10 (c) that the  
289 recirculation area has a similar but weak influence on the WSS distribution in a rounded junction  
290 comparing those observed in a sharp junction for a confluence angle of  $125^\circ$  shown. In order to  
291 quantify how much the recirculation affects the WSS for different branching angles, a two  
292 dimensional plot of WSS on a surface centerline is presented in Figure 11.



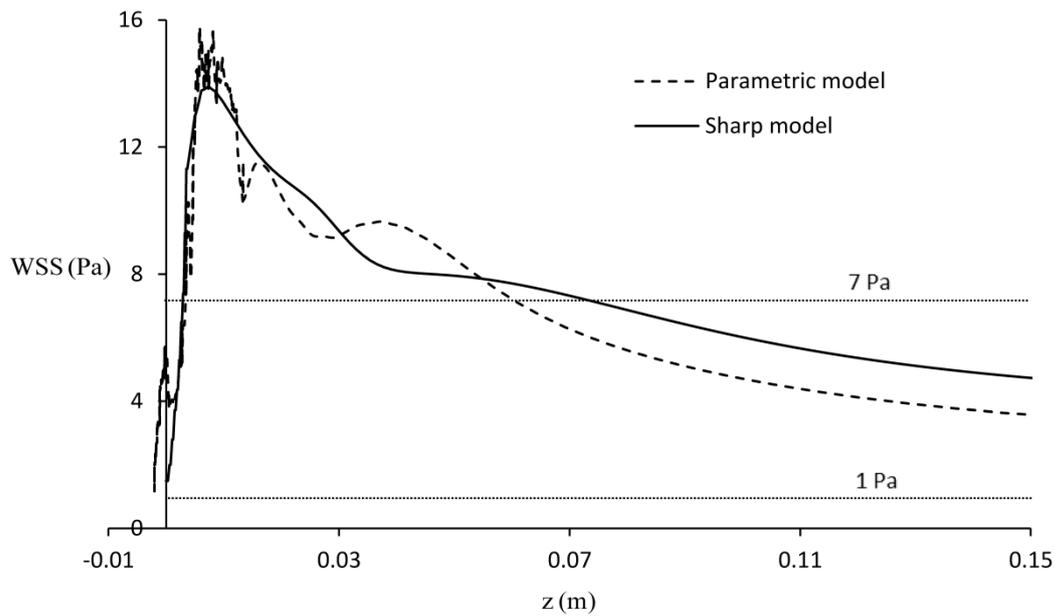
(a) 45°

293



(b) 85°

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(c) 125°

295

296 Figure 11. WSS at the surface center line versus  $z$  of parametric model and sharp model, respectively  
 297 for branching angles (a) 45°, (b) 85° and (c) 125° .

298

299 The dashed lines and solid lines in Figure 11 show the WSS as a function of  $z$  for the smoothed and  
 300 sharp junctions respectively for different branching angles of (a) 45°, (b) 85° and (c) 125° . Here, a

301 ‘healthy window’ is defined as the area along  $z$ -direction within which the WSS is in the range of 1 Pa  
 302 to 7 Pa. The healthy window is an important factor to assess a design of a junction. In the figure, an  
 303 area with healthy WSS is indicated using two dashed lines, representing 1 Pa and 7 Pa, respectively.

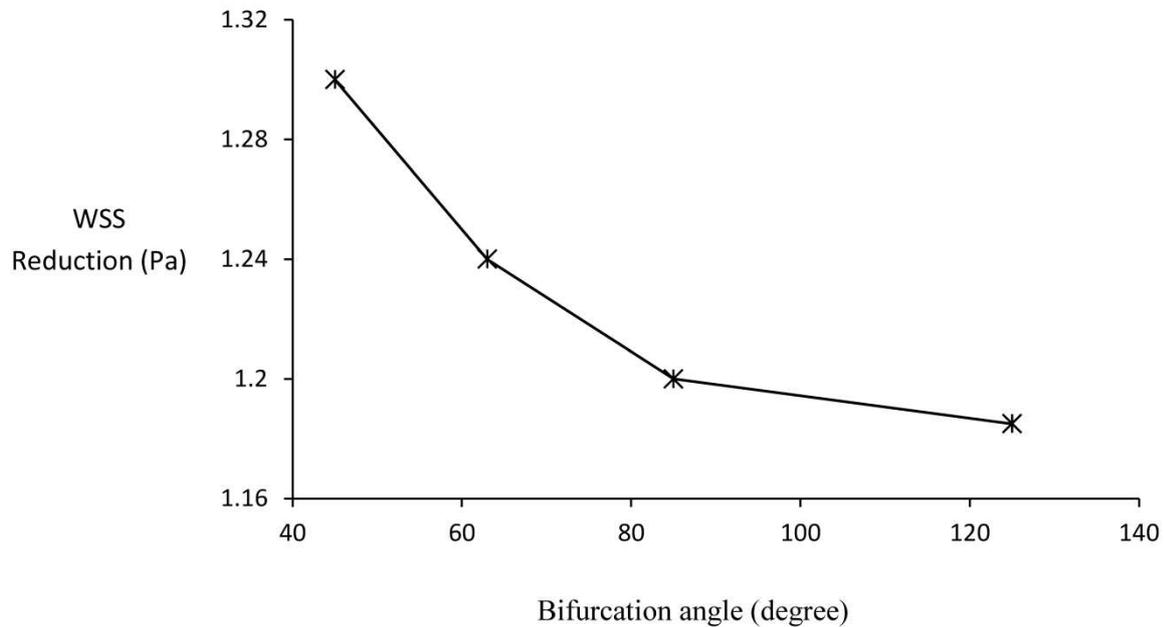
304 In Figure 11 (a), the maximum WSS of the smoothed junction is smaller than that of the sharp one.

305 The WSS values decrease for both junctions downstream. The healthy windows are  $z \in [0\text{m}, 0.005\text{m}]$   
 306 and  $z \geq 0.06\text{m}$  for the sharp junction as seen in Figure 11(a). Those windows expanded to  
 307  $z \in [0\text{m}, 0.01\text{m}]$  and  $z \geq 0.035\text{m}$  for the smoothed junction from as seen in Figure 11(a). The smoothed

308 design of the junction improves the WSS distribution with wider healthy windows. Healthy window

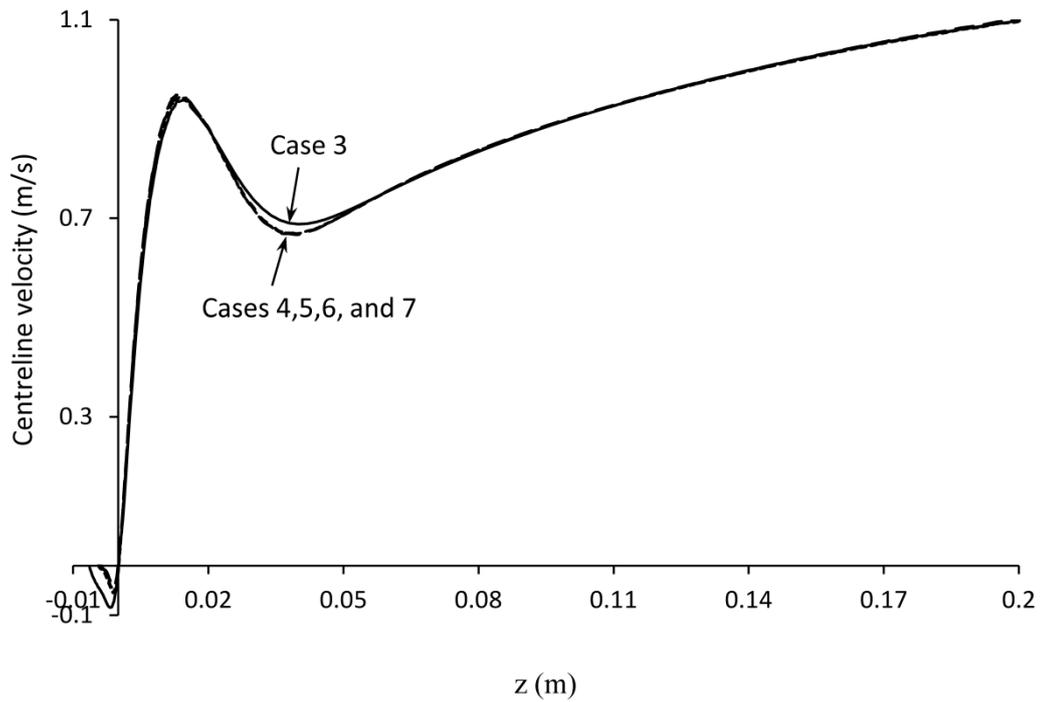
309 expansion can also be observed in Figures 11(b) and 11(c). All the smoothed junctions show a reduced

310 WSS and a more uniform WSS distribution at  $z=150\text{mm}$  downstream compared to the sharp junctions.  
 311 The reduced amount is however different for the three branching angles. A correlation between the  
 312 WSS reduction and the bifurcation angle is shown in figure 12.



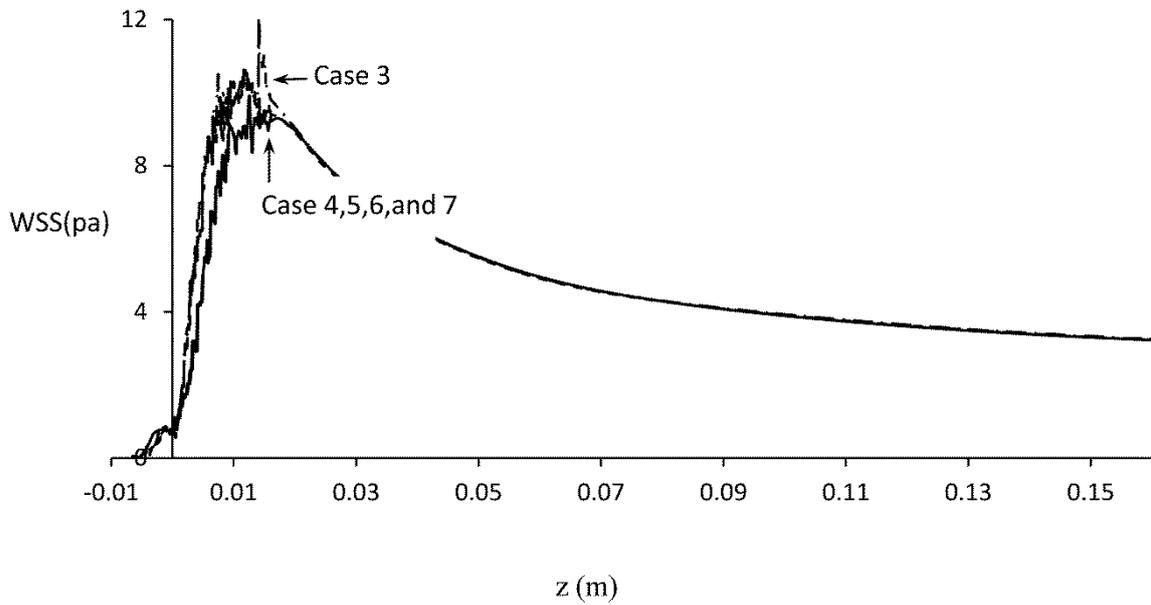
313  
 314 Figure 12. WSS reduction at  $z=0.15$  (m) using parametric model compared with sharp bifurcations  
 315 versus bifurcation angles.  
 316  
 317 From Figure 12, it is seen that junctions with larger bifurcation angle results in a smaller WSS  
 318 reduction. Further increase the bifurcation angle, the WSS reduction decreases more slowly. This  
 319 indicates that the smoothed design has less effect on WSS reduction downstream for bifurcations with  
 320 larger angle.

321 In the parametric design model of a branch junction,  $\bar{C}_{\max}$  and its corresponding  $\bar{V}$  are the most  
 322 important geometric parameters. A larger  $\bar{C}_{\max}$  leads to a smaller  $\bar{V}$ , thus a smaller branching area.  
 323 Further increasing  $\bar{C}_{\max}$ , however, has a limited effect on the branching area as  $\bar{V}$  will decrease more  
 324 slowly. In this section, CFD simulations are presented for smoothed junctions with different  $\bar{C}_{\max}$  for a  
 325 branching angle of  $45^\circ$  (cases 3 - 7). Figure 13(a) shows the central axial velocity as a function of  $z$   
 326 while Figure 13(b) shows the WSS along the centerline as a function of  $z$  coordinate for those cases.



(a)

327



(b)

328

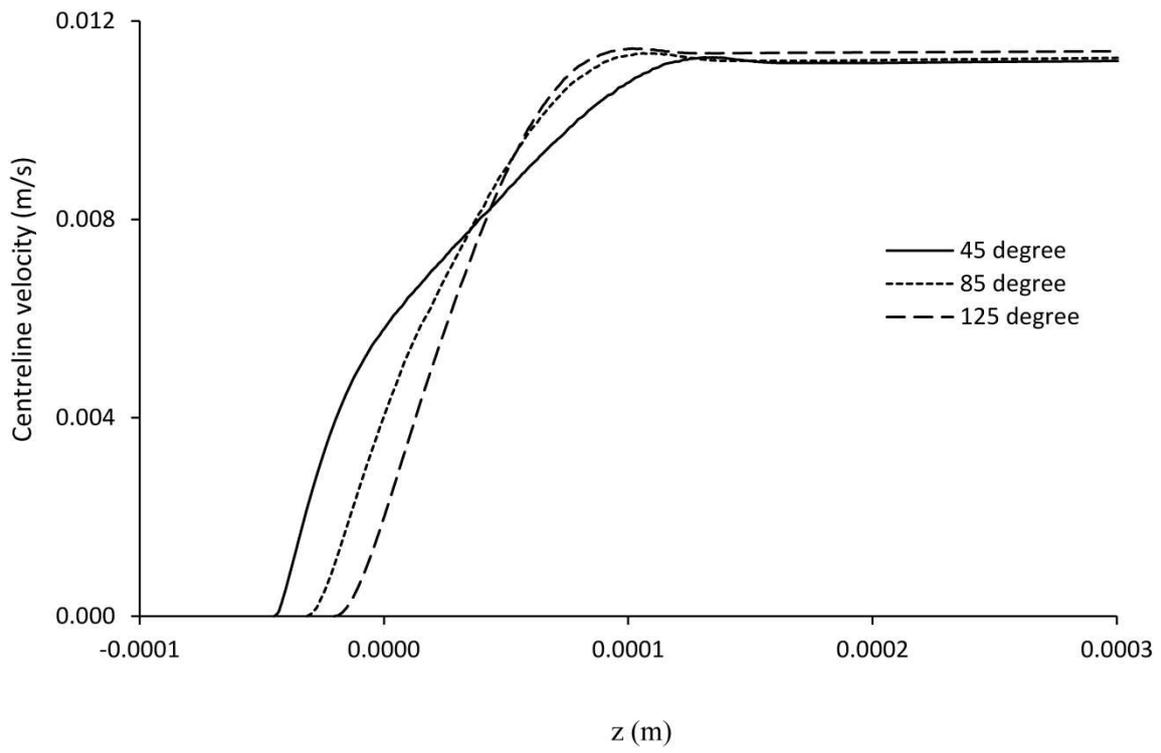
329 Figure 13. (a) Central axial velocities along center line for cases 3-7. (b) WSS at the surface center  
 330 line versus  $z$  for cases 3-7.  
 331

332 In Figure 13(a), cases 4, 5, 6, and 7 show nearly identical velocity distributions. Only case 3 illustrates  
 333 small deviations from the other cases in the area of  $z < 0$  and  $z \in (0.02\text{m}, 0.05\text{m})$ . Figure 13(b) indicates  
 334 that WSS is different for the different junctions with different values of  $\bar{C}_{\max}$  in the area of  $z < 0.02\text{m}$ .

335 Unsurprisingly case3 has the largest deviation. The WSS overlap with each other thereafter for all the  
336 cases. It can be seen from table 2 that case 3 has the largest value of  $\bar{v}$  and smallest value of  $\bar{C}_{\max}$   
337 while cases 4 to 7 have different values of  $\bar{C}_{\max}$  but similar values of  $\bar{v}$ . This is why cases 4-7 have  
338 similar central axial velocities and WSS. Although the velocities and WSS of case 3 show deviations  
339 from the other cases, these deviations disappear downstream. These observations confirm that the  
340 value of  $\bar{C}_{\max}$  and its corresponding  $\bar{v}$  have effect on the distributions of central axial velocities and  
341 WSS. The effects are, however, only in the local areas in contract to the influence from the branching  
342 angle.

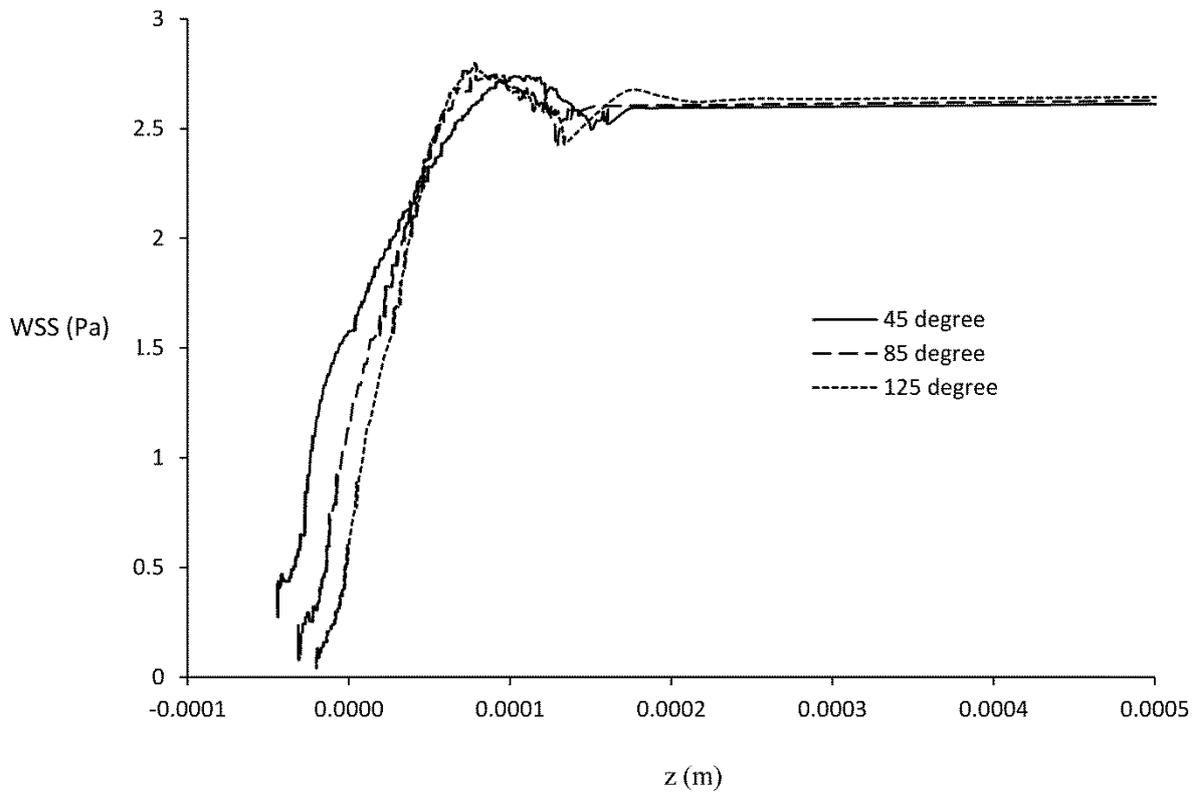
### 343 3.2. CFD analysis for micro-vascular vessels

344 Another important application of the parametric model is the design of micro-vascular networks for  
345 supplying nutrients and oxygen in organ tissue engineering. In these micro-vascular networks, the  
346 vessel diameters are in the range of tens of microns and the blood velocities are low which leads to a  
347 small Reynolds number. In the simulations for the micro-vascular networks, all the diameters and inlet  
348 velocities were scaled down by a factor of 0.01 from those in the previous section. The corresponding  
349 Reynolds number is calculated as 0.06. Figure 14(a) compares the centerline velocities for cases with  
350 branching angles of  $45^\circ$ ,  $85^\circ$  and  $125^\circ$  while Figure 14(b) shows the WSS at the surface centerline as  
351 functions of  $z$  for the same cases.



(a)

352



(b)

353

354 Figure 14. (a) Centerline velocities for branching angles 45°, 85° and 125° and (b) WSS at the surface  
355 center line versus z of parametric model branching angles 45°, 85° and 125°, using a low Reynolds  
356 number.  
357

358 It can be observed from figure 14(a) that the centerline velocities fully develop immediately after its  
359 maximum value regardless of the branching angles. Surprisingly, in the area where z is negative, no  
360 negative velocity can be observed for all the cases which means that backflow and flow recirculation  
361 do not exist. It can also be observed from Figure 14(b) that the WSS are only different in a small area  
362 of  $z < 0.0002$  for all the cases of different branching angles.

#### 363 4. Conclusions

364 In this paper, a set of CFD simulations were carried out to analyse the influence of the different  
365 geometric parameters in the branch design on flow behaviour. The geometric parameters include the  
366 branching angle, maximum curvature at the apex and volume of the branch. Firstly, a large Reynolds  
367 number of 600 was used in the simulations to represent arterial/venous flows. Then a small Reynolds  
368 number of 0.06 was applied to represent micro-vascular flows in a smaller lumen diameter. In the  
369 cases of large Reynolds number, it was found that in addition to manufacturing benefit, the smoothed  
370 junction has two advantages and one drawback. The advantages include: 1) it helps mixing the two  
371 flows better than the sharp junctions; and 2) they have large healthy windows of low WSS. The  
372 drawback is that the recirculation flow may occur in these smoothed junctions. The simulation results  
373 also revealed that different geometric parameters affect the flow behaviour differently: 1) a large  
374 branching angle tends to induce less recirculation flow and WSS reduction in smoothed junctions; and  
375 2) the maximum curvature only affects the flow behaviour locally. For cases with low Reynolds  
376 number, the drawback does not exist because no recirculation was found in such cases. Therefore,  
377 designers are free to choose any  $\bar{C}_{\max}$  for desired branch angles without considering the effect of the  
378 smoothed apex.

#### 379 Acknowledgements

380 This work is part of the project ArtiVasc 3D (<http://www.artivasc.eu/>). It is financially supported by  
381 the European Union's Seventh Framework Programme (FP/2007-2013) under grant agreement No.  
382 263416 (ArtiVasc 3D).

383 **Conflict of interest statement**

384 We have no conflicts of interest to report.

385 **References**

- 386 Caro, C. G., 2008. Discovery of the role of wall shear in atherosclerosis. *Arteriosclerosis, Thrombosis,*  
387 *and Vascular Biology.* 29, 158-161.
- 388 Cliff, W. J., 1976. *Blood vessels.* Cambridge: Cambridge University Press.
- 389 Coppola, G.; Caro, C. G., 2008. Arterial geometry, flow pattern, wall shear and mass transport:  
390 potential physiological significance. *Journal of The Royal Society Interface.* 1-10 p. 2008.
- 391 Edelman, E. R., 1999. *Vascular Tissue Engineering - Designer Arteries.* *Circulation Research.* 85, 1115-  
392 1117.
- 393 Friedman, M. H. et. al., 1983. Arterial geometry affects hemodynamics: a potential risk factor for  
394 atherogenesis. *Atherosclerosis.* 46, 225–231.
- 395 Friedman, M. H.; Seed, W. A., 1993, Hemodynamic implications of a relation between coronary  
396 geometry and pathology: *Atherosclerosis.* 98, 193–199.
- 397 Gibson, I.; Rosen, D. W.; Stucker, B., 2010. *Additive manufacturing technologies: Rapid prototyping*  
398 *to direct digital manufacturing.* Springer 2010.
- 399 Han, X.; Bibb, R.; Harris, R., 2015. Design of bifurcation junctions in artificial vascular vessels  
400 additively manufactured for skin tissue engineering. *Journal of Visual Languages and Computing.* 28,  
401 238-249.
- 402 Hoch, E.; Tovar, G.; Borchers, K., 2014. Bioprinting of artificial blood vessels: current approaches  
403 towards a demanding goal. *European Journal of Cardio - Thoracic Surgery.* 46, 767-778.
- 404 Kamel, R. A. et al., 2013. *Tissue Engineering of Skin.* *Journal of the American College of Surgeons.*  
405 217, 533–555.
- 406 Kannan, R. Y. et al., 2005. The roles of tissue engineering and vascularisation in the development of  
407 micro-vascular networks: a review. *Biomaterials.* 26, 1857-1875.
- 408 Kohler, U. et al., 2001. MRI measurement of wall shear stress vectors in bifurcation models and  
409 comparison with CFD predictions. *Journal of Magnetic resonance imaging.* 14, 563-573.
- 410 Kolesky, D. et al., 2014, 3D Bioprinting of Vascularized, Heterogeneous Cell-Laden Tissue Constructs:  
411 *Adv.Mater.* 26, 3124-3130.
- 412 Kucukgul, C. et al., 2013 3D hybrid bioprinting of macrovascular structures. *Procedia Engineering.* 59,  
413 183-192.
- 414 Liu, G.; Wu, J.; Ghista, D. N. H., Wenhua Kelvin K.L. Wong., 2015. Hemodynamic characterization of  
415 transient blood flow in right coronary arteries with varying curvature and side-branch bifurcation  
416 angles. *Computers in Biology and Medicine.* 64, 117-126.
- 417 Marshall, I. et al., 2004. MRI and CFD studies of pulsatile flow in healthy and stenosed carotid  
418 bifurcation models. *Journal of Biomechanics.* 37, 679-687.
- 419 Miller, J. S., Kelly. et al., 2012. Rapid casting of patterned vascular networks for perfusable  
420 engineered three-dimensional tissues. *Nature Materials.* 11.
- 421 Naito, Y. et al., 2011. *Vascular tissue engineering: towards the next generation vascular grafts.*  
422 *Advanced Drug Delivery Reviews.* 65, 312-323.

- 423 Papaioannou, T. G.; Stefanadis, C., 2004. Vascular wall shear stress: basic principles and methods.  
424 Hellenic Journal of Cardiology. 46, 9-15.
- 425 Patrick, JR, C. W., 2000. Adipose tissue engineering: The future of breast and soft tissue  
426 reconstruction following tumor resection. Seminars in Surgical Oncology. 19: 302-311 p. 2000.
- 427 Rose, F., Oreffo, R., 2002. Bone Tissue Engineering: Hope vs Hype. Biochemical and Biophysical  
428 Research Communications. 292, 1-7.
- 429 Rabinovitz, R. S. Levesque, M. J. Nerem, R. M. 1987. Effects of branching angle in the left main  
430 coronary bifurcation: Circulation. 76: IV-387.
- 431 Ravensbergen, J. et al., 1995. Merging flows in an arterial confluence: the vertebra-basilar junction.  
432 Journal of Fluid Mechnaics. 304: 119-141.
- 433 Ravensbergen, J. et al., 1997. The influence of the blunting of the apex on the flow in a Vertebro-  
434 Basilar junction model. Journal of Biomechanical Engineering. 119, 195-205.
- 435 Wu, W.; Deconinck, A.; Lewis, J., 2011. Omnidirectional Printing of 3D Microvascular Networks: Adv.  
436 Mater. 23, H178-H183.